



Chiral Properties of (2+1)-Flavor QCD in Magnetic Fields at Zero Temperature

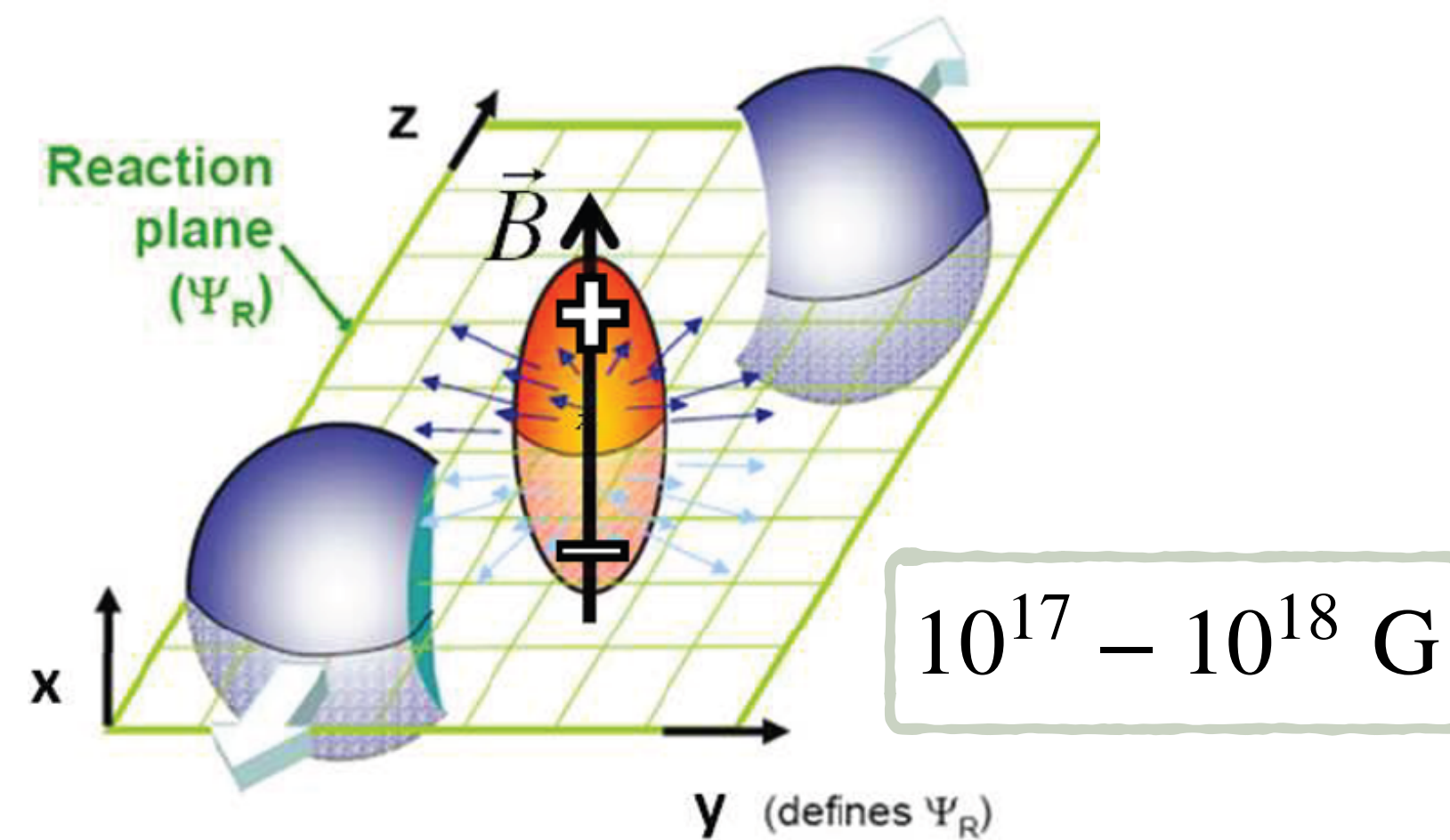
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第五届中国格点量子色动力学研讨会,
10.11, 2025@ 惠州

Strong Magnetic Fields: Strength Comparable to Λ_{QCD}^2

Heavy Ion-Collision



$$\text{RHIC} \sim 5m_\pi^2 \quad \text{LHC} \sim 70m_\pi^2$$

V. Skokov, *et al.*, IJMPA 24, 5925 (2009)
W.-T. Deng, X.-G. Huang PRC 85, 044907 (2012)

Magnetars



$$\text{surface} \sim 10^{14} - 10^{15} \text{ G}$$

$$\text{Kernel} \sim 10^{18} \text{ G}$$

R.Turolla *et al.*, arXiv:1507.02924 [astro-ph.HE]

Chiral Properties in Magnetic Fields

How do chiral properties change in comparable external magnetic fields?

❖ (Inverse) Magnetic catalysis

- Magnetic catalysis (MC): Chiral condensate $\uparrow$ with eB
- Chiral condensate: Order parameter for spontaneous chiral symmetry breaking (SSB).
- $T = 0$: No thermal fluctuations

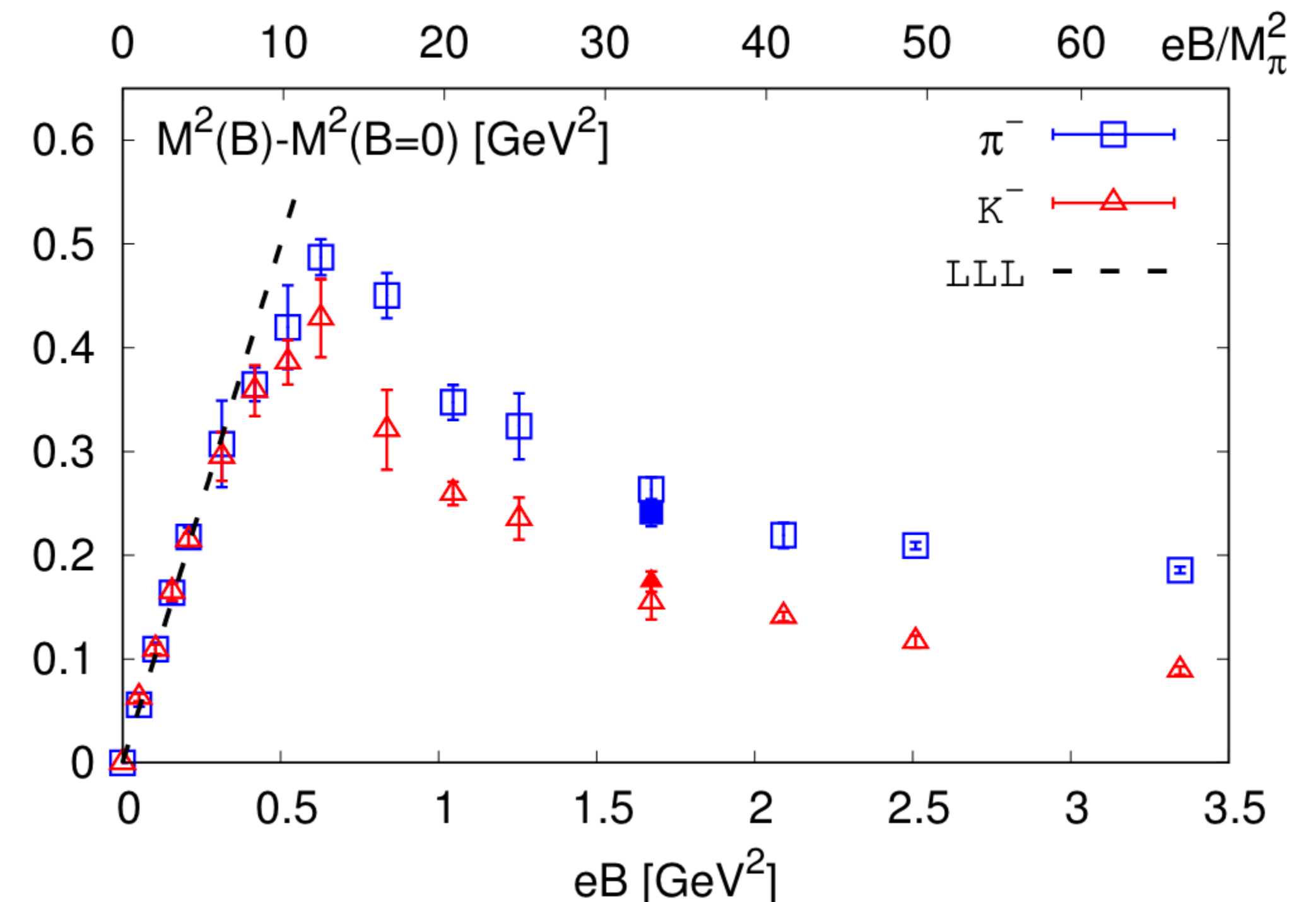
G. S. Bali *et al.*, Phys. Rev. D 86, 071502 (2012)

