BPCS 2025

Enhancing Phase Transition Calculations through Fitting and Neural Network

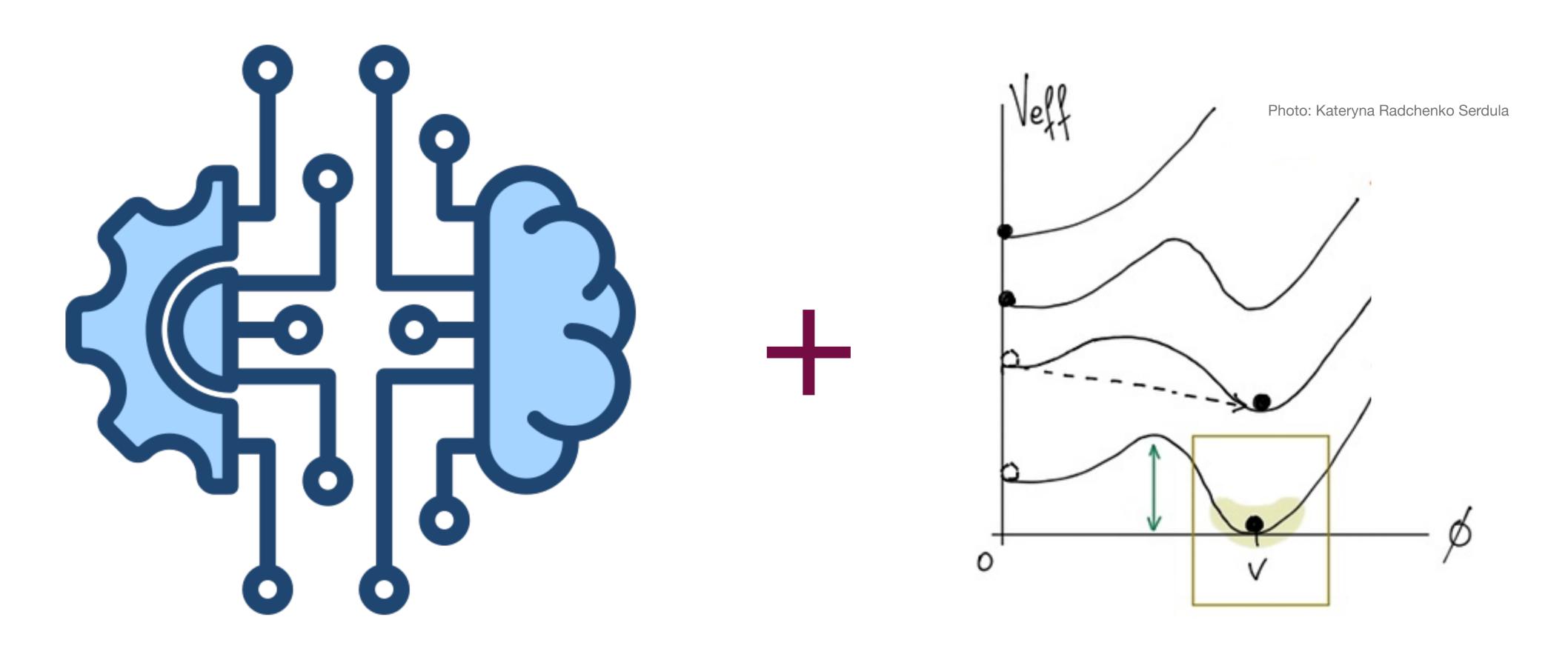
Yang Zhang, Henan Normal University

In collaboration with Ligong Bian, Hongxin Wang, Yang Xiao, Ji-Chong Yang, Jin Min Yang

Based on arXiv:2510.XXXXX



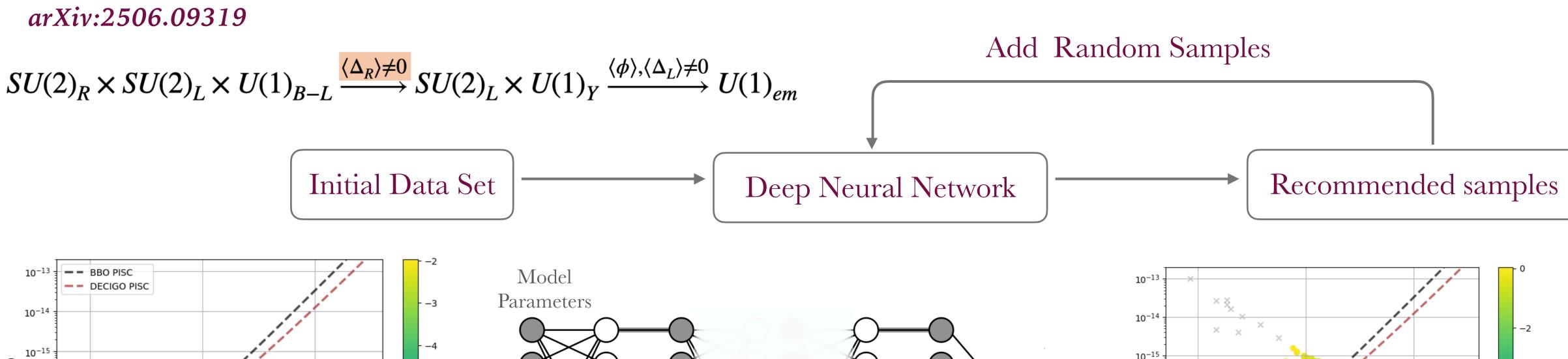
Motivation

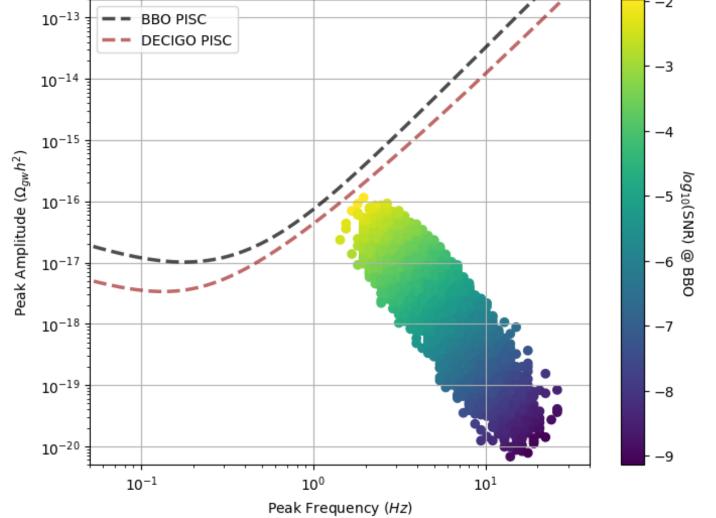


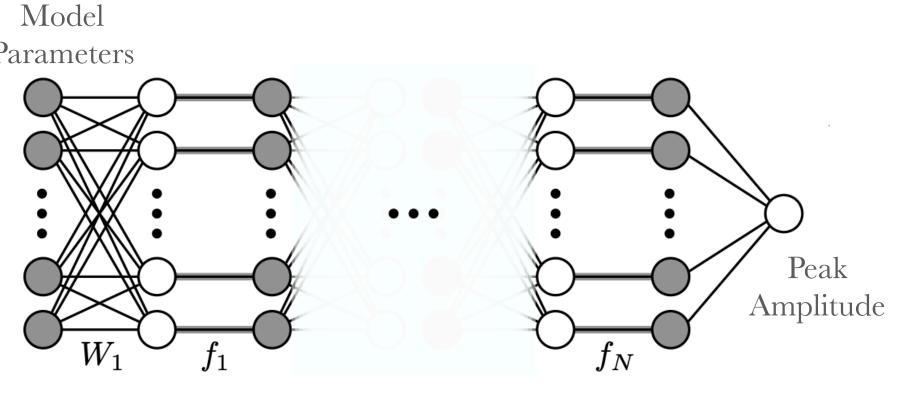
Machine Learning

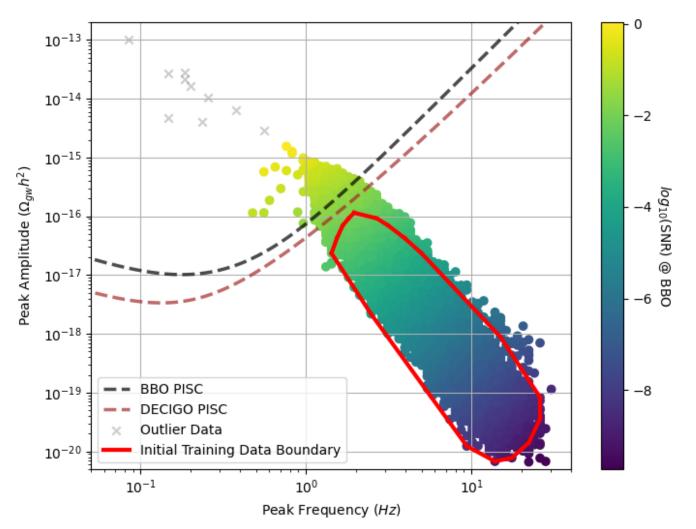
Cosmological Phase Transition

Machine Learning Left-Right Breaking from Gravitational Waves, William Searle, Csaba Balázs, Yang Xiao, Yang Zhang,









 $64 \rightarrow 128 \rightarrow 128 \rightarrow 128 \rightarrow 128 \rightarrow 64$

Solving differential equations with neural networks: Applications to the calculation of cosmological phase transitions Maria Laura Piscopo, Michael Spannowsky, Philip Waite, arXiv:1902.05563

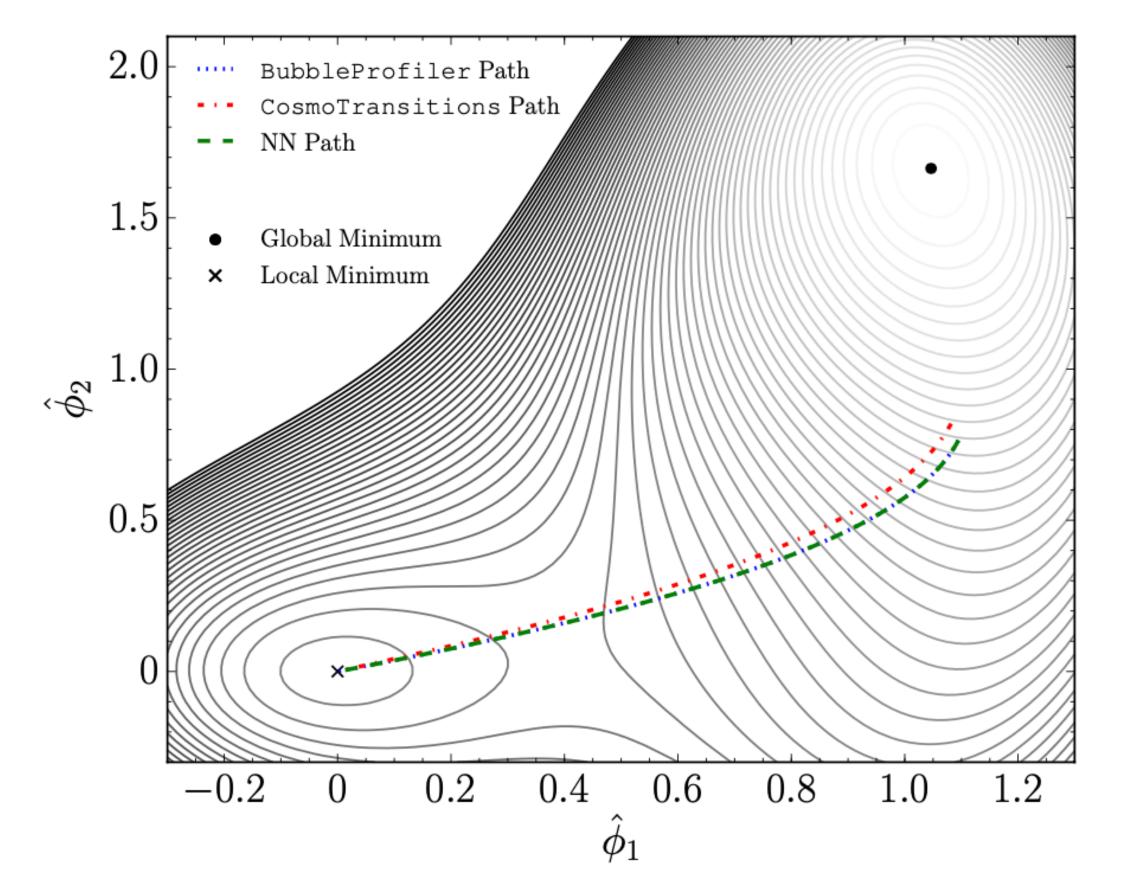
➤ A set of *m* coupled *j*th order differential equations:

$$\mathcal{F}_m(\vec{x}, \phi_m(\vec{x}), \nabla \phi_m(\vec{x}), \cdots, \nabla^j \phi_m(\vec{x})) = 0$$

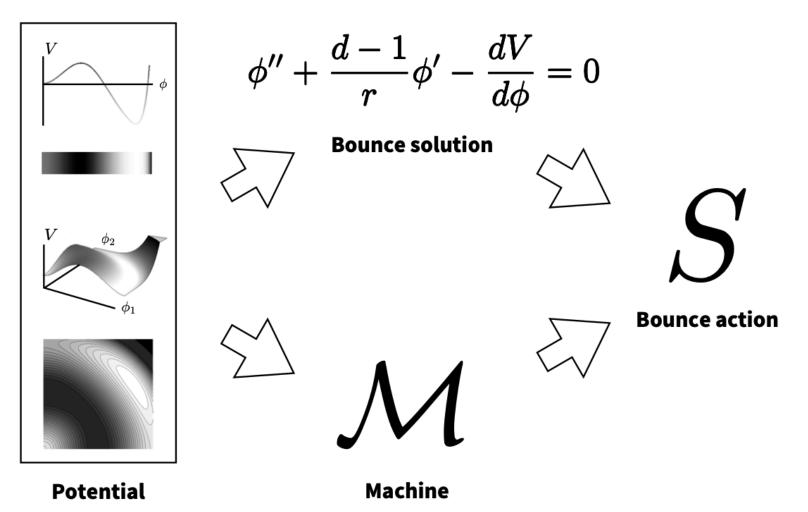
➤ Convert the problem of finding a solution into an optimization problem:

$$egin{aligned} \mathcal{L}(\{w,ec{b}\}) &= rac{1}{i_{ ext{max}}} \sum_{i,m} \hat{\mathcal{F}}_m(ec{x}^i, \hat{\phi}_m(ec{x}^i), \cdots,
abla^j \hat{\phi}_m(ec{x}^i))^2 \ &+ \sum_{ ext{B.C.}} (
abla^p \hat{\phi}_m(ec{x}_b) - K(ec{x}_b))^2 \ , \end{aligned}$$

➤ Minimize the loss function using a neural network.



