

Finite Volume Hamiltonian method for two-particle system

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based on JHEP04(2025)108



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conventional paradigm



possible challenges for Lüscher formula

1' Poisson formula : (for a regular function) summation – integration = $O(e^{-\alpha L})$, $\alpha \sim m_\pi$

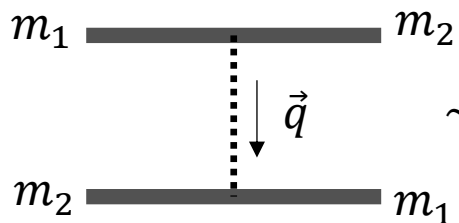
nonnegligible exponential ?

2' separation of on-shell contribution and analytical continuation below threshold

$$\frac{f(q)}{e^{-q^2/2\mu}} = \underbrace{\frac{f(q) - f(\sqrt{2\mu e})}{e^{-q^2/2\mu}}}_{\text{regular}} + \underbrace{\frac{f(\sqrt{2\mu e})}{e^{-q^2/2\mu}}}_{\text{on-shell}}$$

additional analytical properties below threshold ?

both issues matter when One-Pion-Exchange is important!

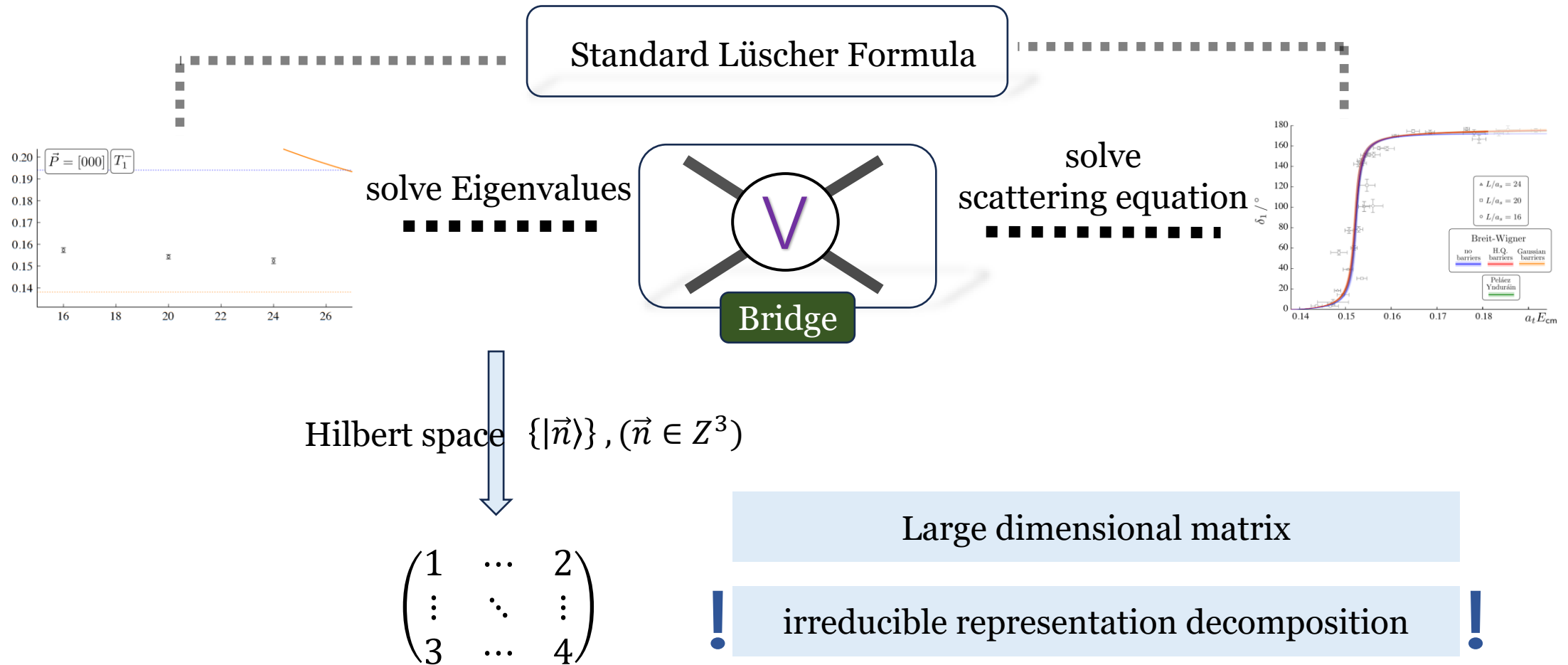


$$\sim \frac{1}{\vec{q}^2 + m_{\text{eff}}^2}, \quad m_{\text{eff}}^2 = m_\pi^2 - (m_2 - m_1)^2$$