I. A. Shovkovy, Lect. Notes Phys. 871, 13 (2013)

G. S. Bali *et al.*, JHEP 02, 044 (2012)

❖ Masses and decay constants of mesons

- Goldstone bosons arising from SSB
- The previous work at $M_\pi \approx 220 \text{ MeV}$ at a finite lattice cutoff H.-T. Ding *et al.*, Phys. Rev. D 104, 014505 (2021)
 - Opposite mass behavior of neutral mesons and charged mesons in magnetic fields



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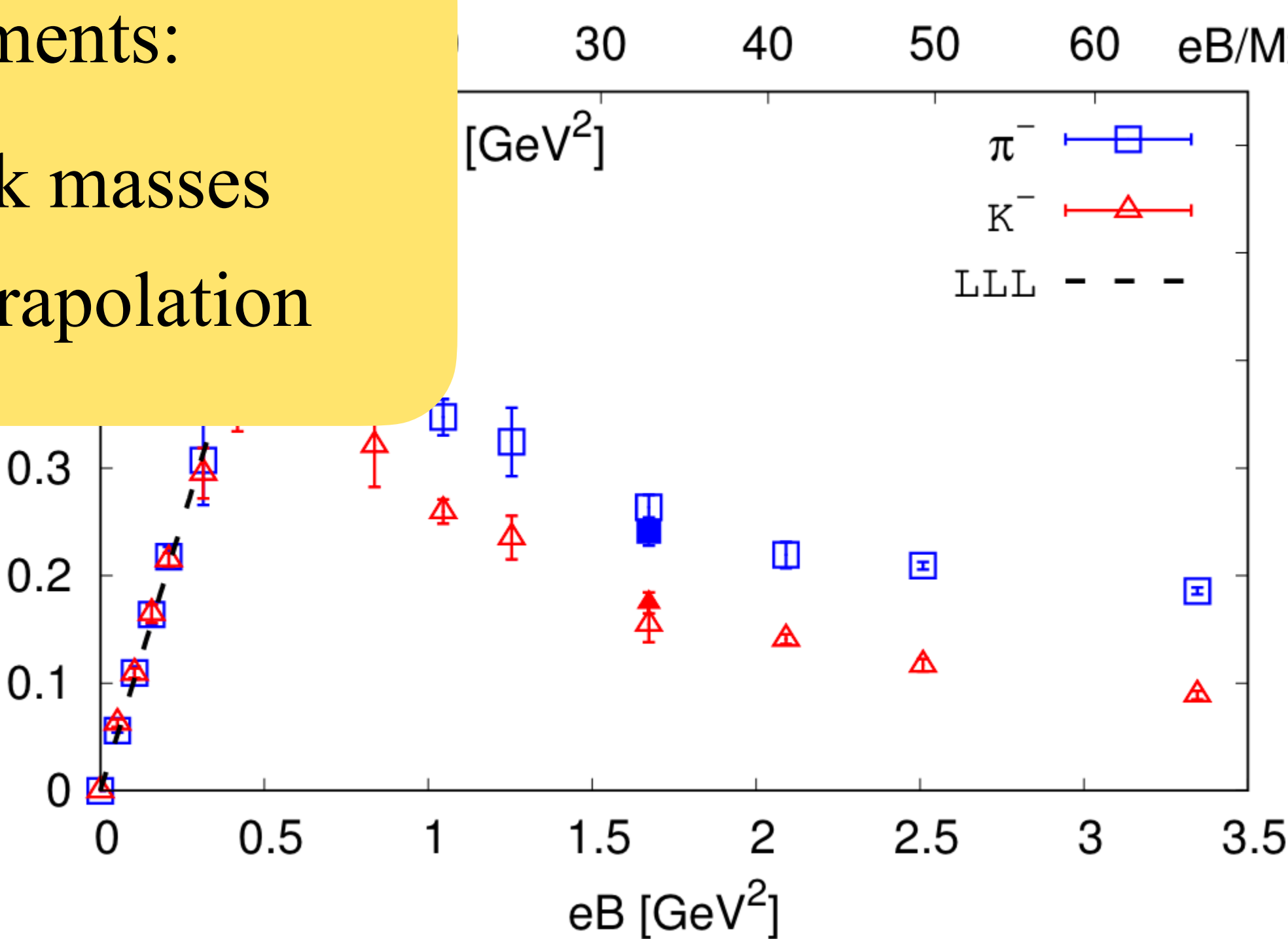
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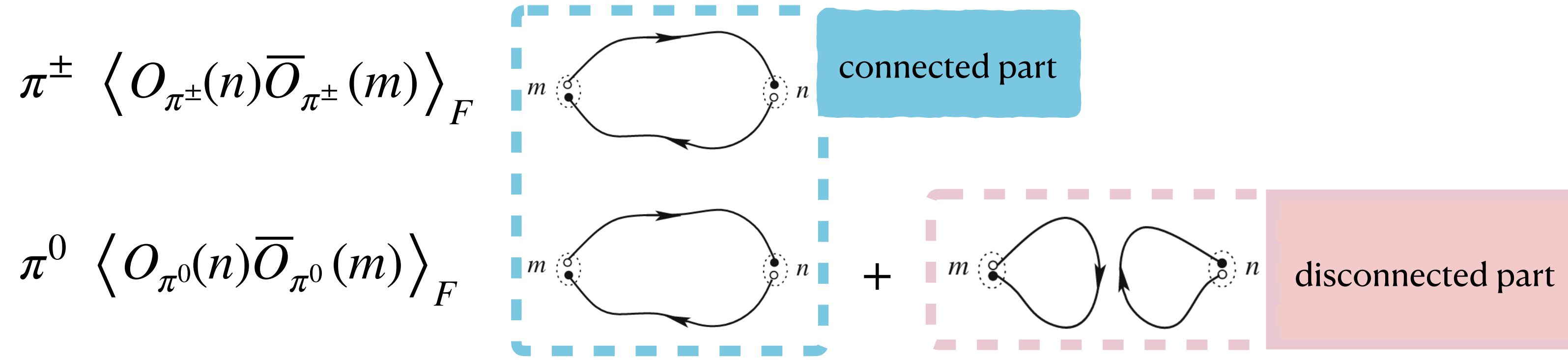
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Our Improvements:

- ✓ Physical quark masses
- ✓ Continuum extrapolation



Mesonic Two Point Correlation Functions



► The connected part of the correlation function,

$$G(\tau) = - \sum_{x,y,z} \zeta(\vec{n}) \text{Tr} \left[\left(M^{-1}(\vec{x}, \tau; \vec{0}, 0) \right)^\dagger M^{-1}(\vec{x}, \tau; \vec{0}, 0) \right],$$

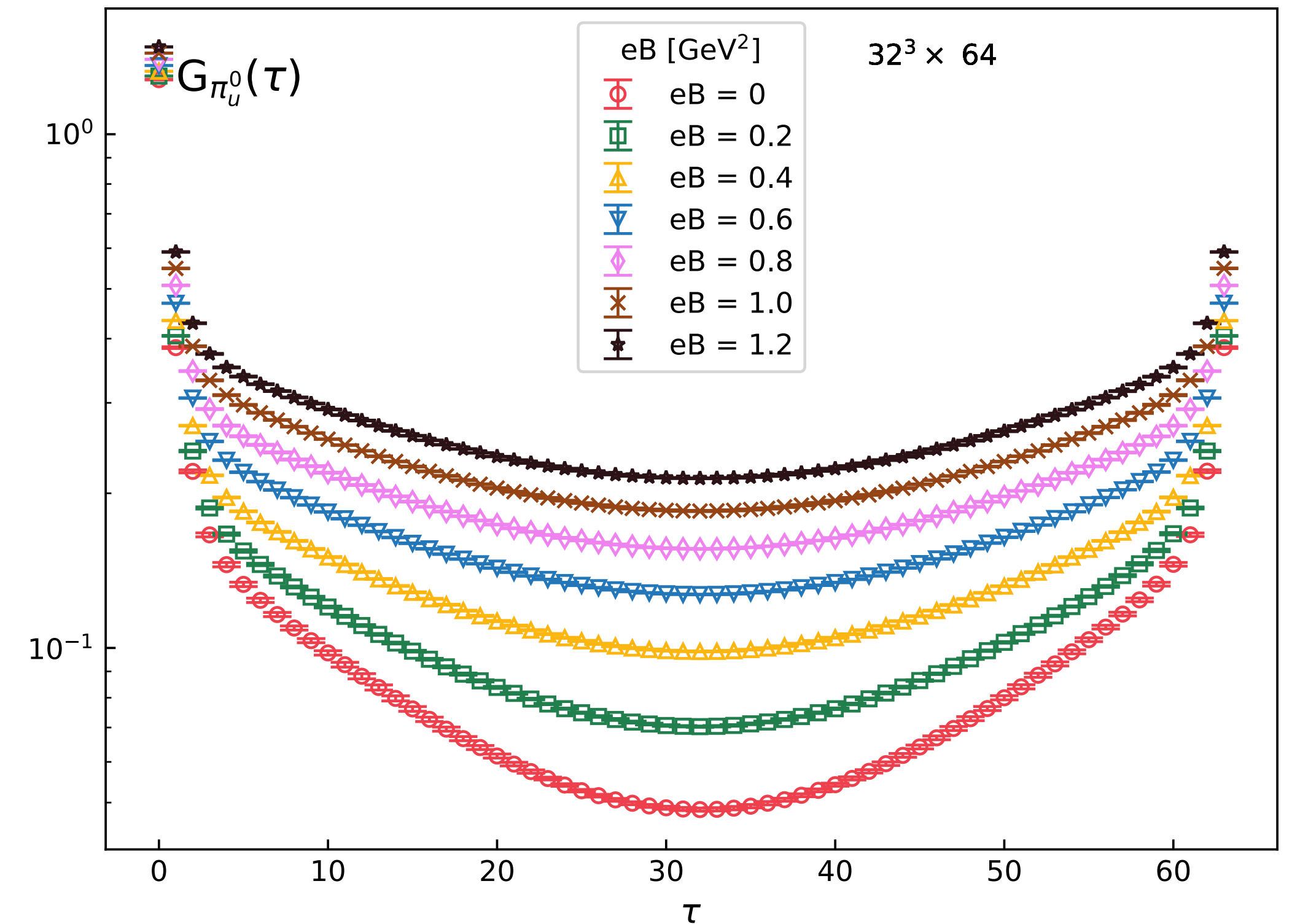
$M^{-1}(\vec{x}, \tau; \vec{0}, 0)$: Staggered fermion propagator from $(\vec{0}, 0)$ to $(\vec{x}, \tau)$

$\zeta(\vec{n})$: Phase factor, in the pseudoscalar channel $\zeta(\vec{n}) = 1$

► Energy Eigenstate Expansion:

$$G(\tau) = \sum_n |\langle 0 | \mathcal{M}(0) | n \rangle|^2 e^{-E_n \tau}$$

► At large distances τ : $\lim_{\tau \rightarrow \infty} G(\tau) \sim A e^{-m_\Gamma \tau}$



Mass & Decay Constant Extraction: Correlated Fitting

► Correlated fits: Account for the full covariance matrix in the fit.

► Fitting ansatz for different number of states:

$$G(n_\tau) = \sum_{i=1}^{N_{\text{nosc}}} A_{\text{nosc},i} e^{-M_{\text{nosc},i} n_\tau} - (-1)^{n_\tau} \sum_{i=0}^{N_{\text{osc}}} A_{\text{osc},i} e^{-M_{\text{osc},i} n_\tau}$$

Multi-exponential fits: $(N_{\text{nosc}}, N_{\text{osc}})$

Fit windows: $\tau \in [\tau_{\text{min}}, N_\tau/2]$

The physical meson mass is given by $M_{\text{nosc},1}$

The amplitude $A_{\text{nosc},1} \propto f_\Gamma$ (meson's decay constant)

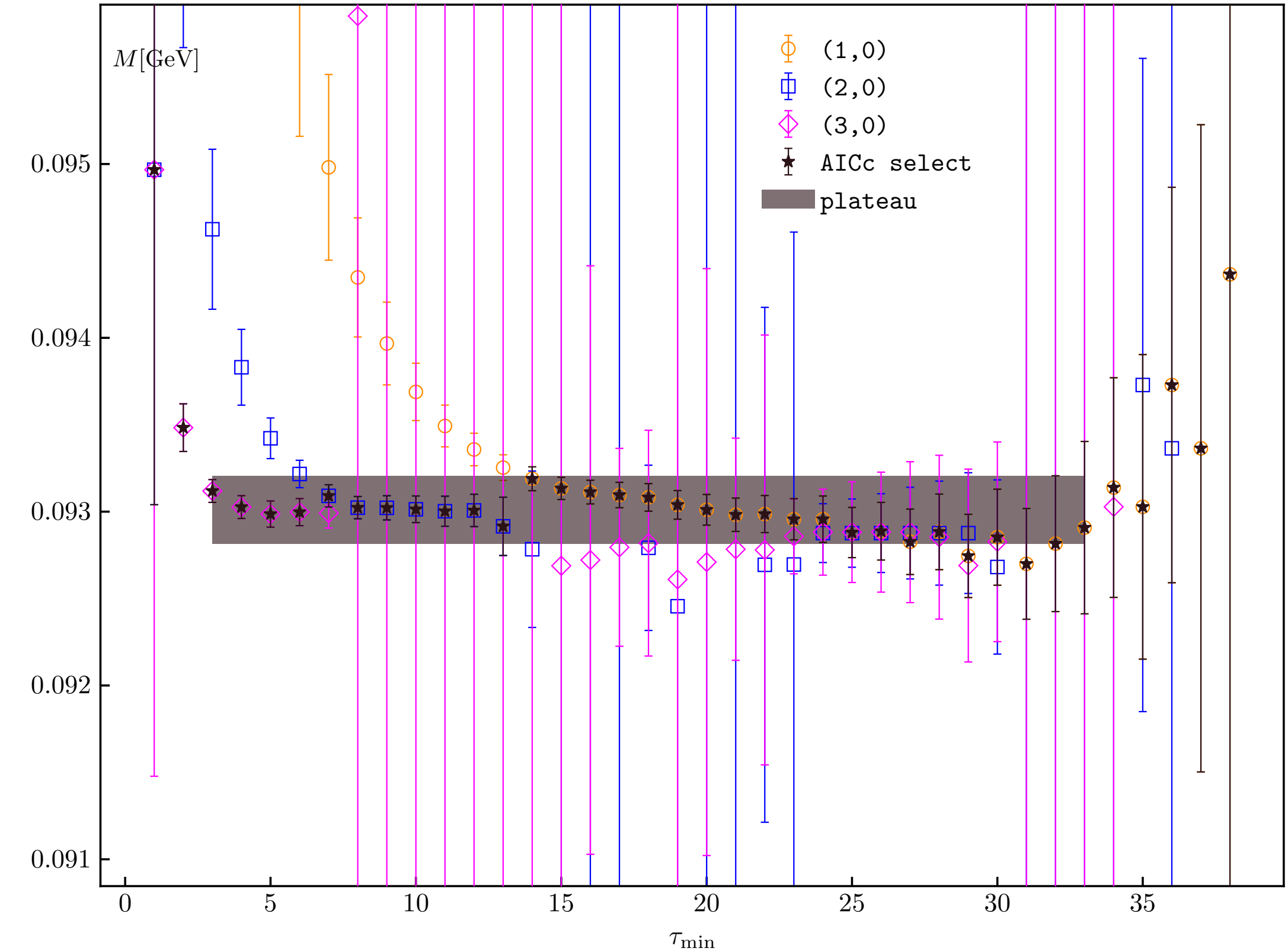
► Select the fit with lowest corrected Akaike Information

