

随机引力波背景的探测

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ICTP-AP

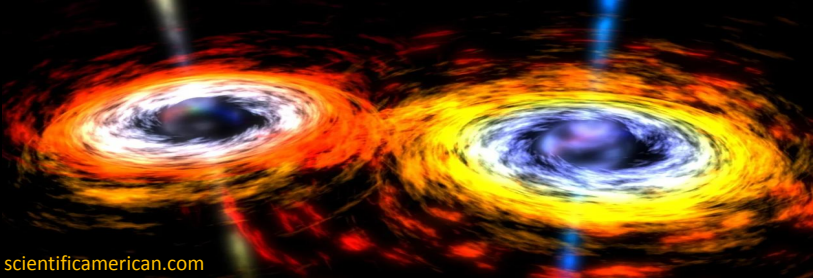
International Centre
for Theoretical Physics Asia-Pacific
国际理论物理中心-亚太地区

第四届新物理数值计算与模拟前沿讲习班

引力波谱



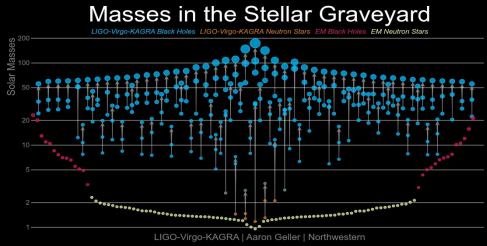
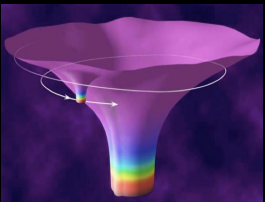
rationalisingtheuniverse.org



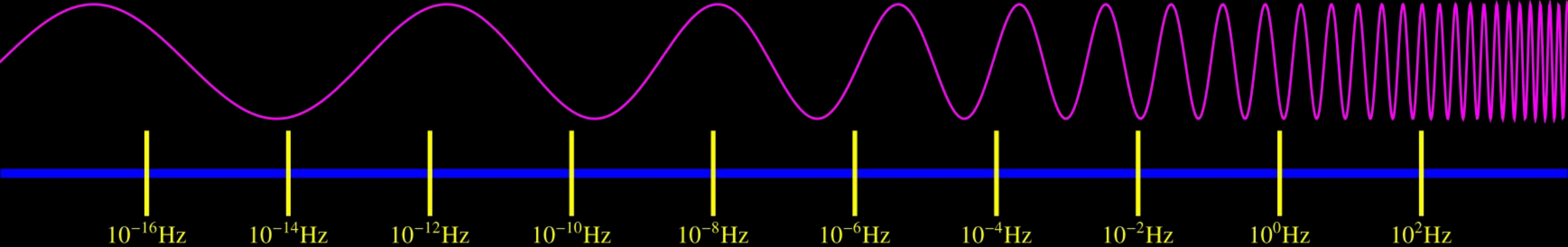
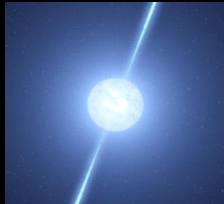
scientificamerican.com



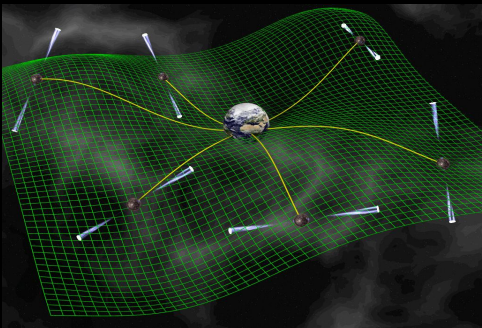
harvard.edu



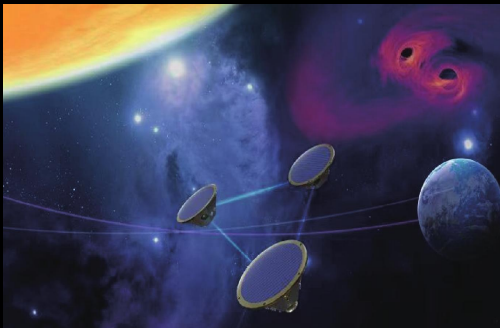
LIGO-Virgo-KAGRA | Aaron Keller | Northwestern



原初引力波探测



脉冲星测时阵列



中国空间引力波探测（太极）



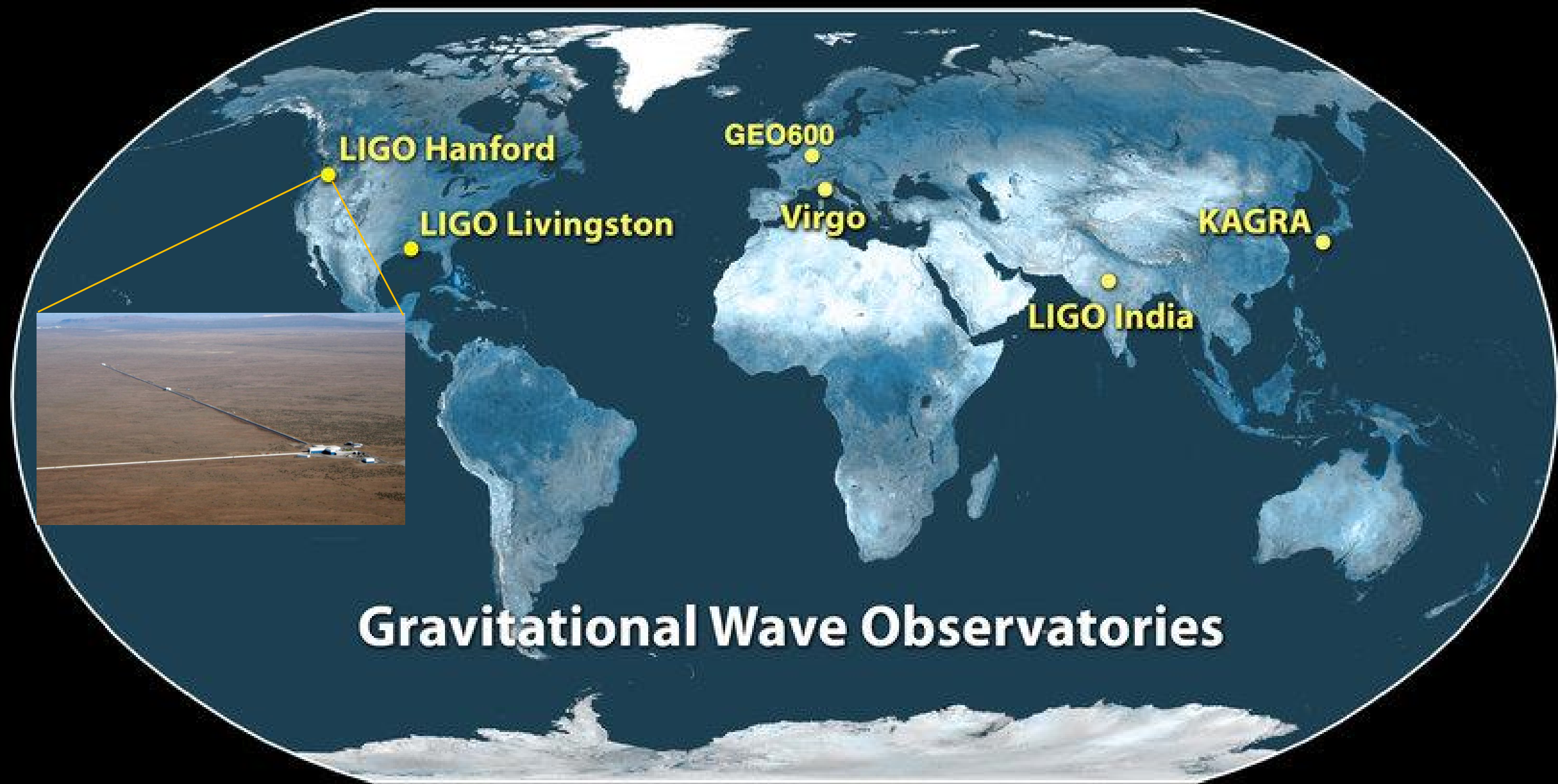
LIGO/Virgo/KAGRA

End-station @ 4 km

Mid-station @ 2 km

4 km (2.5 mi.)





Gravitational Wave Observatories

LIGO Hanford

GEO600

Updated
2025-11-15

01

02

03

04

05

80
Mpc

100
Mpc

100-140
Mpc

150-160+
Mpc

240-325
Mpc

LIGO



30
Mpc

40-60
Mpc

50-60
Mpc

70-130
Mpc

Virgo



0.7
Mpc

1-3
Mpc ≈10
Mpc

25-128
Mpc

KAGRA



G2002127-v33

2016 2017 2018 2019 2020 2021 2022 2023 2024 2025 2026 2027 2028 2029 2030 2031

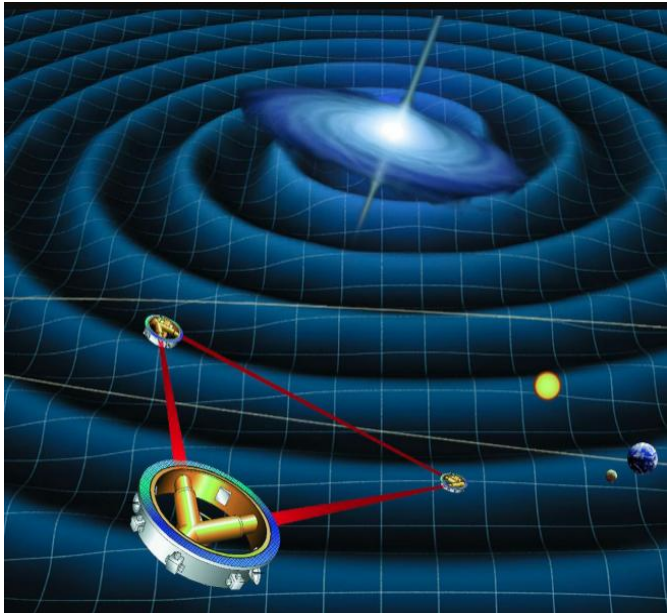
Gravitational Wave Observatories



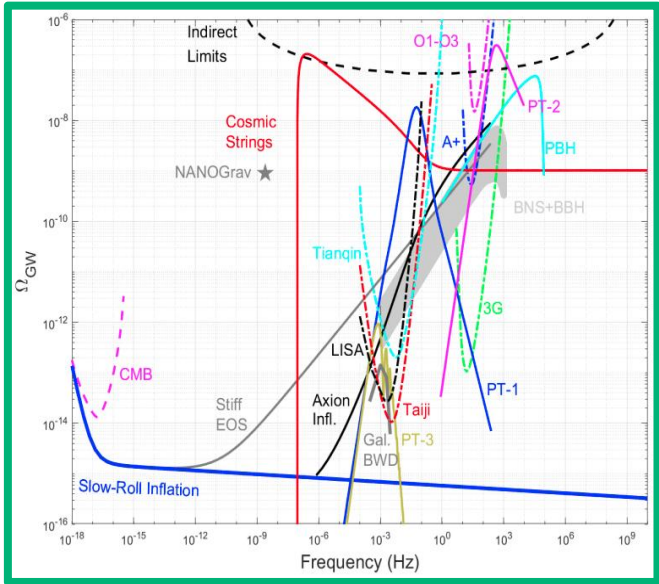
空间引力波探测：太极、天琴、LISA

Example: Phase Transition

Theorist



LIGO, LISA/Taiji/Tianqin, PTA, ...



Gravitational Wave Spectrum

$$\begin{matrix} \alpha \\ \beta \\ v_w \\ T_* \\ g_s \\ \dots \end{matrix}$$

Phase Transition
Parameters

Standard Model of Elementary Particles

three generations of matter (fermions)						interactions / force carriers (bosons)	
I			II			III	
mass $\approx 2.2 \text{ MeV}/c^2$	$\frac{2}{3}$	$\frac{1}{2}$	$\approx 1.28 \text{ GeV}/c^2$	$\frac{2}{3}$	$\frac{1}{2}$	$\approx 173.1 \text{ GeV}/c^2$	$\frac{2}{3}$
u up			c charm			t top	
$\frac{1}{3}$			$\frac{1}{3}$			$\frac{1}{3}$	
$\frac{1}{2}$			$\frac{1}{2}$			$\frac{1}{2}$	
$\approx 4.7 \text{ MeV}/c^2$	$-\frac{2}{3}$	$-\frac{1}{2}$	$\approx 96 \text{ MeV}/c^2$	$-\frac{2}{3}$	$-\frac{1}{2}$	$\approx 4.18 \text{ GeV}/c^2$	$-\frac{2}{3}$
d down			s strange			b bottom	
$\frac{1}{3}$			$\frac{1}{3}$			$\frac{1}{3}$	
$\frac{1}{2}$			$\frac{1}{2}$			$\frac{1}{2}$	
$\approx 0.511 \text{ MeV}/c^2$	-1	$\frac{1}{2}$	$\approx 105.66 \text{ MeV}/c^2$	-1	$\frac{1}{2}$	$\approx 1.7768 \text{ GeV}/c^2$	-1
e electron			μ muon			τ tau	
$\frac{1}{2}$			$\frac{1}{2}$			$\frac{1}{2}$	
$< 1.0 \text{ eV}/c^2$	0	$\frac{1}{2}$	$< 0.17 \text{ MeV}/c^2$	0	$\frac{1}{2}$	$< 18.2 \text{ MeV}/c^2$	0
ν_e electron neutrino			ν_μ muon neutrino			ν_τ tau neutrino	
$\frac{1}{2}$			$\frac{1}{2}$			$\frac{1}{2}$	
$\approx 91.19 \text{ GeV}/c^2$	0	1	$\approx 80.39 \text{ GeV}/c^2$	± 1	1		
Z Z boson			W W boson				
$\frac{1}{2}$			$\frac{1}{2}$				

BSM

SCALAR BOSONS

GAUGE BOSONS
VECTOR BOSONS

Particle Physics Model

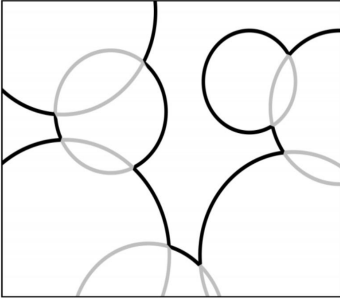
Experimentalist

Questions to Answer

- Set limits when signal is absent
- Parameter estimation when signal is discovered
 - What is precise shape of the signal spectrum?
 - What are the values of α , β , v_w , T^* , etc (PT as an example)?
 - What is the underlying particle physics model?
 - ❖ What are the values of the model parameters?
 - ❖ What this implies for collider experiments?
 - ...