$m_{\text{eff}} \ll m_\pi$ such that $m_{\text{eff}}L \ll 1$
left-hand cut below threshold

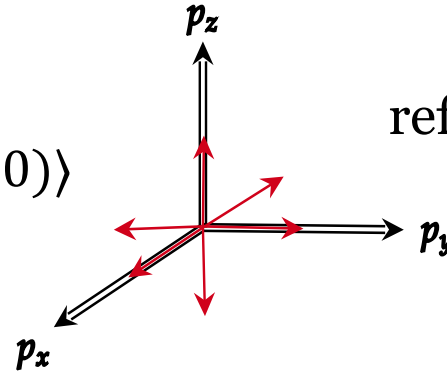
Finite Volume Hamiltonian method

a model-independent relationship if potential is SHORT-RANGED



Sketch : spinless system at rest

subspace: $|(0,0,\pm 1)\rangle, |(0,\pm 1,0)\rangle, |(\pm 1,0,0)\rangle$



reference momentum ,e.g. $\vec{n}_{\text{ref}} = |(0,0,1)\rangle$

irrep basis

$$|\vec{n}_{\text{ref}}; T_1^-, \mu = 1, r = 1\rangle \propto \sum_{g \in O_h} \sum_{v=1}^3 [c_{\vec{n}_{\text{ref}}}^{r=1}]_v D_{\mu=1,v}^{T_1^-}(g) |g\vec{n}_{\text{ref}}\rangle$$

$$I_{v'v} = \sum_{g \in LG(\vec{n}_{\text{ref}})} D^{T_1^-}(g) \text{ eigenvcs } [c^{r=1}], [c^{r=2}], \dots \text{ associated to nonzero eigenvals}$$

Hamiltonian element

$$\langle \vec{n}'_{\text{ref}}, r' | V_L^{T_1^-} | \vec{n}_{\text{ref}}, r \rangle = \left(\frac{2\pi}{L} \right)^3 \sqrt{\frac{|LG(\vec{n}_o)|}{|LG(\vec{n}'_o)|}} \sum_{g \in LC(\vec{n}_o)} \underbrace{[c_{\vec{n}'_{\text{ref}}}^{r'}] \cdot D^{T_1^-}(g) \cdot [c_{\vec{n}_{\text{ref}}}^r]}_{\text{kinematical}} \underbrace{V(\vec{n}', g\vec{n})}_{\text{dynamical}}$$

Sketch : PV system at rest

Hilbert space $\{|\vec{n}, \sigma\rangle\}$, $(\vec{n} \in Z^3, \sigma = 0, \pm 1)$ polarization would mix.

convert to helicity representation $|\vec{n}, \lambda\rangle = \sum_{\sigma} (e^{-iJ_z\phi_{\vec{n}}} e^{-iJ_y\theta_{\vec{n}}})_{\sigma\lambda} |\vec{n}, \sigma\rangle$

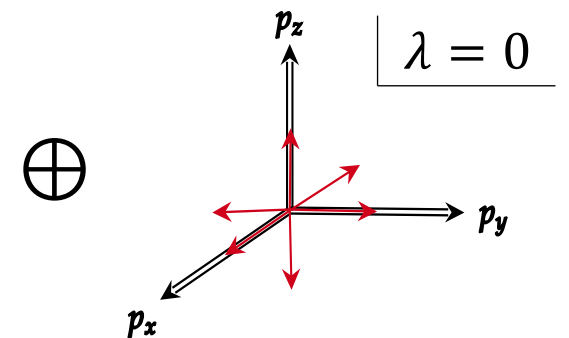
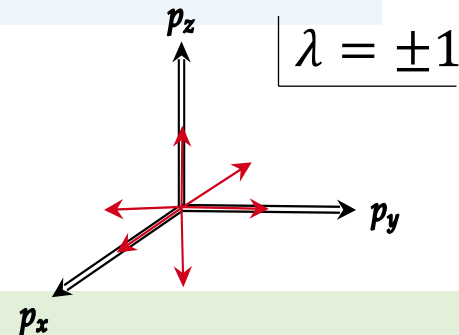
$$g|\vec{n}, \lambda\rangle = e^{-i\lambda\phi_w} |g\vec{n}, \lambda\rangle, \quad g \in O_h^+$$

