

LQCD Determination of Quark Spin in Octet Baryons and SU(3)-Flavor Symmetry

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- ◆ **Chapter 1: Introduction**
- ◆ **Chapter 2: Theoretical Framework**
- ◆ **Chapter 3: Numerical result of Quark Spin and Discussion**
- ◆ **Chapter 4: Conclusion and Outlook**



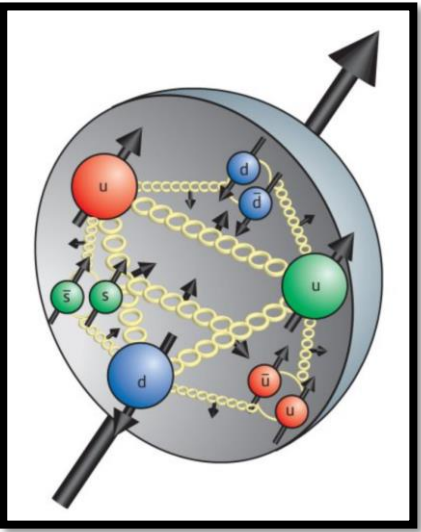
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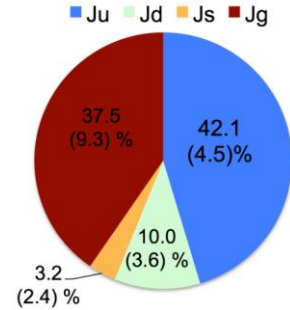
1. Introduction: Decomposition of Proton Spin



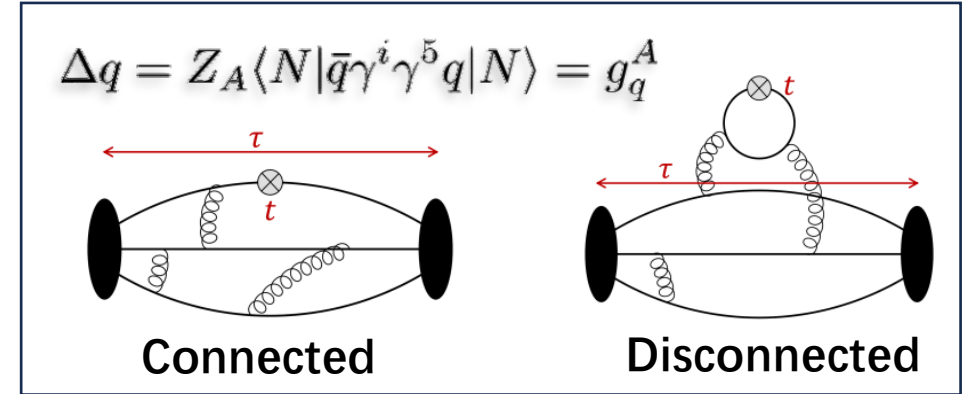
Ji's sum rule: $\frac{1}{2} = \sum_{q=u,d,s,\dots} \left(\frac{1}{2} \Delta q + L_q \right) + J_g$

X.D. Ji, PRL 78, 610(1997)

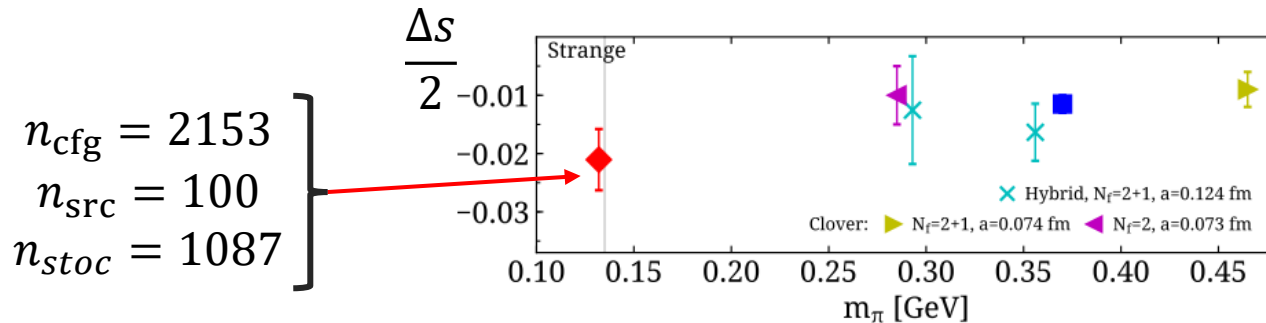
□ Total **intrinsic quark spin** contribute ~30-40% to nucleon;



Nature Rev.Phys. 3, 27-38(2021)

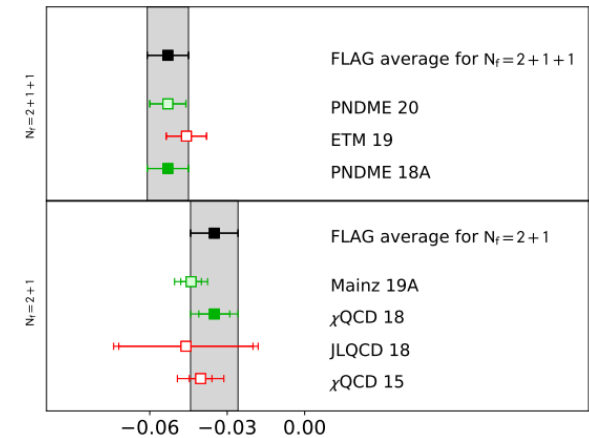


□ LQCD faces **large challenge** on calculate **quark-disconnected insertion** contribution;