Key Observables

bubble collision

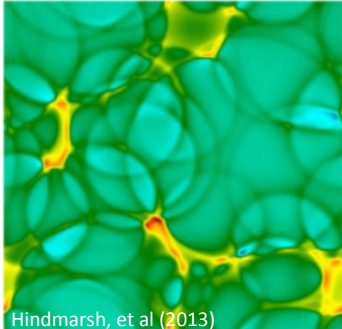


$$\Omega_{\text{coll}}(f)h^2 = 1.67 \times 10^{-5} \Delta \left(\frac{H_{\text{pt}}}{\beta} \right)^2 \left(\frac{\kappa_{\phi} \alpha}{1 + \alpha} \right)^2 \times \left(\frac{100}{g_*} \right)^{1/3} S_{\text{env}}(f),$$

Energy density Spectrum

$$\Omega_{\text{GW}}(f) = \frac{d\rho_{\text{GW}}}{\rho_c d \log f}$$

sound waves



Hindmarsh, et al (2013)

$$\Omega_{\text{sw}}(f)h^2 = 2.65 \times 10^{-6} \left(\frac{H_{\text{pt}}}{\beta} \right) \left(\frac{\kappa_{\text{sw}} \alpha}{1 + \alpha} \right)^2 \left(\frac{100}{g_*} \right)^{1/3} \times v_w \left(\frac{f}{f_{\text{sw}}} \right)^3 \left(\frac{7}{4 + 3(f/f_{\text{sw}})^2} \right)^{7/2} \Upsilon(\tau_{\text{sw}}),$$

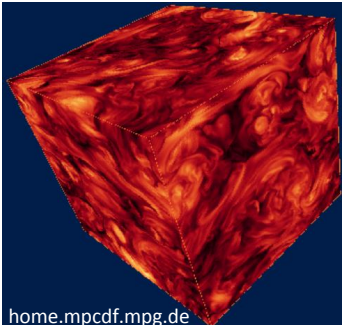
$$\Upsilon = 1 - (1 + 2\tau_{\text{sw}} H_{\text{pt}})^{-1/2} \quad (\text{RD})$$

HG, Sinha, Vagie, White, JCAP [2007.08537]

$$\Upsilon = \frac{2[1 - y^{3(w-1)/2}]}{3(1-w)}$$

Yang Xiao, HG, Jia-Hang Hu,
Jin Min Yang, Yang Zhang [2410.23666]

MHD



home.mpcdf.mpg.de

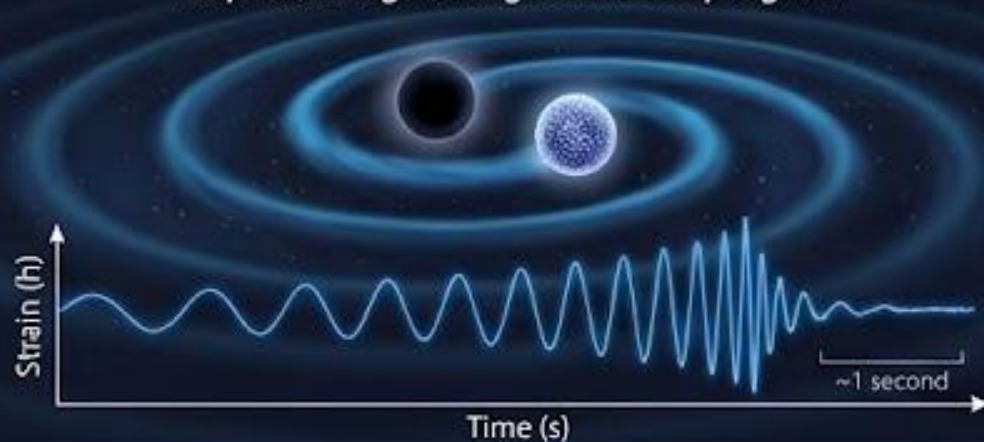
$$h^2 \Omega_{\text{turb}}(f) = 3.35 \times 10^{-4} \left(\frac{H_*}{\beta} \right) \left(\frac{\kappa_{\text{turb}} \alpha}{1 + \alpha} \right)^{\frac{3}{2}} \left(\frac{100}{g_*} \right)^{1/3} v_w S_{\text{turb}}(f)$$

Chiara Caprini et al JCAP [1512.06239]

FOUR TYPES OF GRAVITATIONAL WAVES DETECTED BY LIGO

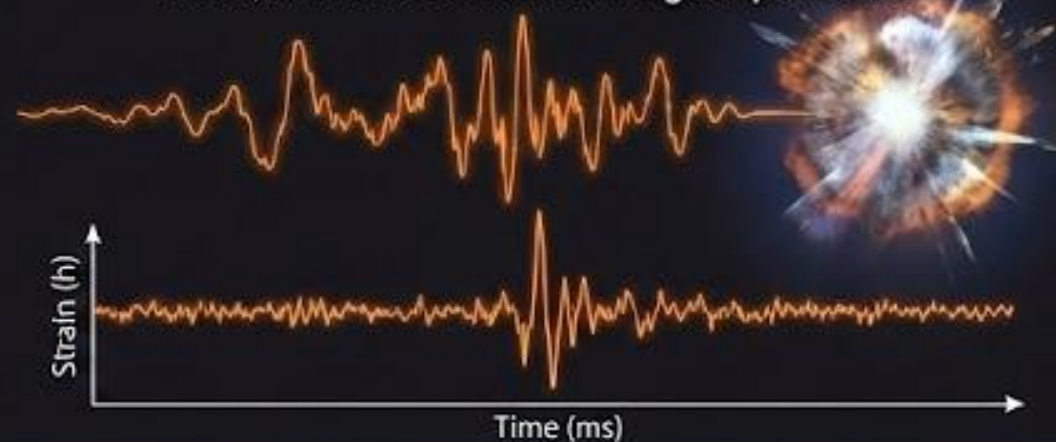
COMPACT BINARY COALESCENCE (CBC)

Inspiral, Merger, Ringdown. Chirp signal.



BURST

Short, unmodeled transient. E.g., Supernovae.



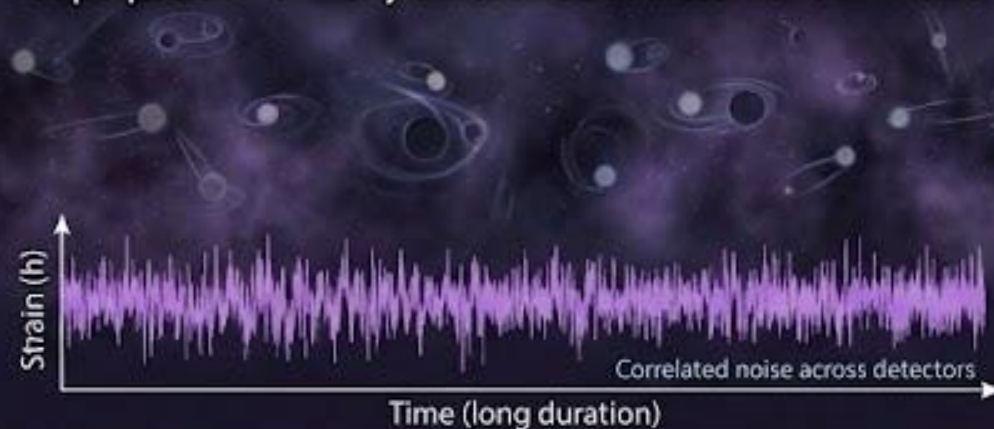
CONTINUOUS WAVES

Constant frequency & amplitude. Spinning neutron stars.



STOCHASTIC BACKGROUND

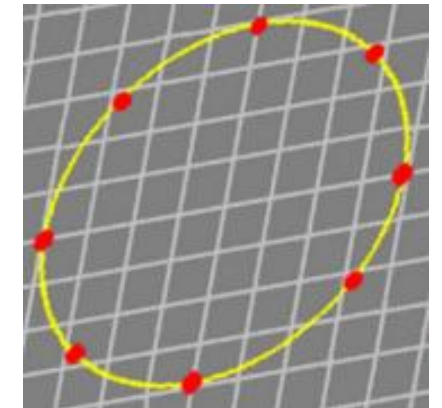
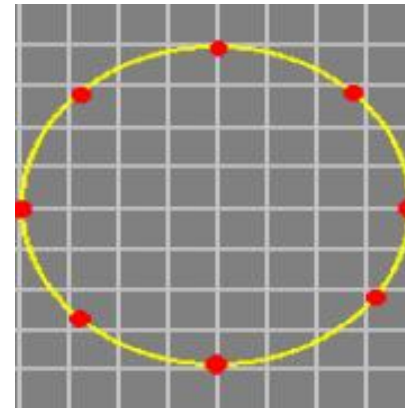
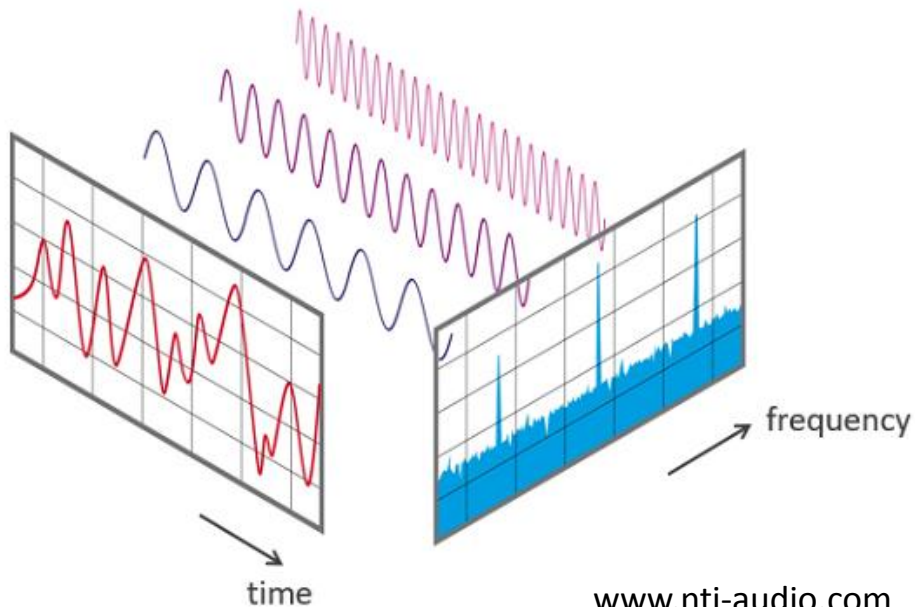
Superposition of many unresolved sources. Random noise.



Plane Wave

- A special gauge, the TT gauge, can be chosen in vacuum (caution: extended sources).
- Only 2 degrees of freedom (+, x).

$$h_{ij}^{\text{TT}}(t, z) = \begin{pmatrix} h_+ & h_\times & 0 \\ h_\times & -h_+ & 0 \\ 0 & 0 & 0 \end{pmatrix}_{ij} \cos[\omega(t - z/c)]$$



Detector Response

- From GW to detector outputs

$$h_{ij}(t, \mathbf{x}) = \sum_{A=+, \times} \int_{-\infty}^{\infty} df \int_{S^2} d^2 \hat{\mathbf{n}} \tilde{h}_A(f, \hat{\mathbf{n}}) e_{ij}^A(\hat{\mathbf{n}}) e^{-2\pi i f(t - \hat{\mathbf{n}} \cdot \mathbf{x} / c)}$$



$$h_I(t) = \sum_A \int df \int d^2 \hat{\mathbf{n}} \tilde{h}_A(f, \hat{\mathbf{n}}) F_I^A(\hat{\mathbf{n}}) e^{-2\pi i f(t - \hat{\mathbf{n}} \cdot \mathbf{x}_I / c)}$$



$$s_I(t) = h_I(t) + n_I(t)$$

different detectors/channels

- Detectors: LIGO (Hanford, Livingston), Virgo, KAGRA, LIGO-India...
- Space-based detectors: Taiji, Tianqin, LISA; TDI channels (A, E, T)

$$h_I(t) = d_I^{ij} h_{ij}(t, \mathbf{x}_I)$$



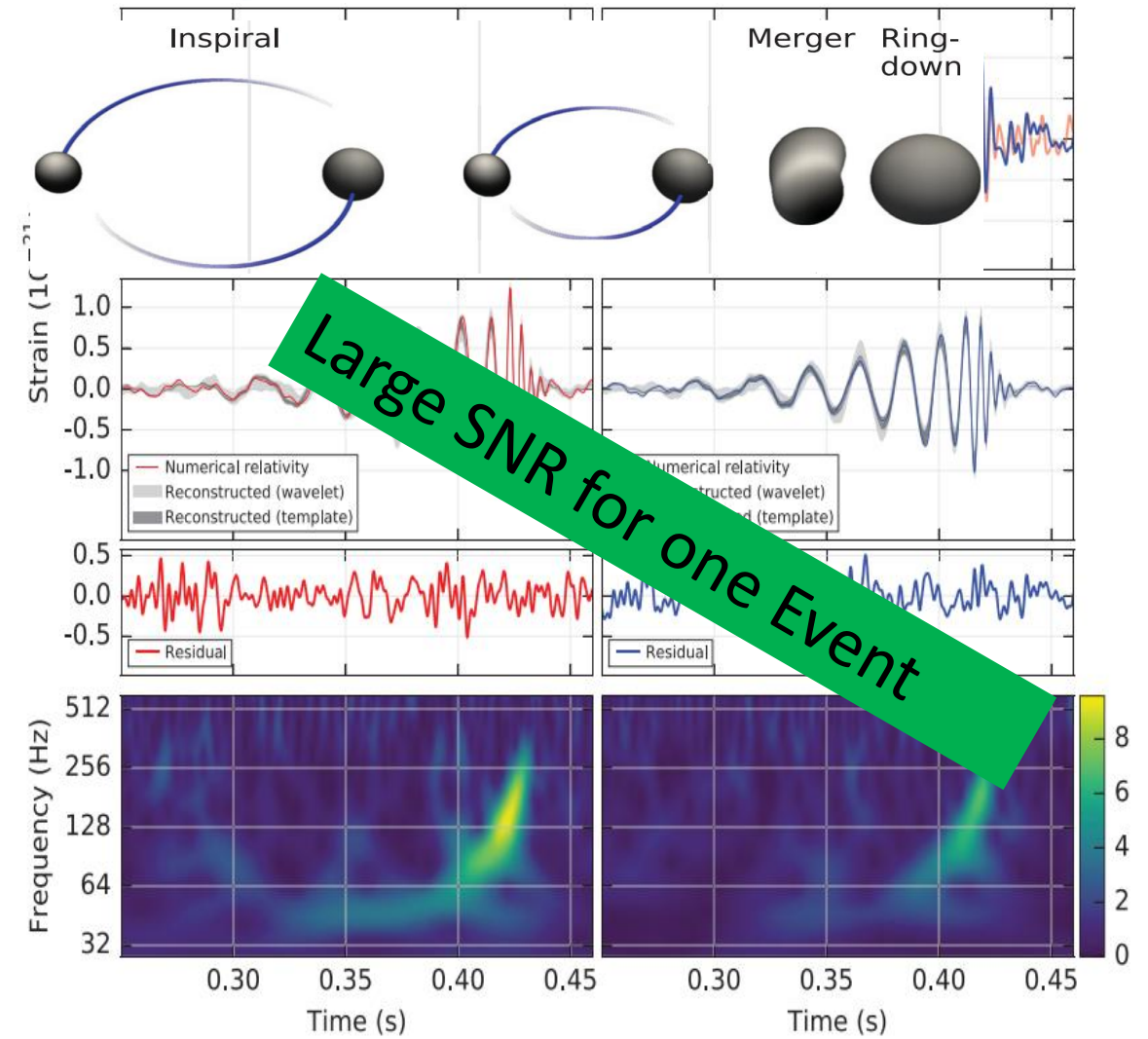
$$F_I^A(\hat{\mathbf{n}}) = d_I^{ij} e_{ij}^A(\hat{\mathbf{n}})$$

$$\begin{aligned} \langle n_I^2(t) \rangle &= \int_0^\infty df P_I(f) \\ \langle \tilde{n}_I^*(f) \tilde{n}_J(f') \rangle &= \frac{\tau}{2} \delta(f - f') \delta_{IJ} P_I(f) \end{aligned}$$