Criterion (AICc):

$$\text{AICc} = 2k - 2 \ln \hat{L} + \frac{2k(k+1)}{n-k-1},$$

► Use Gaussian bootstrapping to estimate final value and error

► We also use the Oblique Lanczos Method to cross-check the extracted energies.



Lattice Setup

- ◆ **Flavors:** $N_f = 2 + 1$
- ◆ **Action:** Highly Improved Staggered Quarks (HISQ) and Tree-level Improved Symanzik action
- ◆ **Quark masses:** Physical values ($m_u = m_d, m_s$); $m_s/m_u \approx 27$
- ◆ **Lattice sizes:** $24^3 \times 48$ (0.112 fm), $32^3 \times 64$ (0.084 fm), $40^3 \times 80$ (0.067 fm), $48^3 \times 96$ (0.056 fm)
- ◆ **eB ranges:** $0 \leq eB \leq 1.22 \text{ GeV}^2$ ($0 \sim 66m_\pi^2$),
 - ❖ **Quantized:** $eB = \frac{6\pi N_b}{N_x N_y} a^{-2}$, $N_b = 0, 2, 4, 6, 8, 10, 12$ (magnetic fluxes)

$N_\sigma^3 \times N_\tau$	a [fm]	β	am_s	am_l	# of conf at eB [GeV ²]						
					0.0	0.2	0.41	0.61	0.81	1.02	1.22
$24^3 \times 48$	0.112	6.722	0.0484	0.001793	600	701	1401	1350	1350	1350	1350
$32^3 \times 64$	0.084	7.016	0.0353	0.001307	600	800	800	800	800	800	800
$40^3 \times 80$	0.067	7.254	0.0274	0.001015	475	555	555	555	555	321	316
$48^3 \times 96$	0.056	7.455	0.0222	0.000822	150	266	250	250	220	306	215

Continuum Extrapolation

❖ Obtain physical results by taking the **continuum limit** ($a \rightarrow 0$)

❖ Fit lattice observable $O(a^2)$ using different ansätze:

- Linear: $O(a^2) = O_0^{\text{linear}} + a^2 A$

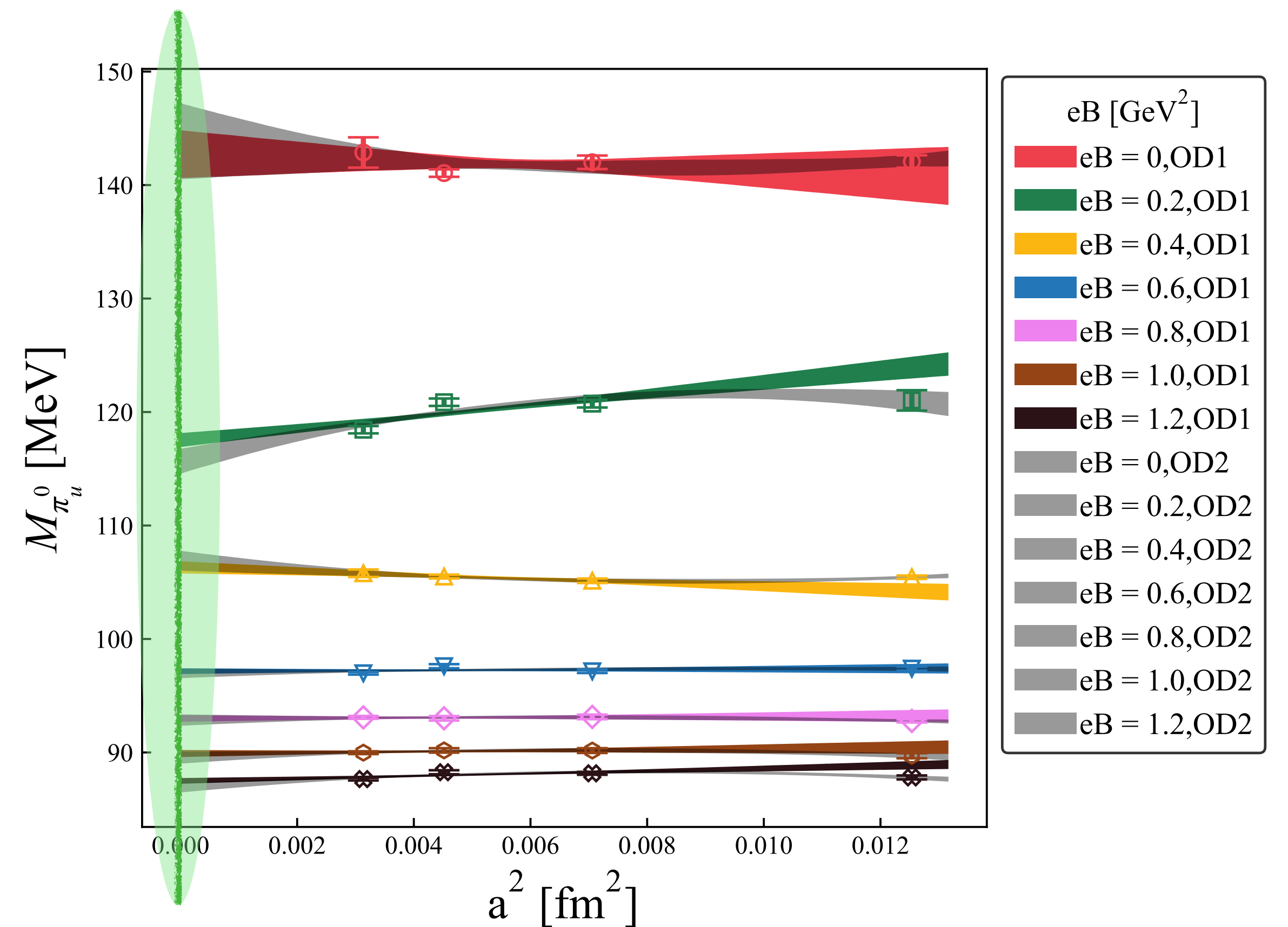
Applied to the three finest lattices.