Machine learning for bounce calculation, Ryusuke Jinno, arXiv:1805.12153



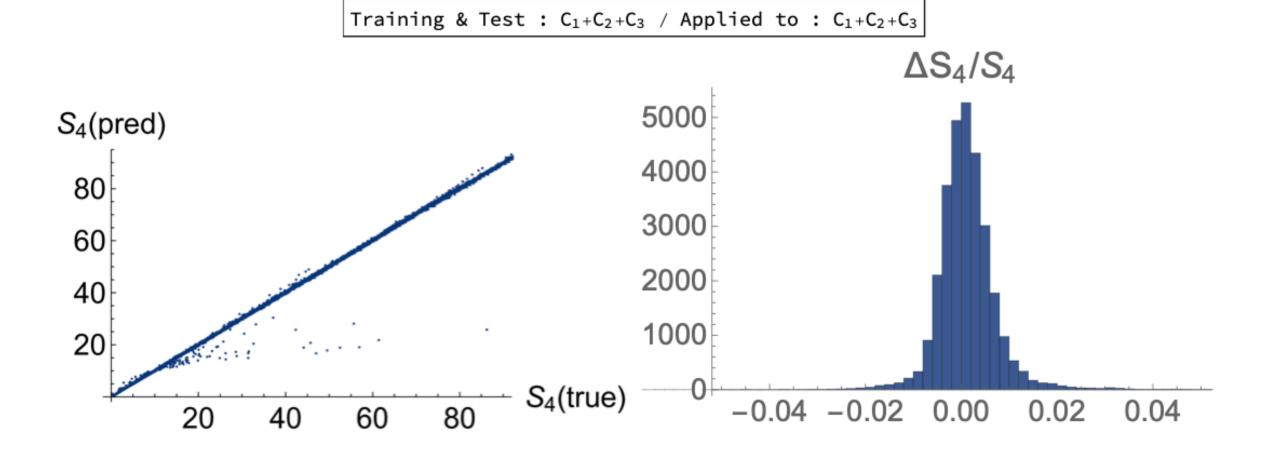
Class 1
$$(C_1)$$
: $V(\phi) = \sum_{n=1}^{7} a_n^{(1)} \phi^{n+1}$,
Class 2 (C_2) : $V(\phi) = \sum_{n=1}^{7} a_n^{(2)} \phi^{2n}$,
Class 3 (C_3) : $V(\phi) = a_1^{(3)} \phi^2 + \sum_{n=2}^{7} a_n^{(3)} \phi^{2n-1}$.

➤ Inputs: values of the potential and its derivatives

$$x_{\text{in}} = \left\{ V(\phi_{\text{sample}}) \middle| \phi_{\text{sample}} = \frac{1}{16}, \frac{2}{16}, \cdots, \frac{15}{16} \right\}$$

$$\oplus \left\{ \frac{dV}{d\phi}(\phi_{\text{sample}}) \middle| \phi_{\text{sample}} = \frac{1}{16}, \frac{2}{16}, \cdots, \frac{15}{16} \right\}$$

$$\oplus \left\{ \frac{d^2V}{d\phi^2}(\phi_{\text{sample}}) \middle| \phi_{\text{sample}} = \frac{0}{16}, \frac{1}{16}, \cdots, \frac{16}{16} \right\}.$$



Transition parameters

➤ Bubble nucleation rate

$$\frac{\Gamma}{V} = A(T)e^{-S_E(T)/T}[1 + \mathcal{O}(\hbar)], A(T) = \left(\frac{S_E(T)}{2\pi T}\right)^3 D(T)$$

➤ Nucleation temperature

$$N(T) = \int_{T_{\text{nuc}}}^{T_{\text{tra}}} \frac{dT}{T} \frac{\Gamma(T)}{H(T)^4} = 1$$

➤ Percolation temperature

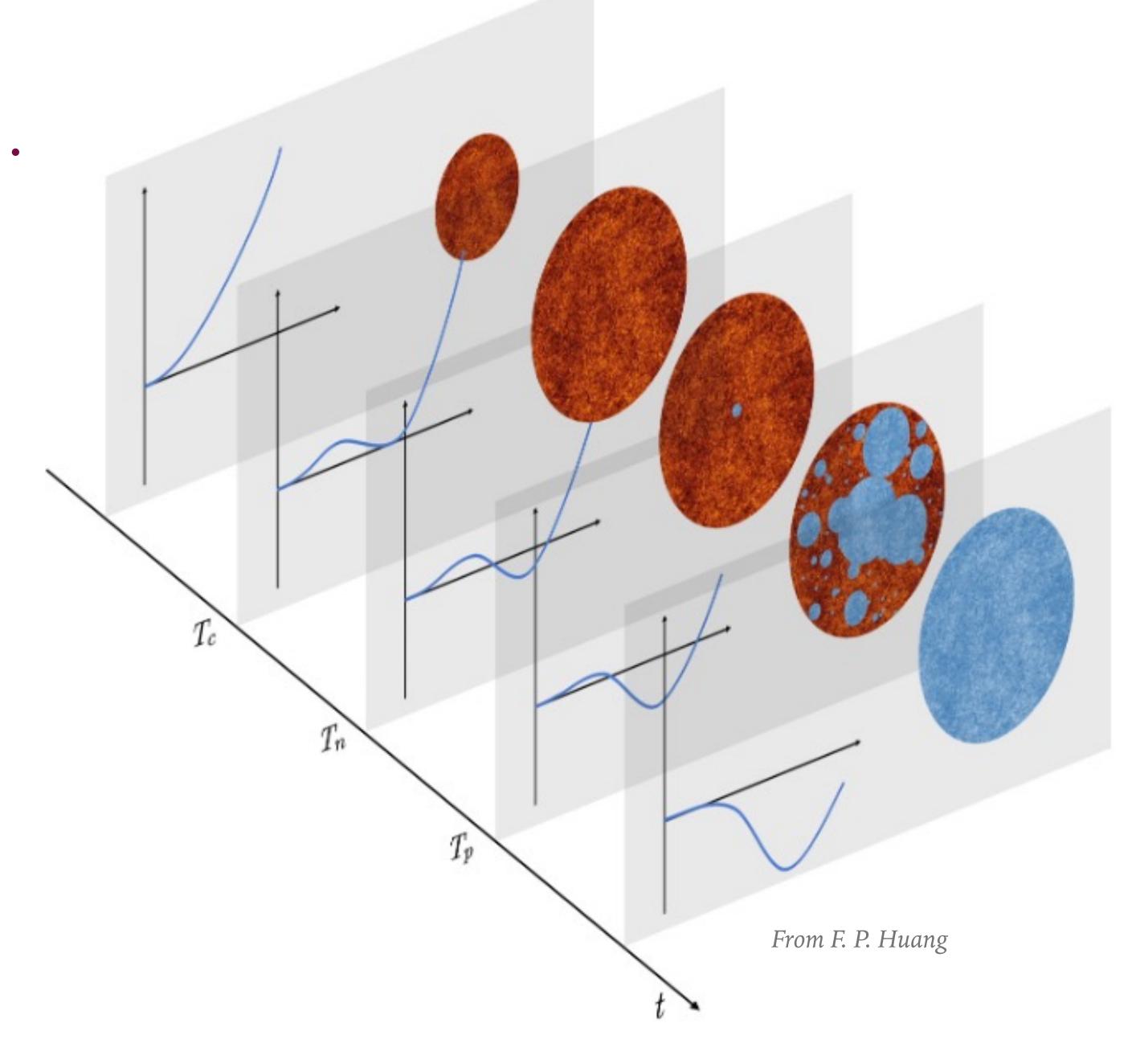
$$P(T_{\text{per}}) = \exp \left[-\frac{64\pi}{3} \xi^4 \int_{T_{\text{per}}}^{T_{\text{tra}}} dT' \frac{\Gamma(T')}{T'^6} \left(\frac{1}{T_{\text{per}}} - \frac{1}{T'} \right)^3 \right] = 70 \%$$

➤ Inverse duration of the phase transition

$$\beta(T) = \frac{d}{dt} \left[\frac{S_E(T)}{T} \right] = TH(T) \frac{d}{dT} \left[\frac{S_E(T)}{T} \right]$$

➤ Ratio of latent heat to radiation density

$$\alpha = \frac{D\theta}{\pi^2 g_* T_*^4 / 30} = \frac{1}{\pi^2 g_* T_*^4 / 30} \left(V(\phi) - \frac{T}{4} \frac{\partial V(\phi, T)}{\partial T} \right) \Big|_{\phi_t}^{\phi_f}$$



Bounce action

➤ Bounce action

$$S_E = S_{d-1} \int_0^\infty \rho^{d-1} \left(\frac{1}{2} \dot{\phi}^2 + V(\phi) \right)$$

where ϕ satisfies

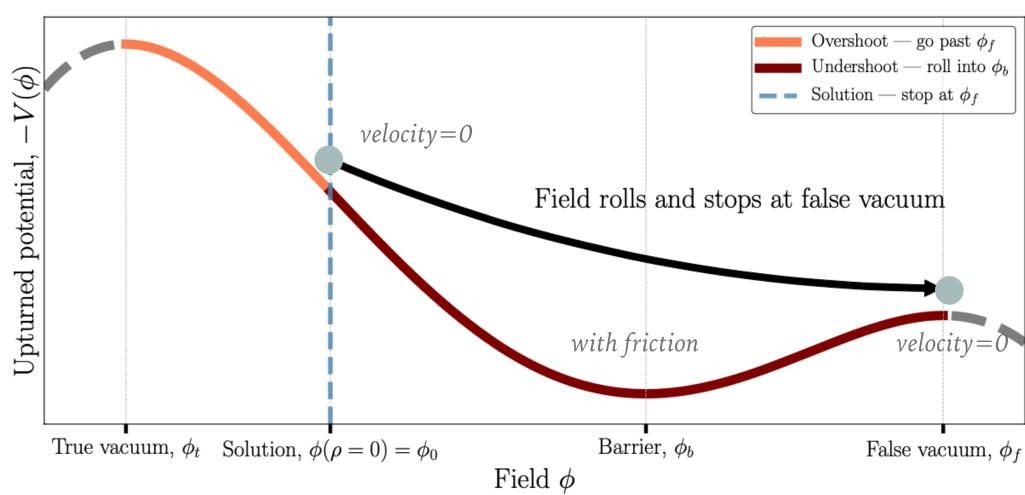
$$\frac{\mathrm{d}^2 \phi(\rho)}{\mathrm{d}\rho^2} + \frac{\alpha}{\rho} \frac{\mathrm{d}\phi(\rho)}{\mathrm{d}\rho} = \Delta V(\phi)$$

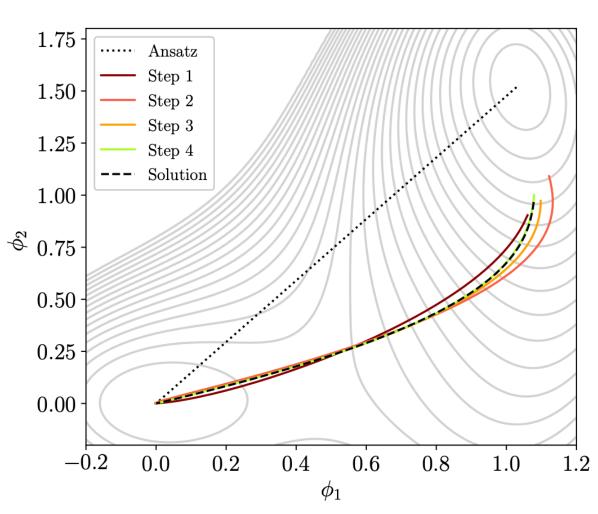
with boundary conditions

$$\frac{\mathrm{d}\phi(\rho)}{\mathrm{d}\rho} \bigg|_{\rho=0} = 0, \quad \frac{\mathrm{d}\phi(\rho)}{\mathrm{d}\rho} \bigg|_{\rho=\infty} = 0$$

$$\phi(\rho \to \infty) = \phi_f$$

The under/over-shooting method





From Bubbleprofiler

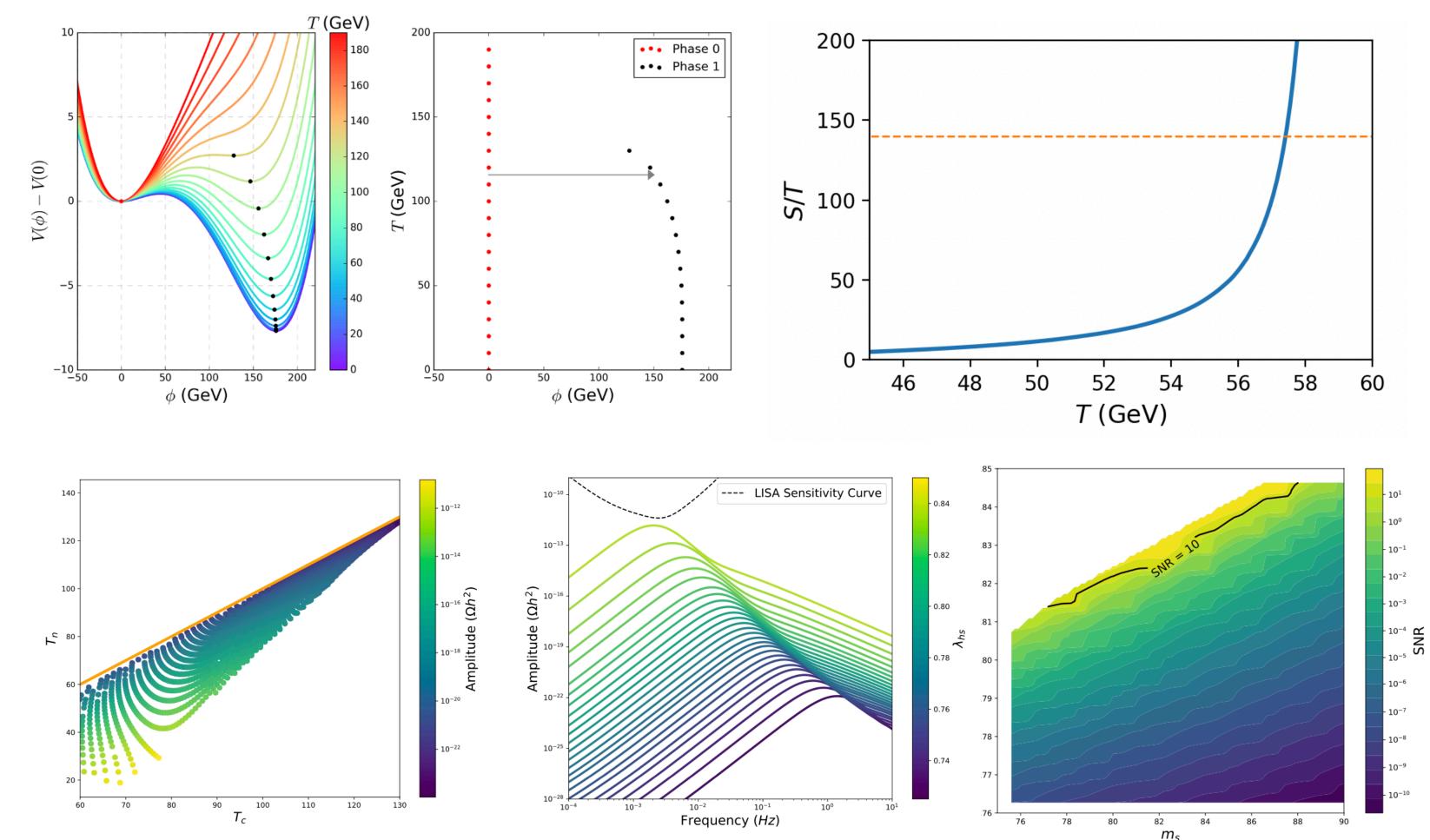
| | Action | | | Time (s) | | | |
|----------|--------|---------------------|------|----------|---------------------|----------|--|
| # fields | BP | CT | AB | BP | CT | AB | |
| 1 | 54.1 | 52.6 | 52.4 | 0.051 | 0.066 | 1.285 | |
| 2 | 20.8 | 21.1 | 20.8 | 0.479 | 0.352 | 7.473 | |
| 3 | 22.0 | 22.0 | 22.0 | 0.964 | 0.215 | 25.209 | |
| 4 | 55.9 | 56.4 | 55.9 | 1.378 | 0.255 | 54.258 | |
| 5 | 16.3 | 16.3 | 16.3 | 2.958 | 0.367 | 305.531 | |
| 6 | 24.5 | 24.5 | 24.4 | 4.853 | 0.337 | 830.449 | |
| 7 | 36.7 | 36.6 | 36.7 | 6.754 | 0.375 | 1430.892 | |
| 8 | 46.0 | 46.0 | 46.0 | 10.014 | 0.409 | 1805.713 | |