Stochastic GW Background

SGB

$\neq$



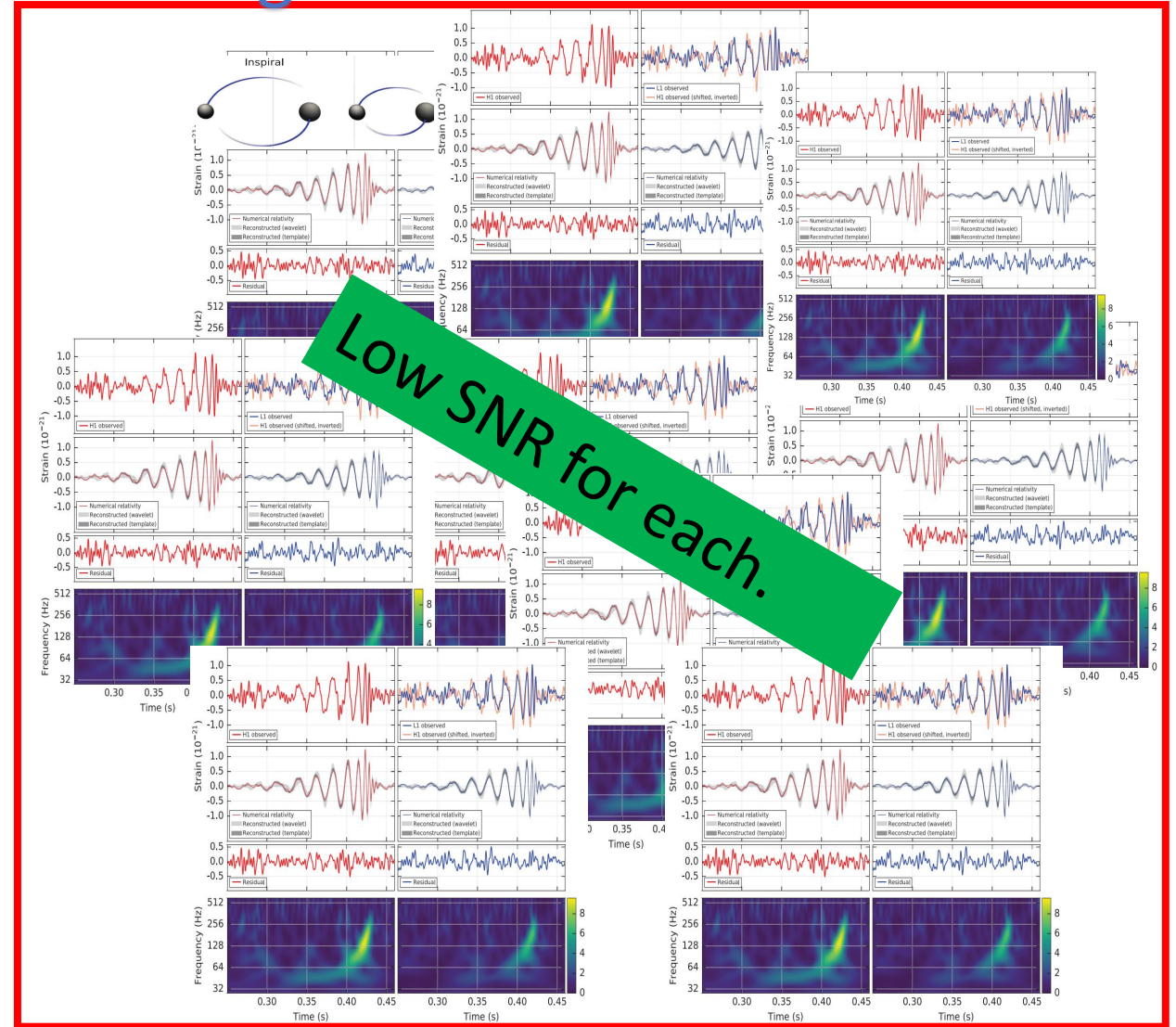
Stochastic GW Background

SGB

~

Similar to a noise.

But how to detect a noise?



LIGO, Phys. Rev. Lett. 116, 061102

Stochastic GW Background

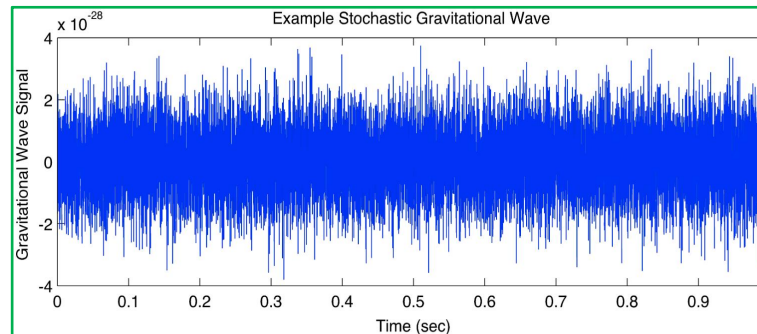
- Assumptions (applies for cosmic sources)

$$\langle \tilde{h}_A^*(f, \hat{\mathbf{n}}) \tilde{h}_{A'}(f', \hat{\mathbf{n}}') \rangle = \delta(f - f') \frac{\delta^2(\hat{\mathbf{n}}, \hat{\mathbf{n}}')}{4\pi} \delta_{AA'} \frac{1}{2} S_h(f)$$

- Stationary
- Gaussian
- Isotropic
- Unpolarized

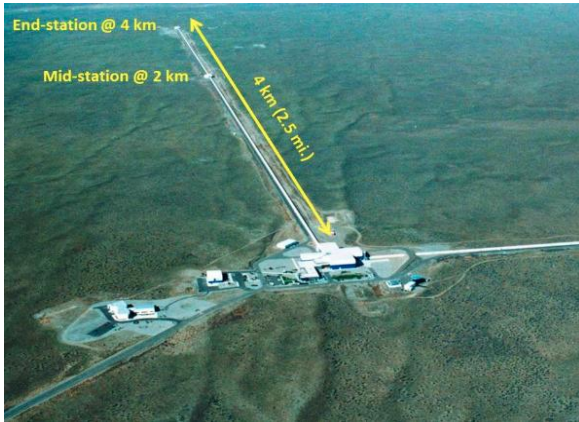
$$S_h(f) = \frac{3H_0^2}{4\pi^2} \frac{\Omega_{\text{GW}}(f)}{f^3}$$

$$\Omega_{\text{GW}}(f) \equiv \frac{f}{\rho_c} \frac{d\rho_{\text{GW}}}{df}$$



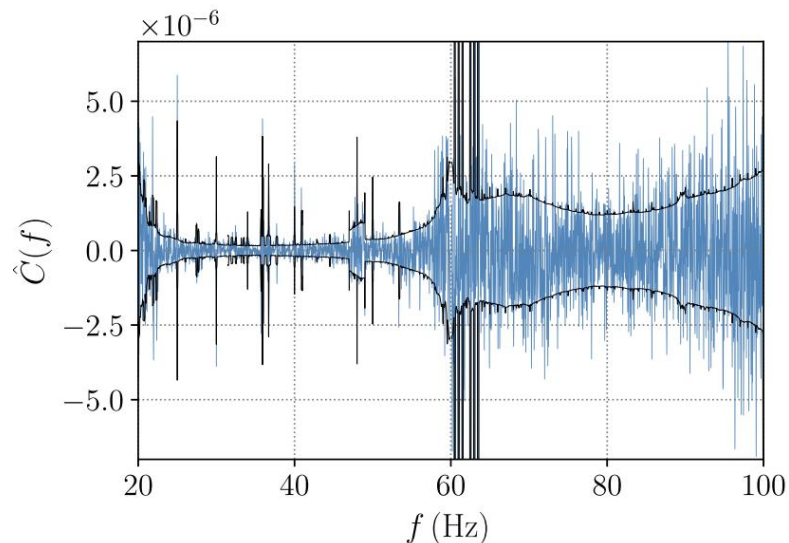
Question: how to simulate GW?

LIGO-Virgo-KAGRA



$$\hat{C}^{IJ}(f) = \frac{2}{T} \frac{\text{Re}[\tilde{s}_I^*(f)\tilde{s}_J(f)]}{\gamma_{IJ}(f)S_0(f)}$$

$$\langle \hat{C}^{IJ}(f) \rangle = \Omega_{\text{GW}}(f)$$



- Glitches and gating
- Narrow spectral artifacts
- Non-Gaussian noise
- Magnetic noise
- Calibration uncertainty

LVK collaborations, PRD [2101.12130]

PT searches:

Cosmological and HEP implications: [2510.26848]
Romero, ..., HG, ..., PRL [2102.01714]

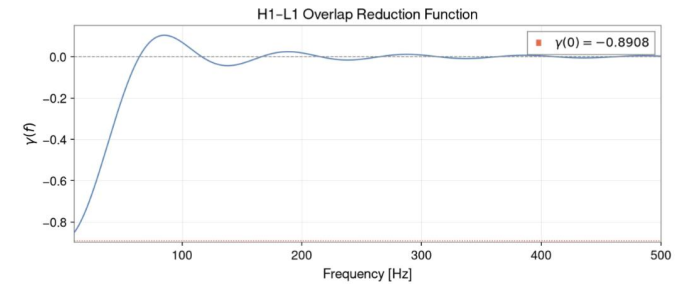
Cross-Correlation

- Single detector/channel output

$$\tilde{h}_I(f) = \sum_A \int_{S^2} d^2 \hat{\mathbf{n}} \tilde{h}_A(f, \hat{\mathbf{n}}) F_I^A(\hat{\mathbf{n}}) e^{2\pi i f \hat{\mathbf{n}} \cdot \mathbf{x}_I / c}$$

- cross correlation of outputs

$$\langle \tilde{h}_I^*(f) \tilde{h}_J(f') \rangle = \frac{1}{2} \delta(f - f') \Gamma_{IJ}(f) S_h(f)$$



overlap reduction function

$$\gamma_{IJ}(f) = \frac{\Gamma_{IJ}(f)}{\Gamma_{II}(0)} = \frac{\Gamma_{IJ}(f)}{2/5}$$

- orf describes the reduction of cross-correlation from mis-aligned, separated detectors

Cross-Correlation

For a detector pair (IJ), segment t:

$$\hat{C}_t^{IJ}(f_k) = \frac{2}{T_{\text{seg}}} \frac{\text{Re}[\tilde{s}_{I,t}^*(f_k) \tilde{s}_{J,t}(f_k)]}{\gamma_{IJ}(f_k) S_0(f_k)}$$



$$\hat{C}_k^{IJ} \equiv \hat{C}^{IJ}(f_k) = \frac{\sum_t \hat{C}_t^{IJ}(f_k) [\sigma_t^{IJ}(f_k)]^{-2}}{\sum_t [\sigma_t^{IJ}(f_k)]^{-2}}$$

$$[\sigma^{IJ}(f_k)]^2 = \left[\sum_t \frac{2T_{\text{seg}} \Delta f \gamma_{IJ}^2(f_k) S_0^2(f_k)}{P_{I,t}(f_k) P_{J,t}(f_k)} \right]^{-1}$$



$$\langle \hat{C}_k^{IJ} \rangle = \Omega_M(f_k | \Theta)$$

$$\text{Var}[\hat{C}_k^{IJ}] = [\sigma_k^{IJ}]^2$$

SNR

- cross-correlation estimator

$$Y = \int_{-T/2}^{T/2} dt \int_{-T/2}^{T/2} dt' s_1(t) s_2(t') \boxed{Q(t-t')} \quad \text{filter}$$

- SNR

$$\frac{S}{N} \equiv \frac{\langle Y \rangle}{\sqrt{\langle Y^2 \rangle - \langle Y \rangle^2}} = T^{1/2} \frac{\int_{-\infty}^{\infty} df S_h(f) \Gamma(f) \tilde{Q}(f)}{\left[\int_{-\infty}^{\infty} df |\tilde{Q}(f)|^2 P_1(f) P_2(f) \right]^{1/2}}$$

- Optimized SNR

$$\frac{S}{N} = \left[2T \int_0^{\infty} df \Gamma^2(f) \frac{S_h^2(f)}{P_1(f) P_2(f)} \right]^{1/2}$$

optimal filter

$$\tilde{Q}_{\text{opt}}(f) \propto \frac{\Gamma(f) S_h(f)}{P_1(f) P_2(f)}$$

Bayesian Analysis

We can now construct the likelihood

$$p(\hat{C}_k^{IJ} | \Theta) \propto \exp \left[-\frac{1}{2} \sum_{IJ} \sum_k \left(\frac{\hat{C}_k^{IJ} - \Omega_M(f_k | \Theta)}{\sigma_{IJ}(f_k)} \right)^2 \right]$$

