

2022年7月31日 - 8月7日

轻子生成 理论框架与最新进展

周也铃, HIAS, 2022-08-02







ICTP-AP International Centre for Theoretical Physics Asia-Pacific 国际理论物理中心-亚太地区

Baryon asymmetry



Parameter(s)	$\Omega_{ m b} h^2$	$\Omega_{ m c} h^2$	$100\theta_{\rm MC}$	H_0	n _s	$\ln(10^{10}A_{\rm s})$
Base ACDM	0.02237 ± 0.00015	0.1200 ± 0.0012	1.04092 ± 0.00031	67.36 ± 0.54	0.9649 ± 0.0042	3.044 ± 0.014
<i>r</i>	0.02237 ± 0.00014	0.1199 ± 0.0012	1.04092 ± 0.00031	67.40 ± 0.54	0.9659 ± 0.0041	3.044 ± 0.014
$dn_s/d\ln k$	0.02240 ± 0.00015	0.1200 ± 0.0012	1.04092 ± 0.00031	67.36 ± 0.53	0.9641 ± 0.0044	3.047 ± 0.015
$dn_s/d\ln k, r$	0.02243 ± 0.00015	0.1199 ± 0.0012	1.04093 ± 0.00030	67.44 ± 0.54	0.9647 ± 0.0044	3.049 ± 0.015
$d^2n_s/d\ln k^2$, $dn_s/d\ln k$.	0.02237 ± 0.00016	0.1202 ± 0.0012	1.04090 ± 0.00030	67.28 ± 0.56	0.9625 ± 0.0048	3.049 ± 0.015
$N_{\rm eff}$	0.02224 ± 0.00022	0.1179 ± 0.0028	1.04116 ± 0.00043	66.3 ± 1.4	0.9589 ± 0.0084	3.036 ± 0.017

Baryogenesis



Baryogenesis via leptogenesis



Baryogenesis via leptogenesis



● 几种常见机制的简介

- ◎ 理论进展
- ◎ 实验验证
- 当前的重要问题

● 几种常见机制的简介

● 理论进展

- ◎ 实验验证
- 当前的重要问题

Sakharov conditions for baryogenesis

B violation

C/CP violation

Out of equilibrium dynamics

在满足洛伦兹不变、局域、厄密的标准场 论框架下,Sakharov三条件是必要条件

Sakharov conditions for leptogenesis

SM L/B-L violation

C/CP violation

Out of equilibrium dynamics

依然是必要条件吗? $L_{SM} = L_e + L_\mu + L_\tau$

Origin of neutrino masses



Minimal seesaw: I = 1,2; General: I = 1,2,3

Classical thermal leptogenesis

Recipes of thermal leptogenesis (in type I seesaw)