$$P|\vec{n}, \lambda\rangle = |-\vec{n}, -\lambda\rangle$$

different helicity $|\lambda|$ would not mix, as if spinless !

Sketch : others

- moving system : only substitution $O_h \rightarrow C_{4v}, C_{3v}, C_{2v}$ is needed
- fermionic system: take care of double cover group O_h^2
- isospin symmetry : incorporation of permutation group



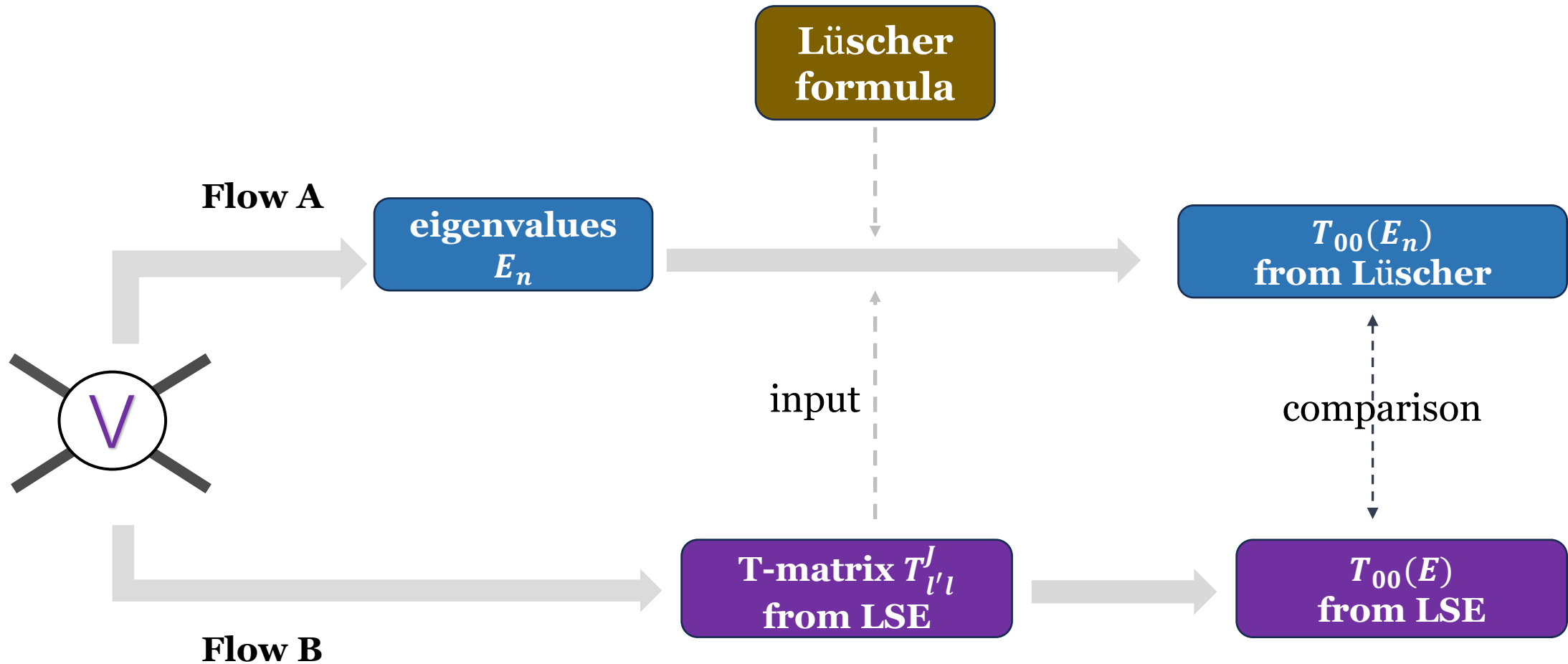
Application: Toy Model and issues from left-hand cut

$$\mathbf{PV} \rightarrow \mathbf{PV} \quad V_{\lambda'\lambda}(\vec{p}, \vec{k}) = \frac{1}{(2\pi)^3} \delta_{\lambda'\lambda} V(\vec{p}, \vec{k})$$