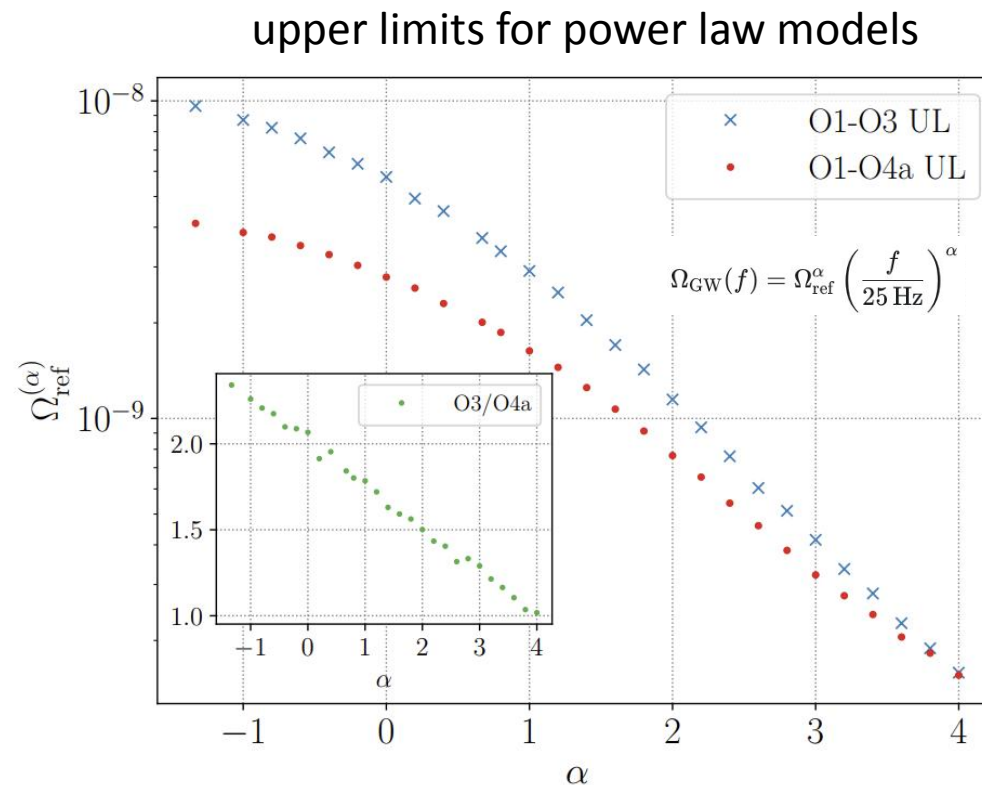
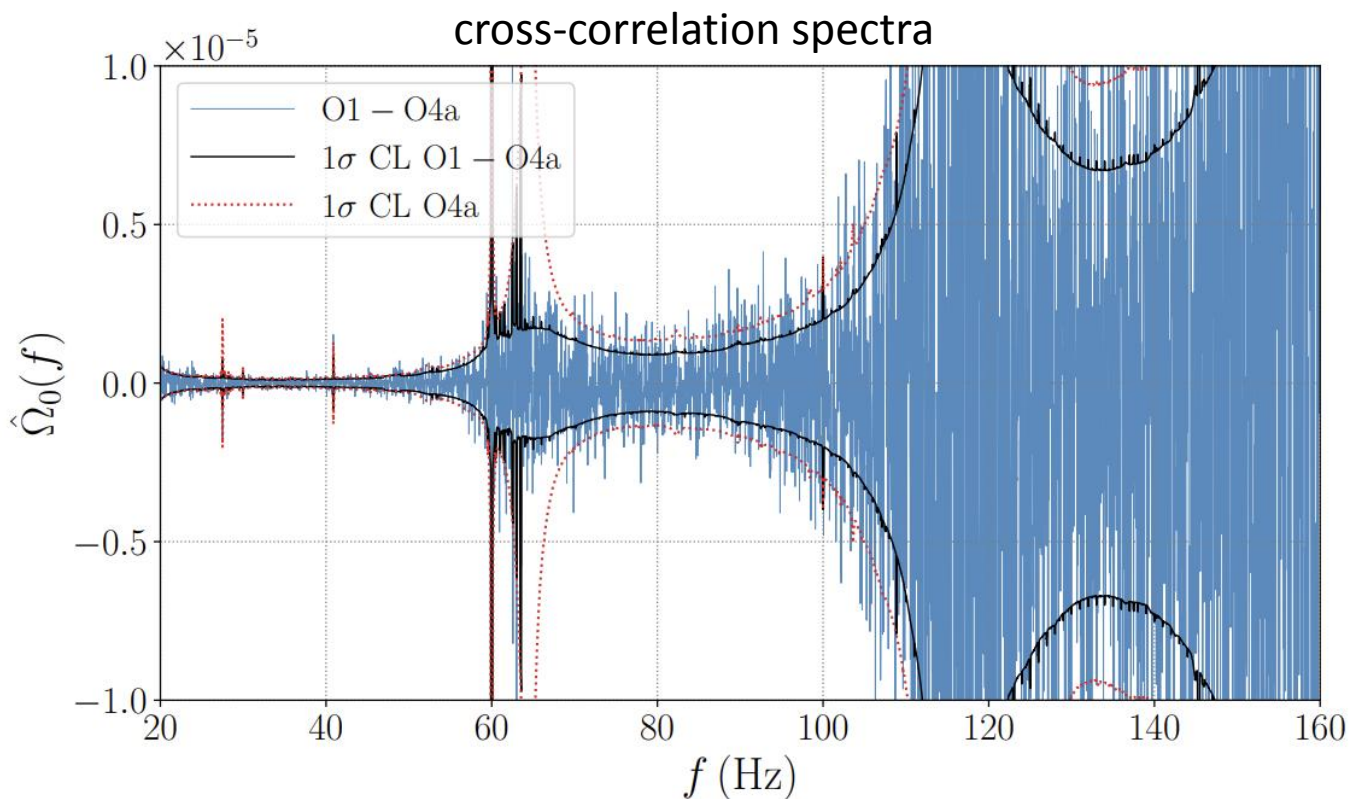
- Specify a prior
- Specify a model (astro, cosmo)
- Sample the posterior
- Model comparison (Bayes factor)
- Parameter estimation

Bayes's theorem

$$p(H_i | D, I) = \frac{p(H_i | I) p(D | H_i, I)}{p(D | I)}$$

Cross-Correlation: O4a

LVK [2508.20721]



Consistent with noise only scenario. Exclusion for power law models obtained

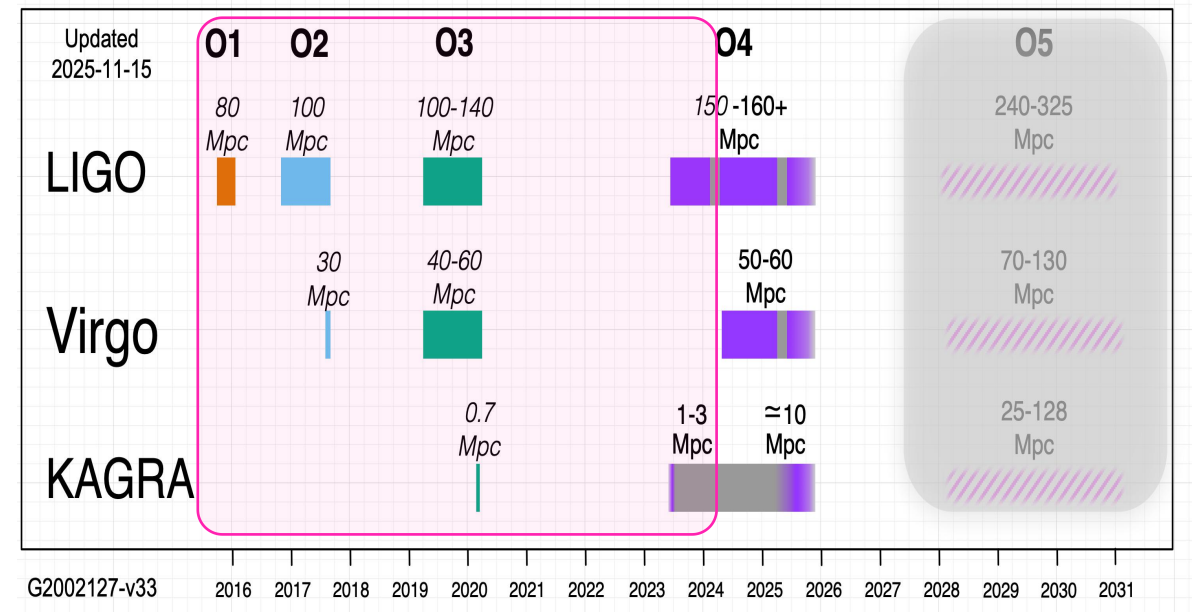
Cosmological Sources in O4a Analysis

Cosmological and High Energy Physics implications from gravitational-wave background searches in LIGO-Virgo-KAGRA's O1-O4a runs

The LIGO Scientific Collaboration, The Virgo Collaboration, and The KAGRA Collaboration*

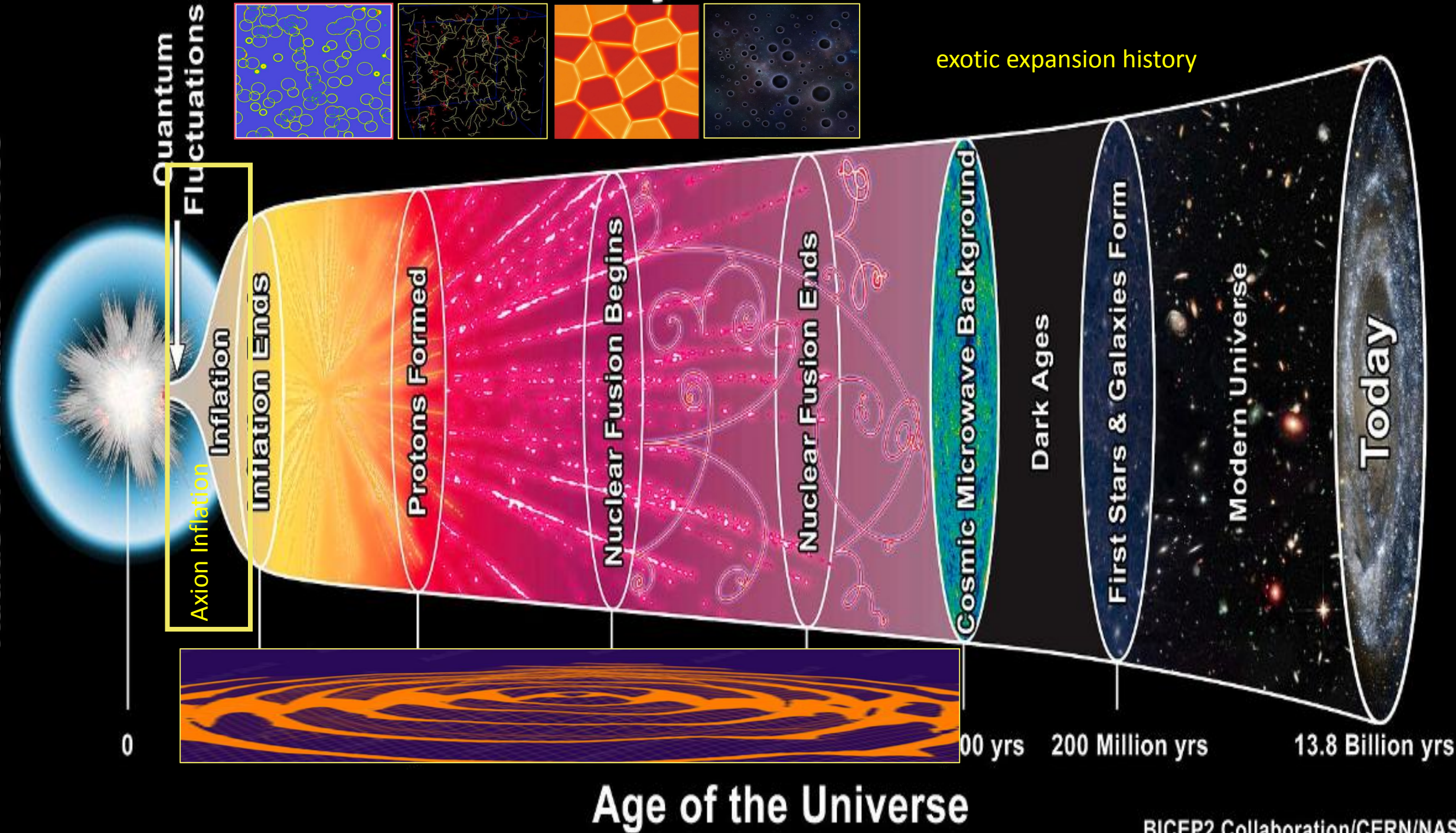
arxiv:2510.26848

- Cosmological phase transitions
- Cosmic strings
- Domain walls
- Stiff equation of state
- Axion Inflation
- Second order scalar perturbations
- Primordial black holes
- Parity violation (not covered here)

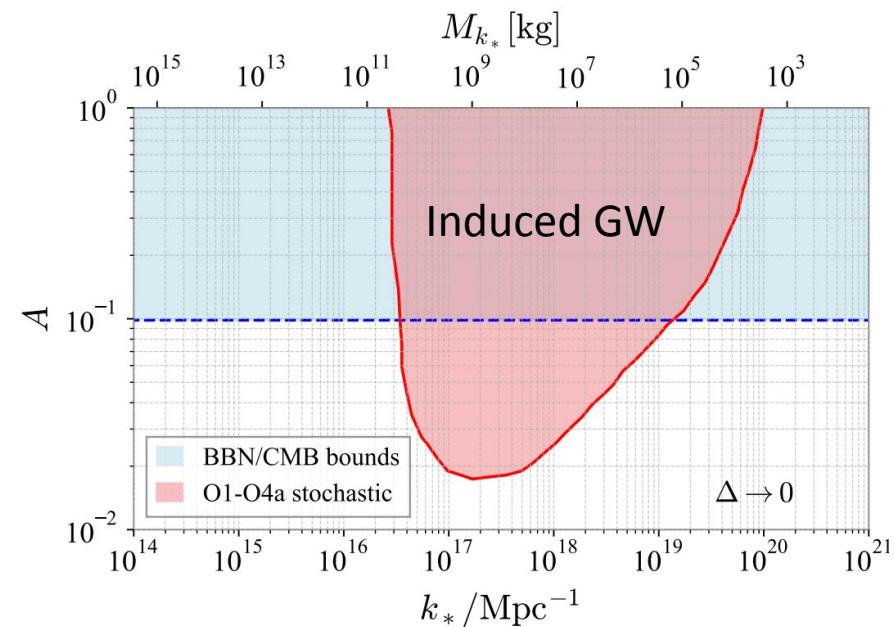
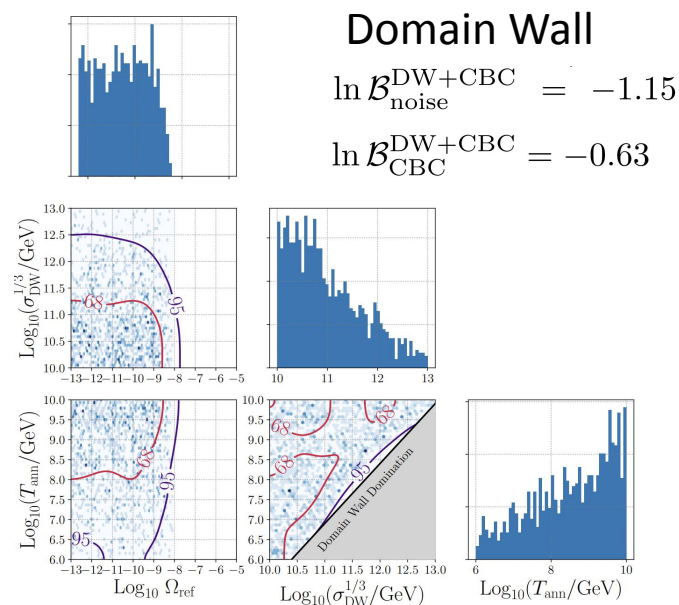
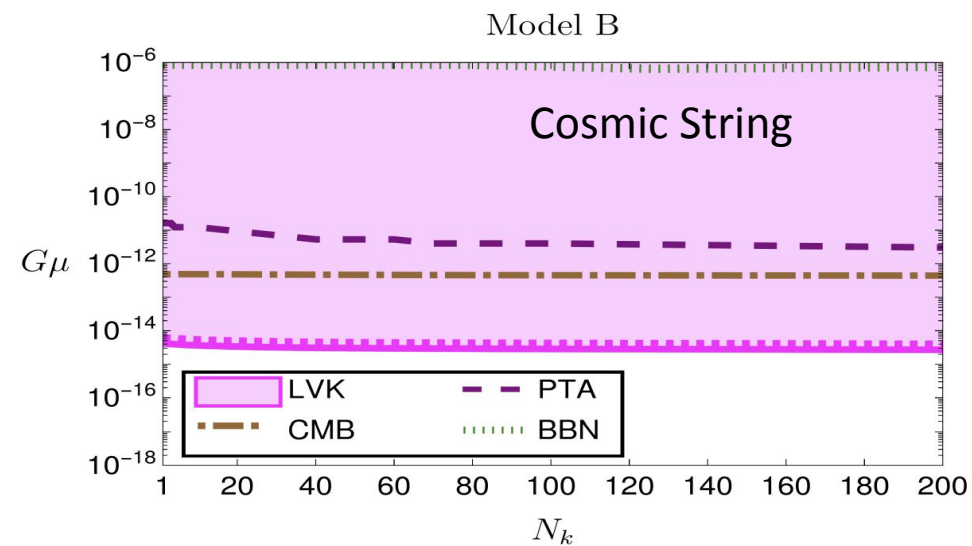
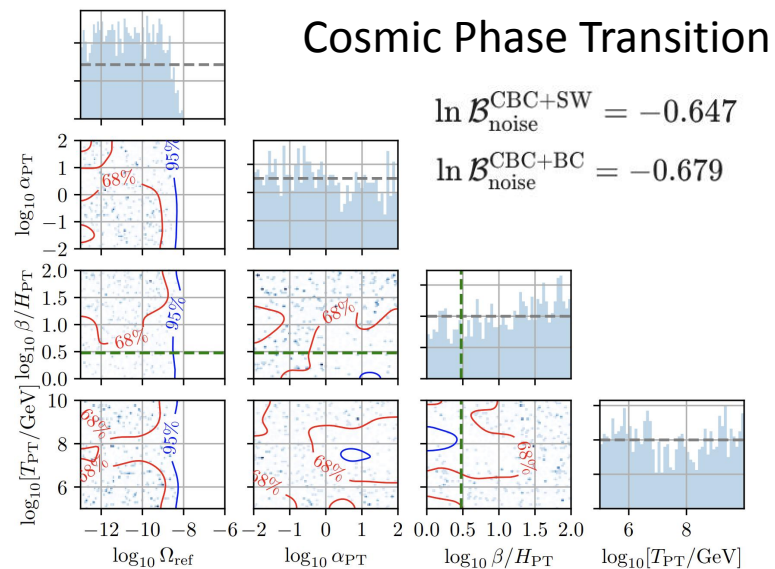