Covi, Roulet, Vissani, hep-ph/9605319

An example



Evolution of leptogenesis works



500+ citations

Discrete gauge symmetries and the origin of baryon and lepton number conservation in supersymmetric versions of the standard model Luis E. Ibanez (CERN), Graham G. Ross (Oxford U.) (May, 1991) Published in: <i>Nucl.Phys.B</i> 368 (1992) 3-37	Largest temperature of the radiation era and its cosmological implications#2Gian Francesco Giudice (CERN), Edward W. Kolb (Fermilab and Chicago U., EFI and Chicago U., Astron. Astrophys. Ctr.), Antonio Riotto (Pisa, Scuola Normale Superiore and INFN, Pisa) (May, 2000)Published in: Phys.Rev.D 64 (2001) 023508 • e-Print: hep-ph/0005123 [hep-ph]
Baryogenesis via leptogenesis M.A. Luty (Chicago U., EFI) (1992) Published in: <i>Phys.Rev.D</i> 45 (1992) 455-465	Thermal production of gravitinos#24M. Bolz (DESY), A. Brandenburg (DESY), W. Buchmuller (DESY) (Dec, 2000)Published in: Nucl.Phys.B 606 (2001) 518-544, Nucl.Phys.B 790 (2008) 336-337 (erratum) • e-Print: hep-ph/0012052 [hep-ph]
Baryogenesis from a lepton asymmetric universe Marion Flanz (Dortmund U.), Emmanuel A. Paschos (Dortmund U.), Utpal Sarkar (Dortmund U 1994) Published in: <i>Phys Lett B</i> 345 (1995) 248-252, <i>Phys Lett B</i> 384 (1996) 487-487 (erratum), <i>E</i>	A Lower bound on the right-handed neutrino mass from leptogenesis#2Sacha Davidson (Durham U.), Alejandro Ibarra (Oxford U.) (Feb, 2002)Published in: Phys.Lett.B 535 (2002) 25-32 • e-Print: hep-ph/0202239 [hep-ph]
Thermal inflation and the moduli problem David H. Lyth (Lancaster U.), Ewan D. Stewart (Tokyo U.) (Sep, 1995) Published in: Phys.Rev.D 53 (1996) 1784-1798 • e-Print: hep-ph/9510204 [hep-ph]	Resonant leptogenesis#20Apostolos Pilaftsis (Manchester U.), Thomas E.J. Underwood (Manchester U.) (Sep, 2003)Published in: Nucl.Phys.B 692 (2004) 303-345 • e-Print: hep-ph/0309342 [hep-ph]
CP violating decays in leptogenesis scenarios Laura Covi (SISSA, Trieste), Esteban Roulet (SISSA, Trieste), Francesco Vissani (ICTP, Trieste) 1996)	Towards a complete theory of thermal leptogenesis in the SM and MSSM#19G.F. Giudice (CERN), A. Notari (Pisa, Scuola Normale Superiore), M. Raidal (NICPB, Tallinn), A.Riotto (INFN, Padua), A. Strumia (Pisa U. and INFN, Pisa) (Oct, 2003)Published in: Nucl.Phys.B 685 (2004) 89-149 • e-Print: hep-ph/0310123 [hep-ph]Published Interpretation
CP violation and baryogenesis due to heavy Majorana neutrinos Apostolos Pilaftsis (Munich, Max Planck Inst.) (May, 1997) Published in: Phys.Rev.D 56 (1997) 5431-5451 • e-Print: hep-ph/9707235 [hep-ph]	Big-Bang nucleosynthesis and hadronic decay of long-lived massive particles#1Masahiro Kawasaki (Tokyo U., ICRR), Kazunori Kohri (Osaka U., Dept. Earth Space Sci.), TakeoMoroi (Tohoku U.) (Aug, 2004)Published in: Phys.Rev.D 71 (2005) 083502 • e-Print: astro-ph/0408426 [astro-ph]Earth Space Sci.)
Neutrino masses and leptogenesis with heavy Higgs triplets Ernest Ma (UC, Riverside), Utpal Sarkar (Ahmedabad, Phys. Res. Lab) (Feb, 1998) Published in: <i>Phys.Rev.Lett.</i> 80 (1998) 5716-5719 • e-Print: hep-ph/9802445 [hep-ph]	The nuMSM, dark matter and neutrino masses #1 Takehiko Asaka (IPT, Lausanne), Steve Blanchet (IPT, Lausanne), Mikhail Shaposhnikov (IPT, Lausanne) (Mar, 2005)
Baryogenesis via neutrino oscillations Evgeny K. Akhmedov (ICTP, Trieste and Kurchatov Inst., Moscow), V.A. Rubakov (Moscow, INI Trieste and Tokyo U., ICRR), A.Yu. Smirnov (ICTP, Trieste and Moscow, INR) (Mar, 1998) Published in: Phys.Rev.Lett. 81 (1998) 1359-1362 • e-Print: hep-ph/9803255 [hep-ph]	Published in: Phys.Lett.B 631 (2005) 151-156 • e-Print: hep-ph/0503065 [hep-ph] The νMSM, dark matter and baryon asymmetry of the universe #1 Takehiko Asaka (IPT, Lausanne), Mikhail Shaposhnikov (IPT, Lausanne) (May, 2005) #1 Published in: Phys.Lett.B 620 (2005) 17-26 • e-Print: hep-ph/0505013 [hep-ph] ● 915 citation