- Quadratic: $O(a^2) = O_0^{\text{quad}} + a^2 A + a^4 B$

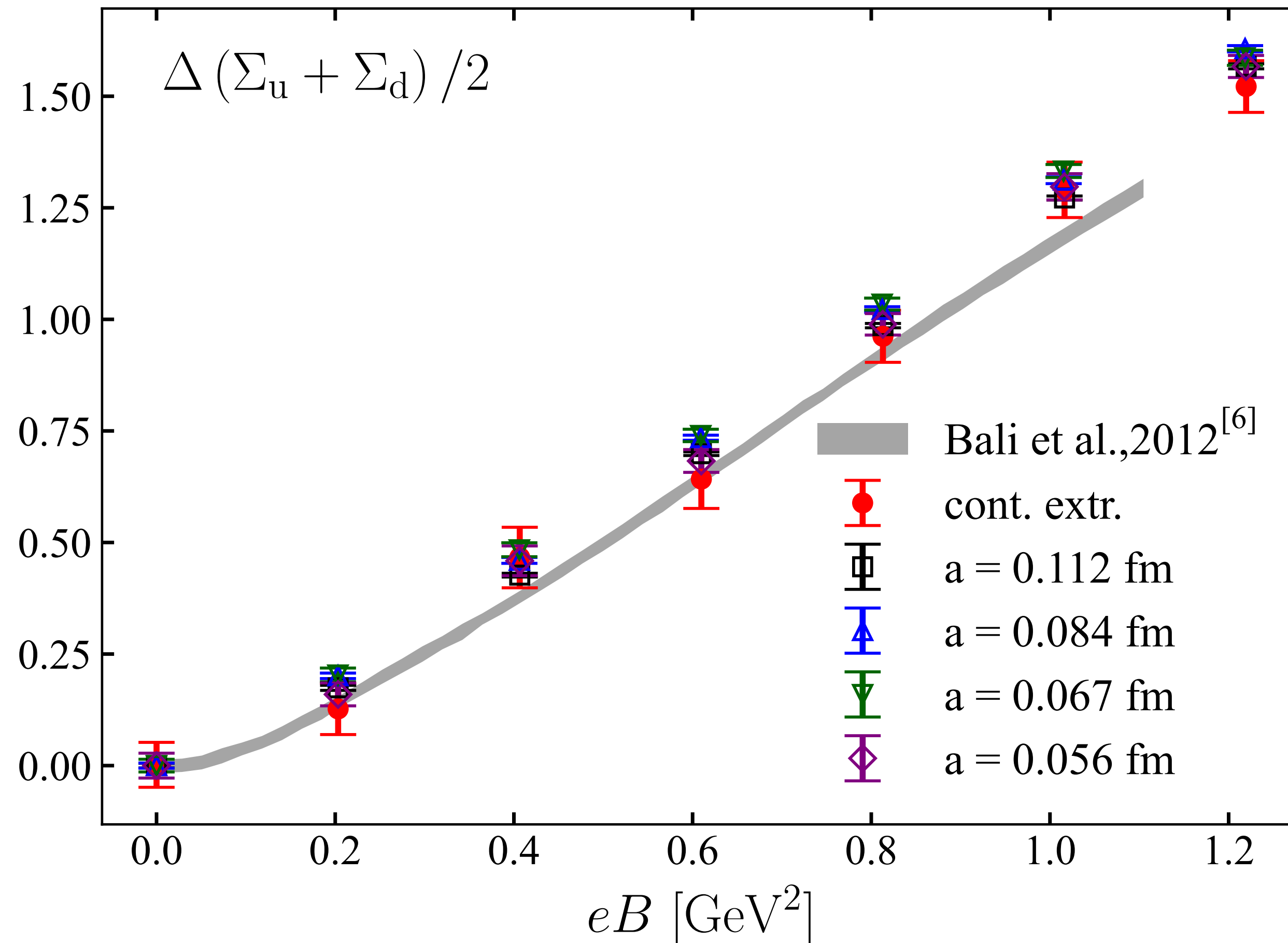
Applied to all four lattices.

❖ Final Continuum Value:

$$O_{\text{cont}} = \frac{1}{2} \left(O_0^{\text{linear}} + O_0^{\text{quad}} \right)$$



Light Quark Chiral Condensates in Magnetic Fields



*G. S. Bali *et al.*, Phys. Rev. D **86**, 071502 (2012)

Chiral Condensate

$$\langle \bar{\psi} \psi \rangle_q(B, T) \equiv \frac{T}{V} \frac{\partial \ln \mathcal{Z}(B, T)}{\partial m_f}$$

$\mathcal{Z}$: Partition function

V : Spatial volume

$T = 1/(aN_\tau)$

$q = u, d$

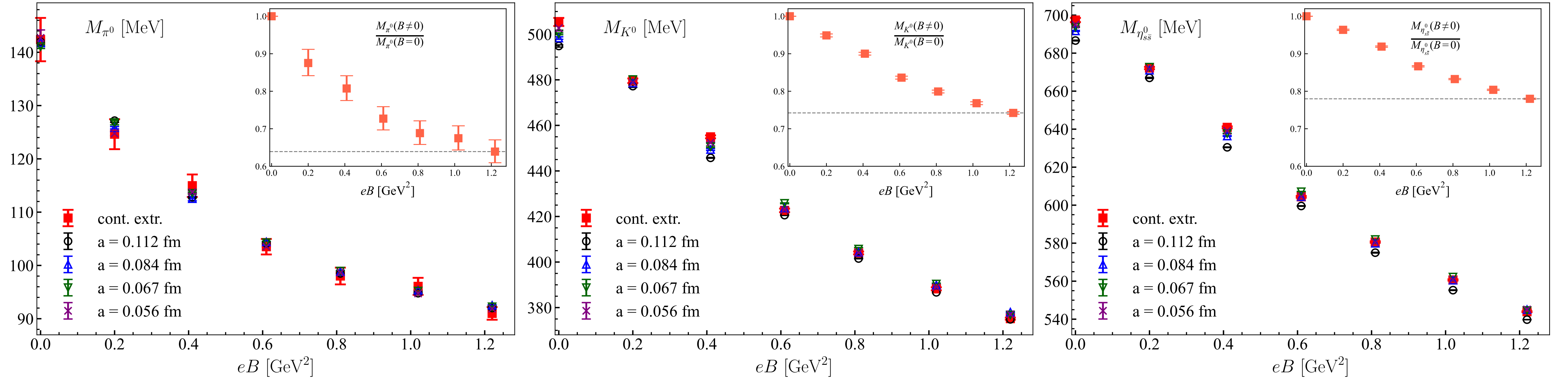
To remove additive and multiplicative divergences

Renormalized light quark chiral condensate

$$\Delta \Sigma_q(B) = \frac{2m_q}{m_\pi^2 F_\pi^2} \left(\langle \bar{\psi} \psi \rangle_q(B, T=0) - \langle \bar{\psi} \psi \rangle_q(B=0, T=0) \right)$$

☑ Magnetic Catalysis: the light quark condensates increase with the magnetic field at $T = 0$