O4a: 2023-05-24 to 2024-01-16 16:00:00 UTC

Radius of the Visible Universe

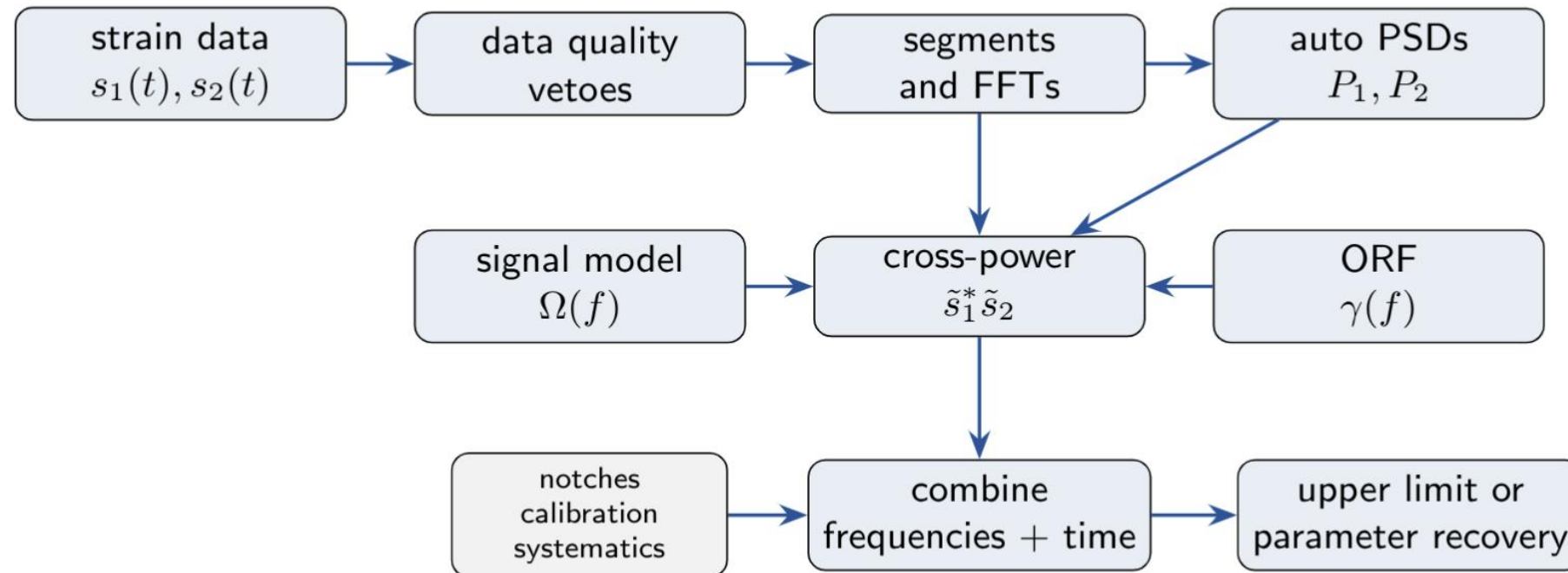


Representative Results



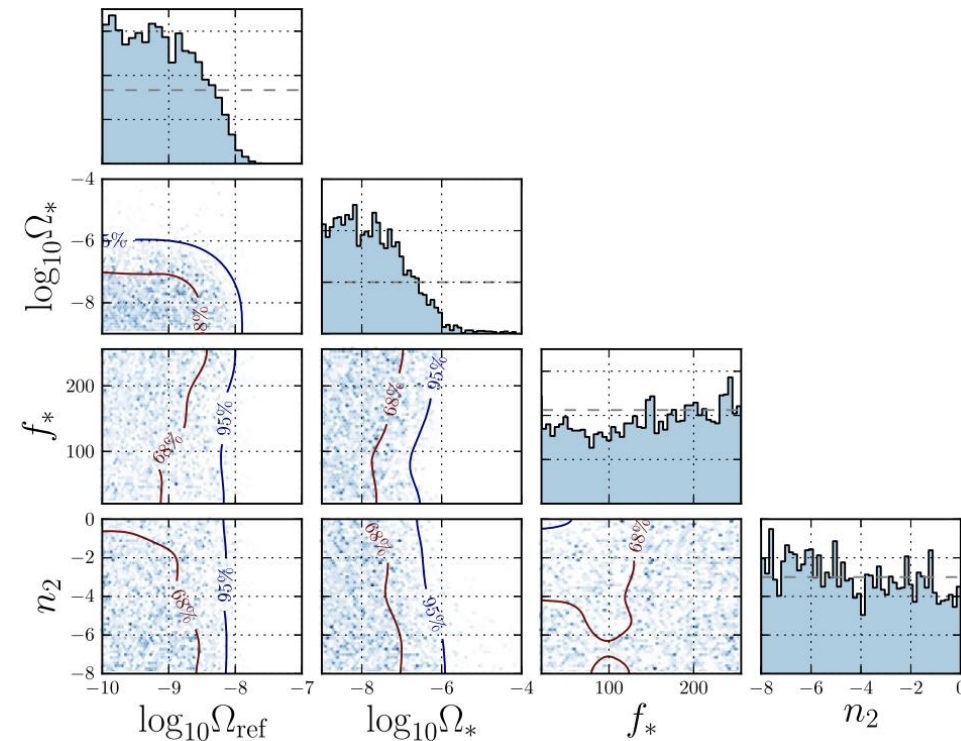
pyGWB

<https://github.com/a-renzini/pygwb>



Reproducing O3 PT Studies

- Phase transitions + CBC





空间引力波探测：太极、天琴、LISA

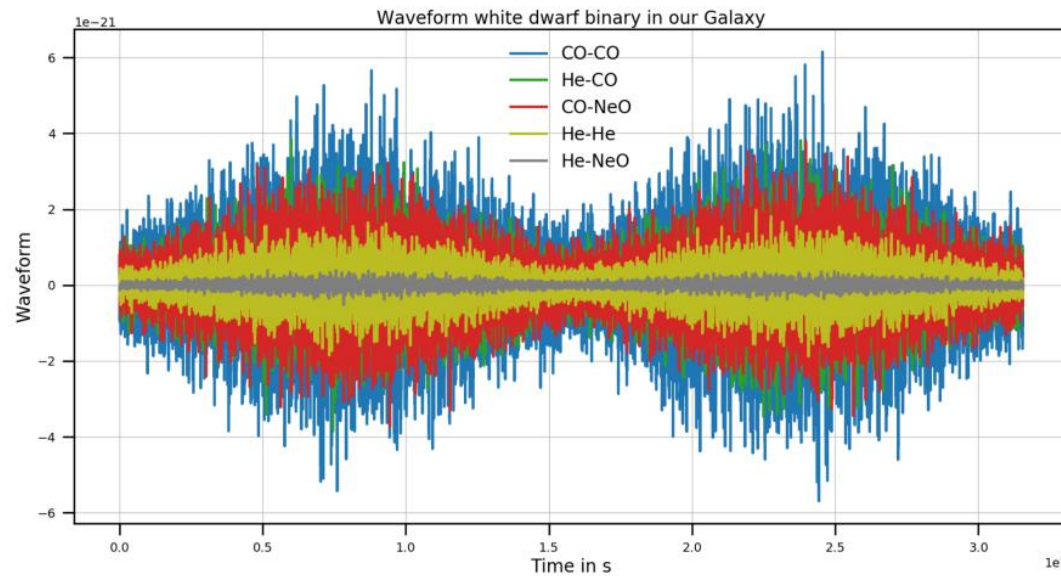
Sources

Cosmological

- Inflation
- Phase transitions (EW, GUT, etc)
- Cosmic strings, Domain walls
- SIGW
- PBH

Astrophysical

- Galactic DWD
- BBH, BNS, etc
- EMRI
- SMBH binary



Detection in Space

Stochastic GW detection in space:

- Complicated, and correlated noise
- Complications from time-delay interferometry
- Solution: **null channel method**, or with a **detector network**

Studies on PT detection in space:

Gowling, Hindmarsh, Hooper, Torrado, JCAP [2209.13551]

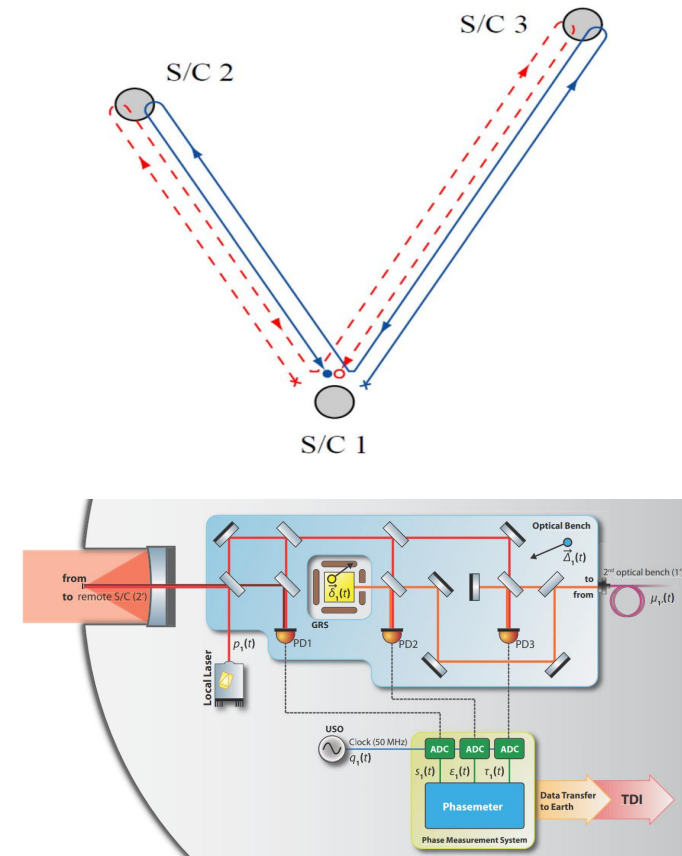
Gowling, Hindmarsh, JCAP [2106.05984]

Boileau, et al, JCAP [2209.13277]

Huang, Bi, Cao, Huang, CPC [2309.14045]

Lewicki, et al, PRD [2403.03769]

Caprini, et al, JCAP [2403.03723]

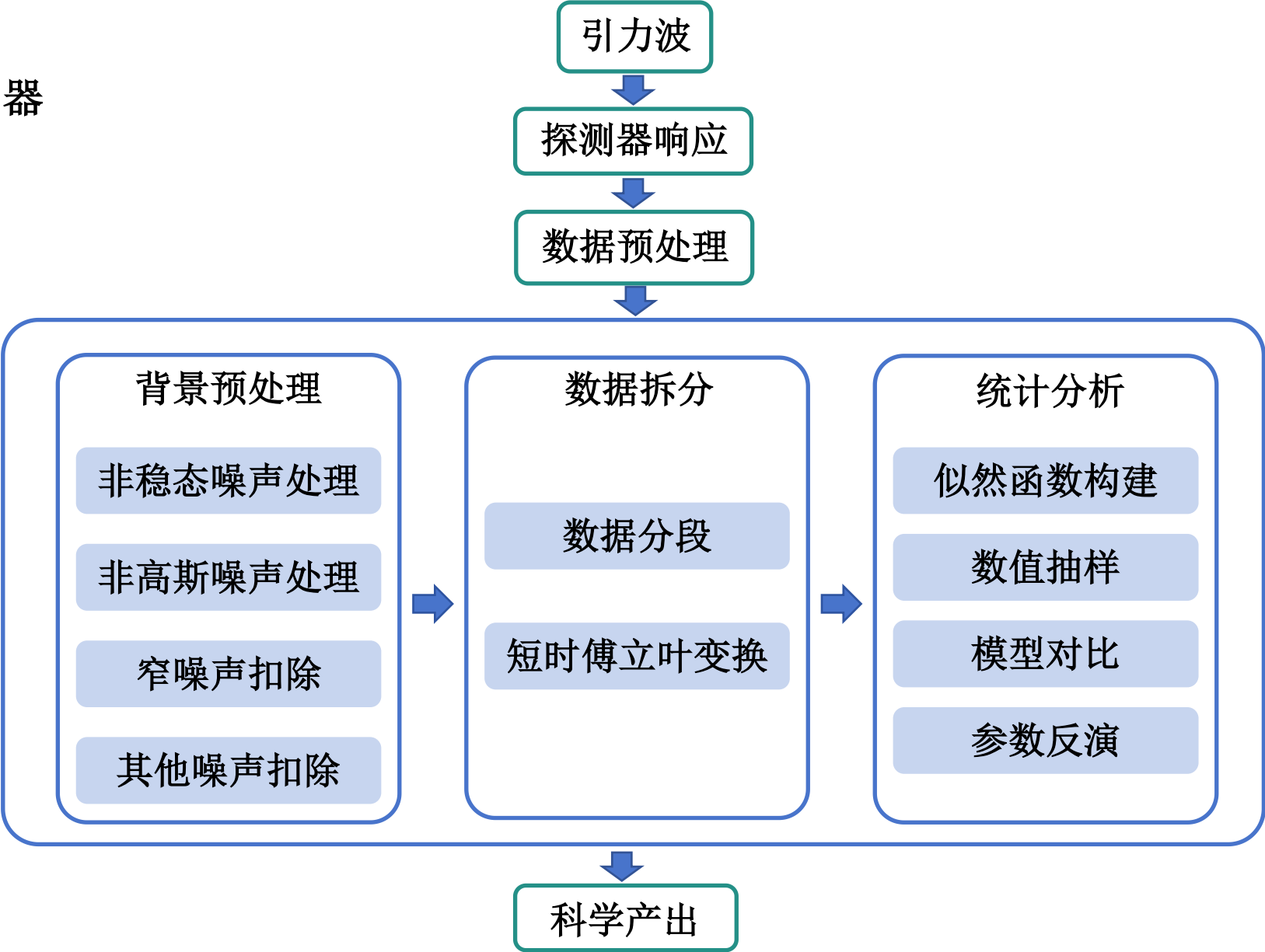


Cosmo SGB detectable down to $\Omega_{GW} \sim O(10^{-13})$

Boileau et al, MNRAS [2105.04283]