C. Alexandrou et al., PRL 119, 142002 (2017)

g_A^s

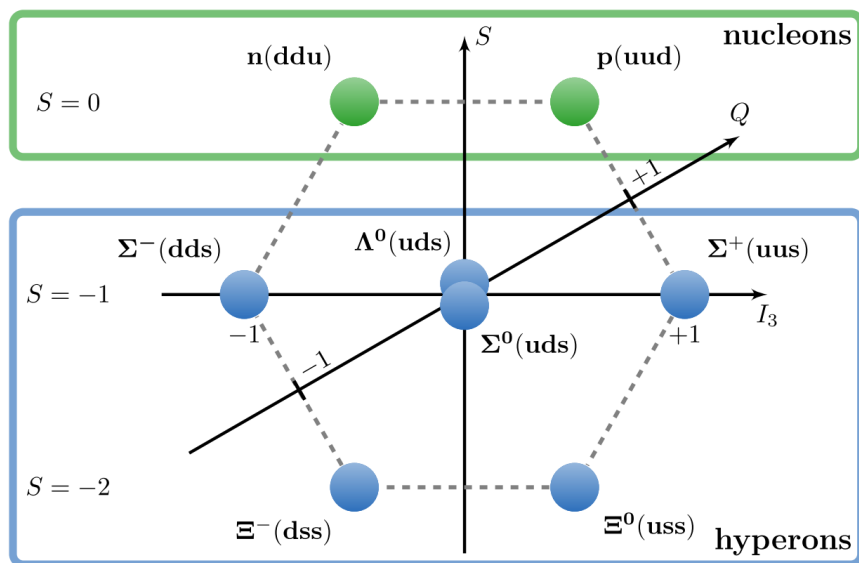


PNDME, PRD 98, 094512 (2018)

χ QCD, PRD 98, 074505 (2018)

Mainz, arXiv 1911.01177 (2019)

1. Introduction: SU(3) flavour symmetry

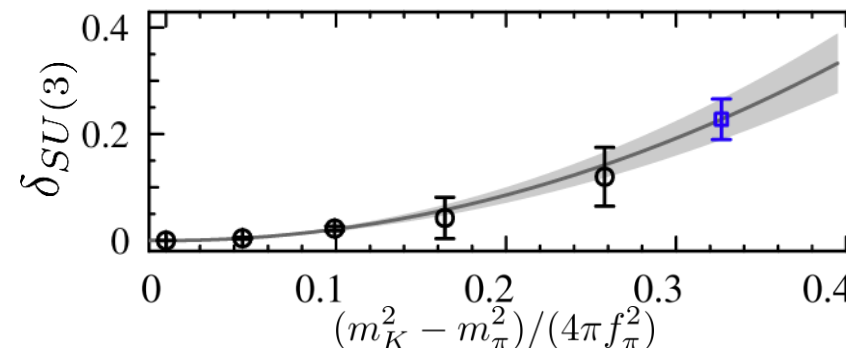


$$\begin{aligned} m_u^{\overline{\text{MS}}(2\text{GeV})} &\simeq 2.1\text{MeV} \\ m_d^{\overline{\text{MS}}(2\text{GeV})} &\simeq 4.7\text{MeV} \\ m_s^{\overline{\text{MS}}(2\text{GeV})} &\simeq 93\text{MeV} \end{aligned}$$

FLAG(2024), arXiv:2411.04268

- There is moderate **SU(3) symmetry breaking** effects for the iso-vector axial charges, $\delta_{SU(3)}/g_A^n \sim 9\%$;

$$\delta_{SU(3)} = g_A^n(u-d) - g_A^\Sigma(u-d) + g_A^\Xi(u-d)$$

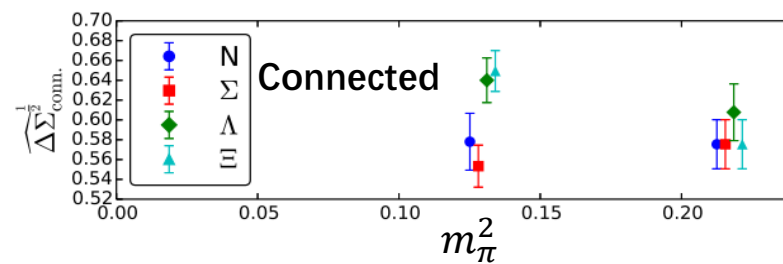


H.W. Lin et al, PRD **79**, 034507 (2009)

A. Savanur, PRD **102**, 014501 (2020)

RQCD, PRD **108**, 034512 (2023)

- How does SU(3) flavour symmetry and breaking affect **Hyperons Spin Structure**? A. J. Chambers et al, PRD **90**, 014510 (2014)



✓ In this study we use a novel ***Blending method*** to study **“Quark Spin in Octet Baryons”** (including quark DI contribution) at 4 ensembles (include 3 pion mass and 2 lattice spacing).



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2: Theoretical Framework: Blending Method



- ◆ ***Distillation*** can significantly improve the **signal-to-noise ratio** and suppress **excited-state contamination**;

