# Chemical Potential Formalism for Entropic Effects in Confined Polymers

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#### **Experimental motivations**





Both tug-of-war and recoiling processes are caused by entropic forces.

A DNA is pulled into and escapes from a nanoslit. Chia-Fu Chou's lab (Academia Sinica, Taiwan)

### The entropic effect for confined polymers



#### **Experiments for entropic forces**



 $\approx 0.17 \sim 0.51 \text{ pN}$ 

# Entropy-Driven Single Molecule Tug-of-War of DNA at Micro-Nanofluidic Interfaces

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	Outside the channel	Inside the channel
coarse-grained BU	<i>C</i> <sub>0</sub>	CI
number of BU	$n_{0}$	$n_{ m I}$
configuration number of BU	$\Omega_{n_{\mathrm{O}}}$	$\Omega_{n_{\mathrm{I}}}$
excess configuration number caused by adding an extra BU outside/inside the channel	$\Phi_{n_0} = \frac{\Omega_{n_0+1}}{\Omega_{n_0}}$	$\Phi_{n_{\rm I}} = \frac{\Omega_{n_{\rm I}+1}}{\Omega_{n_{\rm I}}}$

# **Entropic recoiling force**

If  $C_{\rm I} = C_0$ 

Entropic free energy

 $-\tilde{T}S = -k_B\tilde{T}\ln(\Omega_{n_I}\Omega_{n_O})$   $k_B = \text{Boltzmann's constant}$  $\tilde{T} = \text{Kelvin temperature}$ 



The entropic recoiling force describes the tendency of a chain to escape from a channel,

$$\tilde{f}_R = \frac{-\tilde{T}\Delta S}{l} = \left(\frac{k_B\tilde{T}}{l}\right) \ln\left(\frac{\Omega_{n_I-1}\Omega_{n_O+1}}{\Omega_{n_I}\Omega_{n_O}}\right),\tag{1}$$

when the chain is shifted a distance / (BU size) out of the channel.

A dimensionless version of force

$$f_R = \frac{\tilde{f}_R l}{\varepsilon} = T \ln\left(\frac{\Omega_{n_I - 1} \Omega_{n_O + 1}}{\Omega_{n_I} \Omega_{n_O}}\right) = T \ln\left(\frac{\Phi_{n_o}}{\Phi_{n_I - 1}}\right)$$
(2)

with  $T = \tilde{T} / \tilde{T}_{room}$  and  $\varepsilon \equiv k_B \tilde{T}_{room}$ .

If  $f'_R$  is the force of a second chain of the same form as (2), then

$$f_R = f_R' + T \ln(\alpha_R) \tag{3}$$

with 
$$\alpha_R \equiv \frac{\Phi'_{n_I-1}\Phi_{n_O}}{\Phi_{n_I-1}\Phi'_{n_O}}$$
.

If 
$$C_{I} \neq C_{0}$$
  
$$f = T \ln \left( \left( \tilde{\Phi}_{n_{0}} \right)^{n_{0}} \right)$$

$$f_R = T \ln \left( \frac{\left( \tilde{\Phi}_{n_I - 1} \right)^{n_i}}{\left( \tilde{\Phi}_{n_I - 1} \right)^{n_i}} \right) \quad (4)$$



$$\begin{split} \widetilde{\Phi}_{n_{I}} &\equiv \frac{g_{n_{I}+1}\Omega_{n_{I}+1}}{g_{n_{I}}\Omega_{n_{I}}} = \Psi_{n_{I}}\Phi_{n_{I}} \\ \widetilde{\Phi}_{n_{O}} &\equiv \frac{g_{n_{O}+1}\Omega_{n_{O}+1}}{g_{n_{O}}\Omega_{n_{O}}} = \Psi_{n_{O}}\Phi_{n_{O}} \end{split}$$

$$\begin{cases} \Psi_{n_0} \equiv \frac{g_{n_0+1}}{g_{n_0}} \\ \Psi_{n_I} \equiv \frac{g_{n_I+1}}{g_{n_I}} \end{cases}$$

Two forces of this form are related by

difference btw  $C_{I}$  and  $C_{0}$ 

$$f_{R} = f_{R}' + T \ln(\tilde{\alpha}_{R}) \qquad \tilde{\alpha}_{R} = G \bar{\alpha}_{R} = \gamma \alpha_{R} \qquad (5)$$

$$G \equiv \frac{g_{n_{I}-1}'^{n_{i}} g_{n_{0}}^{n_{0}}}{g_{n_{I}-1}^{n_{i}} g_{n_{0}}'^{n_{0}}} \qquad \bar{\alpha}_{R} \equiv \frac{\Phi_{n_{I}-1}'^{n_{i}} \Phi_{n_{0}}^{n_{0}}}{\Phi_{n_{I}-1}^{n_{i}} \Phi_{n_{0}}'^{n_{0}}} \qquad \gamma \equiv G \frac{\bar{\alpha}_{R}}{\alpha_{R}}.$$