数据处理流程

基于单探测器



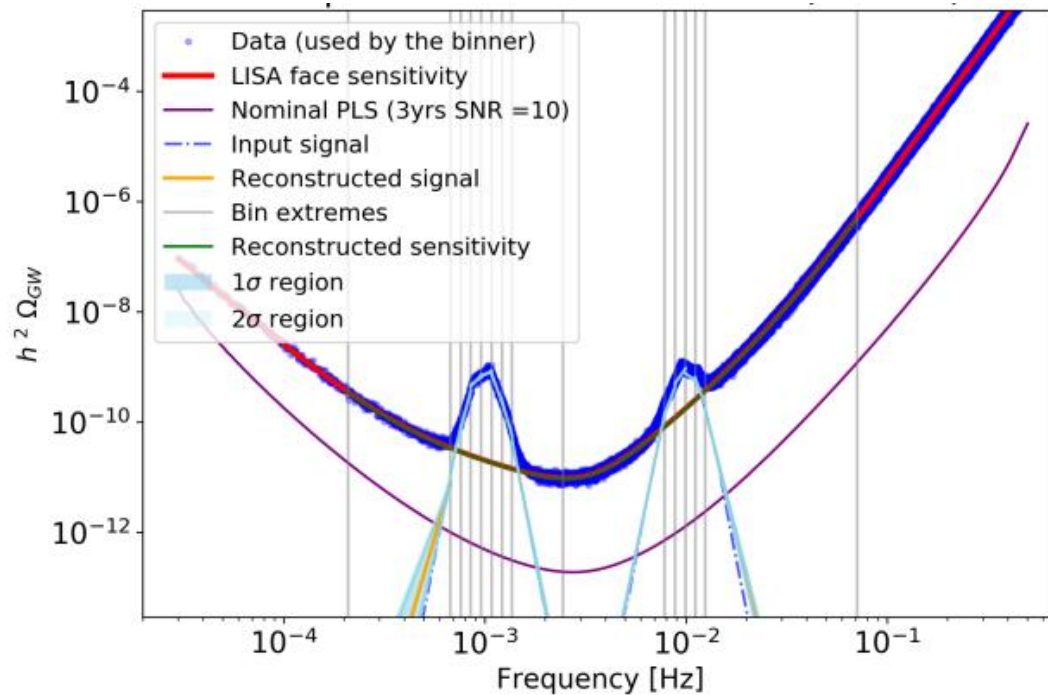
Agnostic Detection

Fit to data bin by bin:

SGWBinner

$$h^2 \Omega_{\text{GW}}(f, \vec{s}_i) = 10^{\alpha_i} \left(\frac{f}{\sqrt{f_{\min,i} f_{\max,i}}} \right)^{p_i} \theta(f - f_{\min,i}) \theta(f_{\max,i} - f)$$

$$h^2 \Omega_{\text{GW}}(f, \vec{s}_i) = 10^{\alpha_i} \theta(f - f_{\min,i}) \theta(f_{\max,i} - f)$$

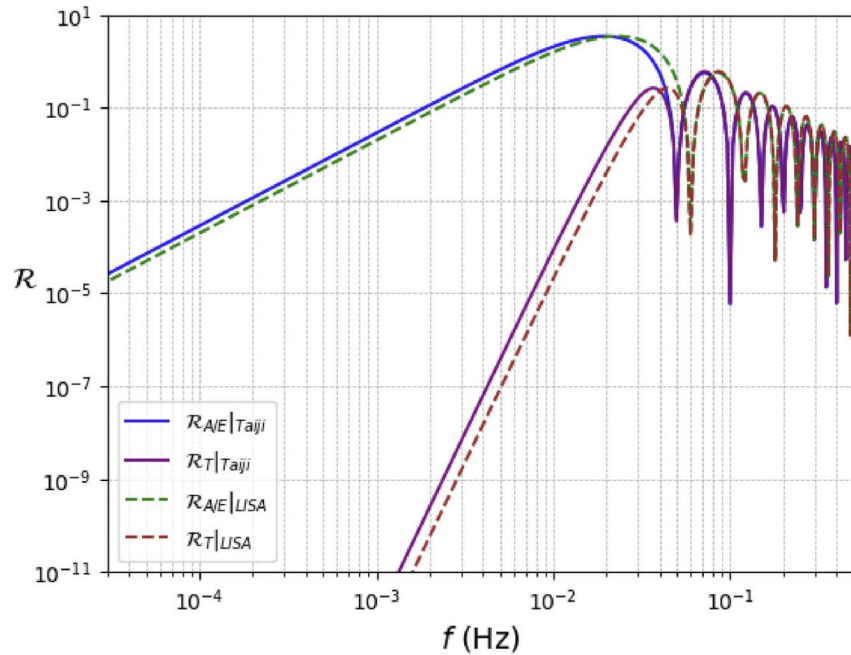


Other Methods

- Excess power method: Karnesis, Petiteau, CQG [1906.09027]
- Principal component analysis: Pieroni, Barausse, JCAP [2004.01135]
- ...

Null Channel Method

Adams, Cornish, PRD [1002.1291]



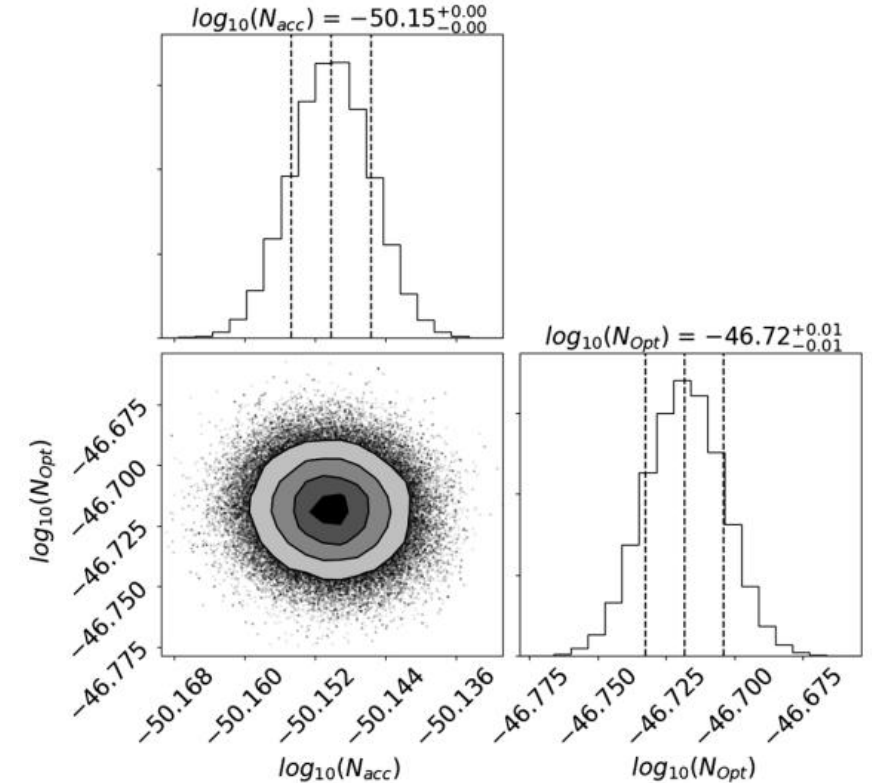
TDI: cancel dominant noise from laser frequency fluctuation

$$\begin{cases} A = \frac{1}{\sqrt{2}}(Z - X) \\ E = \frac{1}{\sqrt{6}}(X - 2Y + Z) \\ T = \frac{1}{\sqrt{3}}(X + Y + Z). \end{cases}$$

$$P_a(f) = S_a(f) + N_a(f)$$

$$S_a(f) = \frac{3H_0^2}{4\pi^2} \frac{\Omega_{\text{GW}}(f)}{f^3} \mathcal{R}_a(f)$$

Shuo Guan, [HG](#), Dian Jiao, Qingyuan Liang, Lei Wu, Yang Zhang, PRD [2511.00996]



Boileau, Christensen, Meyer, Cornish, PRD [2011.05055]

TDI optimization: Wang, Li, Xu, Fan, PRD [2201.10902], ...

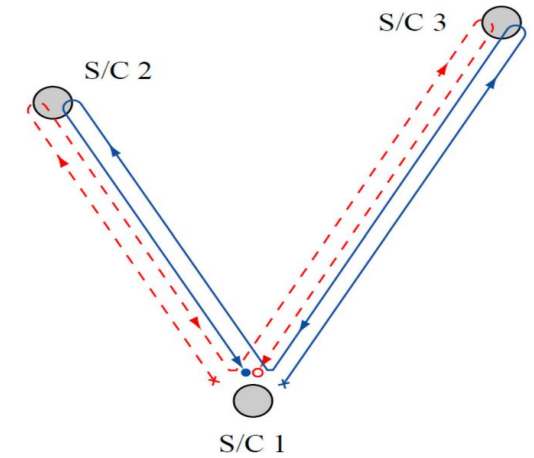
Detector Response

3 orthogonal TDI channels

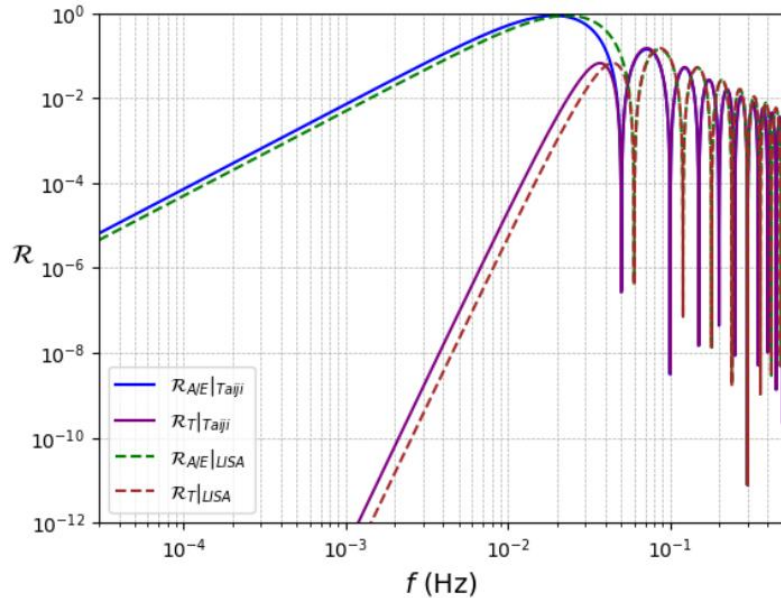
TDI: time-delay interferometry
Tinto, Dhurandhar, LRR, 2021

$$\langle \tilde{d}_a(f) \tilde{d}_b^*(f') \rangle = \frac{1}{2} P_a(f) \delta_{ab} \delta(f - f')$$

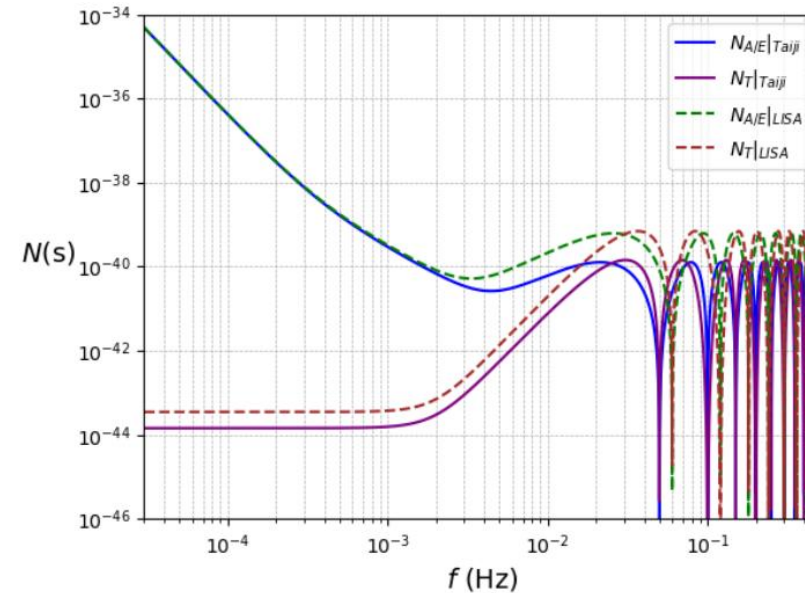
$$P_a(f) = \frac{3H_0^2}{4\pi^2} \frac{\Omega_{\text{GW}}}{f^3} R_a(f) + N_a(f)$$



Response functions



Noise spectra



Noise Sources

There are many noise sources. The two usually considered:

$$N_1(f) = \frac{1}{L^2} \left\{ 4S_{\text{OMS}}(f) + 8 \left[1 + \cos^2 \left(\frac{f}{f_*} \right) \right] \right. \\ \left. \times \frac{1}{(2\pi f)^4} S_{\text{acc}}(f) \right\} |W(f)|^2, \\ N_2(f) = -\frac{1}{L^2} \left[2S_{\text{OMS}}(f) + \frac{8}{(2\pi f)^4} S_{\text{acc}}(f) \right] \\ \times \cos \left(\frac{f}{f_*} \right) |W(f)|^2.$$

$$N_A = N_E = N_1 - N_2,$$

$$N_T = N_1 + 2N_2,$$

- Optical path-length fluctuation

$$\sqrt{S_{\text{OMS}}(f)} = \delta x \sqrt{1 + \left(\frac{2 \text{ mHz}}{f} \right)^4} \left(\frac{\text{m}}{\sqrt{\text{Hz}}} \right)$$

- Test mass acceleration noise

$$\sqrt{S_{\text{acc}}(f)} = N_{\text{acc}} \sqrt{1 + \left(\frac{0.4 \text{ mHz}}{f} \right)^2} \sqrt{1 + \left(\frac{f}{8 \text{ mHz}} \right)^4} \\ \times \left(\frac{\text{m}}{\text{s}^2 \sqrt{\text{Hz}}} \right),$$