$$-\nabla^2|\phi\rangle = \lambda|\phi\rangle,$$

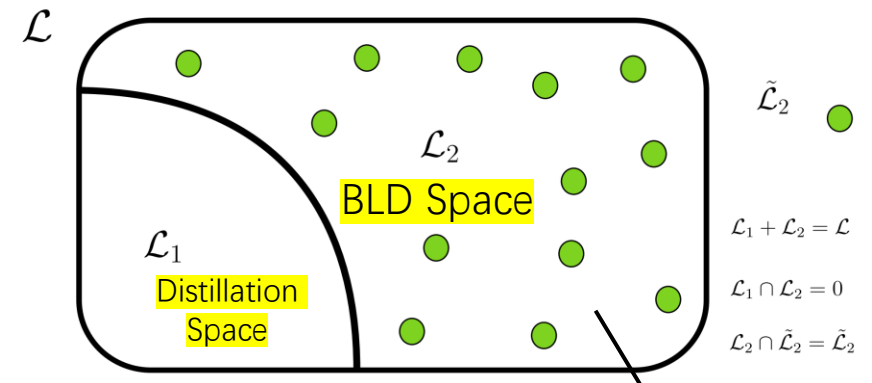
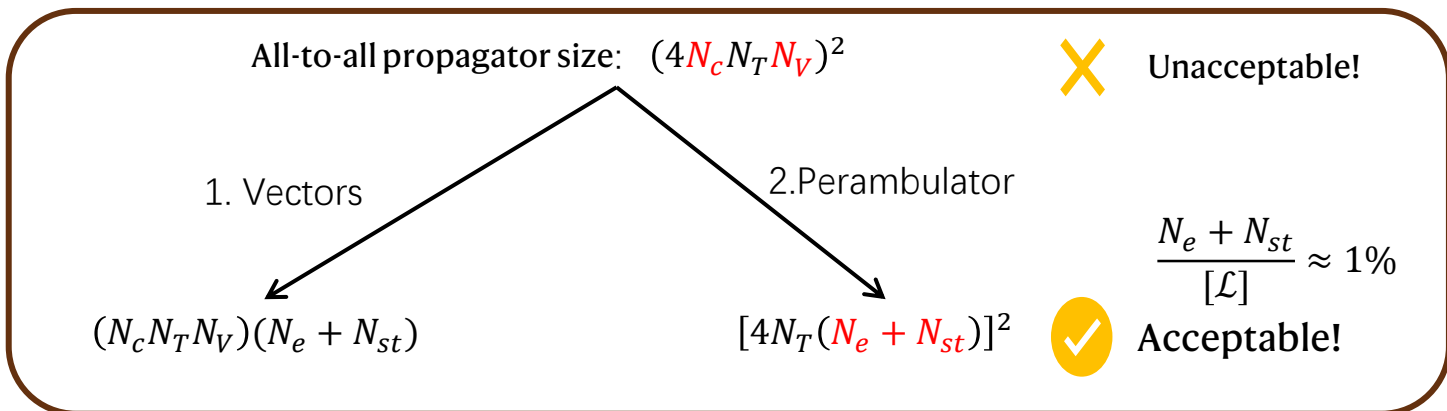
$$L = \{|\lambda\rangle; -\nabla^2|\phi\rangle = \lambda|\phi\rangle\}. \quad \text{PRD. 80, 054506 (2009)}$$

- Distillation space $\mathcal{L}_1$ is not complete, resulting in **bias of local current ME**;
- ***Blending* method**(BLD) can provide an unbiased estimation of identity matrix, based on samples in $\mathcal{L}_2$;

[Z.C. Hu et al. , arXiv: 2505.01719\(2025\)](#)

- ✓ **All-to-All fermion Propagator could be stored and reused.**

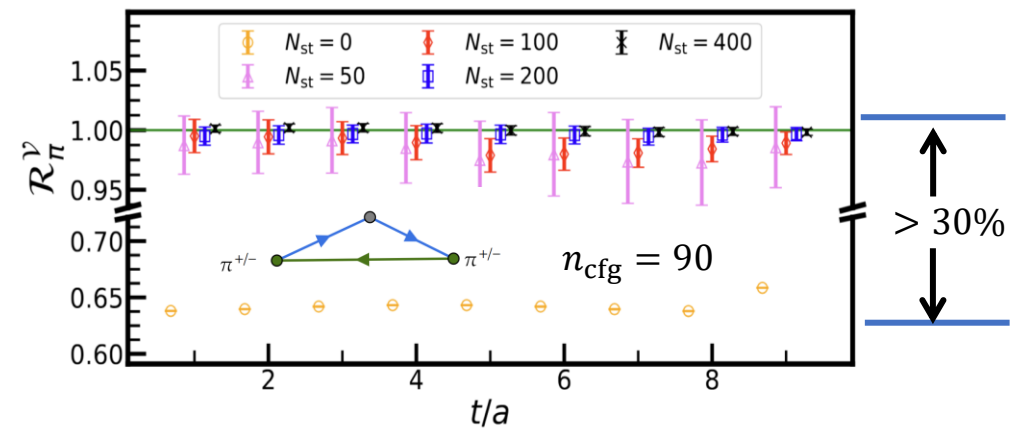
$$S_{ij}(t_1, t_2) = \int d^3x d^3y \langle \phi_i(\vec{x}, t_1) | S(\vec{x}, t_1; \vec{y}, t_2) | \phi_j(\vec{y}, t_2) \rangle,$$



Samples by Noise vector

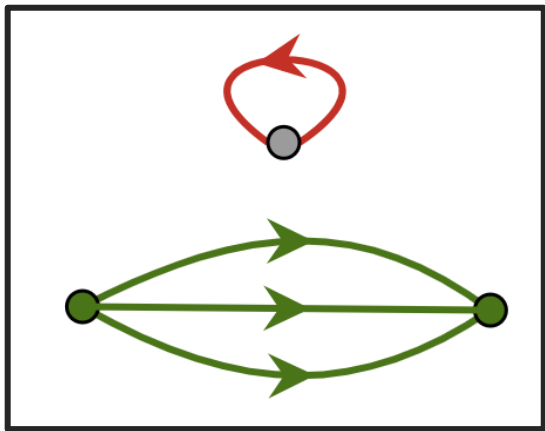
Conserved Current Test:

$$\langle H | V^c | H \rangle / \langle H | H \rangle = 1$$

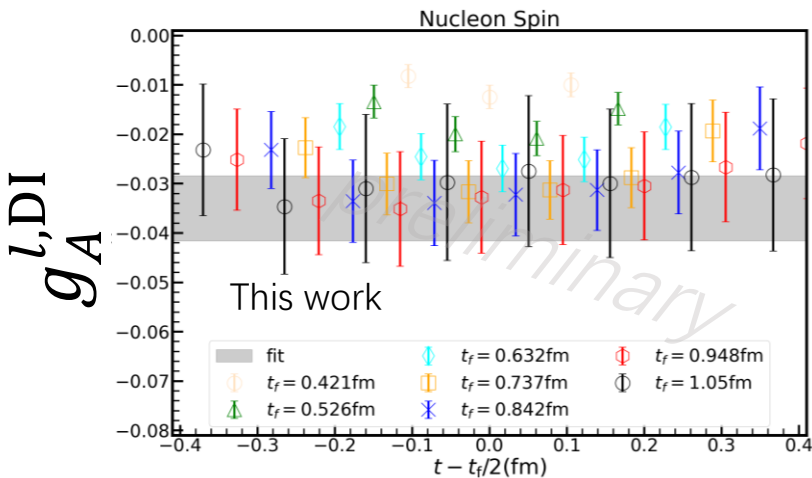
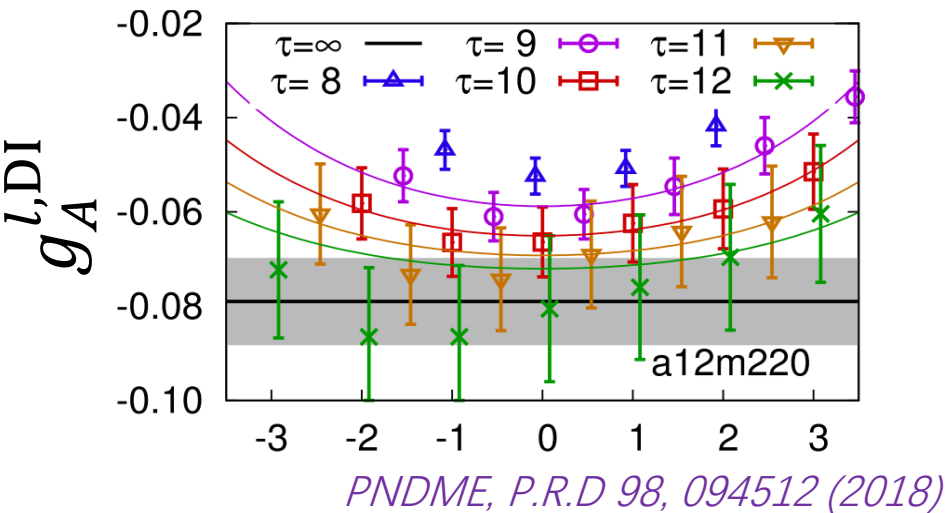


- ✓ $\langle \pi^+ | V^c | \pi^+ \rangle = 1.0002(29)$ with $N_{st} = 400$, $m_\pi \approx 290\text{MeV}$;

2: Theoretical Framework: Blending Method



- a) A better ground-state convergence!
- b) The cost could be control very well.
- c) Data of fermion Propagator could be stored and reused.



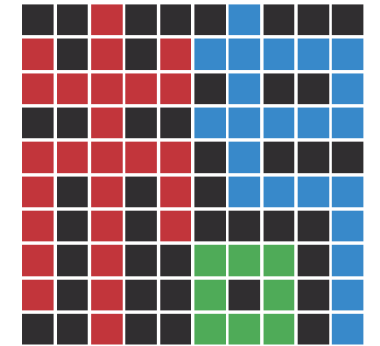
	Ensemble	m_π	$L^3 \times T$	n_{cfg}	n_{inv}	$g_A^{l,DI}$	$g_A^{s,DI}$
PNDME Collaboration	a12m220	228MeV	32x64	958	11000	-0.075(9)	-0.037(6)
CLQCD(This work)	C32P23	227MeV	32x64	80	3200 [†]	-0.035(7)	-0.021(2)

2: Theoretical Framework: Lattice Setup



- 2+1 flavor tadpole-improved **Clover Fermion** with Symanzik gauge action;

	$a(\text{fm})$	$m_\pi(\text{MeV})$	$L_x \times L_t$	$m_\pi L$	n_{cfg}
C32P29	0.105	292	32x64	5.0	200
C32P23	0.105	227	32x64	3.9	80
C48P14	0.105	136	48x96	3.8	45
F48P30	0.077	303	48x96	5.7	40