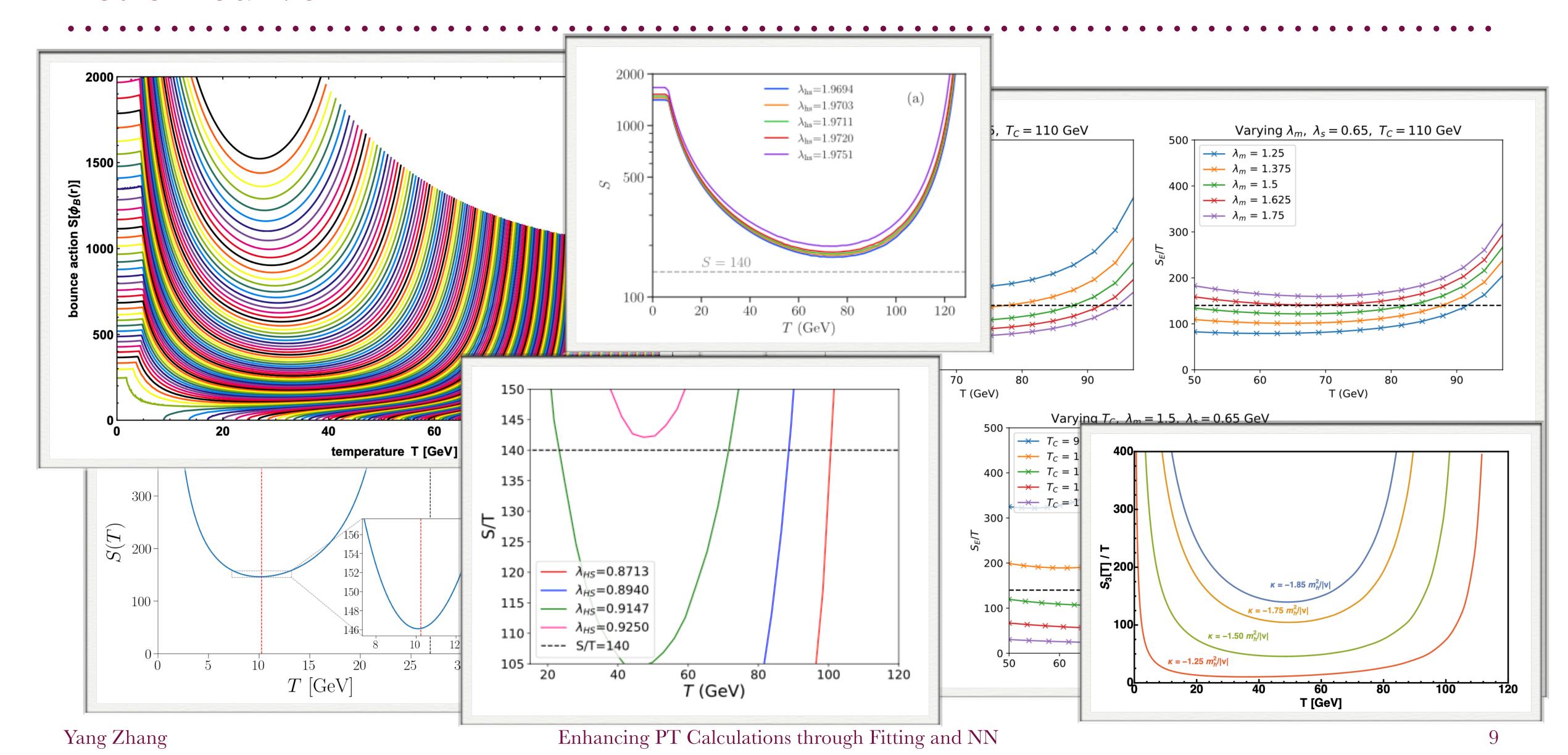
Bounce action

➤ Tools for calculating bounce action

- CosmoTransitions
- PhaseTracer2
- BSMPT3
- PT2GWFinder
- AnyBubble
- BubbleProfiler
- FindBounce
- SimpleBounce

• • • • • •





Model GW Parameters Potential Action at a fixed T Model Parameters temperature T [GeV]

Beyond the Standard Model Cocktail, Yann Gouttenoire, arXiv:2207.01633

For a polynomial potential of the type

$$V(\phi) = a\phi^2 - b\phi^3 + \frac{\lambda}{4}\phi^4$$

> Semi-analytical approximations of the bounce action:

$$S_3 = \frac{8\pi b}{\lambda^{3/2}} \frac{8\sqrt{\delta}}{81(2-\delta)^2} (\beta_1 \delta + \beta_2 \delta^2 + \beta_3 \delta^3), \qquad S_3 = \frac{13.72 a^{3/2}}{b^2} f(\delta/2),$$

$$\delta \equiv \frac{2\lambda a}{b^2},$$

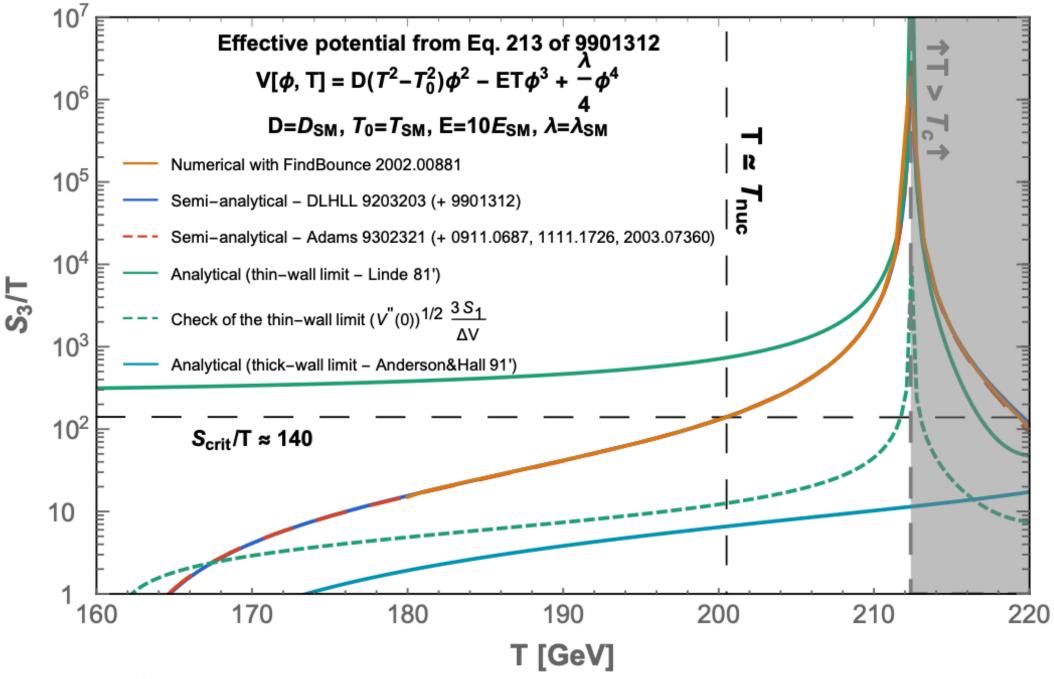
$$\beta_1 = 8.2938$$
, $\beta_2 = -5.5330$, $\beta_3 = 0.8180$

$$\alpha_1 = 13.832$$
, $\alpha_2 = -10.819$, $\alpha_3 = 2.0765$

$$S_3 = \frac{13.72 a^{3/2}}{b^2} f(\delta/2),$$

$$f(x) = 1 + \frac{x}{4} \left[1 + \frac{2.4}{1 - x} + \frac{0.26}{(1 - x)^2} \right]$$

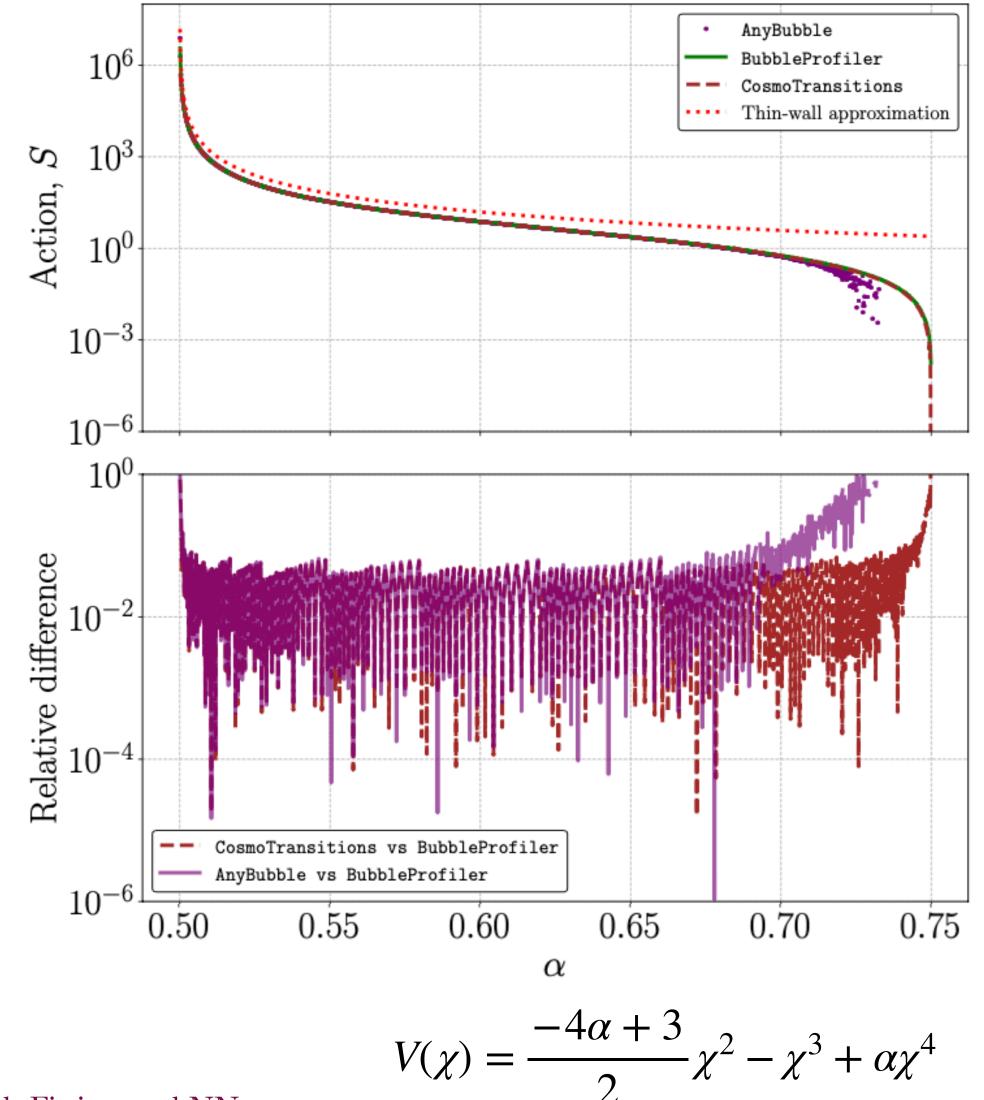
Different methods for computing the O_3 bounce action



arXiv:2404.17632, Marco Matteini, Miha Nemevšek, Yutaro Shoji, Lorenzo Ubaldi

$$\beta_{1} = 8.2938, \quad \beta_{2} = -5.5330, \quad \beta_{3} = 0.8180, \quad f(x) = 1 + \frac{x}{4} \left[1 + \frac{2.4}{1 - x} + \frac{0.26}{(1 - x)^{2}} \right] \qquad V = \frac{1}{2} m^{2} \phi^{2} + \eta \phi^{3} + \frac{\lambda}{8} \phi^{4}, \quad \varepsilon_{\alpha} \equiv 1 - \lambda \frac{m^{2}}{4\eta^{2}}, \quad \alpha_{1} = 13.832, \quad \alpha_{2} = -10.819, \quad \alpha_{3} = 2.0765.$$

- ➤ In general cases, especially for highdimensional scenarios, the action curve has no analytical expression.
- ➤ Meanwhile, there are two main issues with numerical calculation of the action:
 - unavoidable numerical errors
 - excessive computational time
- They will lead to
 - \bullet large error on the β parameter
 - ullet extremely slow to get the $T_{
 m P}$



Bubbleprofiler

Transition parameters

➤ Bubble nucleation rate

$$\frac{\Gamma}{V} = A(T)e^{-S_E/T}[1 + \mathcal{O}(\hbar)]$$

➤ Nucleation temperature

$$N(T) = \int_{T_{\text{nuc}}}^{T_{\text{tra}}} \frac{dT}{T} \frac{\Gamma(T)}{H(T)^4} = 1$$

➤ Percolation temperature

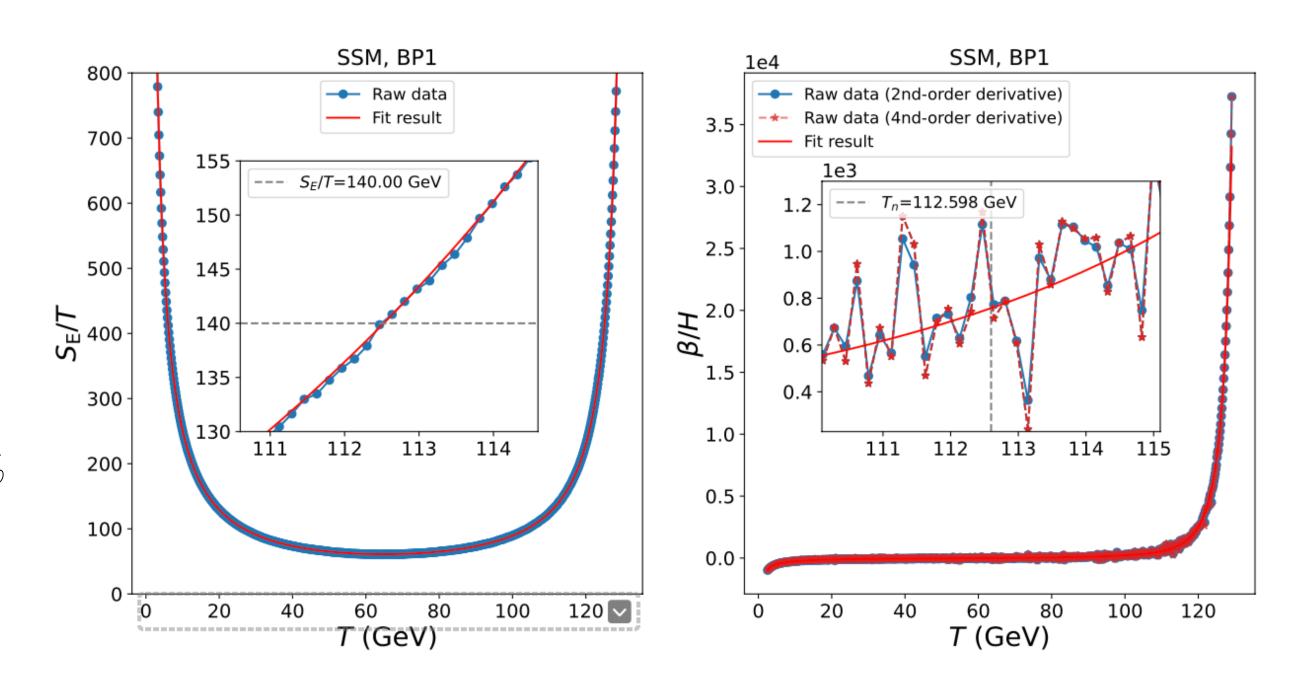
$$P(T_{\text{per}}) = \exp \left[-\frac{64\pi}{3} \xi^4 \int_{T_{\text{per}}}^{T_{\text{tra}}} dT' \frac{\Gamma(T')}{T'^6} \left(\frac{1}{T_{\text{per}}} - \frac{1}{T'} \right)^3 \right] = 70 \%$$