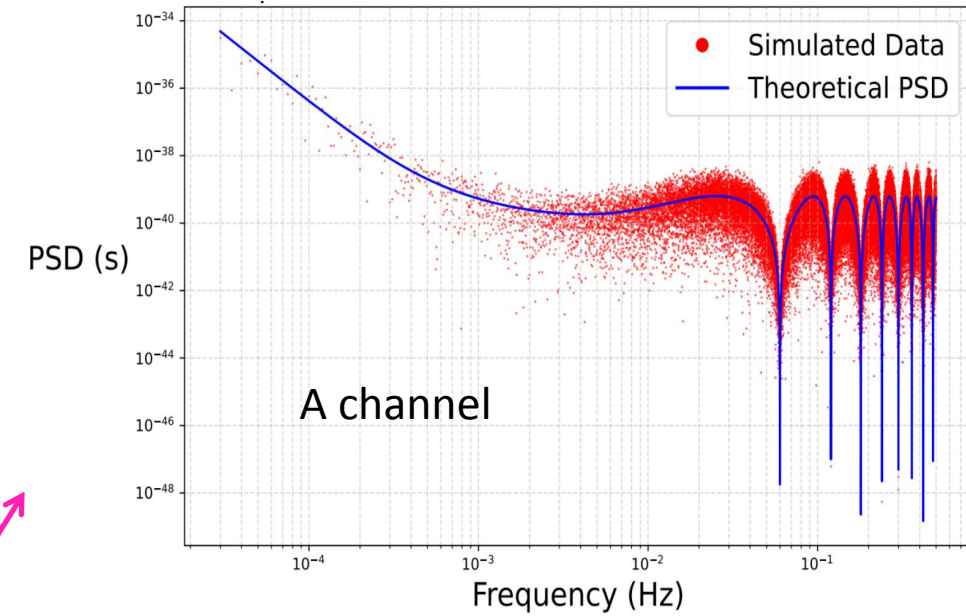
$$\delta x = 8 \times 10^{-12}$$

$$N_{\text{acc}} = 3 \times 10^{-15}$$

Likelihood

The core of the statistical analysis: likelihood

both signal and noise behave as random variables

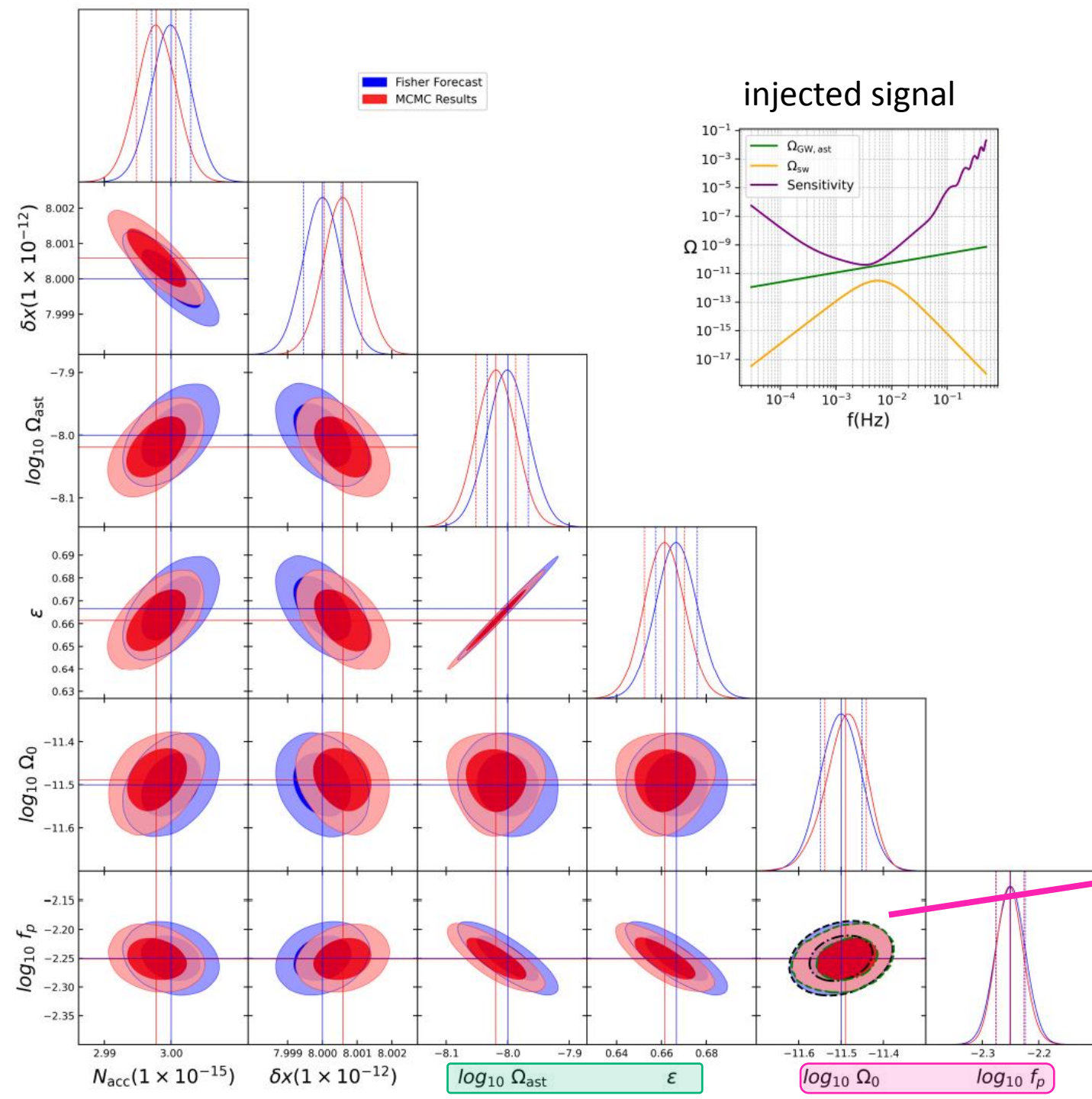


$$\mathcal{L} = \prod_{\kappa=1}^{N_0} \prod_{k=1}^{N/2} \frac{1}{8\pi^3 \sigma_A^2 \sigma_E^2 \sigma_T^2} \exp \left[- \sum_{a=A,E,T} \frac{|\tilde{d}_a^\kappa(f_k)|^2}{2\sigma_a^2} \right]$$

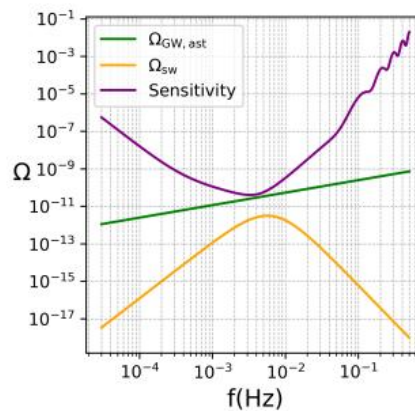
$N_0 = 126$: number of segments
 $N = 10^6$: frequency bins in one segment
 $T = 1.26 \times 10^8$ s : observation time
 $f_s = 1$ Hz : sampling frequency

$$\sigma_a^2 = \frac{T f_s^2}{4} \left[\frac{3H_0^2}{4\pi^2} \frac{\Omega_{\text{GW}}}{f^3} R_a(f) + N_a(f) \right]$$

astrophysical background (extra-galactic BBH, BNS)
 astrophysical foreground (double white dwarf)
 cosmological background (PT)



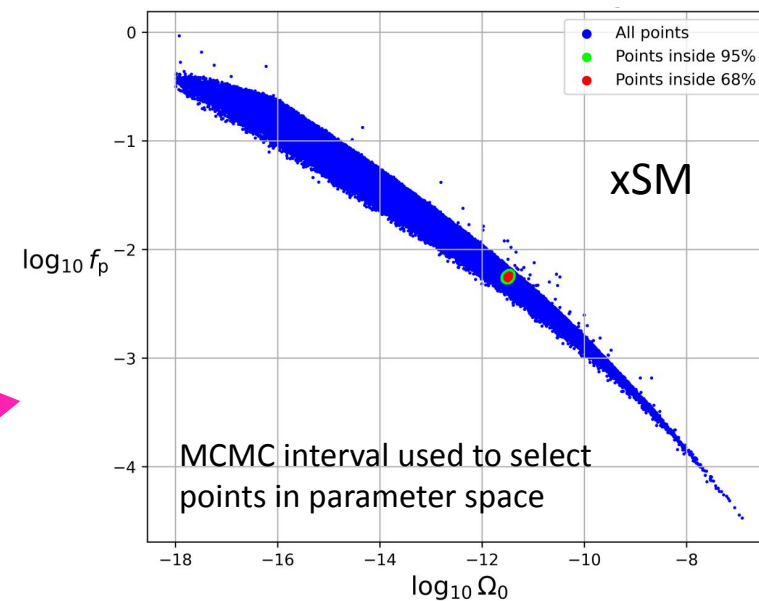
injected signal

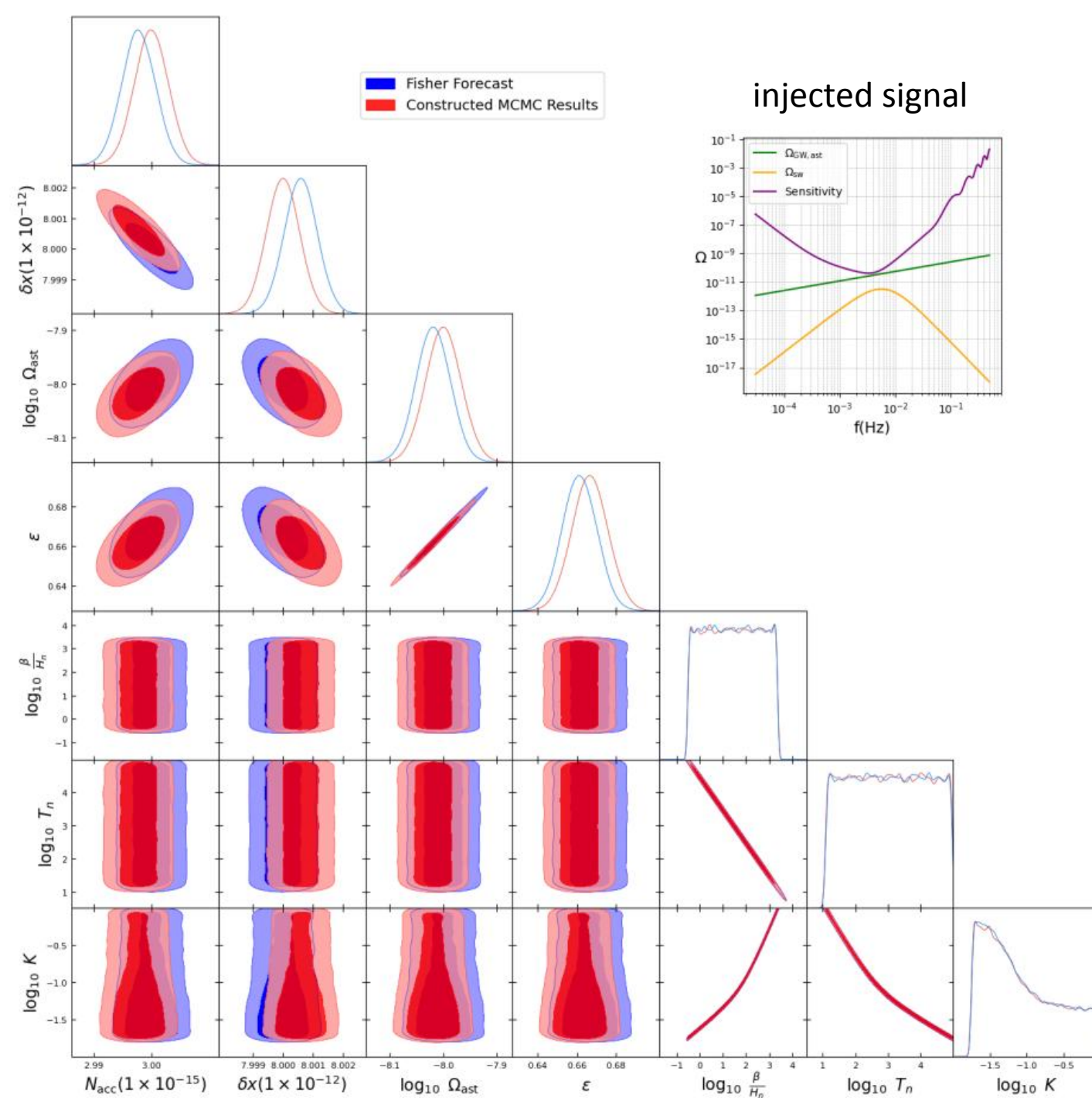


Simplified SW spectrum with no parameter degeneracy

$$\Omega_{\text{sw}}(f) = \Omega_0 \left(\frac{f}{f_p} \right)^3 \left(\frac{7}{4 + 3(f/f_p)^2} \right)^{7/2}$$

$$\Omega_{\text{astro}}(f) = \Omega_{\text{ast}} \left(\frac{f}{f_{\text{ref}}} \right)^\epsilon$$





$$\Omega_{\text{sw}}(f)h^2 = 2.65 \times 10^{-6} \left(\frac{H_n}{\beta} \right) \left(\frac{\kappa_{\text{sw}} \alpha}{1 + \alpha} \right)^2 \left(\frac{100}{g_*} \right)^{1/3} v_w \left(\frac{f}{f_p} \right)^3 \left(\frac{7}{4 + 3(f/f_p)^2} \right)^{7/2} \Upsilon(\tau_{\text{sw}}),$$

$$f_p = \frac{19}{v_w} \left(\frac{\beta}{H_n} \right) \left(\frac{T_n}{100 \text{ GeV}} \right) \left(\frac{g_*}{100} \right)^{1/6} 10^{-6} \text{ Hz}$$

Parameter degeneracy forbids a determination of phase transition parameters.