- Model 1: $V(\vec{p}, \vec{k}) = c_S \left(\frac{\Lambda^2}{\vec{k}^2 + \Lambda^2} \right)^2 \left(\frac{\Lambda^2}{\vec{p}^2 + \Lambda^2} \right)^2$ S-wave short
- Model 2: $V(\vec{p}, \vec{k}) = c_D \left(\frac{1}{3} (\vec{p} \cdot \vec{k})^2 - \vec{p}^2 \vec{k}^2 \right) \frac{1}{\Lambda^4} \left(\frac{\Lambda^2}{\vec{k}^2 + \Lambda^2} \right)^4 \left(\frac{\Lambda^2}{\vec{p}^2 + \Lambda^2} \right)^4$ D-wave short
- Model 3: $V(\vec{p}, \vec{k}) = c_S \left(\frac{\Lambda^2}{\vec{k}^2 + \Lambda^2} \right)^2 \left(\frac{\Lambda^2}{\vec{p}^2 + \Lambda^2} \right)^2 + \frac{1}{2} \int d \cos \theta \frac{-g m_{\text{eff}}^2}{(\vec{p} + \vec{k})^2 + m_{\text{eff}}^2}$ S-wave short + S-wave long
- Model 4: $V(\vec{p}, \vec{k}) = c_S \left(\frac{\Lambda^2}{\vec{k}^2 + \Lambda^2} \right)^2 \left(\frac{\Lambda^2}{\vec{p}^2 + \Lambda^2} \right)^2 + \frac{-g m_{\text{eff}}^2}{(\vec{p} + \vec{k})^2 + m_{\text{eff}}^2}$ S-wave short + full long