Masses of Neutral Mesons: $\pi^0, k^0, \eta_{s\bar{s}}^0$



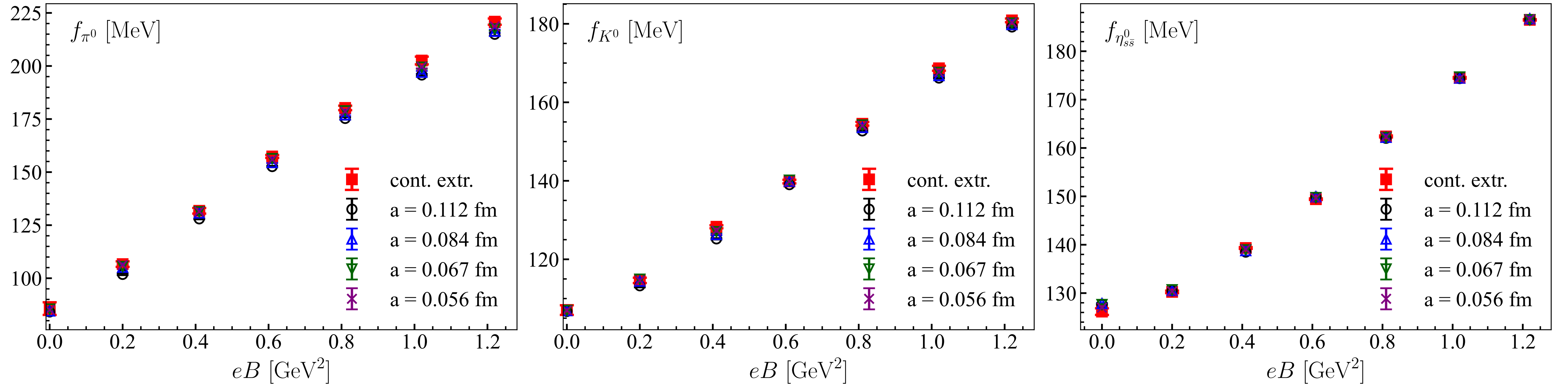
M_{π^0} using $G_{\pi^0} = (G_{\pi_u^0} + G_{\pi_d^0})/2$;

$M_{K^0(s\bar{d})}$

$M_{\eta_{s\bar{s}}^0}$

- ☑ The masses of neutral mesons decrease as eB increases, with the neutral pion mass reaching about 65% of its original value at $eB \sim 1.2 \text{ GeV}^2$.
- ☑ The heavier the neutral meson, the less it is affected by the magnetic field.

Decay Constants of $\pi^0, K^0, \eta_{s\bar{s}}^0$



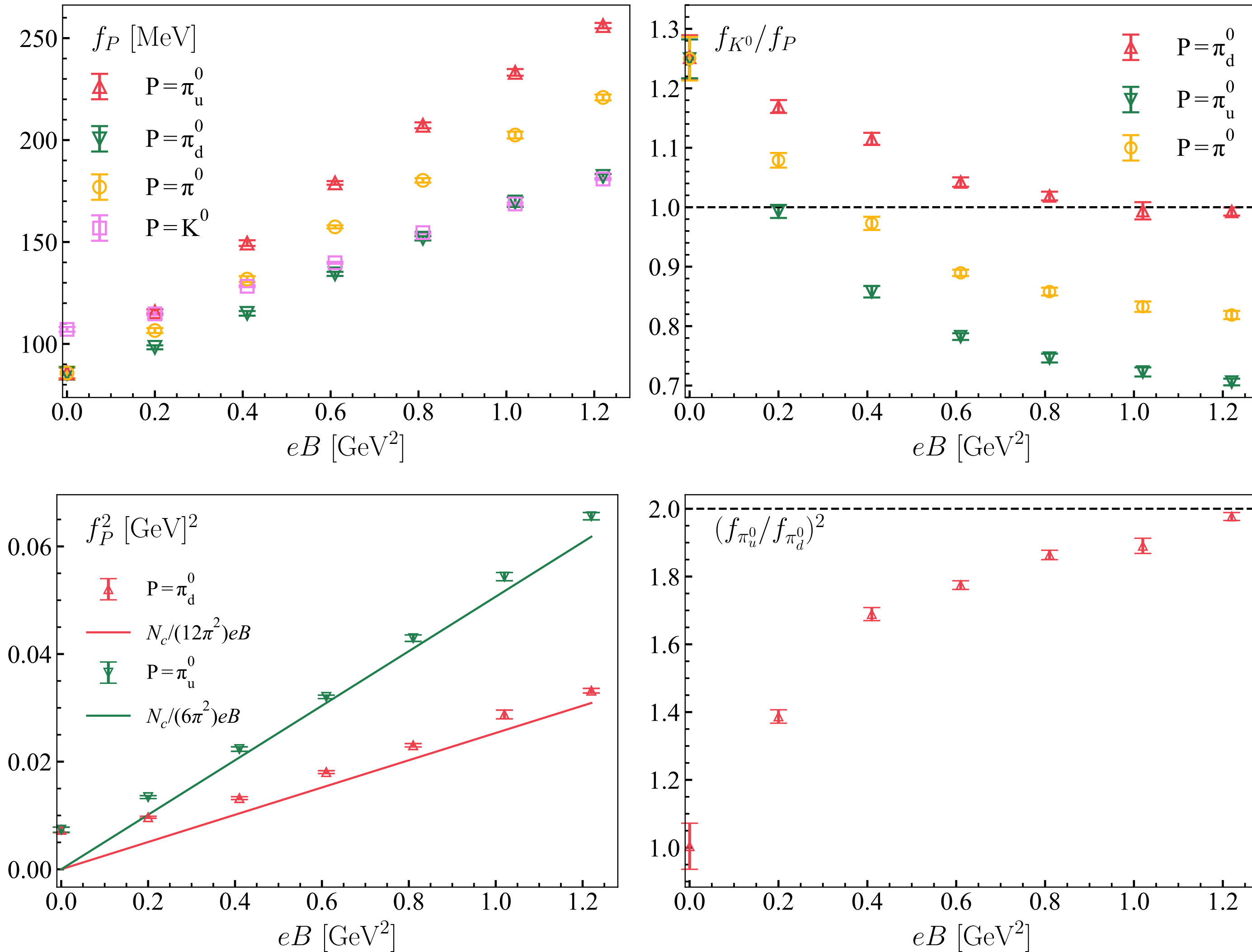
$$f_{\pi^0} = 2m_l \sqrt{\frac{A_{\pi^0;\text{nosc},1}}{4M_{\pi^0}^3}}$$

$$f_{K^0} = (m_l + m_s) \sqrt{\frac{A_{K^0;\text{nosc},1}}{4M_{K^0}^3}}$$

$$f_{\eta_{s\bar{s}}^0} = 2m_s \sqrt{\frac{A_{\eta_{s\bar{s}}^0;\text{nosc},1}}{4M_{\eta_{s\bar{s}}^0}^3}}$$

☑ All three decay constants increase with the magnetic field strength.