Future solutions: more observables

xSM Model

Barger, Langacker, McCaskey, Ramsey-Musolf, Shaughnessy,
PRD [0706.4311]

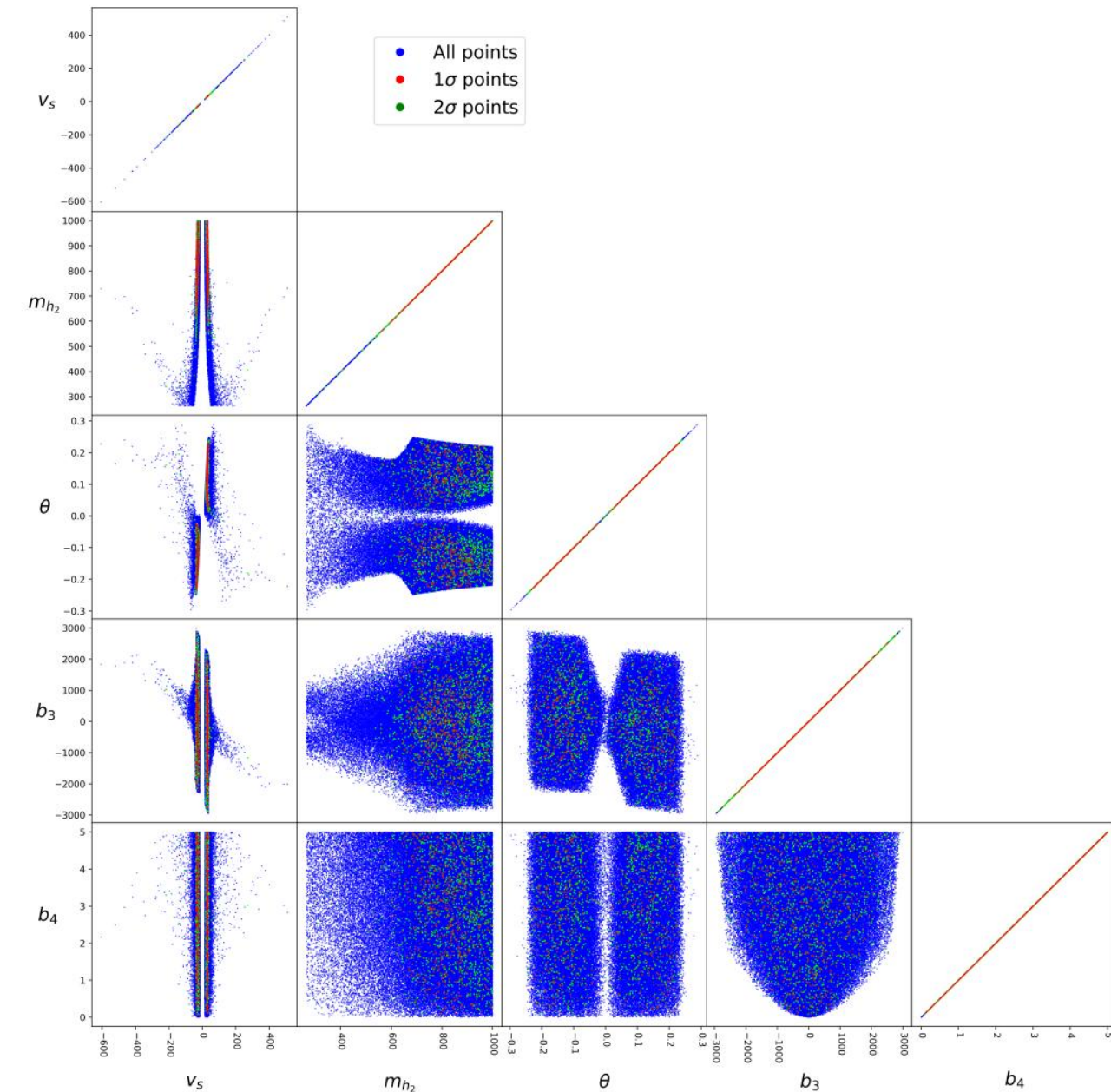
$$V(\phi, S) = -\mu^2 H^\dagger H + \lambda (H^\dagger H)^2 + \frac{a_1}{2} H^\dagger H S + \frac{a_2}{2} H^\dagger H S^2 + \frac{b_2}{2} S^2 + \frac{b_3}{3} S^3 + \frac{b_4}{4} S^4$$

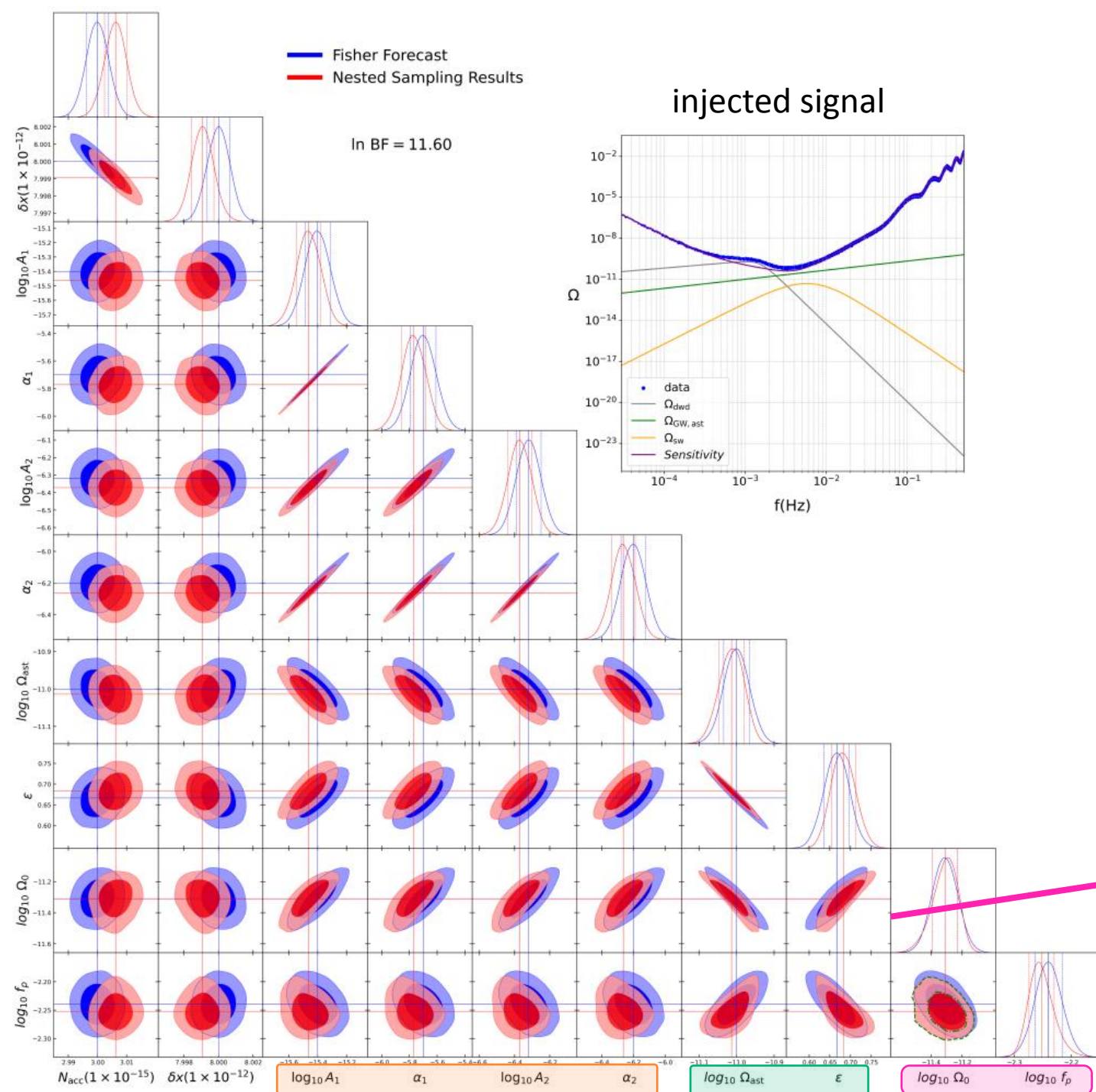


$v_s, \quad m_{h_2}, \quad \theta, \quad b_3, \quad b_4$

Points survive pheno constraints.

Shuo Guan, [HG](#), Dian Jiao, Qingyuan Liang, Lei Wu, Yang Zhang,
PRD [2511.00996]



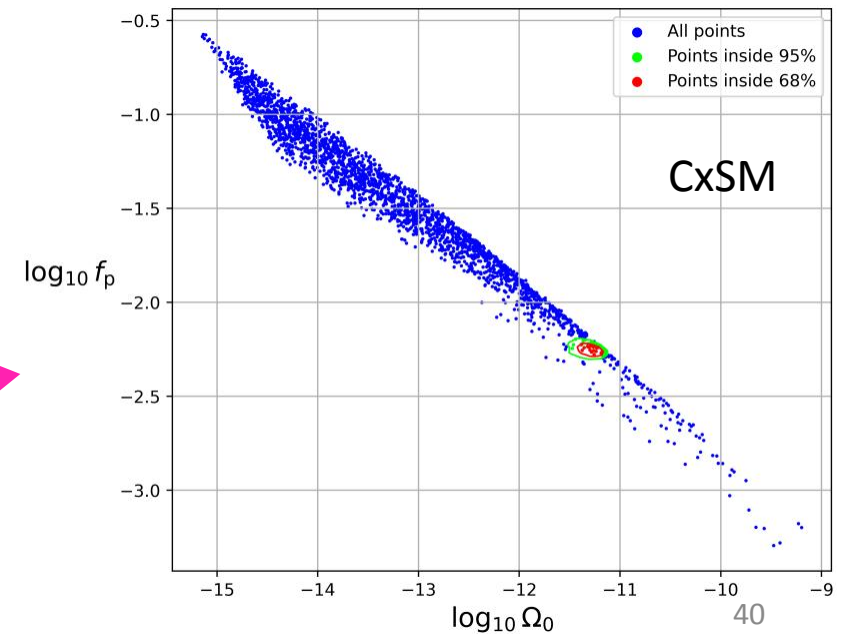


Simplified SW spectrum with no parameter degeneracy

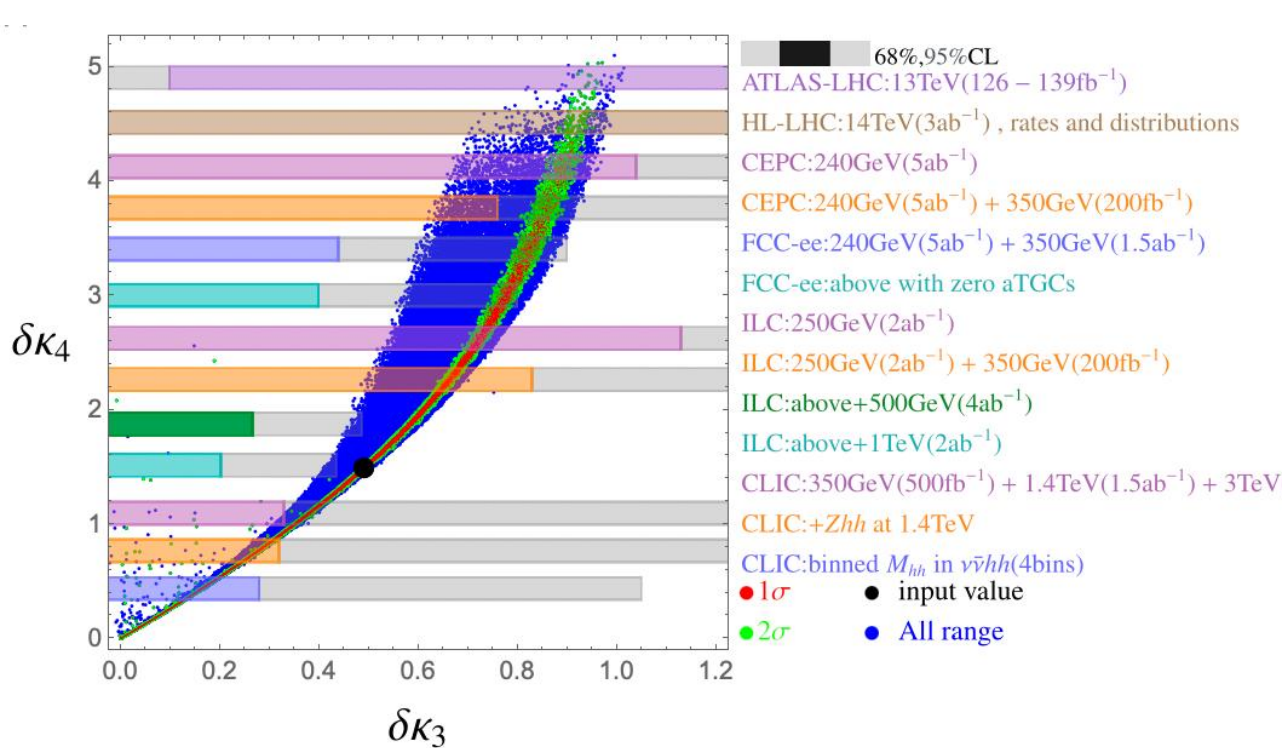
$$\Omega_{\text{sw}}(f) = \Omega_0 \left(\frac{f}{f_p} \right)^3 \left(\frac{7}{4 + 3(f/f_p)^2} \right)^{7/2}$$

$$\Omega_{\text{astro}}(f) = \Omega_{\text{ast}} \left(\frac{f}{f_{\text{ref}}} \right)^\epsilon$$

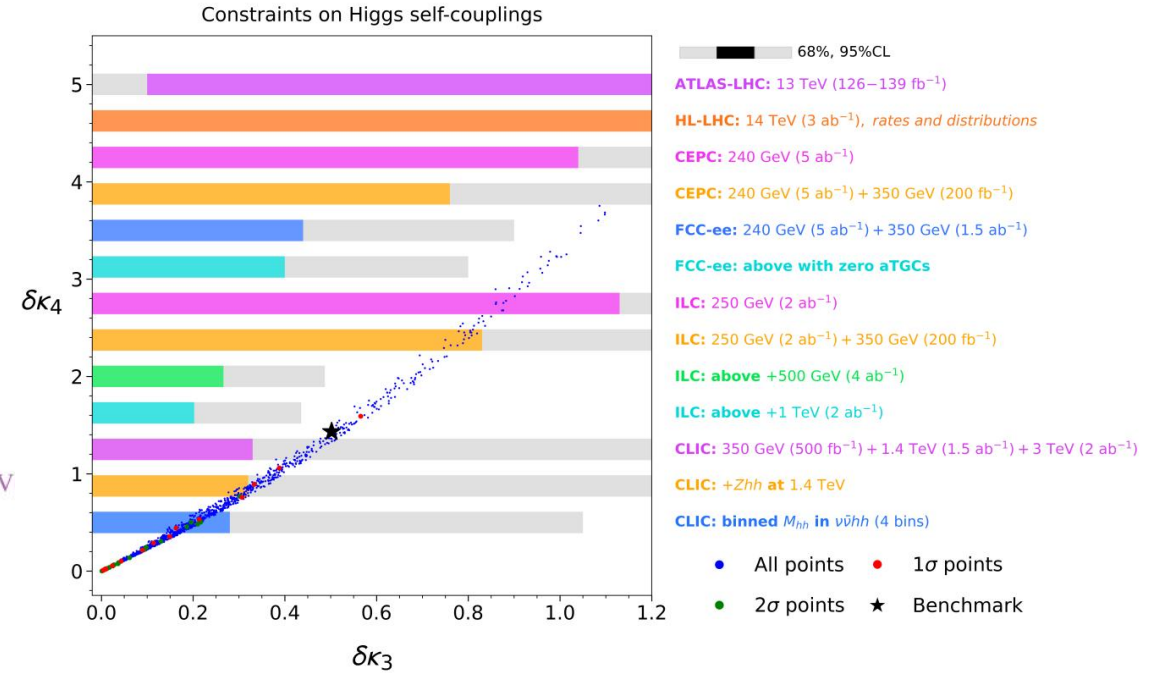
$$\Omega_{\text{dwd}}(f) = \frac{A_1 (f/f_*)^{\alpha_1}}{1 + A_2 (f/f_*)^{\alpha_2}}$$



Higgs Self-Couplings



Shuo Guan, [HG](#), Dian Jiao, Qingyuan Liang, Lei Wu, Yang Zhang,
PRD [2511.00996]



Qingyuan Liang, Ligong Bian, HG, Yongcheng Wu,
PRD [2511.21488]

未来挑战