味效应是否需要考虑 --取决于 e, μ, τ 的Yukawa耦合是否在 sphaleron导致的重子与轻子的转化过程中发生作用

•
$$\Gamma_{L_{\tau}\leftrightarrow\tau_{R}(+\text{Higgs})} \simeq 5 \cdot 10^{-3} \times y_{\tau}^{2} T \text{ vs } H \simeq 1.66g_{\star} T^{2}/M_{\text{pl}}$$

• $T > 10^{12} \text{ GeV}, \tau_{R} \nleftrightarrow L_{\tau} \leftrightarrow B, \mu_{R} \nleftrightarrow L_{\mu} \leftrightarrow B$
• $10^{9} < T < 10^{12} \text{ GeV}, \tau_{R} \leftrightarrow L_{\tau} \leftrightarrow B, \mu_{R} \nleftrightarrow L_{\mu} \leftrightarrow B$
• $T > 10^{6} \text{ GeV}, e_{R} \leftrightarrow L_{e} \leftrightarrow B$

Resonant leptogenesis Pilaftsis, 9707235; 9702393; 9812256

CP asymmetry used in classical thermal leptogenesis

$$\begin{split} \varepsilon_{N_{1}} &= \frac{\Gamma(N_{1} \to lH) - \Gamma(N_{1} \to \bar{l}H^{\dagger})}{\Gamma(N_{1} \to lH) + \Gamma(N_{1} \to \bar{l}H^{\dagger})} & \text{fill} \| \text{bll} \text{fill} \\ &= -\frac{1}{8\pi} \sum_{\beta \neq 1} \frac{\Im \left[\left(\lambda^{\dagger} \lambda \right)_{\beta 1}^{2} \right]}{\left(\lambda^{\dagger} \lambda \right)_{11}} \left\{ \left[-\frac{M_{\beta}}{M_{1}} \left(1 - \left(1 + \frac{M_{\beta}^{2}}{M_{1}^{2}} \right) \right) \ln \left(1 + \frac{M_{1}^{2}}{M_{\beta}^{2}} \right) \right] + \left[\frac{M_{1}/M_{\beta}}{1 - \left(M_{1}^{2}/M_{\beta}^{2} \right)} \right] \right\} \\ &\simeq -\frac{3}{16\pi} \sum_{\beta \neq 1} \frac{\Im \left[\left(\lambda^{\dagger} \lambda \right)_{\beta 1}^{2} \right]}{\left(\lambda^{\dagger} \lambda \right)_{11}} \frac{M_{1}}{M_{\beta}} & \text{(For } M_{1} \ll M_{\beta}) \\ & \text{[L. Covi, E. Roulet and F. Vissani, PLB384 (1996)]} \end{split}$$

• CP asymmetry in the limit $M_1 \approx M_2$

$$\epsilon_i = \frac{1}{(\lambda^{\dagger}\lambda)_{ii}} \frac{\text{Im}[(\lambda^{\dagger}\lambda)_{21}^2]}{8\pi} \frac{M_1 M_2 (M_2^2 - M_1^2)}{(M_2^2 - M_1^2)^2 + A^2} \qquad A = M_i \Gamma_i - M_j \Gamma_j$$

Enhanced in the resonant region

These results are all obtained at zero temperature

Leptogenesis via RH neutrino oscillation



The "generalised" lepton number $\mathbf{L} = L + L_N$ is conserved.

$$P(N_{\alpha} \to N_{\beta}) - P(\overline{N}_{\alpha} \to \overline{N}_{\beta}) \propto \operatorname{Im} \{ \exp(-i \int_{0}^{t} \frac{\Delta M_{ij}^{2}}{2E} a(t) dt) \} \times \operatorname{Im} \{ Y_{\alpha i} Y_{\beta i}^{*} Y_{\alpha j}^{*} Y_{\beta j} \}$$