Decay Constants of $\pi^0, K^0, \eta_{s\bar{s}}$

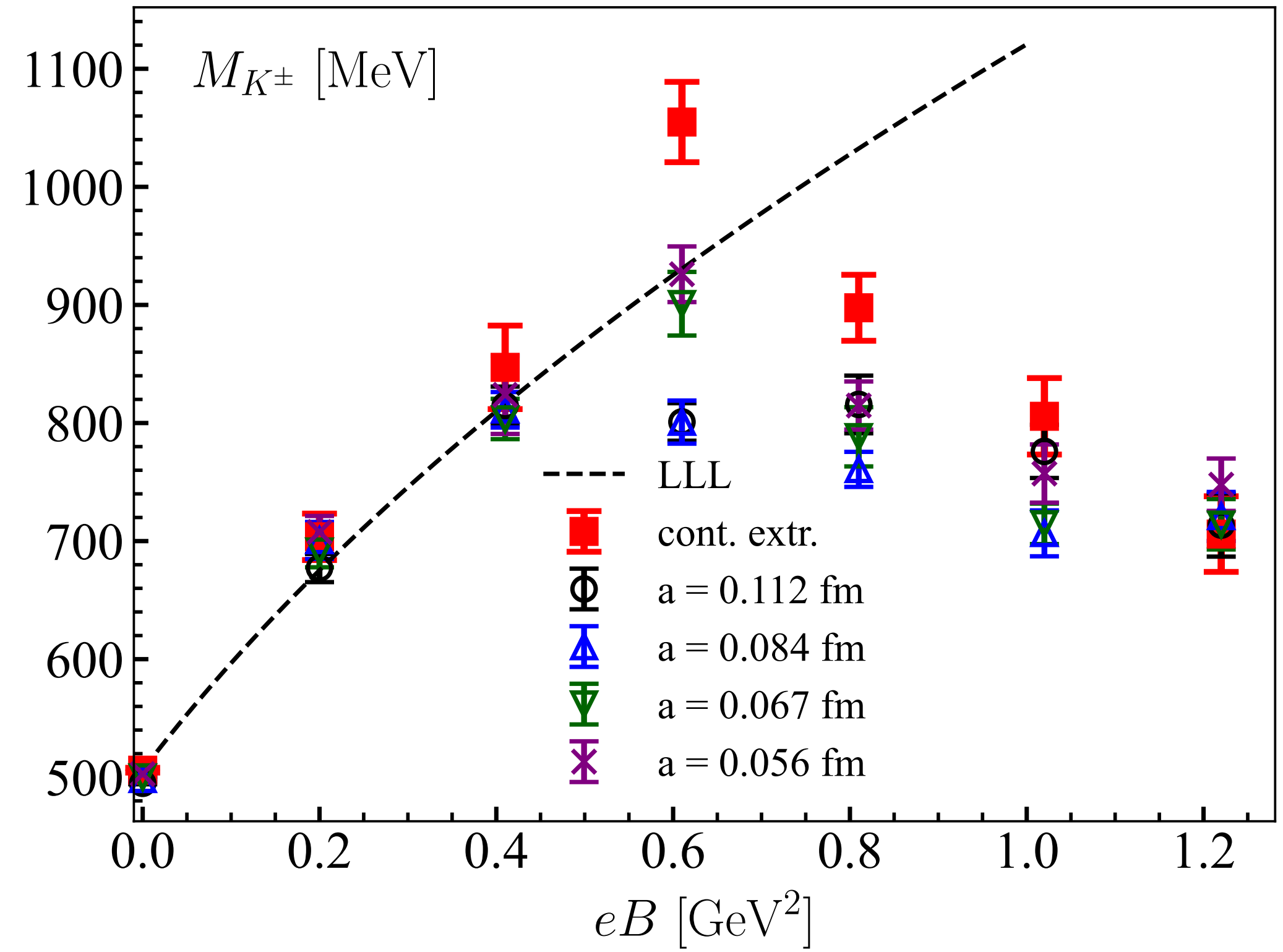
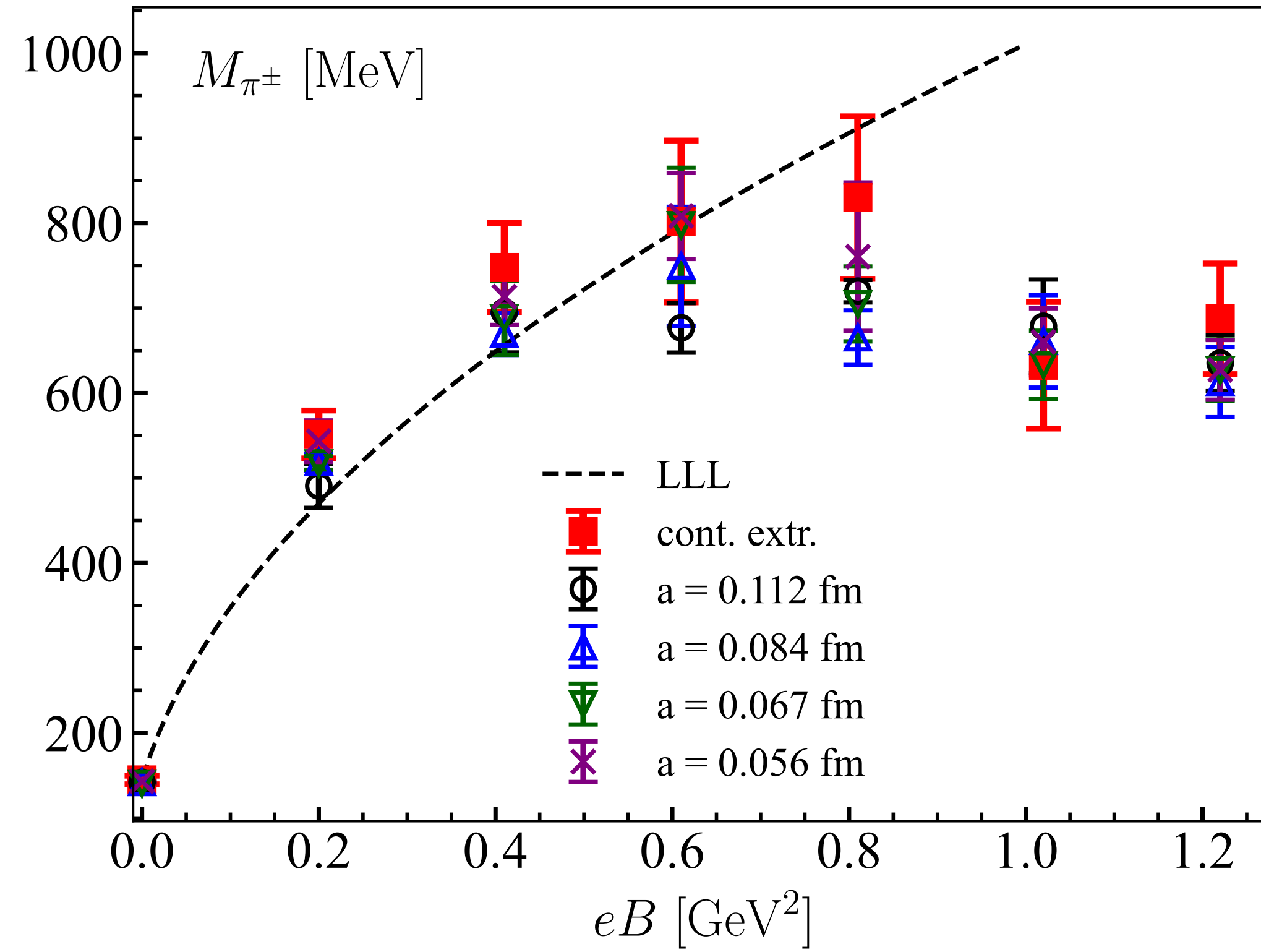


- ☑ $eB \uparrow$, $f_{\pi_u^0}$ increase fastest, f_{K^0} increase slowest
- ☑ f_{K^0}/f_π decreases with eB , especially for the up-quark component

$$\tilde{f}^2 = \frac{(N_d f_u + N_u f_d)^2}{N_f^2} = \frac{(\sqrt{2}N_d + N_u)^2 N_c}{12\pi^2 N_f^2} |eB|.$$