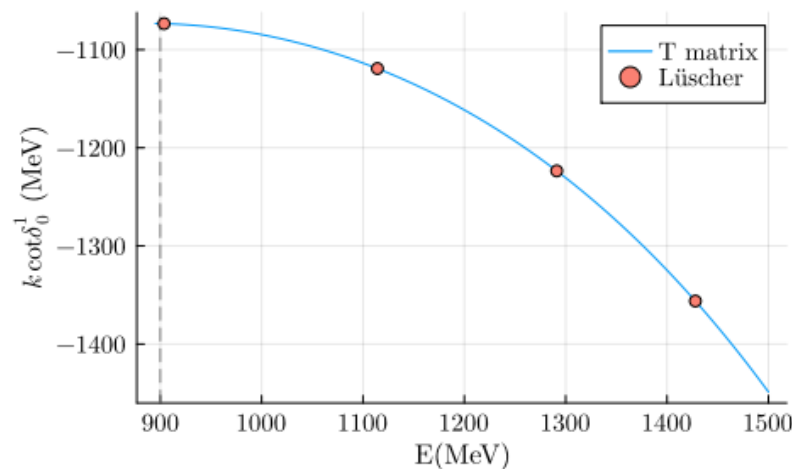
$$m_{\text{eff}} L = 2$$

Application: Toy Model and issues from left-hand cut

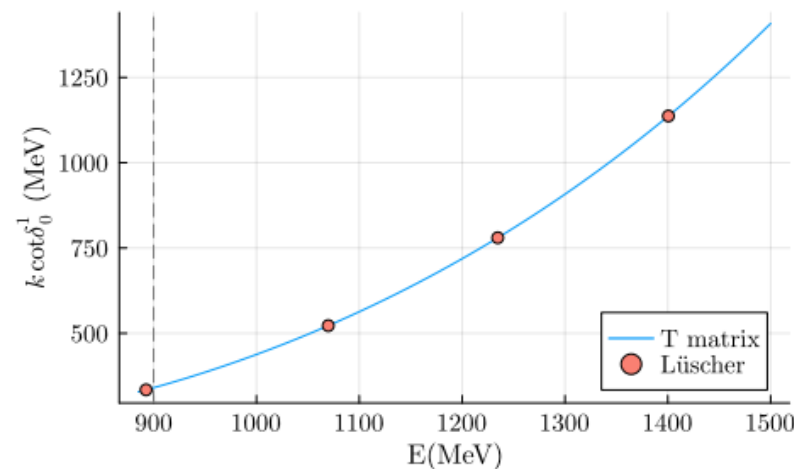


Application: Toy Model 1 and Model 2

pure S-wave short-range

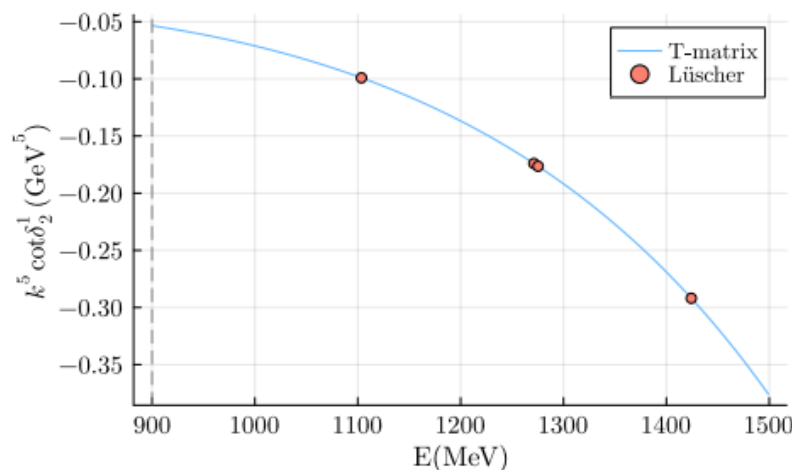


(a) Model I: repulsive potential: $c_S > 0$