CP不对称产生于震荡过程中,而非衰变过程中,这是量子力学效应

 ν MSM (neutrino minimal Standard Model): 2 GeV RHN + 1 keV vDM

● 几种常见机制的简介

◎ 理论进展

- ◎ 实验验证
- 当前的重要问题

Theory development I: density matrix approach

Oscillations of interaction states → 宏观的量子相干效应

e.g., oscillation of active neutrinos $i\frac{d}{dt}|\nu_{\alpha}\rangle = \sum_{\beta=e,\mu,\tau} H_{\nu,\alpha\beta}|\nu_{\beta}\rangle$ $H_{\nu} = \frac{1}{2E}M_{\nu}M_{\nu}^{\dagger} + H_{\text{matter}}$

- Damping and Boltzmann evolution → 经典的粒子数密度的演化
 e.g., evolution of RHN and SM lepton number densities in classical leptogenesis
- Combined together \rightarrow **Density Matrix** formalism Raffelt, Sigl, 93
- Density matrix approach in ARS leptogenesis
 Hernández, Kekic, López-Pavón, Racker, Salvado, 1508.03676; 1606.06719

$$\rho_{N,ij} = |N_i\rangle f_{ij} \langle N_j| \qquad H_N = \frac{1}{2E} M_N M_N^{\dagger} + \frac{T^2}{8E} \lambda^{\dagger} \lambda$$
$$\frac{d\rho_N}{dt} = -i[H_N, \rho_N] + \{\Gamma_N^a, \rho_N\} + \{\Gamma_N^p, \rho_N^{eq} - \rho_N\}$$
$$\frac{d}{dt} \rightarrow aH \frac{d}{da} \qquad \text{oscillation annihilation production}$$

Theory development II: closed-time-path approach

case

QFT in non-equilibrium

 QFT at zero temperature or in thermal equilibrium



Reviews e.g., Quiros, hep-ph/9901312; Berges, hep-ph/0409233

Outline

References

Show full outline \checkmark

Cited By (927)



Physics Reports Volume 118, Issues 1–2, February 1985, Pages 1-131

5 87 88 AV 88 78 78 78 78 78 78 78 78 78 78 78 78
REPORTS
ADVIDU DEVOLUTION

Equilibrium and nonequilibrium formalisms made unified

Kuang-chao Chou, Zhao-bin Su *, Bai-lin Hao, Lu Yu

Show more \checkmark

Theory development II: closed-time-path approach

Propagators



Applying CTP approach in EW baryogenesis and phase transition

周五晁伟的报告

Classical formalism vs CTP formalism

Applying CTP in leptogenesis

Higgs H

Leptons L & N

$$\begin{split} \Delta_{q}^{<,>} &= \frac{-2\varepsilon(q^{0})\mathrm{Im}\Pi_{q}^{R}}{[q^{2} + \mathrm{Re}\Pi_{q}^{R}]^{2} + [\mathrm{Im}\Pi_{q}^{R}]^{2}} \Big\{ \vartheta(\mp q^{0}) + f_{B,|q^{0}|}(x) \Big\} \,, \\ S_{k}^{<,>} &= \frac{-2\varepsilon(k^{0})\mathrm{Im}\Sigma_{k}^{R\,2}}{[k^{2} + \mathrm{Re}\Sigma_{q}^{R\,2}]^{2} + [\mathrm{Im}\Sigma_{q}^{R\,2}]^{2}} \Big\{ \vartheta(\mp k^{0}) - f_{F,|q^{0}|}(x) \Big\} P_{L} \not \!\!\!\!\! k P_{R} \,, \end{split}$$

Lepton number density obtained from Wightman operator

$$n_L \sim \operatorname{tr}(\gamma^0 S_L^+)$$
 $S_L^+ = \frac{1}{2}(S_L^> + S_L^<)$

- CP violation source ——Either in oscillation or collision terms
- CPV source in RHN decay



Anisimov, Buchmuller, Drewes, Mendizabal, 1012.5821

Comparison with the classical calculation

Dev, Garny, Klaric, Millington, Teresi, 1711.02863.