CLQCD

*CLQCD, PRD **109**, 054507 (2024)*

*CLQCD, PRD **111**, 054504 (2025)*



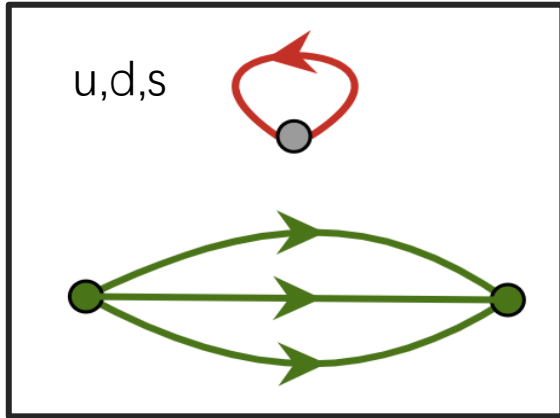
◆ Chapter 1: Introduction

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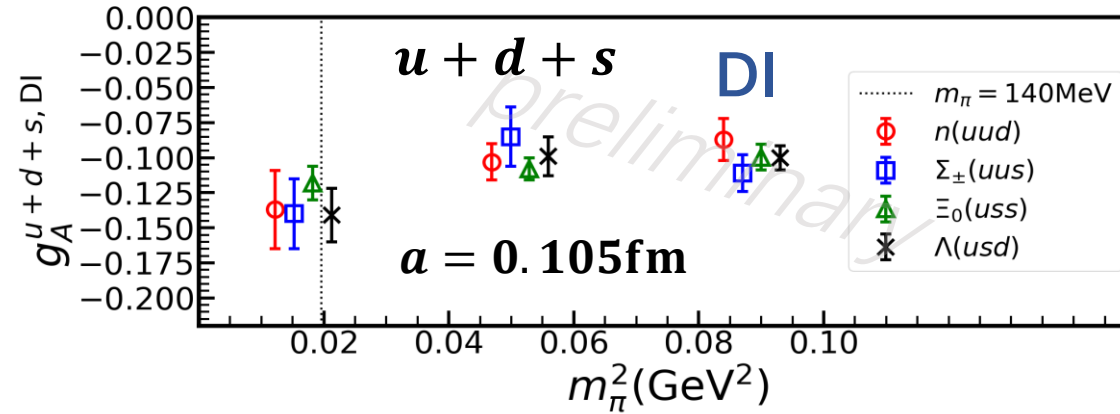
◆ **Chapter 3: Numerical result of Quark Spin and Discussion**

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3.1 Disconnected Insertion Contribution



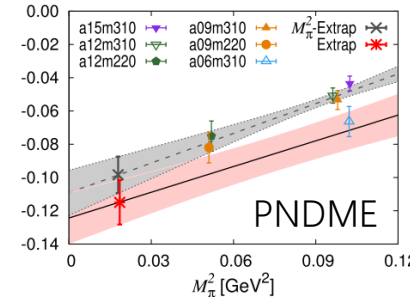
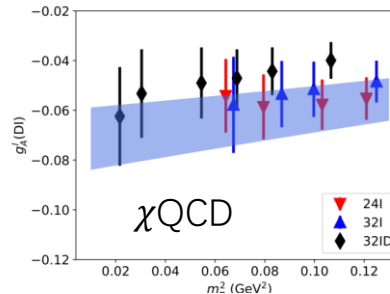
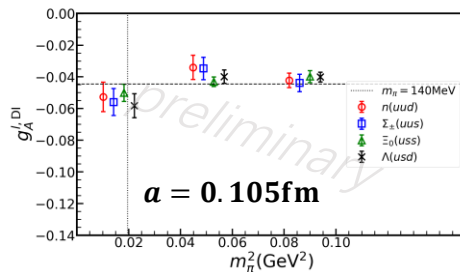
✓ The contribution of sea quarks is basically consistent in all octet baryons;



■ How about SU(3) symmetry for gluon helicity ΔG ?

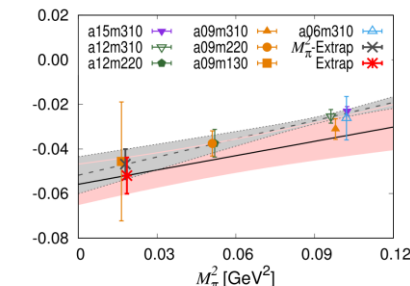
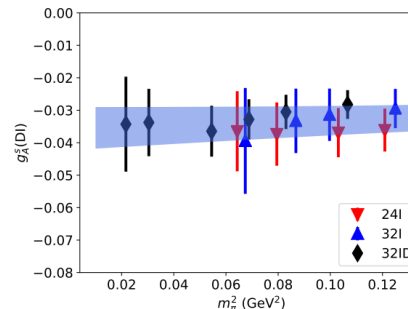
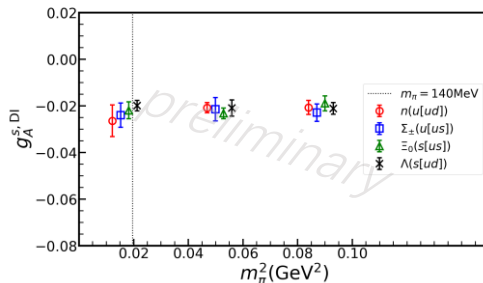
➤ We didn't see remarkable m_π dependence at for **disconnected quark insertion** contribution;

up/down



■ CLQCD result is more closed χ QCD collaboration, which is support for **weaker m_π dependence**.

strange



PNDME, PRD **98**, 094512 (2018)

χ QCD, PRD **98**, 074505 (2018)

3.2 The total singlet quark spin and SU(3)-flavour symmetry



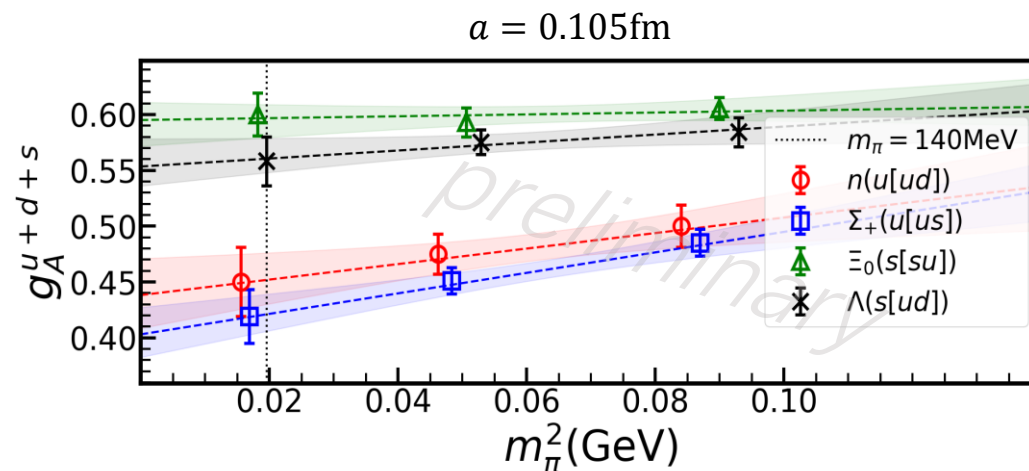
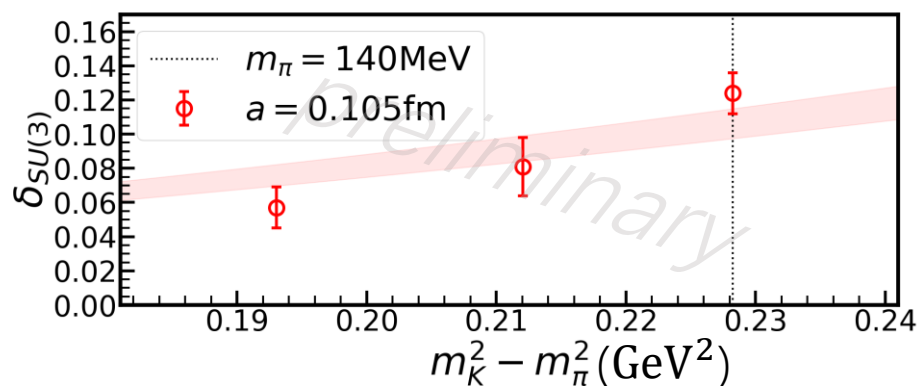
up/down dominant