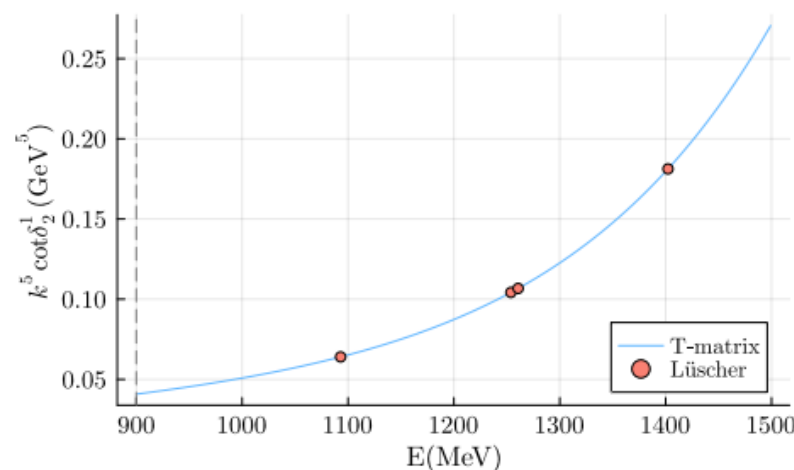


(b) Model I: attractive potential: $c_S < 0$

pure D-wave short-range



(c) Model II: repulsive potential: $c_D > 0$

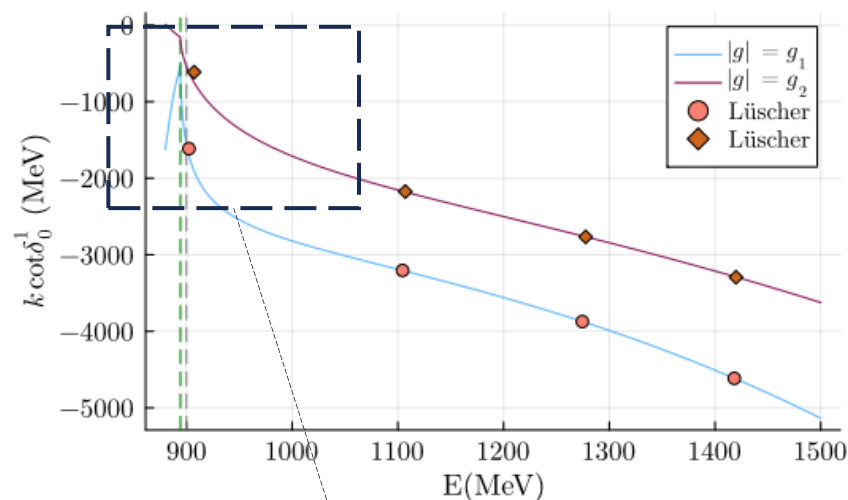


(d) Model II: attractive potential: $c_D < 0$

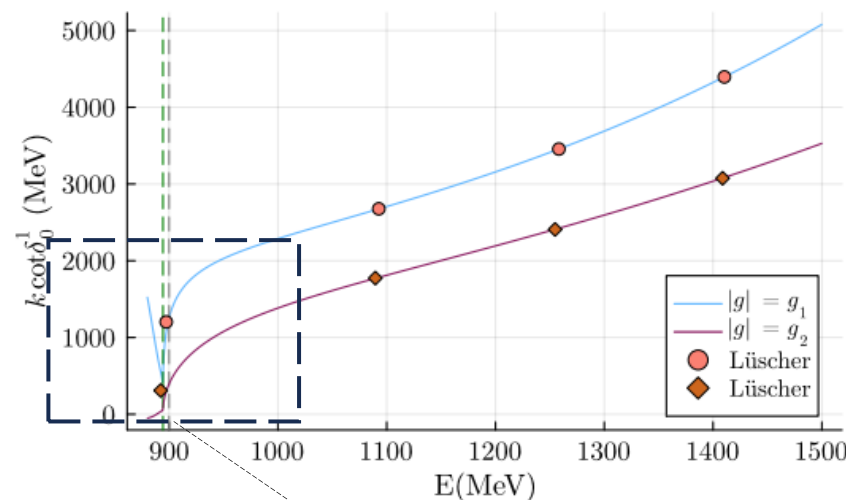
Application: Toy Model 3