KB: a general Kadanoff-Baym approach, two-time evolution Wigner: one-time treatment of KB approach as follows

$$iS_{Nij}(p,X) = \int d^4r \ e^{ip \cdot r} S_{Nij}(X+r/2,X-r/2)$$

● 几种常见机制的简介

● 理论进展

◎ 实验验证

• 当前的重要问题

Direct test of leptogenesis



Leptogenesis vs charged LFV



 $BR(\mu \to eee) < 1.0 \times 10^{-12} (90\% \text{ C.L.}),$

 $\begin{aligned} & \operatorname{CR}(\mu_{22}^{\,48}\mathrm{Ti} \to e_{22}^{\,48}\mathrm{Ti}) \ < \ 4.3 \times 10^{-12} \ (90\% \ \mathrm{C.L.}) \,, \\ & \operatorname{CR}(\mu_{79}^{\,197}\mathrm{Au} \to e_{79}^{\,197}\mathrm{Au}) \ < \ 7.0 \times 10^{-13} \ (90\% \ \mathrm{C.L.}) \,. \end{aligned}$

Leptogenesis vs neutrino oscillation experiments

N M 站地图、联系我们、English、中国科学院、北京分院 Institute of High Energy Physics Chinese Academy of Sciences					
首页 概况 机构 科研队伍 科研成果 研究生 博士后 院地合作 国际交流 创新文化 科学传播 党群园地 信息公开 网上博展馆 图片库 视频库					
団结 唯实 创新 奉献 「請輸入关键字 站内捜索					
√ 您现在的位置: 首页 > 新闻动态 > 高能新闻 > 2012年高能新闻					
大亚湾中微子实验发现新的中微子振荡 2012-03-08 文章来源: 【大中小】					
 ^{3]其} 如果它(θ₁₃)足够大,我们就能进行下一代实验来测量中微 子振荡中的宇称和电荷对称性(CP)破坏,以理解宇宙中物质 ^{4%±} -反物质不对称现象,即宇宙中'反物质消失之谜'。 					
代起即有迹象,当时称作太阳甲微于之迹和大气甲微于之迹。1998年日本的超级神闪实验止式发现大气甲微于振荡,随后太阳甲微于振荡也被多个实验证实。第三种振荡则一直未被发现,甚至有理论预言其根本不存在(即其振荡几率为零)。					
中国科学院高能物理研究所的科研人员2003年提出设想,利用我国大亚湾核反应堆群产生的大量中微子,来寻找中微子的这第三种振荡,其振荡几率用sin ² 2013表示。中国高能物理学会理事长赵光达院士表示,"013不仅是物理学中的一个基本参数,其数值的大小也对未来中微子物理的发展方向起着决定性的作用。如果它足够大,我们就能进行下一代实验来测量中微子振荡中的宇称和电荷对称性(CP)破坏,以理解宇宙中物质-反物质不对称现象,即宇宙中'反物质消失之谜'。否则我们就不知道如何进行下一代实验"。					

Connecting leptogenesis with low-energy experiments

Casas-Ibarra Parametrisation

$$M_{\nu} = Y M_N^{-1} Y^T v^2 \qquad \Longrightarrow \qquad Y = \frac{1}{v} U \sqrt{\hat{M}_{\nu}} R \sqrt{\hat{M}_N}$$