$$\left\{ \begin{array}{l} n = u(u^T C \gamma^5 d) \\ \Sigma^+ = u(u^T C \gamma^5 s) \end{array} \right.$$

strange dominant

$$\left\{ \begin{array}{l} \Xi_0 = s(s^T C \gamma^5 u) \\ \Lambda_0 = 2s[u^T d] + d[u^T s] - u[d^T s] \end{array} \right.$$

$$\delta_{SU(3)} = g_{A(u-d)}^n - g_{A(u-d)}^\Sigma + g_{A(u-d)}^\Xi$$



- There is moderate **SU(3) symmetry breaking** effects for the **iso-vector axial** charges, $\delta_{SU(3)}/g_A^n \sim 10\%$;

- There are SU(3) symmetry breaking in quark spin at $a = 0.105\text{fm}$;
- SU(3) symmetry keep well **inside di-quark operator**;

$$n(u[ud]) \xrightleftharpoons[d]{s} \Sigma(u[us]) \quad \Xi(s[su]) \xrightleftharpoons[s]{d} \Lambda(s[ud])$$

3.2 The total singlet quark spin and SU(3)-flavour symmetry



- ◆ After **continuum extrapolation**, there is a tendency for the SU(3) symmetry to **partially recover**.

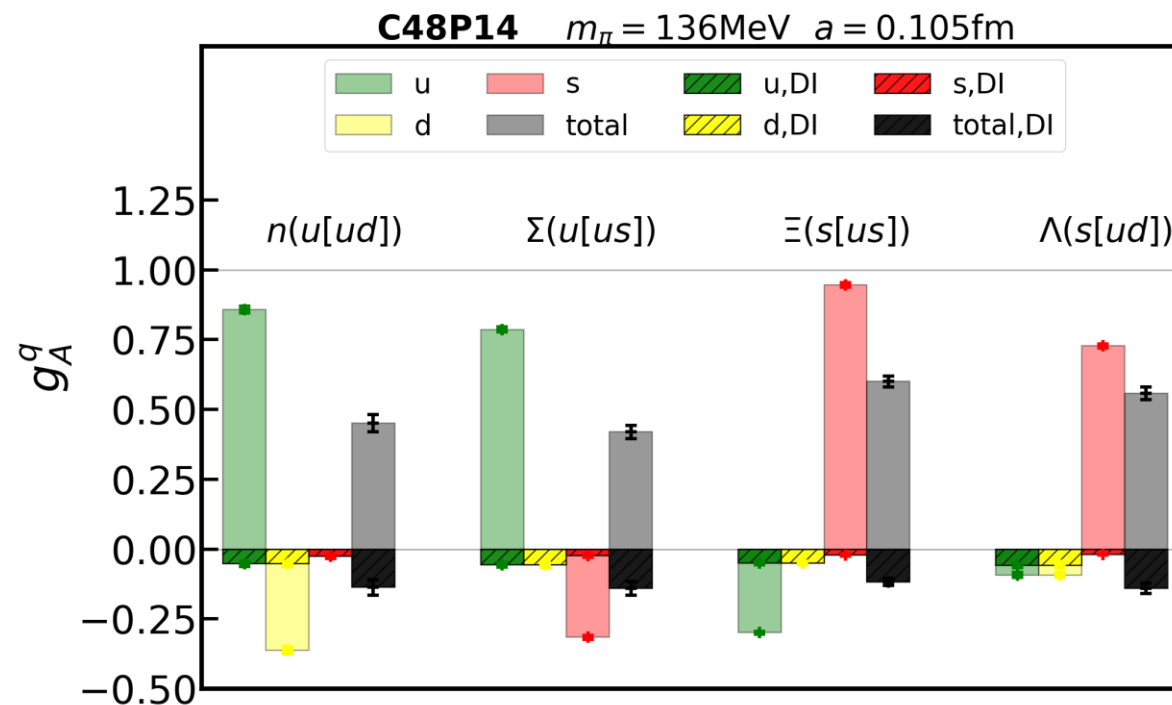
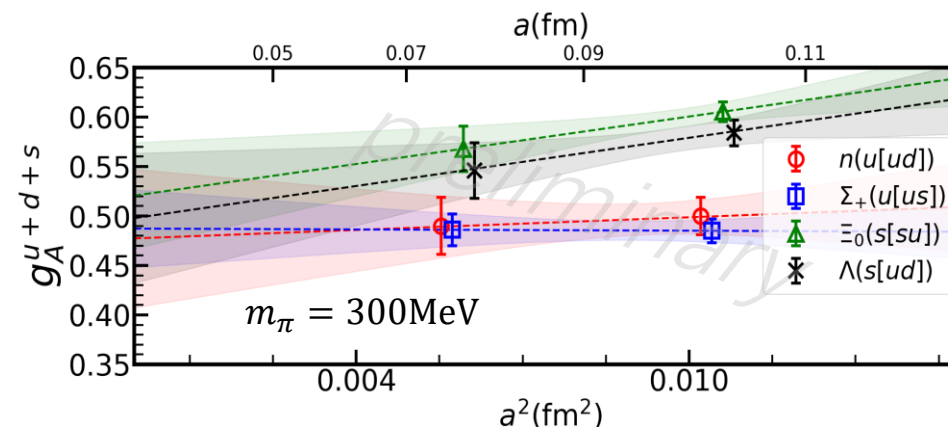
$$n: \frac{1}{2} \Delta q \sim 0.20 - 0.25$$