pure S-wave short + pure S-wave long

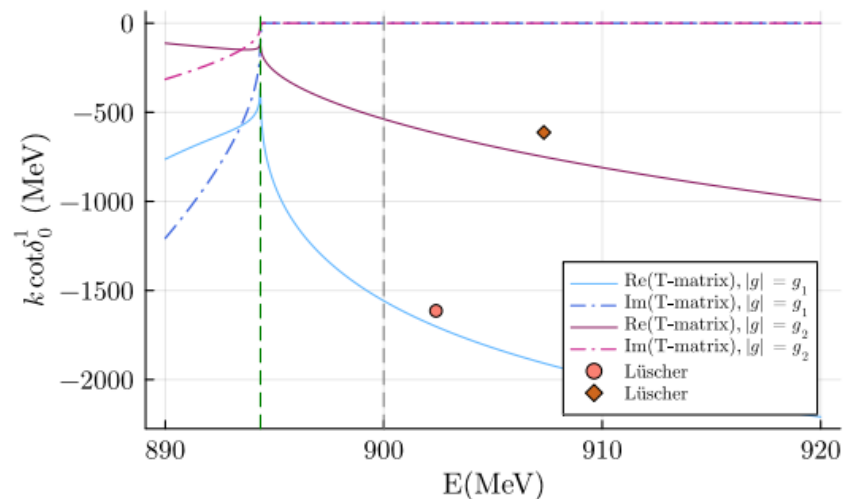
$$m_{\text{eff}}L = 2$$



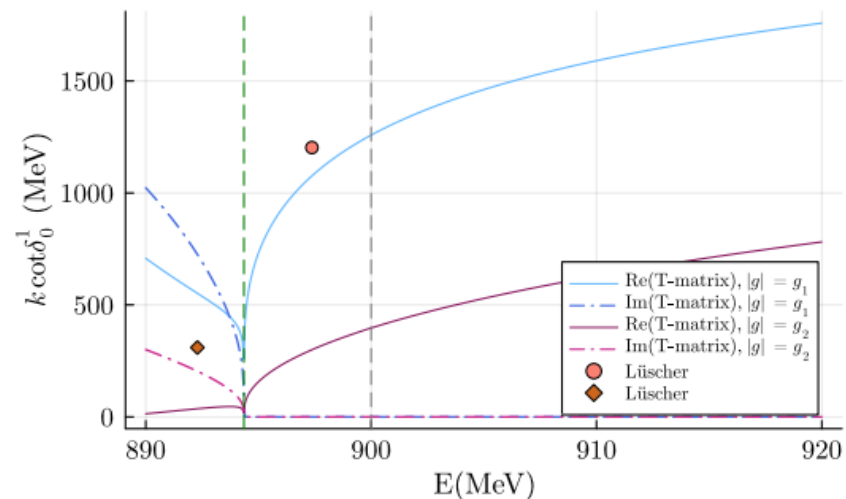
(a) repulsive potential: $c_S > 0, g < 0$



(b) attractive potential: $c_S < 0, g > 0$



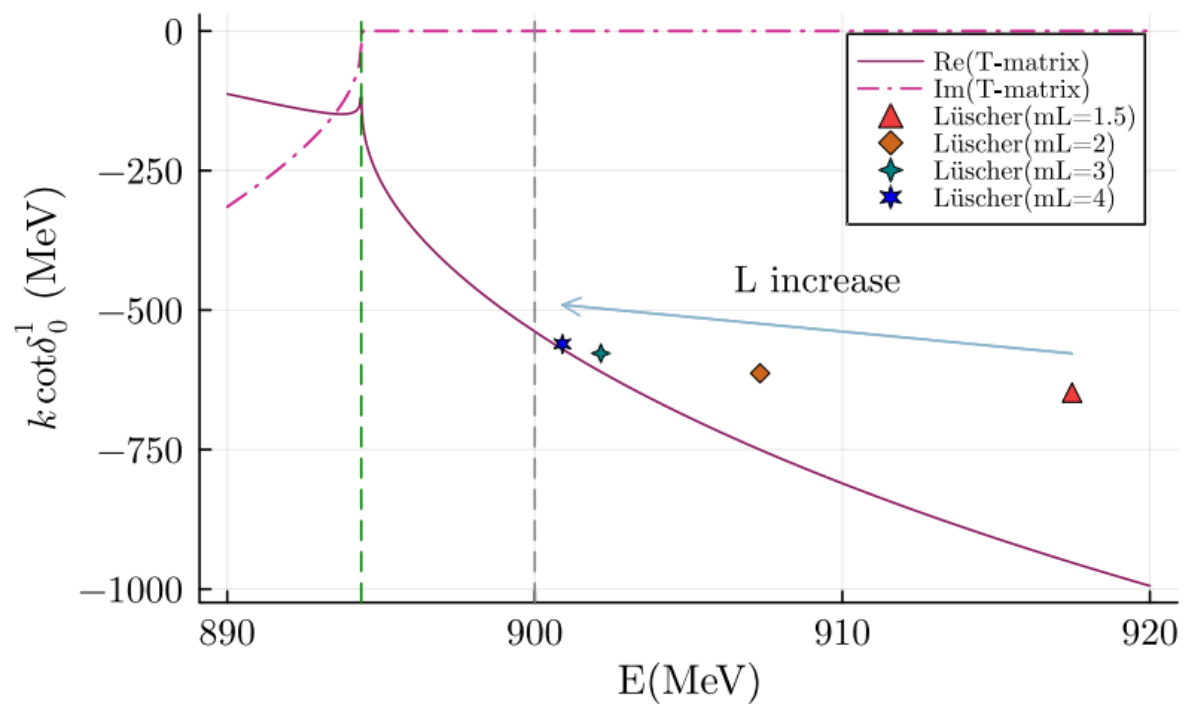
(c) detail near the threshold for (a)