(3 RHN assumed)

$$\begin{array}{lll} \mathsf{PMNS} \\ \mathsf{matrix} \end{array} \quad U = \begin{pmatrix} c_{12}c_{13} & s_{12}c_{13} & s_{13}e^{-i\delta} \\ -s_{12}c_{23} - c_{12}s_{23}s_{13}e^{i\delta} & c_{12}c_{23} - s_{12}s_{23}s_{13}e^{i\delta} & s_{23}c_{13} \\ s_{12}s_{23} - c_{12}c_{23}s_{13}e^{i\delta} & -c_{12}s_{23} - s_{12}c_{23}s_{13}e^{i\delta} & c_{23}c_{13} \end{pmatrix} \begin{pmatrix} 1 & 0 & 0 \\ 0 & e^{i\frac{\alpha_{21}}{2}} & 0 \\ 0 & 0 & e^{i\frac{\alpha_{31}}{2}} \end{pmatrix} \end{array}$$

Complex
orthogonal
$$R = \begin{pmatrix} 1 & 0 & 0 \\ 0 & c_1 & s_1 \\ 0 & -s_1 & c_1 \end{pmatrix} \begin{pmatrix} c_2 & 0 & s_2 \\ 0 & 1 & 0 \\ -s_2 & 0 & c_2 \end{pmatrix} \begin{pmatrix} c_3 & s_3 & 0 \\ -s_3 & c_3 & 0 \\ 0 & 0 & 1 \end{pmatrix}$$
 $R^T R = R R^T = 1$
matrix

$$c_I = \cos(x_I + iy_I)$$
 $s_I = \sin(x_I + iy_I)$

New parameters from the heavy sector: M_I , $x_I + iy_I$ for I = 1,2,3

Leptogenesis from low-energy CP violation

Figure 4.2. In the $m_1 = 0$ case with z being real (i.e., $z = \theta$), the contour lines of $M_1/(10^{10} \text{ GeV})$ for successful flavored leptogenesis are shown (a) in the θ - σ plane by assuming $\delta = 0$ and (b) in the θ - δ plane by assuming $\sigma = 0$.

Xing, Zhao, 2008.12090

Leptogenesis from low-energy CP violation

Figure 9: The two-dimensional projections for intermediate scale leptogenesis with $M_1 = 1.29 \times 10^8$ GeV for $x_1 = 0$, $y_2 = 0$, $x_3 = 180^\circ$, $y_1 = y_2 = 180^\circ$, with *CP* violation provided

Moffat, Pascoli, Petcov, Turner, 1809.08251

Leptogenesis with flavour textures

当前中微子理论家偏向于具有特殊味道结构的费米子质量矩阵或者味混合矩阵

Leptogenesis with flavour textures

Resonant leptogenesis in minimal seesaw with TM1 mixing

Xing, S. Zhou, 0607302

Leptogenesis with modular-type flavour texture

The observed baryon asymmetry allows to determine the flavor scale Λ_1

Right-handed neutrino mass:

 $M_1 = 7.844 \times 10^{10} \text{ GeV}, M_2 = 1.826 \times 10^{11} \text{ GeV}, M_3 = 2.517 \times 10^{11} \text{ GeV}$

ULYSSES: Universal LeptogeneSiS Equation Solver

a python package calculating the baryon and lepton asymmetry in type-I seesaw https://github.com/earlyun

https://github.com/earlyuniverse/ulysses

Model	example input file	Description
1DME	1N3F.dat	DME 1 RHN Density Matrix Equation
2DME	2N3F.dat	DME 2 RHN
3DME	3N3F.dat	DME 3 RHN
1BE1F	1N1F.dat	one-flavoured BE 1 RHN Boltzmann Equation
1BE2F	1N2F.dat	two-flavoured BE 1 RHN
1BE3F	1N3F.dat	three-flavoured BE 1 RHN
2BE1F	2N1F.dat	one-flavoured BE with 2 RHN
2BE2F	2N2F.dat	two-flavoured BE with 2 RHN
2BE3F	2N3F.dat	three-flavoured BE with 2 RHN
3DMEsct	3N3F.dat	DME 3 RHN including scattering effects
1BE1Fsf	1N1F.dat	1BE1F evolving in scale factor
2RES	Res.dat	2BE3F in the resonant regime
2 RESsp	Res.dat	2RES including spectator processes

Table 1: Overview of built-in plugins. We abbreviate density matrix equations, Boltzmann equations and (decaying) right-handed neutrino as DME, BE and RHN respectively. The evolution variable is $z = M_1/T$ for all plugins other than 1BE1Fsf which evolves in the cosmological scale factor.