$$\Sigma: \frac{1}{2} \Delta q \sim 0.20 - 0.25$$

$$\Xi: \frac{1}{2} \Delta q \sim 0.25 - 0.30$$

$$\Lambda: \frac{1}{2} \Delta q \sim 0.25 - 0.30$$

- There is moderate **SU(3) symmetry breaking** effects for the **quark spin** $\sim 10(5)\%$;





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Conclusion

- 1) The **sea quark contribution is consistent** in Octet baryons ranging $m_\pi \in [135, 220, 300]\text{MeV}$;
- 2) SU(3)-flavour symmetry breaking has ~10% influence in **quark spin**;
- 3) The symmetry is better preserved within specific diquark structures, as seen in the pairs like, $(n, \Sigma) (\Lambda, \Xi)$;
- 4) There is a tendency for the SU(3) symmetry to **partially recover** after **continuum extrapolation**.

Outlook

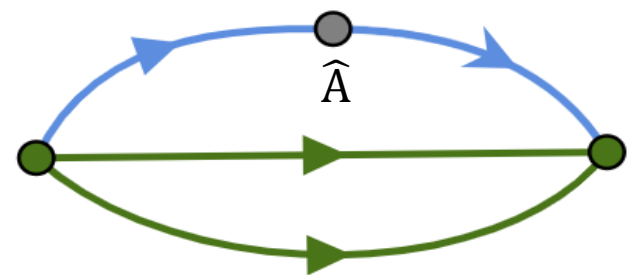
- 1) The chiral behavior of **sea quark helicity** should be checked carefully;
- 2) More systematical computation in different lattice spacing such as $a = 0.077\text{fm}$ and $a = 0.052\text{fm}$ will be done in the near future;

Thanks you!

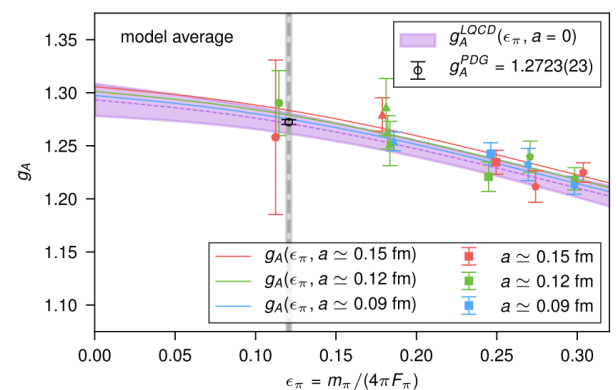


◆ Low-lying space of Laplace (Distillation) can get high precision nucleon external state without momentum;

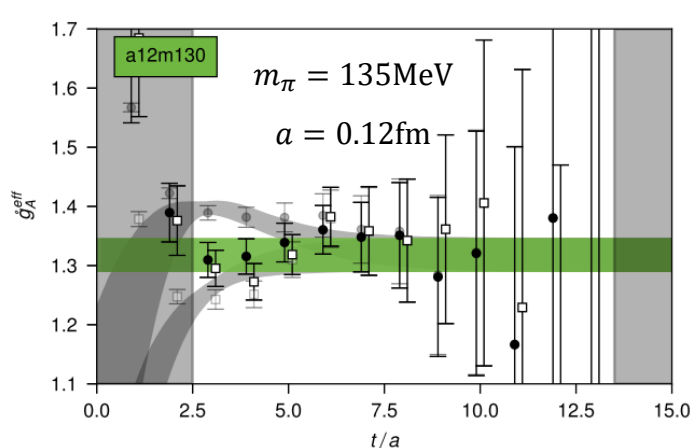
✓ We observe substantially improved results at the physical point($m_\pi = 135\text{MeV}$);



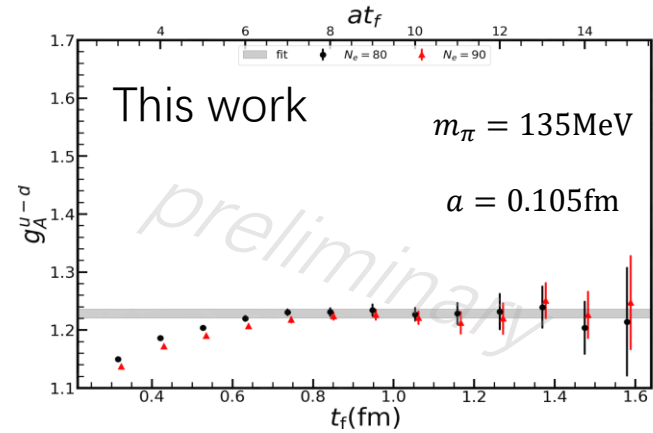
$\langle P|\bar{u}\gamma^5\gamma^i u - \bar{d}\gamma^5\gamma^i d|P\rangle$



Severe model dependent !!!