➤ Inverse duration of the phase transition

$$\beta(T) = \frac{d}{dt} \left(\frac{S_E}{T} \right) = TH(T) \frac{d}{dT} \left(\frac{S_E}{T} \right)$$

➤ Ratio of latent heat to radiation density

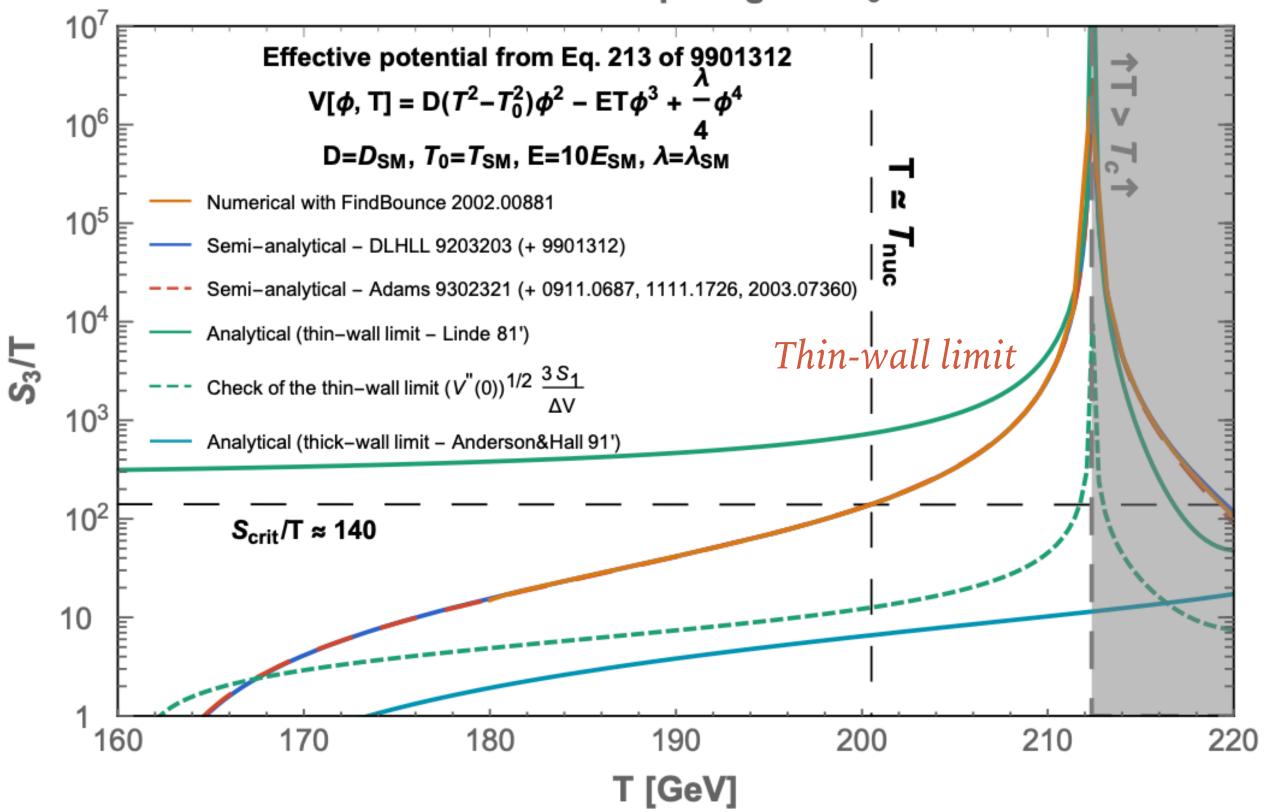
$$\alpha = \frac{D\theta}{\pi^2 g_* T_*^4 / 30} = \frac{1}{\pi^2 g_* T_*^4 / 30} \left(V(\phi) - \frac{T}{4} \frac{\partial V(\phi, T)}{\partial T} \right) \Big|_{\phi_t}^{\phi_f}$$

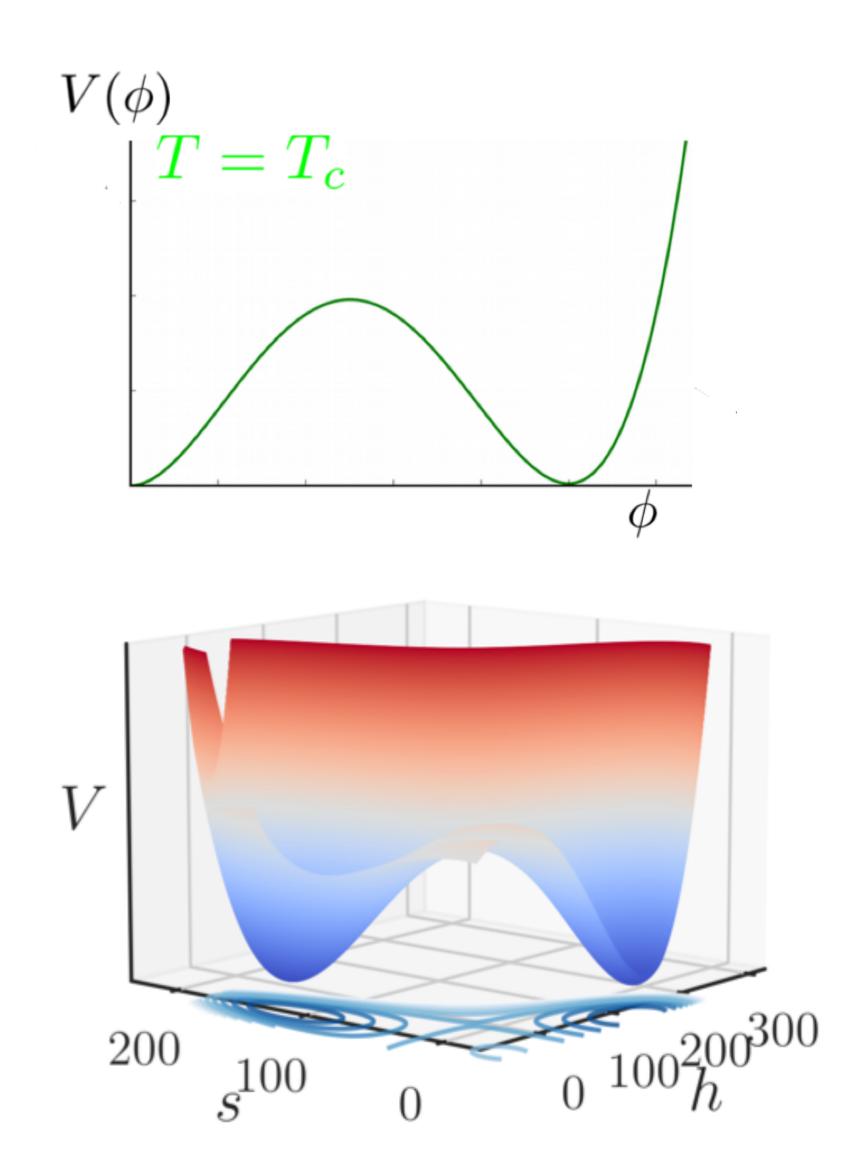


$$f'(x) \approx \frac{f(x+h) - f(x-h)}{2h}$$
, (2nd – order)
 $f'(x) \approx \frac{-f(x+2h) + 8f(x+h) - 8f(x-h) + f(x-2h)}{12h}$, (4nd – order)

The action curve always diverges at $T = T_C$.

Different methods for computing the O_3 bounce action





This divergent behavior is independent of the dimensionality of the potential

► In the thin-wall approximation, the potential difference $\epsilon = V_{\text{eff}}(\phi_f; T) - V_{\text{eff}}(\phi_t; T)$ is much smaller than the height of the barrier, so we neglect the viscous damping term in the bounce equation:

$$\frac{\mathrm{d}^2 \phi}{\mathrm{d}r^2} + \frac{2}{r} \frac{\mathrm{d}\phi}{\mathrm{d}r} = \frac{\partial V_{\mathrm{eff}}(\phi; T)}{\partial \phi} \rightarrow \frac{\mathrm{d}^2 \phi}{\mathrm{d}r^2} = \frac{\partial V_{\mathrm{eff}}(\phi; T)}{\partial \phi} \rightarrow \frac{d\phi}{dr} = \sqrt{2V_{\mathrm{eff}}(\phi; T)}.$$

Define the surface tension of the bubble

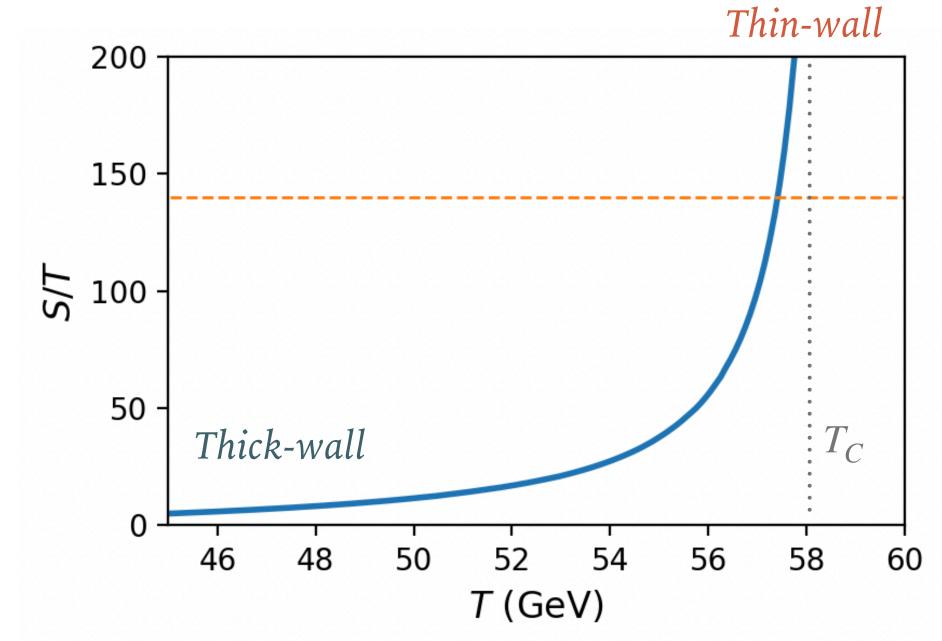
$$\sigma = \int_0^\infty dr \left[\frac{1}{2} \left(\frac{d\phi}{dr} \right)^2 + V_{\text{eff}}(\phi; T) \right] = \int_{\phi_t}^{\phi_f} d\phi \sqrt{2V_{\text{eff}}(\phi; T)},$$

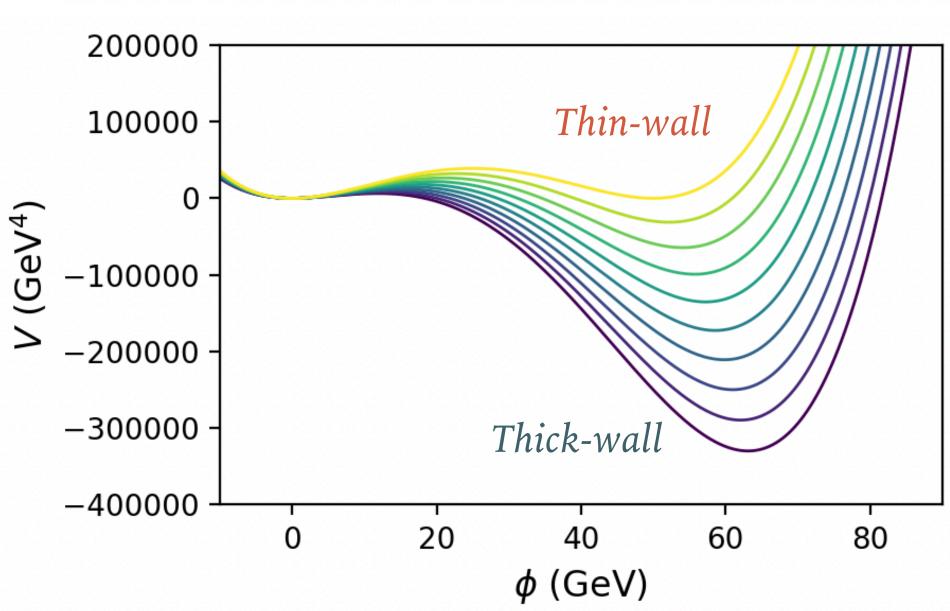
$$S_E = 4\pi \int_0^{+\infty} r^2 dr \left[\frac{1}{2} \left(\frac{\partial \phi}{\partial r} \right)^2 + V_{\text{eff}}(\phi; T) \right] = 4\pi R^2 \sigma - \frac{4}{3}\pi R^3 \epsilon.$$

where R is the radius of the critical bubble and can be calculated by minimization of S_E , $R=\frac{2\sigma}{2}$. As

$$\epsilon = V_{\text{eff}}(\phi_f; T) - V_{\text{eff}}(\phi_t; T) = \left(\frac{\partial V_{\text{eff}}(\phi_f; T)}{\partial T} \middle|_{T = T_c} - \frac{\partial V_{\text{eff}}(\phi_t; T)}{\partial T} \middle|_{T = T_c}\right) (T - T_c)$$