● 几种常见机制的简介

- 最新理论进展
- 实验检验

● 当前的重要问题

New mechanisms?

- In the framework of Type-II seesaw
 - Type-II leptogenesis Ma, Sarkar, 9802445 CP violation from triplet Higgs decay $\Delta \to LL$ and $\Delta^{\dagger} \to \overline{L}\overline{L}$
 - Via Affleck-Dine Mechanism 下午韩成成的报告

Barrie, Han, Murayama, 2106.03381; 2204.08202

- Lepton asymmetry directly generated from phase transition
 - Only a time-varying Weinberg operator

Pascoli, Turner, Zhou, 1609.07969; 1808.00470; 1808.00475

Including RH neutrinos

Long, Tesi, Wang, 1703.04902; Huang, Xie, 2206.04691

CPT violation/Lorentz violation

Barenboim, Salvado; Fujikawa, Tureanu; Sarkar ...

Testability

- Most well-motivated: neutrino oscillation experiments
 test only mixing parameters and low-energy CP violation
- Direct detection

- Indirect constraints by Charged LFV — up to 10⁴ TeV RHN
- New test: Stochastic GW background detection
 - GW generated from phase transitions
 - —— for phase transition temperature up to 10⁸ TeV
 - GWs generated from cosmic strings
 - —— requires U(1) breaking
 - -- for the symmetry-breaking scale up to 10¹⁴ TeV
 - --- requires RHN gain masses from a Spontaneous Symmetry Breaking

$$\lambda_{N,IJ} \phi \,\overline{N}_I N_J^c \Longrightarrow \qquad M_R = \lambda_N v_\phi$$

引力波的产生与探测

An example: GWs in Majoron model

	Inputs			Predictions				
	$m_S/{ m GeV}$	$ ilde{\mu}/{ m GeV}$	$M/{ m GeV}$	$v_0/{ m GeV}$	$T_{\star}/{ m GeV}$	lpha	eta/H_{\star}	a_0
A1	0.06190	0.0005857	0.5361	3.5873	0.6504	0.1248	2966	0.05951
A2	156.2	13.15	465.6	1014	721	0.04139	754.8	0.3886
A3	1036	13.72	7977	44424	9180	0.08012	1975	0.06268
A4	43874	1856	181099	567378	247807	0.05611	809.7	0.1944

GW from cosmic strings

Cui, Lewicki, Morrissey, Wells, 1808.08968

Towards a complete theory — GUT

- RHN with a U(1) symmetry can be easily embedded into a SO(10) GUT
- In GUTs, quark and lepton's masses and mixing are correlated together. RHN neutrino mass spectrum are not free.

$$\begin{split} M_u &= M_{10} + M_{126} + M_{120}, \\ M_d &= r_1 M_{10} + r_2 M_{126} + r_3 M_{120}, \\ M_e &= r_1 M_{10} - 3r_2 M_{126} + Ar_4 M_{120}, \\ M_\nu^D &= M_{10} - 3M_{126} + AM_{120}, \end{split}$$

 $M_R = M_{126} v_R / v_{ew}$ Dutta, Mimura, Mohapatra, 0406262 N_1 leptogenesis or N_{2,3} leptogenesis? — — In some GUT theory, mass and Yukawa coupling of N₁

may not allow N_1 leptogenesis. N_2 leptogenesis main dominate the lepton asymmetry.

Di Bari, Riotto, 1012.2343; Di Bari, King, 1507.06431

Towards a complete theory — GUT

不必要的拓扑缺陷: 包括磁单极和畴壁

$$G_x = G_{3221}$$
 or G_{421}

在任何大统一理论的破缺链 中,都不得不引入暴涨用来 稀释不必要的拓扑缺陷。

King, Pascoli, Turner, **YLZ**, 2005.13549