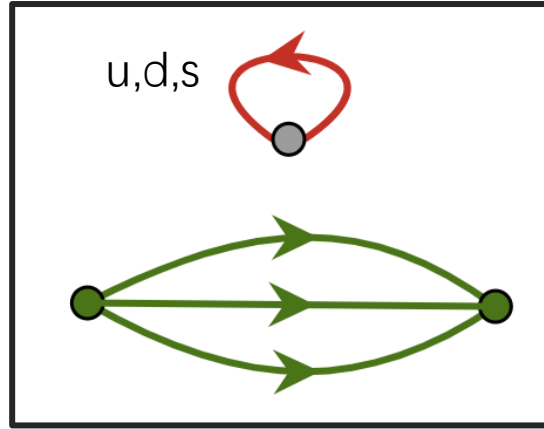
Nature 558,91 (2018)



	Ensemble	n_{cfg}	g_A^{u-d}	$\delta x/ x $
CLQCD(This work)	C48P14	46	1.221(07)	0.57%
CalLat	a15m130	1000	1.270(72)	5.66%
	a12m130	1000	1.292(30)	2.32%
RQCD [PRD 108, 034512 (2023)]	D452	1000	1.17(24)	20.5%

✓ It is expected to achieve experimental-level precision of 0.1% with an internationally comparable level of statistics (e.g., 1000 configurations);

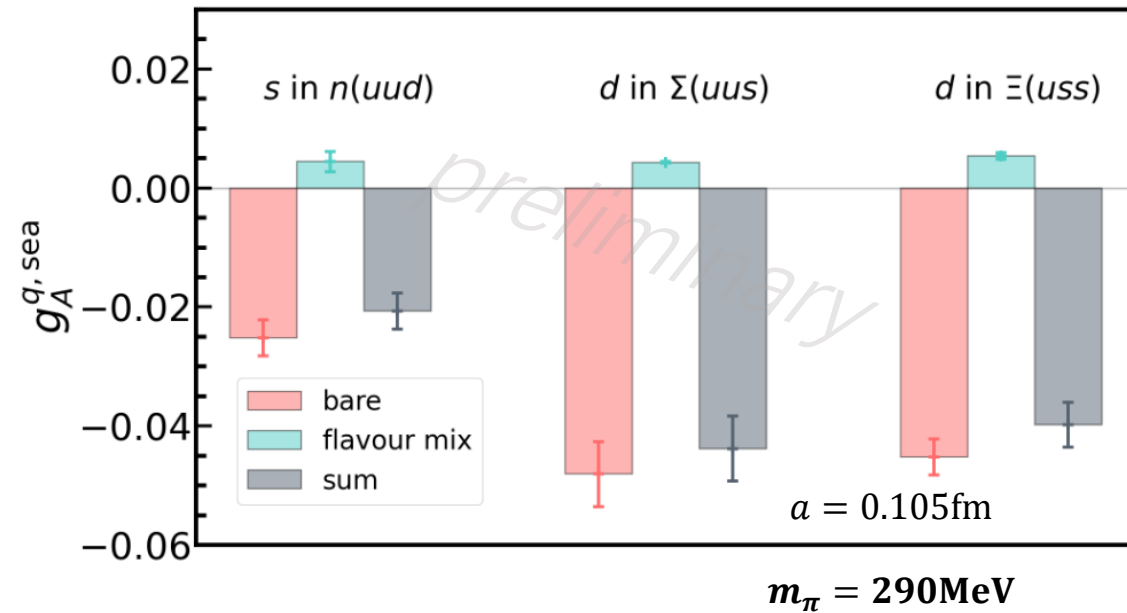
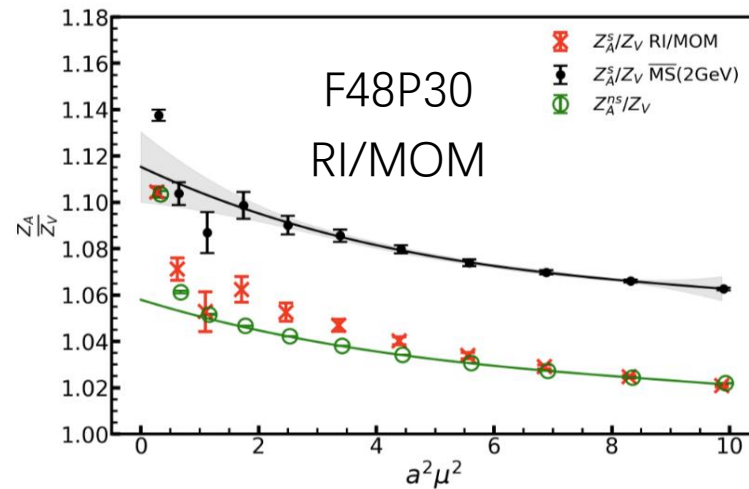
- ◆ Flavour Non-singlet and Singlet renormalized constant;



$$g_A^{q,\text{bare}} \equiv \langle N | \bar{q} \gamma_i \gamma_5 q | N \rangle;$$

Flavor Mixing:
$$g_A^{q,R} = Z_A^{\text{ns}} g_A^{q,\text{bare}} + \frac{Z_A^{\text{s}} - Z_A^{\text{ns}}}{N_f} \sum_{q=u,d,s} g_A^{q,\text{bare}},$$

A. J. Chambers et al. , PRD 92, 114517 (2015)



- ◆ Flavour mixing will lead to **1σ shift** at current precision **globally**;

- ◆ $\delta g_A^{\text{mix}} = \frac{Z_A^{\text{s}} - Z_A^{\text{ns}}}{N_f} \sum_q g_A^{q,B}$ couldn't be ignored for **strange quark in nucleon**, which has 20%~30% contribution; **P17**