$$S_E = \frac{16\pi\sigma^3}{3\epsilon^2} \propto \frac{1}{(T - T_c)^2}$$



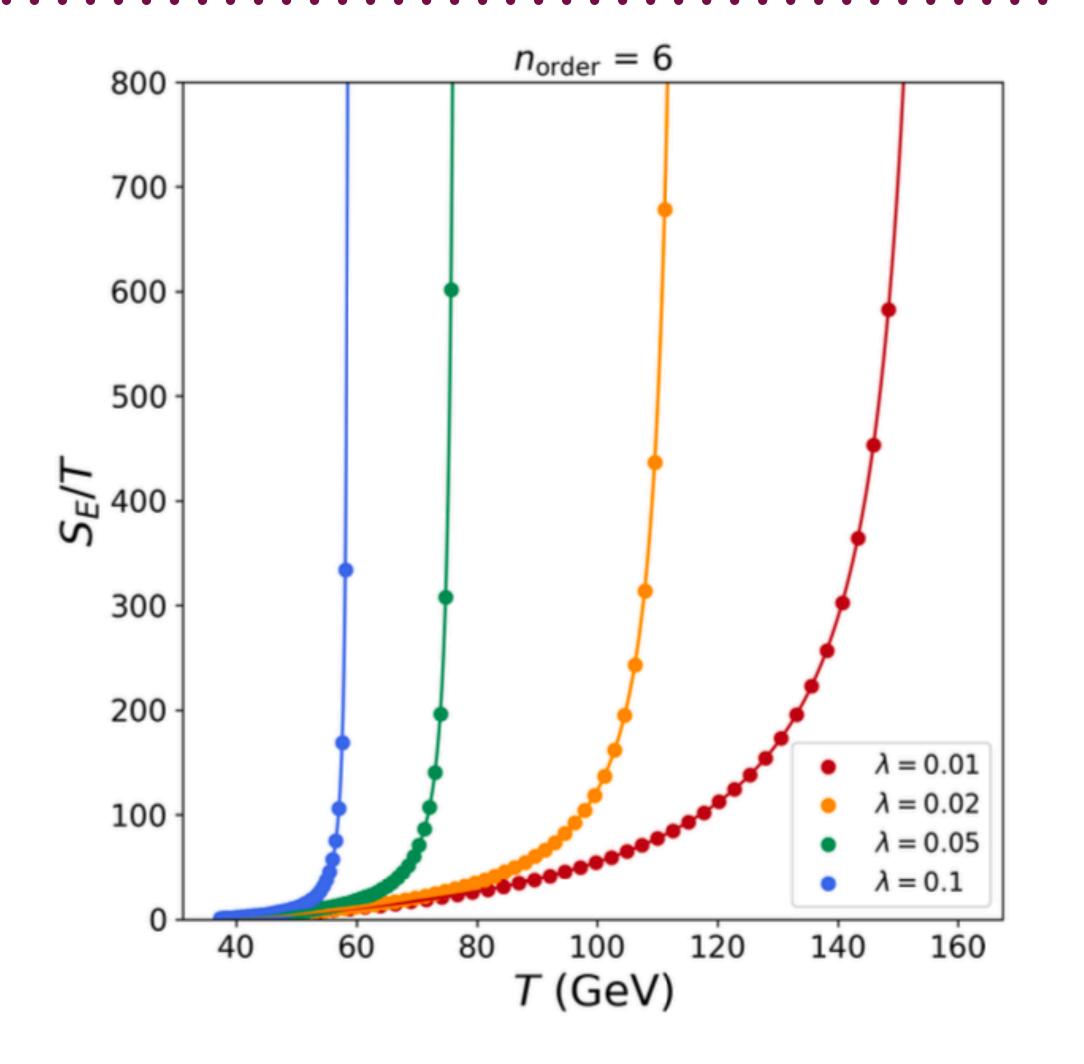


Therefore, it is reasonable to use the polynomial fitting formula

$$S_E = \frac{1}{(T - T_c)^2} \sum_{i=0}^{n_{\text{order}}} q_i T^i$$

➤ And one can get the expression for the inverse phase transition duration time

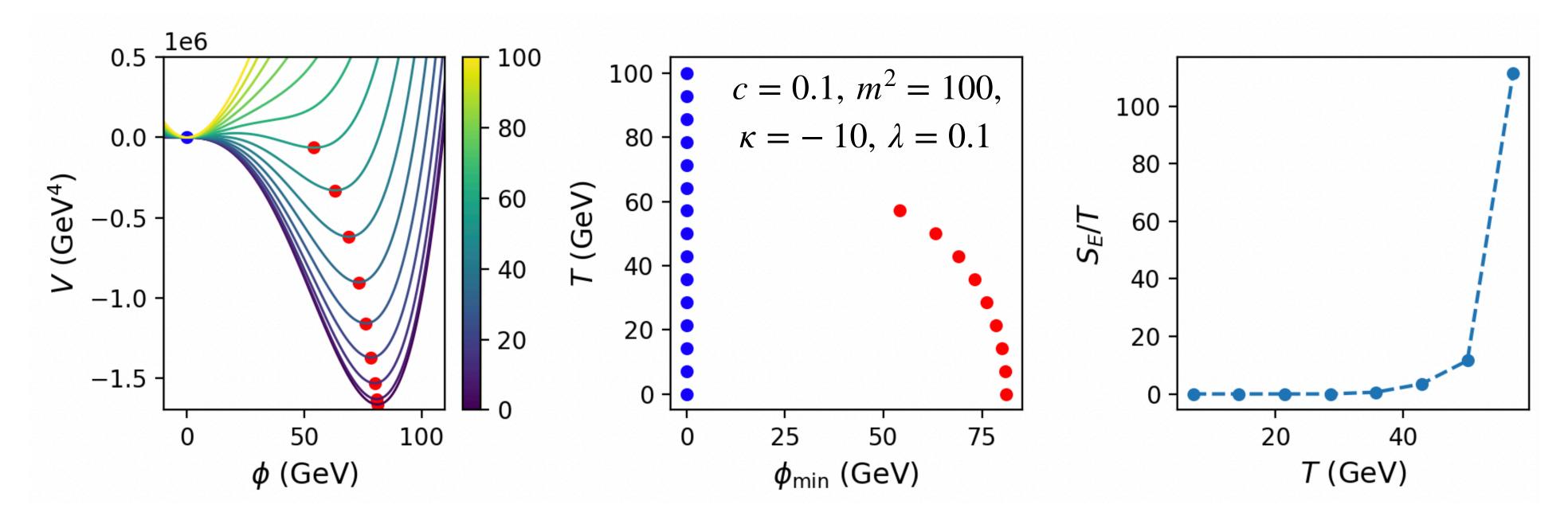
$$\frac{\beta}{H} = \frac{1}{T(T - T_c)^3} \left[\sum_{i=1}^{n_{\text{order}}} q_i i T^i (T - T_c) - \sum_{i=0}^{n_{\text{order}}} q_i T^i (3T - T_c) \right]$$

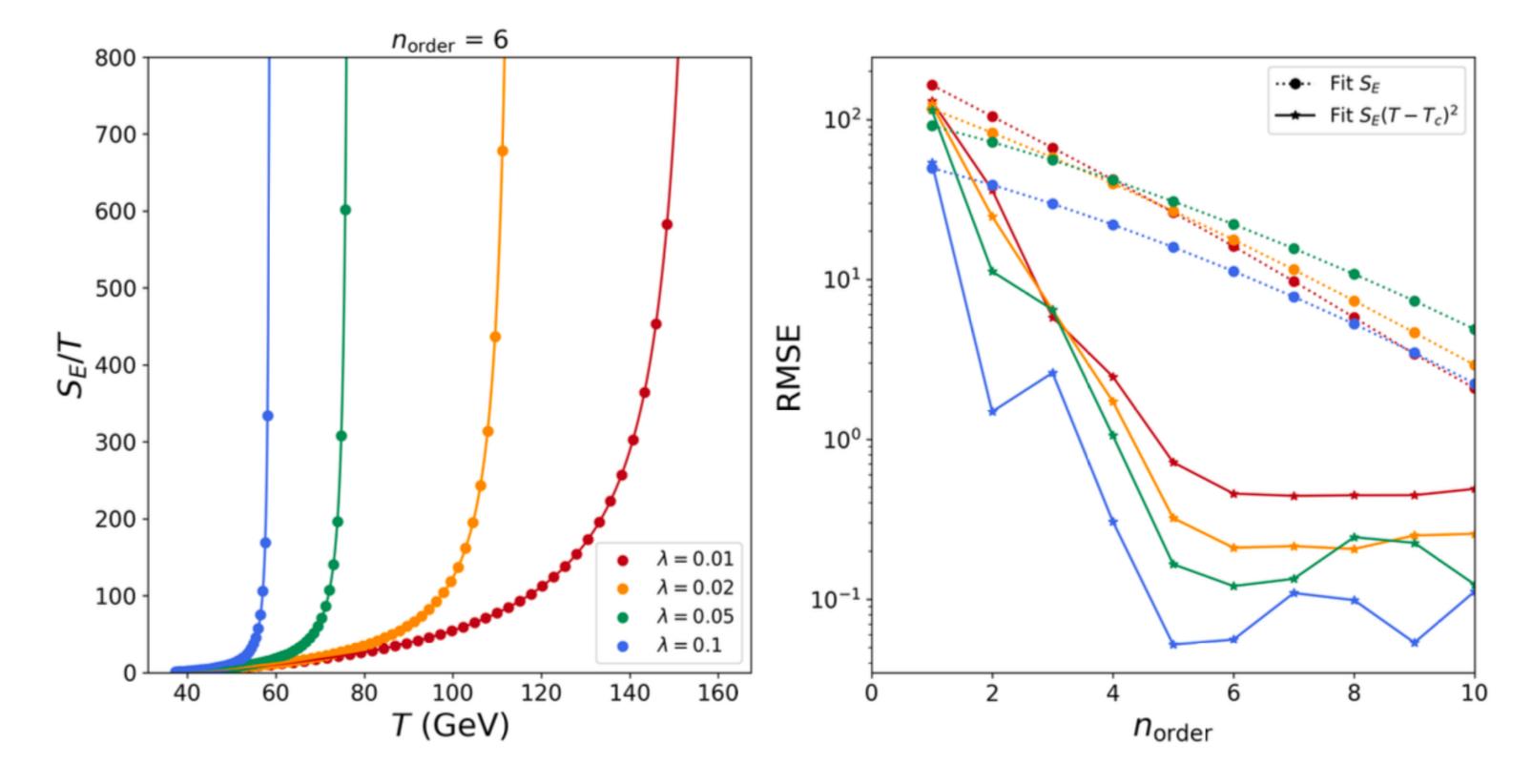


➤ To validate the polynomial fitting approach, we utilize a 1D toy model in which the action can be accurately computed.

$$V_{\text{eff}}(\phi; T) = (cT^2 - m^2)\phi^2 + \kappa\phi^3 + \lambda\phi^4$$

➤ It mimics a simple model that includes high-temperature corrections.

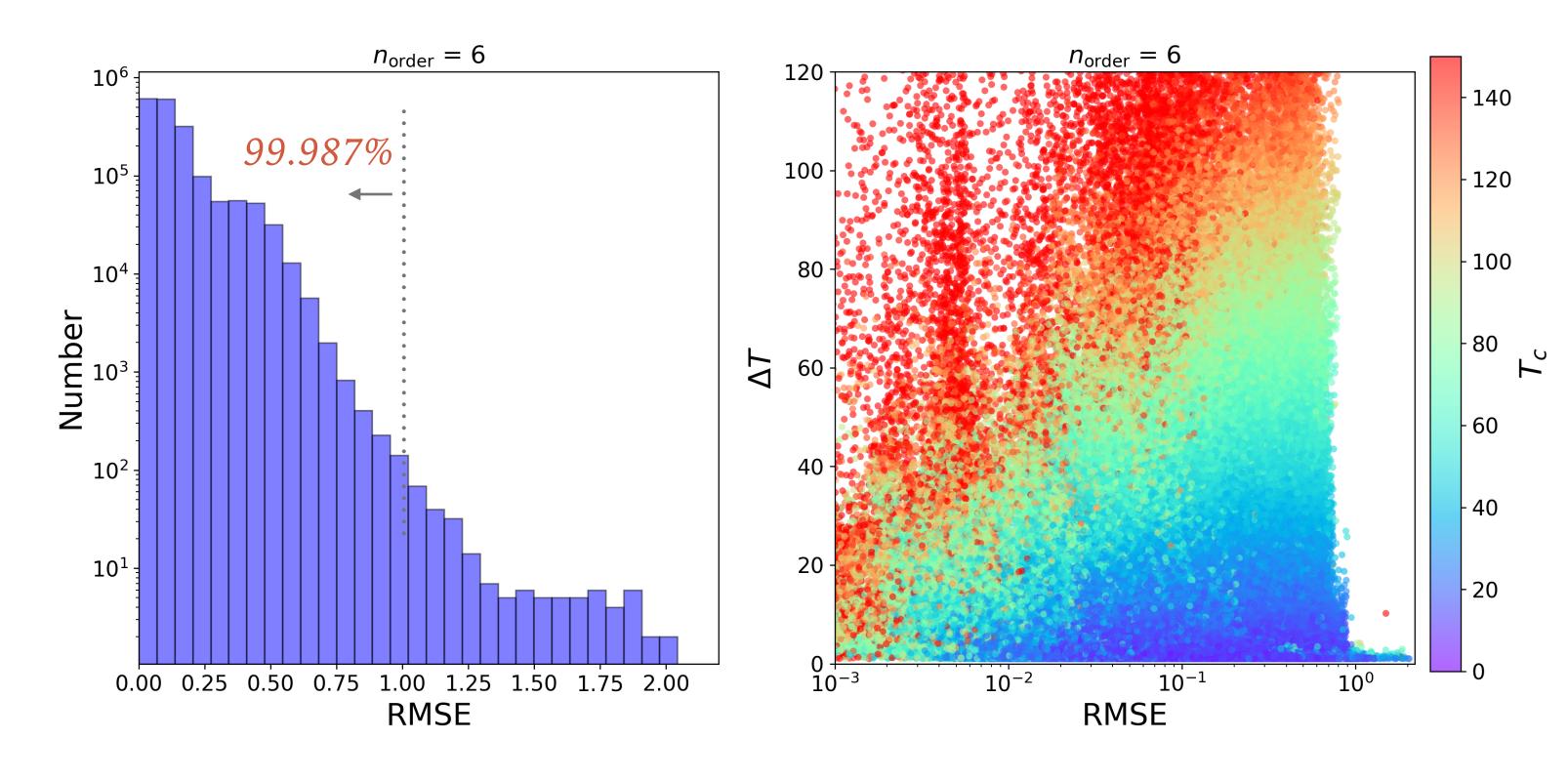




- The fitting results align very well with the raw data.
- ➤ We utilize the root mean square error (RMSE) to quantify the degree of agreement,

RMSE =
$$\sqrt{\frac{1}{n} \sum_{i=1}^{n} \left(\frac{S_E(T_i)}{T_i} - \frac{\hat{S}_E(T_i)}{T_i} \right)^2}$$

➤ With the factor of $(T - T_C)^2$, the MSE drops quickly with the increasing of n_{order} .