#### **Entropic drift & tension forces**

The entropic drift force describes the tendency of the drift of the whole chain from the right to the left space,



$$f_D = T \ln\left(\frac{\Omega_{n_R-1}\Omega_{n_L+1}}{\Omega_{n_R}\Omega_{n_L}}\right) = T \ln\left(\frac{\Phi_{n_L}}{\Phi_{n_R-1}}\right).$$
 (6)

The entropic tension force is the tendency of polymer stretching inside the channel,

$$f_T = T \ln \left( \frac{\Omega_{n_L+1} \Omega_{n_I-1} \Omega_{n_R} + \Omega_{n_L} \Omega_{n_I-1} \Omega_{n_R+1}}{\Omega_{n_L} \Omega_{n_I} \Omega_{n_R}} \right)$$
$$= T \ln \left( \frac{\Phi_{n_L} + \Phi_{n_R}}{\Phi_{n_I-1}} \right).$$
(7)

All these forces are related to a second force of the same type by



Upper & bottom bounds of forces (candidates for reference forces)



A correspondence between the entropic force for confined polymers & the chemical potential for solutions.

$$f = f^{\bullet} + T \ln(\alpha), \quad \alpha = \gamma \beta \qquad \gamma = 1$$
granular chain  

$$f = f^{\bullet} + T \ln(\alpha), \quad \alpha = \gamma \beta \qquad \gamma = 1$$
granular chain  
inter-unit factor  $\beta$   
intra-unit factor  $\gamma$   

$$\mu = \mu^{\bullet} + R\tilde{T} \ln(a), \quad a = rc \qquad r = 1$$
dilute solution  

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# Numerical experiments





Two ways to extract entropic forces:

- Jarzynski equality
- Recursion formula (count microstates)

#### Force relation between different types of polymers



#### Force relation btw different channel widths & stiffnesses



#### Entropic drift & tension forces of tug-of-war



**Entropic forces** calculated by the RF coincide with our analytical formulas.

Our study reveals the overestimated traditional scaling approach.

In a, b, c, monomer number outside the channel N = 4, 20, **100**,  $\xi = n_1 / (N - n_1)$ .

100

# Applications

# **Comparison with existing experiments**

experiments	Polymers & channel	f <sub>R</sub> <sup>exp</sup> (pN)	f <sub>R</sub> <sup>theor</sup> (pN)
Α 200 nm C 1 μm 100 hm	A dsDNA recoiling from a nano channel 90 nm wide, 100 nm deep. (Biophysical J.) Modeled by a strip of width 110 nm.	≈ 0.22	≈ 0.18
	A dsDNA recoiling from an array of nanopillars 35 nm in diameter with a center-to-center spacing of 160 nm, free space width 125 nm (Phys. Rev. Lett.) Modeled by a strip of width 125 nm.	≈ 0.04	≈ 0.12
f	A dsDNA undergoing (i) a tug-of-war scenario and (ii) a recoiling-retraction scenario in a nanoslit of 110 nm in height (Nano Lett.) Modeled by a slit of height 110 nm.	≈ 0.30	≈ 0.28

A simple theory provides simultaneous consistency with several experimentally indirectly "measured" forces.

## Predictions

experiments	Polymers & channel	Entropic forces (pN)
alamy stock photo	A ssDNA, partly confined in a 2D nanochannel or a 3D circular nanotube of width 1.6 nm. Modeled by an SC comprising beads of three nucleotide bases.	f <sub>R</sub> <sup>theor</sup> ≈ 11.32 (nanochannel) 21.76 (nanotube)
	A ssDNA oligomer threaded through an $\alpha$ -hemolysin (10 nm long, inner size 1.4 $\sim$ 4.1 nm. Modeled by an NSC through a channel of the same width as the bead, with a hypothesized effective stiffness = $1 \sim 10^2$ pN/nm.	f <sub>T</sub> <sup>theor</sup> ≈ 1.47~14.09

General force scale: fN ~ pN

Upper bound seems to be  $\approx 20 \text{ pN}$ .

#### Summary

Theoretically (mathematical equivalence)

In practice (justify pure entropic contributions in experimental measurements)









pprox 0.04 pN

pprox 0.22 pN

pprox 0.17  $\sim$  0.51 pN

Upper bound of entropic forces in nature (≈20 pN)



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Hong-Qing Xie





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# Chemical potential formalism for polymer entropic forces

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