V. A. Miransky, I. A. Shovkovy, Phys.Rev. D66, 045006(2002)

The Masses of Charged Meson $\pi^\pm, K^\pm$



The Lowest Landua Level (LLL),

$$M_{\text{ps}}^\pm(B) = \sqrt{\left(M_{\text{ps}}^\pm(B=0)\right)^2 + |eB|}.$$



$M_{\pi^\pm}, M_{K^0}$

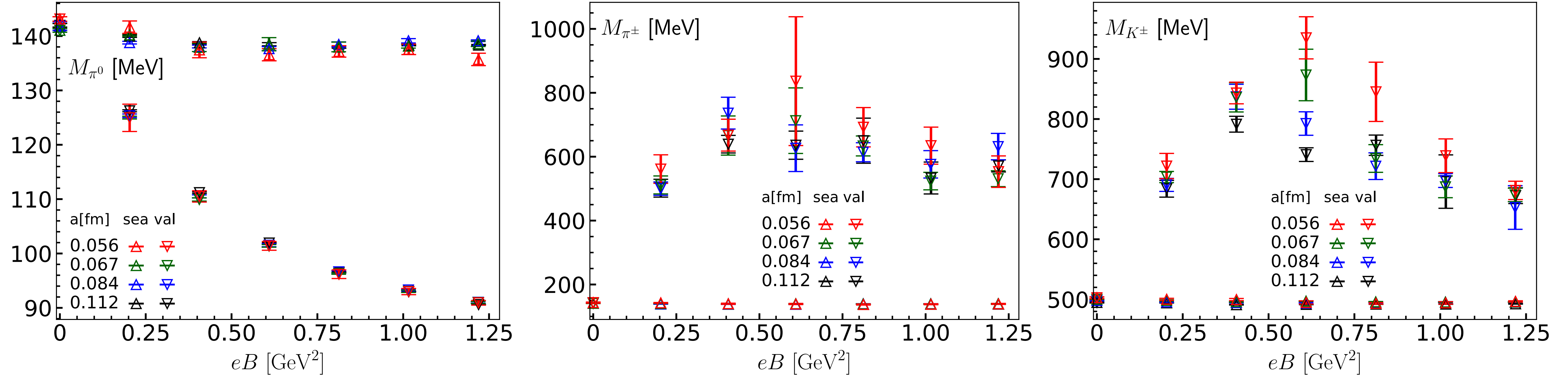


Initially increase with eB ,



Around $eB \approx 0.6 \text{ GeV}^2$, plateau or decrease slightly, deviate from the lowest Landau level.

Masses From Sea and Valence Quark Contributions



- ☑ Sea quark: weak eB dependence, change $< 3\%$
- ☑ Valence quark: follows total mass behavior vs eB

Summary

- 📌 Magnetic catalysis observed up to $eB \sim 1.2 \text{ GeV}^2$.

From Lattice correlator calculations in pseudoscalar channel at zero temperature

- 📌 Neutral meson masses $\downarrow$ as $eB \uparrow$, with heavier mesons less affected.
- 📌 Decay constants of neutral pseudoscalar mesons $\uparrow$ with eB .
- 📌 The masses of charged mesons deviate from Lowest Landau Levels at $eB \approx 0.6 \text{ GeV}^2$, with no further increase.
- 📌 Valence quarks drive the eB dependence of meson masses; sea-quark effects are minor.

Thank You for your attention!

Backup

Oblique Lanczos Method: Fit - Independent Energy Extraction

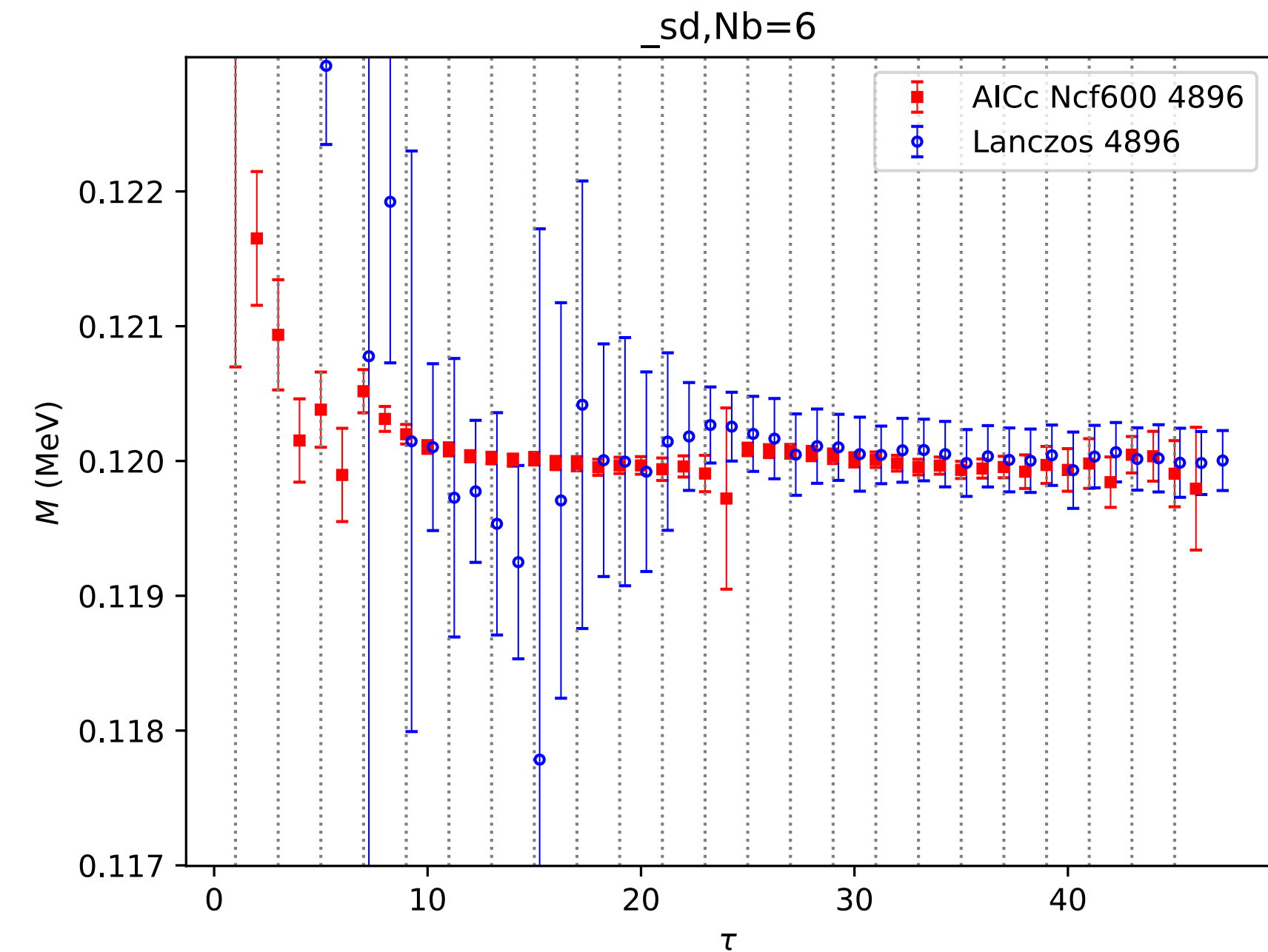
► Why use Oblique Lanczos?

- 🌐 **Cross-check** multi-state fit analysis
- 🌐 Provide a **fit-independent** estimate of ground-state mass

► How it works:

- 🌐 Direct application to transfer matrix $T = e^{-aH}$
- 🌐 Extracts energy levels from Krylov-space projection
- 🌐 Energy levels: $E_k = -(1/a)\ln \lambda_k$

λ_k : Eigenvalues of effective tridiagonal matrix T_m



Ground-state mass for K^0 via Oblique Lanczos (blue points) and Correlated Fit (red point) on a $48^3 \times 96$ lattice at $eB = 0.6 \text{ GeV}^2$

☑ **Agreement** with AICc-preferred fit results