➤ A random scan in

$$c \in [0,2], m^2 \in [0,200],$$

$$\lambda \in [0,2], \kappa \in [-30,0].$$

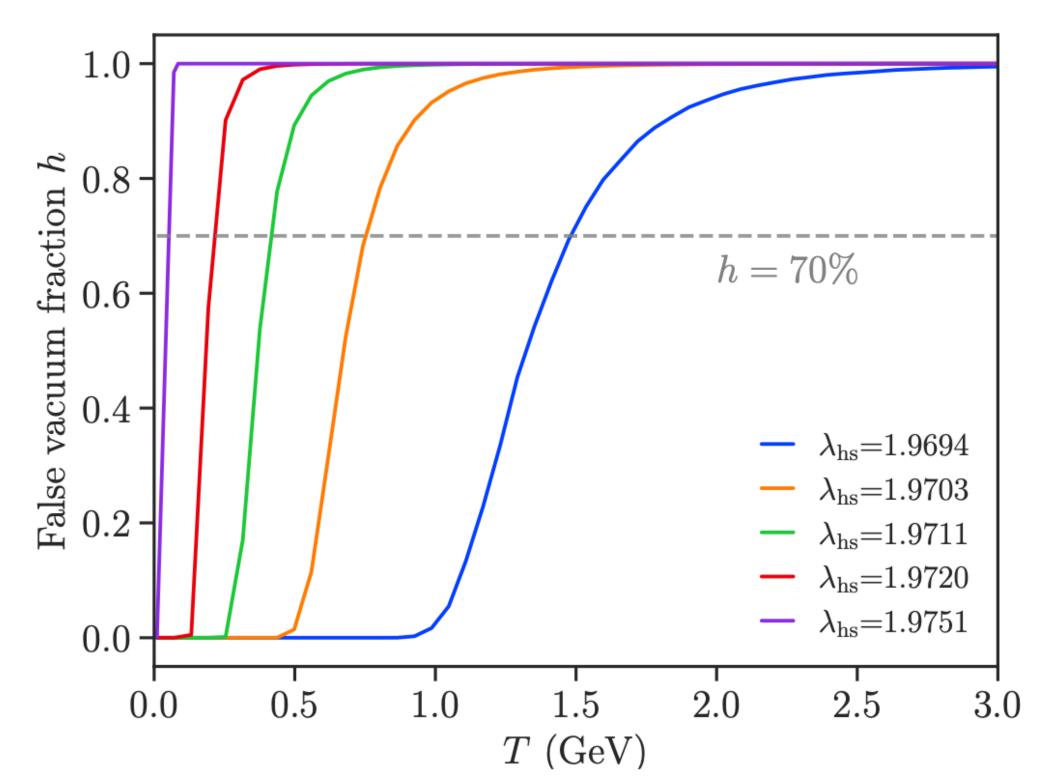
➤ The majority of samples exhibit an MSE below 1, with a maximum value of 4.2.

$$S_E/T_{\rm nuc} \simeq 140$$

The RMSE exceeds 1 only for $\Delta T < 2$ GeV

 \triangleright With the action curve function, we can calculate the percolation temperature T_P fairly quick.

$$P(T_{\rm P}) = \exp\left[-\frac{64\pi}{3}\xi^4 \int_{T_{\rm per}}^{T_{\rm tra}} dT' \frac{\Gamma(T')}{T'^6} \left(\frac{1}{T_{\rm P}} - \frac{1}{T'}\right)^3\right] = 70\%$$



| | | c | m^2 | κ | λ | T_C | T_N | T_P | Time |
|----------------|--------------------|-----|-------|----------|------|-------|-------|-------|---------------------|
| BP1 | without action fit | 0.1 | 100 | -10 | 0.01 | 161.2 | 125.7 | 122.0 | 5.60s |
| | with action fit | | | | | 161.2 | 125.7 | 122.0 | 0.04s |
| BP2 | without action fit | 0.1 | 100 | -10 | 0.02 | 116.2 | 101.5 | 99.6 | 7.54s |
| | with action fit | | | | | 116.2 | 101.4 | 99.6 | 0.05s |
| $\mathbf{BP3}$ | without action fit | 0.1 | 100 | -10 | 0.05 | 77.5 | 73.0 | 72.4 | 10.8s |
| | with action fit | | | | | 77.5 | 73.0 | 72.4 | 0.05s |
| BP4 | without action fit | 0.1 | 100 | -10 | 0.1 | 59.2 | 57.4 | 57.2 | $\overline{11.46s}$ |
| | with action fit | 0.1 | | | | 59.2 | 57.4 | 57.2 | 0.05s |
| | | | | | | | | | |

- ➤ Without action fit: Each integration step requires multiple evaluations of the action, and we need several integration to determine the temperature corresponding to 70%.
- ➤ With action fit: 30 times evaluations of the action to get the data for fit.

➤ Now we turn to a physical model, the singlet scalar extensions of the Standard Model, which is wildly used in instructional studies of phase transition:

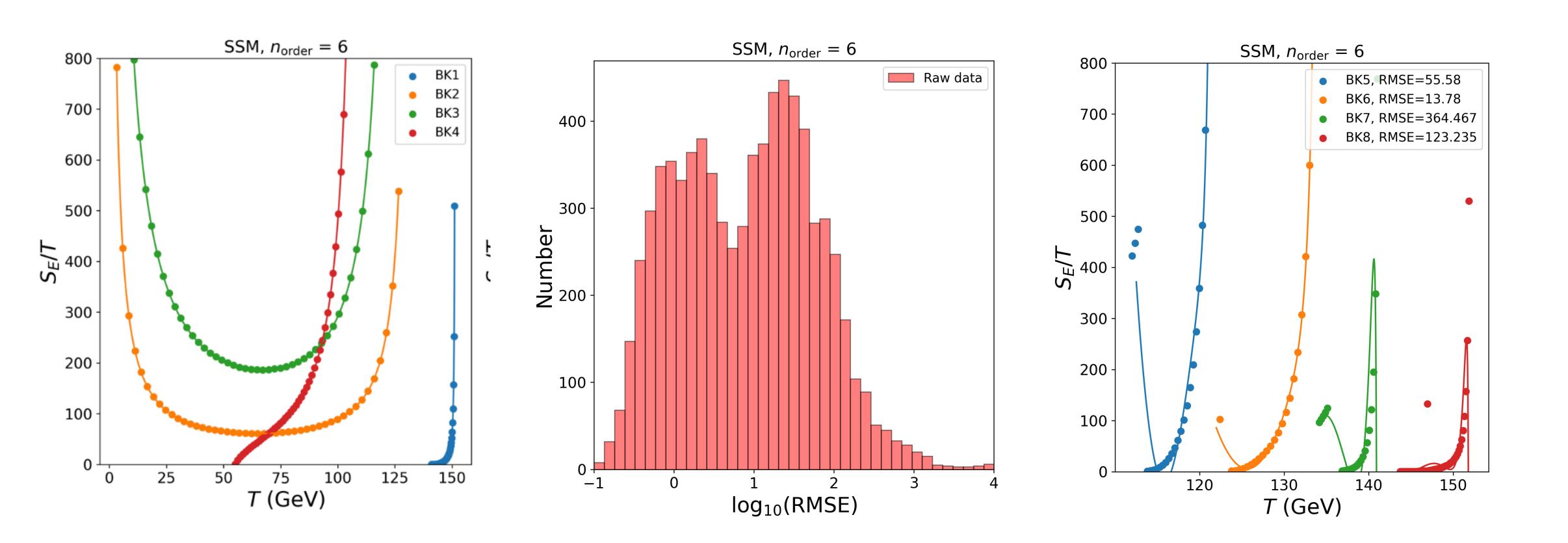
$$V_0(h,s) = -\frac{\mu_H^2}{2}h^2 + \frac{\lambda_H}{4}h^4 - \frac{\mu_S^2}{2}s^2 + \frac{\lambda_S}{4}s^4 + \frac{\lambda_{HS}}{4}h^2s^2$$

$$V_{\text{eff}}(h,s;T) = V_0(h,s) + V_{\text{CW}}(h,s) + V_{\text{CT}}(h,s) + V_{\text{1T}}(h,s;T) + V_{\text{ring}}(h,s;T)$$

We choose the OS-like scheme, the Landau gauge and the Parwani method:

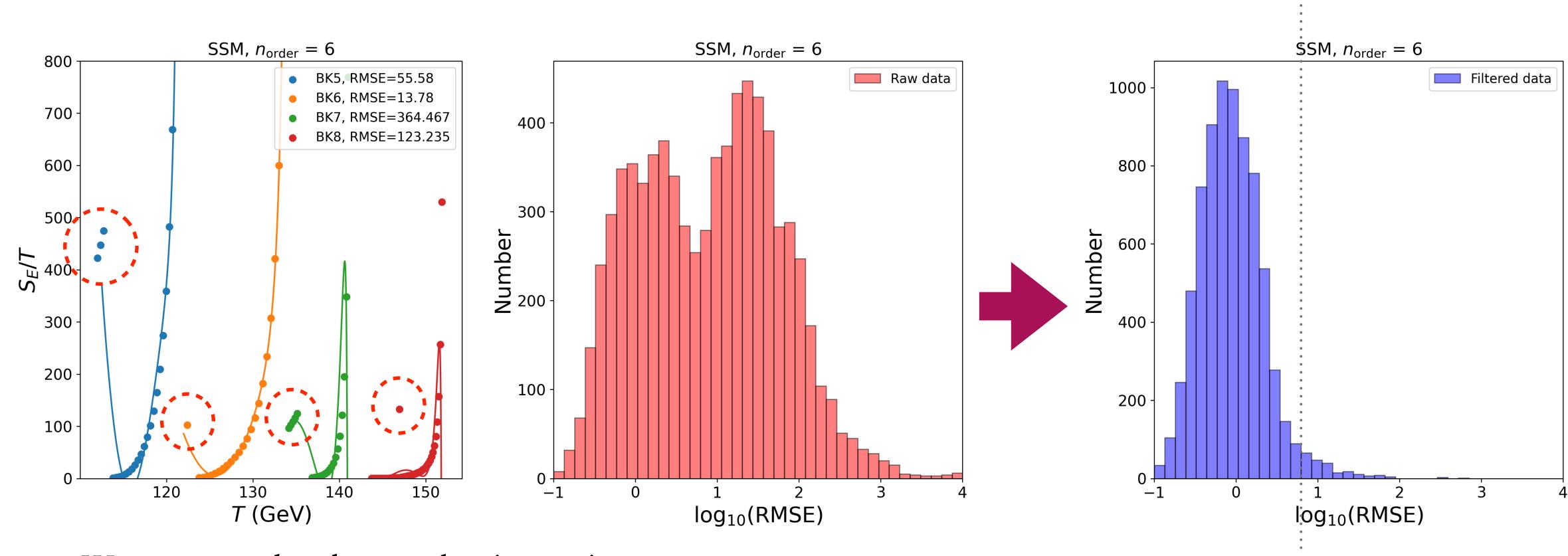
$$V_{\text{CW}}(h,s) + V_{\text{CT}}(h,s) = \sum_{i} (-1)^{s_i} \frac{g_i}{64\pi^2} \left\{ m_i^4(h,s) \left[\log \frac{m_i^2(h,s)}{m_i^2(v_h,v_s)} - \frac{3}{2} \right] + 2m_i^2(h,s) m_i^2(v_h,v_s) \right\}$$

$$V_{1T}(h,s) = \frac{T^4}{2\pi^2} \left[\sum_B g_B J_B \left(\frac{m_B(h,s)}{T} \right) + \sum_F g_F J_F \left(\frac{m_F(h,s)}{T} \right) \right]$$



- For some of the samples, the fitting results also align very well with the calculated action.
- \succ Half of the samples have large RMSE, because of incorrect S_F .





➤ We removes the abnormal points using

$$S(T_i) - S(T_{i-1}) < 0$$
 and $|S(T_i) - S(T_{i-1})| < |S(T_{i-1}) - S(T_{i-2})|$

> As temperature decreases, the action should decrease monotonically, and the rate of decrease should slow.

Calculate one β/H for 10 times with different h

| | h = 1 | h = 0.1 | h = 0.01 | h = 0.001 | Fititng |
|-------------|---------|---------|----------|-----------|---------|
| 1 | 1621.95 | 1573.90 | 1530.38 | 1461.57 | 1500.25 |
| 2 | 1622.96 | 1534.27 | 1915.43 | 901.45 | 1500.24 |
| 3 | 1622.34 | 1533.25 | 1505.82 | 991.253 | 1499.92 |
| 4 | 1618.56 | 1503.12 | 1278.73 | -1502.6 | 1500.33 |
| 5 | 1621.08 | 1466.95 | 1202.39 | 5570.78 | 1500.12 |
| 6 | 1579.43 | 1503.90 | 1635.98 | 845.765 | 1500.00 |
| 7 | 1622.35 | 1506.99 | 1691.52 | 1728.72 | 1499.98 |
| 8 | 1623.08 | 1517.53 | 1077.92 | 2533.38 | 1500.44 |
| 9 | 1620.85 | 1503.17 | 1171.85 | -380.135 | 1500.08 |
| 10 | 1622.14 | 1523.33 | 1812.96 | 1863.13 | 1500.13 |
| Mean | 1617.47 | 1516.64 | 1482.30 | 1401.33 | 1500.15 |
| Uncertainty | 12.74 | 26.46 | 273.30 | 1769.83 | 0.15 |

$$\frac{\mathrm{d}f(x)}{\mathrm{d}x} = \frac{f(x+h) - f(x-h)}{2h}$$

- ➤ It is doable to model the action curve using polynomial fitting.
- ➤ For one benchmark point, we only need to calculate action about 30 times, then we can
 - \bullet Precisely calculate the β
 - Improve the calculation of A(T)
 - \bullet Rapidly calculate the $T_{\rm nuc}$ and $T_{\rm per}$.

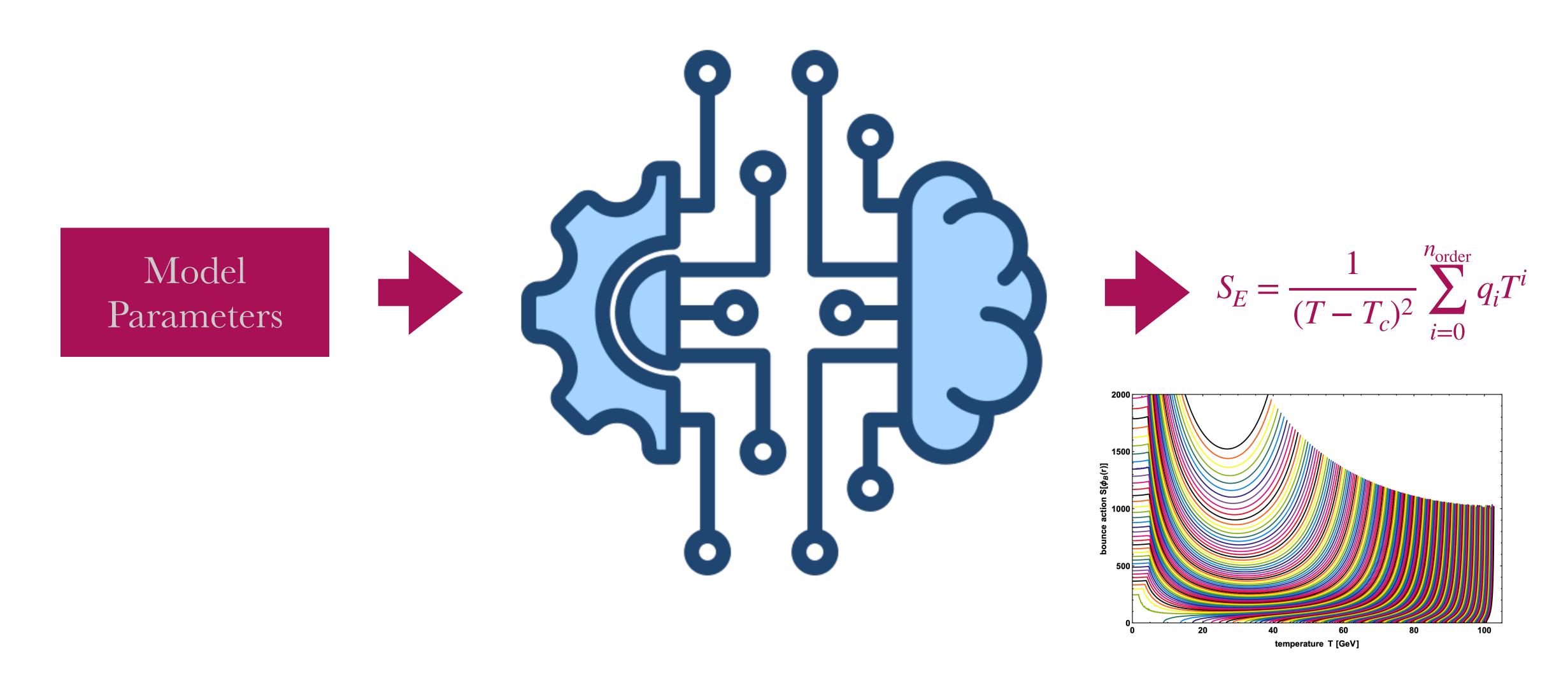
Action curve fit in PhaseTracer2

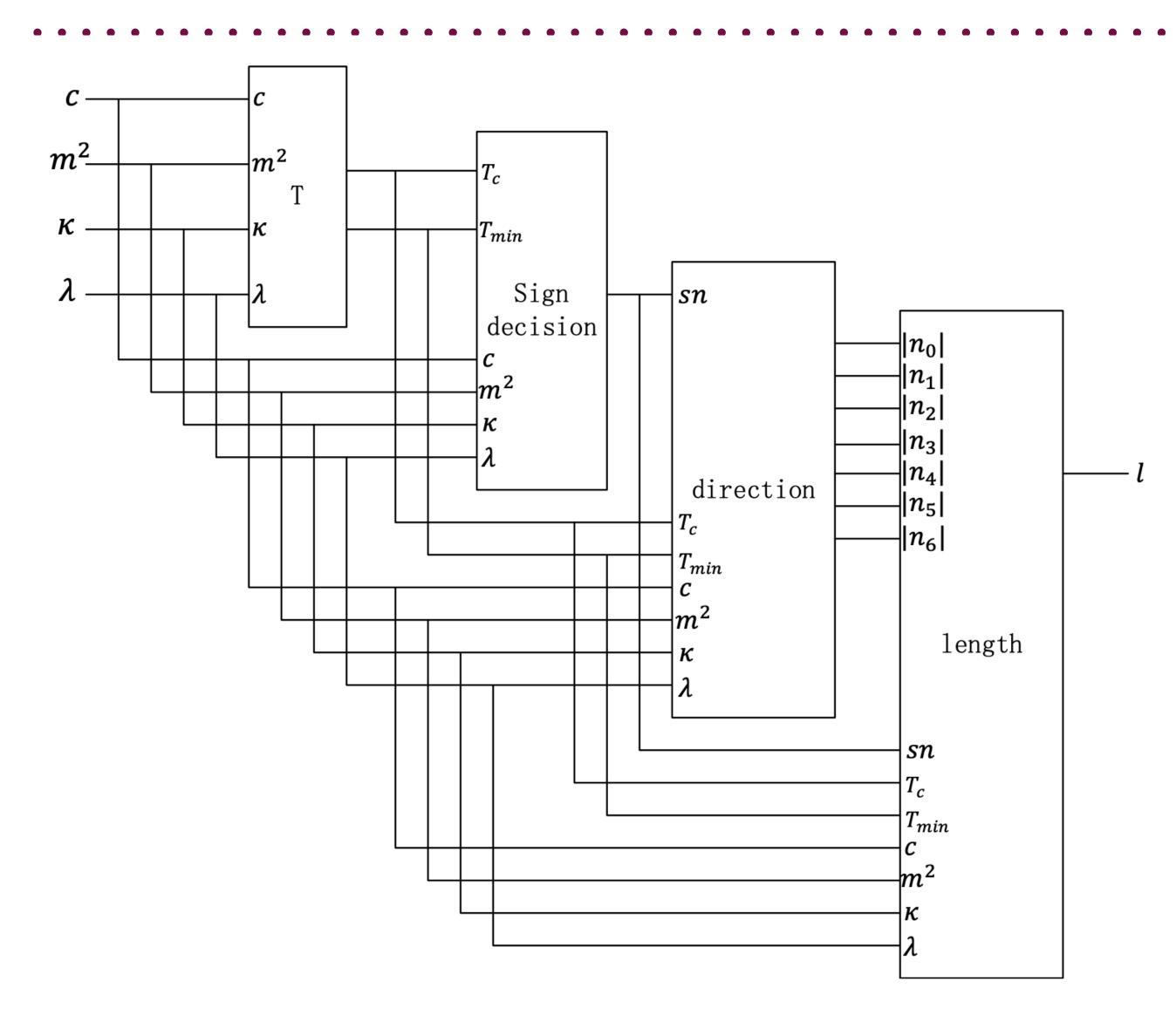
➤ We have added the action fitting in the "fit_action" branch of PhaseTracer2, and will merge to the master branch soon.

```
PhaseTracer /
                      PhaseTracer
                       1 Pull requests 1
                                          Discussions
         • Issues 5
<> Code
 PhaseTracer (Public)
                  우 9 Branches ♥ 6 Tags

property fit_action ▼
                                          Q Go to file
  This branch is 18 commits ahead of master.
// Make TransitionFinder object and find the transitions
PhaseTracer::TransitionFinder tf(pf, ac);
tf.set_fit_action_curve(true);
tf.set_calculate_percolation(true);
tf.find_transitions();
std::cout << tf;</pre>
```

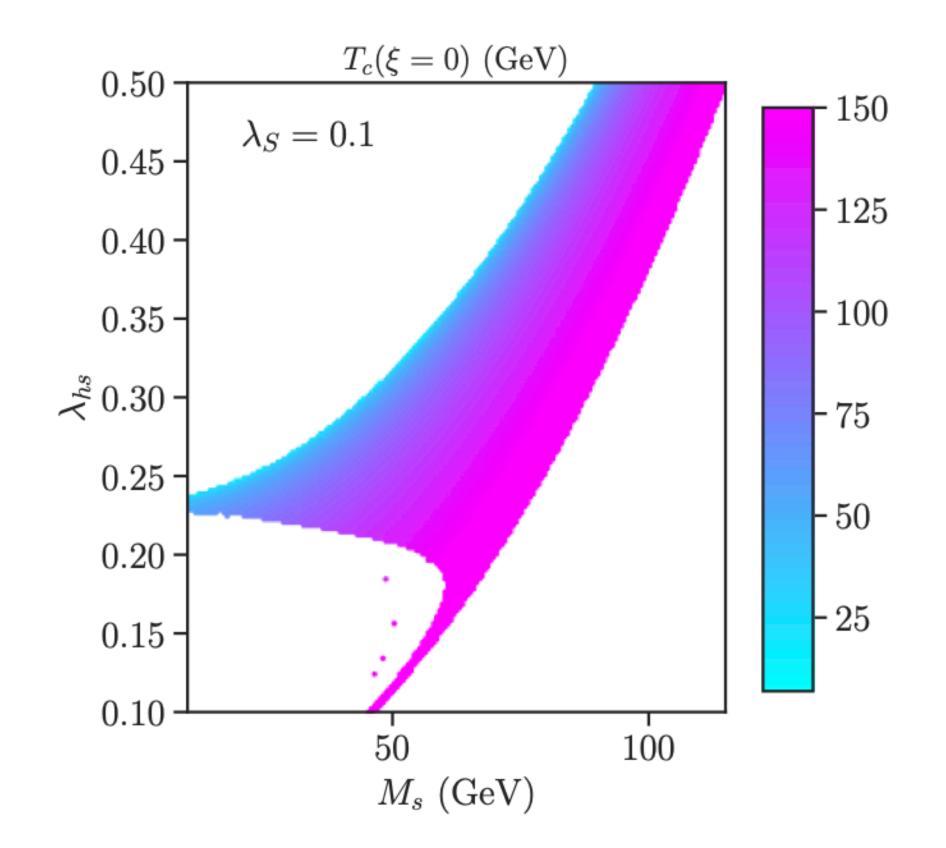
```
PhaseTracer — -zsh — 62×21
=== transition from phase 0 to phase 1 ===
changed = [true]
TC = 59.1608
false vacuum (TC) = [-8.64205e-06]
true vacuum (TC) = [50.0002]
gamma (TC) = 0.845158
delta potential (TC) = 0.00117793
Action curve fitting succeeded with MSE = 0.00145456
TN = 57.4032
false vacuum (TN) = [-9.70809e-06]
true vacuum (TN) = [53.583]
TP = 57.175
transition was not subcritical
=== gravitational wave spectrum generated at T = 57.4032 ===
alpha = 0.00138512
beta over H = 7318.24
peak frequency = 0.124794
peak amplitude = 3.3802e-23
signal to noise ratio for LISA = 6.62623e-13
zy@Yangs-MacBook-Pro-2 PhaseTracer % ./bin/run_1D_test_model
```

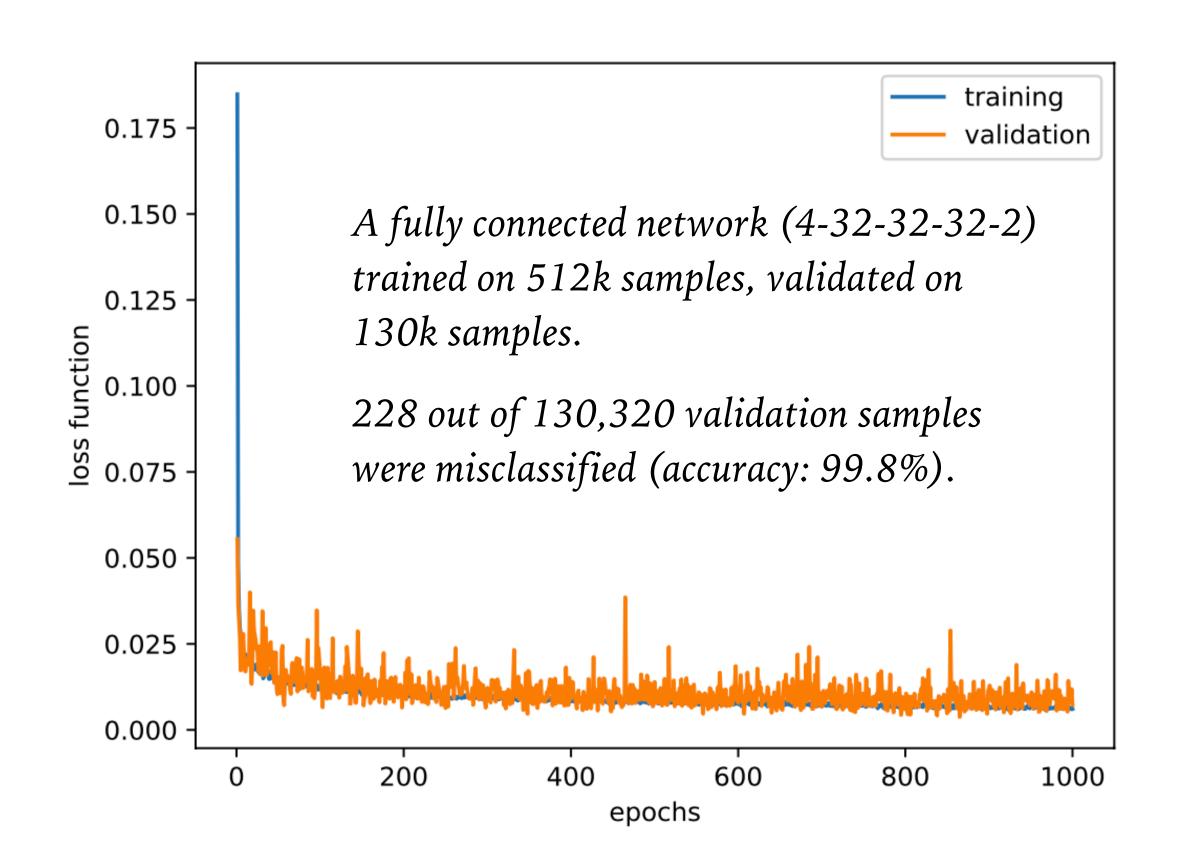




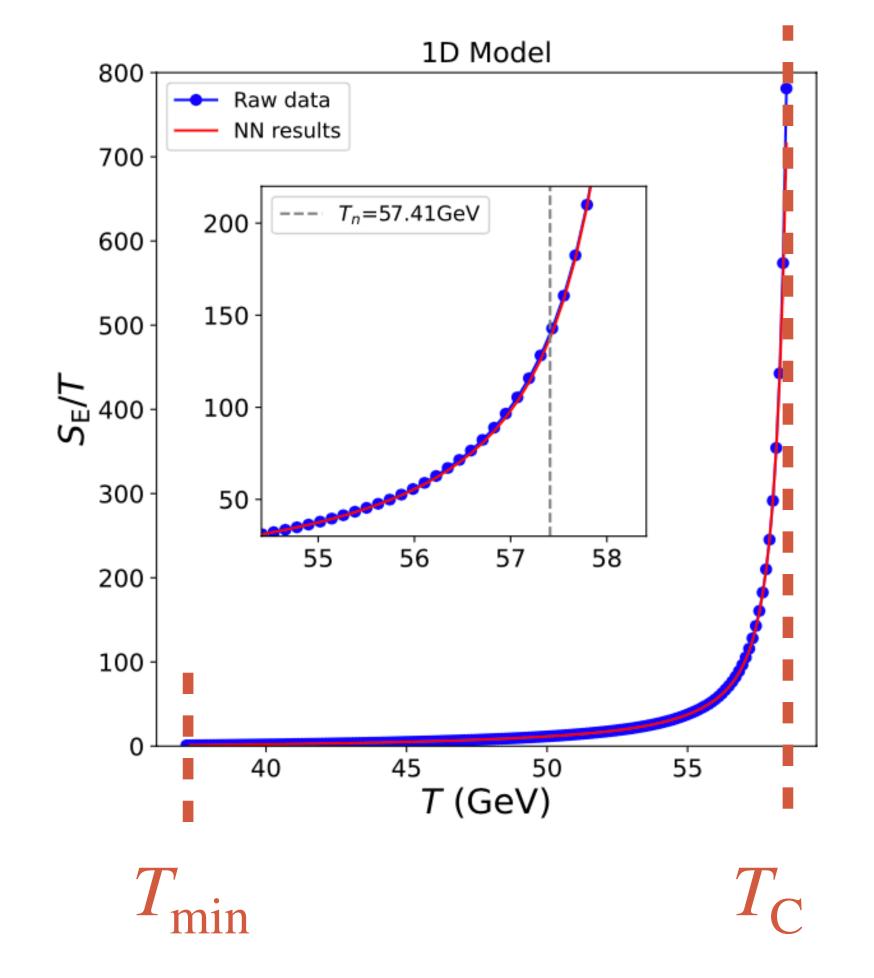
- Step 1: distinguish parameter regions where valid first-order phase transition occurs.
- Step 2: predict the overlap temperature range of the two phases, i.e. T_C and T_{\min} .
- Step 3: predict the polynomial coefficients.
- Step 4: validate the accuracy using a few point in the curve.
- ➤ We employ the conventional fully connected neural network for the first two step, and the Kolmogorov-Arnold Network (KAN) for step 3.

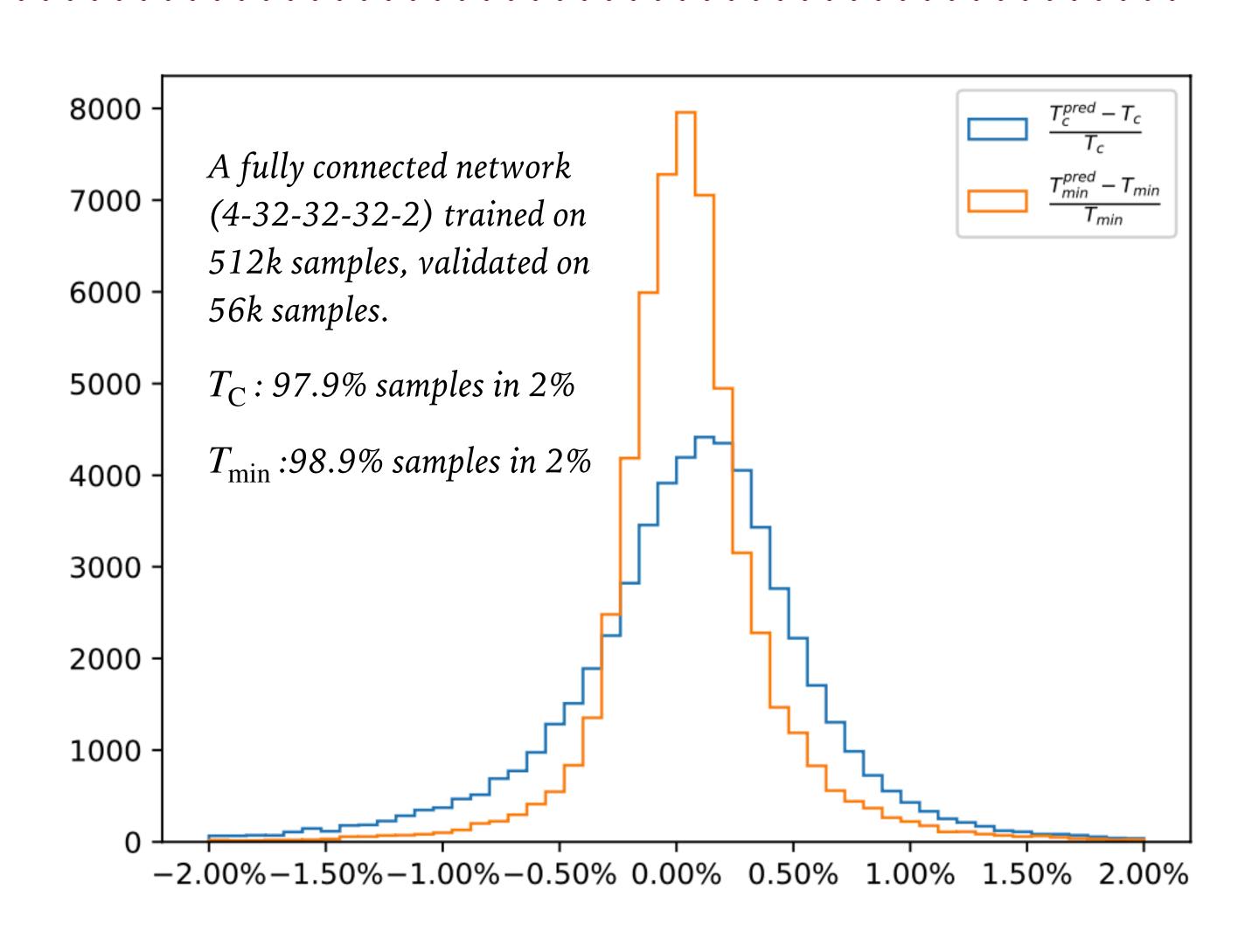
- ➤ Step 1: Distinguish the parameter space that has valid action curve
- ➤ We utilize the 1D toy model to illustrate the performance of machine learning.





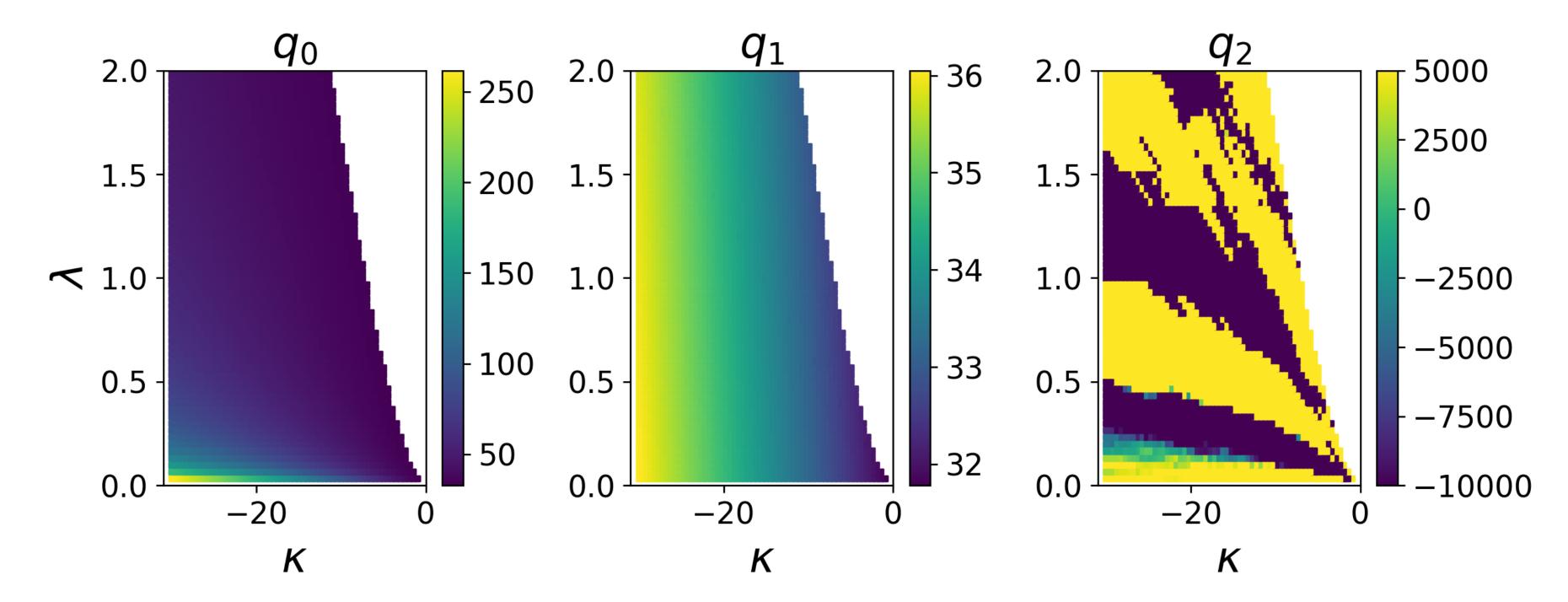
Step 2: Predict $T_{\rm C}$ and $T_{\rm min}$





$$S_E = \frac{1}{(T - T_c)^2} \sum_{i=0}^{n_{\text{order}}} q_i T^i$$

➤ Step 3: Predict the polynomial coefficients



➤ The variation of some coefficients with the input parameters is not smooth, covering a wide range and occasionally switching sign.

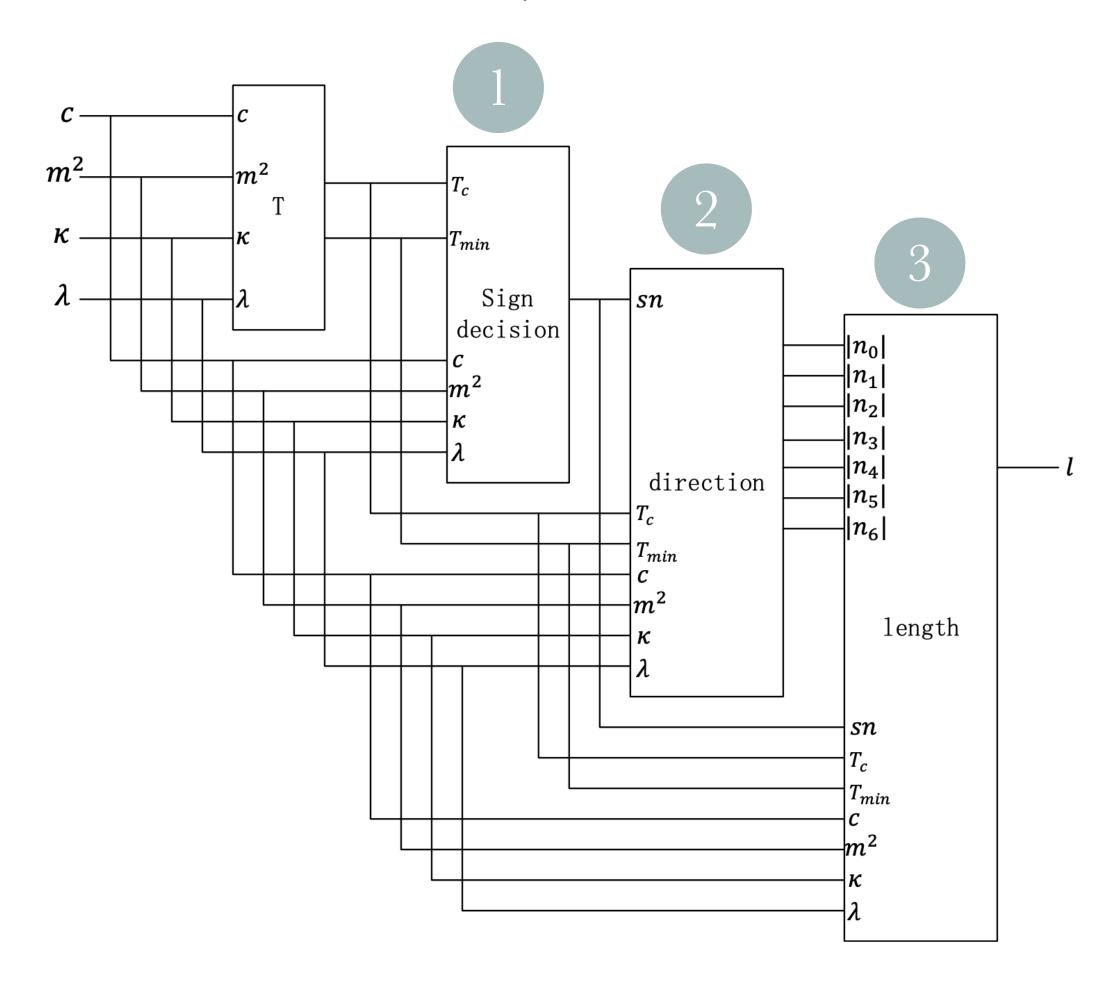
- ➤ Step 3: Predict the polynomial coefficients
 - We have to predict the function as a whole, rather than predicting the coefficients separately.
 - The difference between two functions can be measured using

$$\epsilon = \int_{x_0}^{x_1} dx \left(f_1(x) - f_2(x) \right)^2$$

which is the Euclidean distance in a Hilbert space.

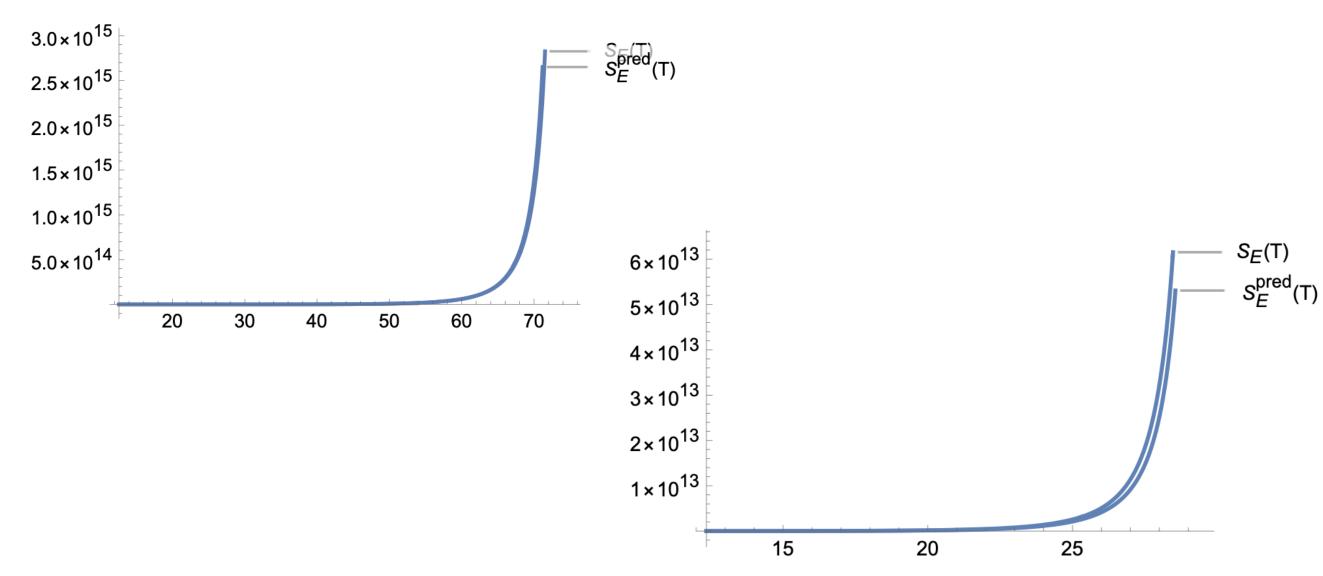
- Thus, we construct an orthogonal function set by subtracting the projection of the original function onto the existing orthogonal basis, following the Gram-Schmidt orthogonalization process.
- We use a fast KAN network with architecture 13-64-64-64-64-64-1

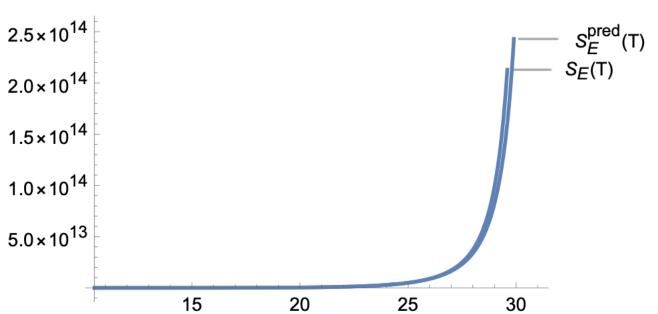
➤ Step 3: Predict the polynomial coefficients



- Network-1: predict the sign of the coefficients
 - → 96.69% accurate
- Network-2: predicts the direction $|n_i|$ of the polynomial in the Hilbert space
 - → 96.59% accurate
- Network-3: predicts the length of the polynomial in the Hilbert space
 - → 71.48% accurate

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Summary

- ➤ Using machine learning to predict action curve function is more practical than predicting isolated action value or observable. It provides the flexibility to perform subsequent calculations derived from the action curve.
- ➤ Neural networks, particularly KANs, show promise in achieving this objective, although their predictive accuracy requires further improvement.
- ► With action curve fitting, we can calculate the T_P and T_N quickly, and the β more accurately. This is independent of machine learning, and can be used in PhaseTracer now.

Thanks!