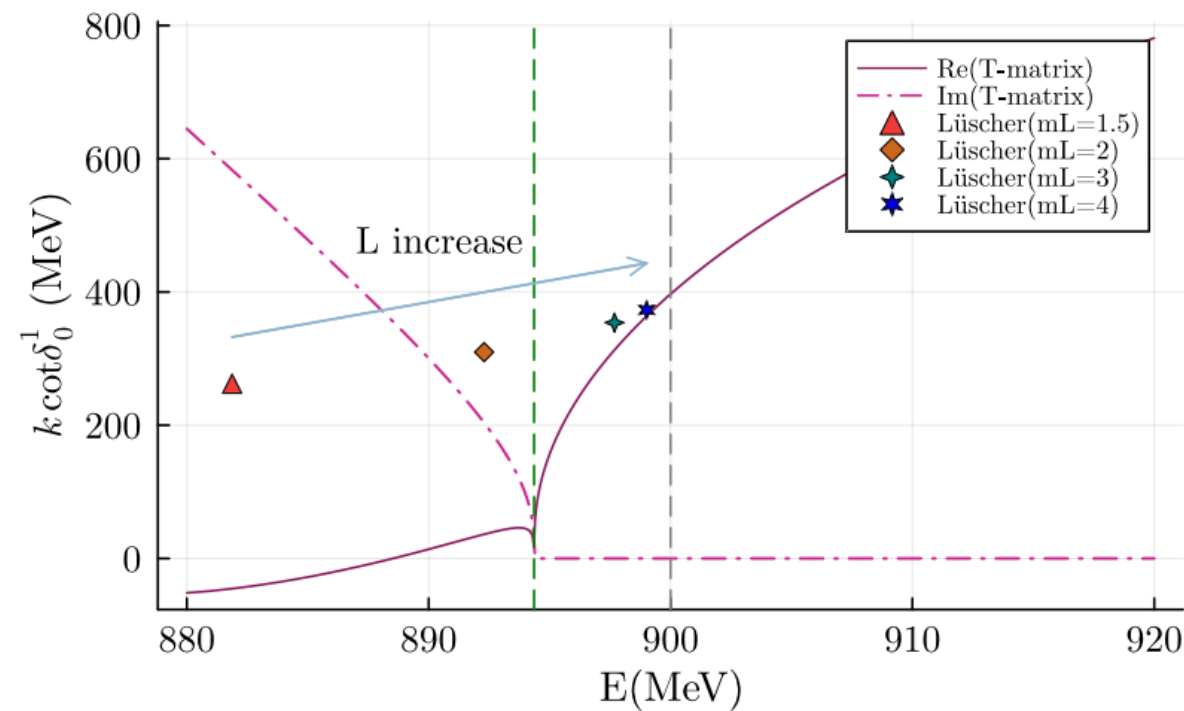
(d) detail near the threshold for (b)

Blue: weaker g
Red: stronger g

Application: Toy Model 3



(a) repulsive potential: $c_S > 0, g < 0$

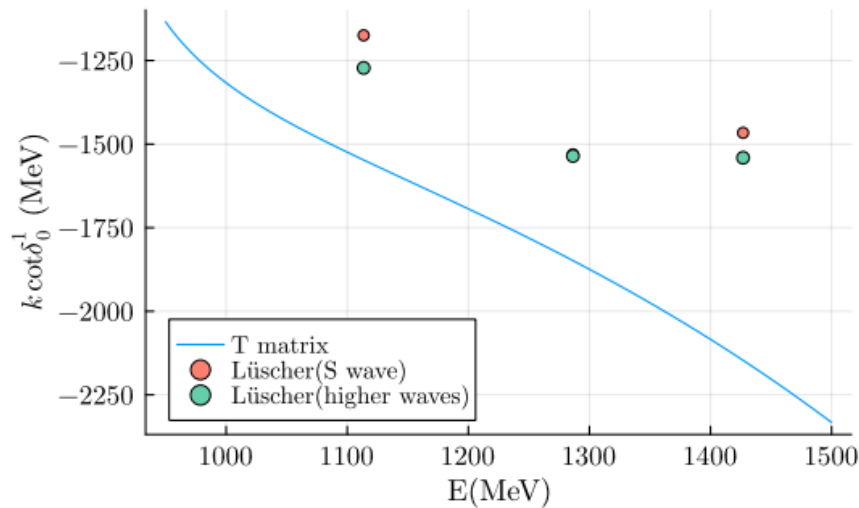


(b) attractive potential: $c_S < 0, g > 0$

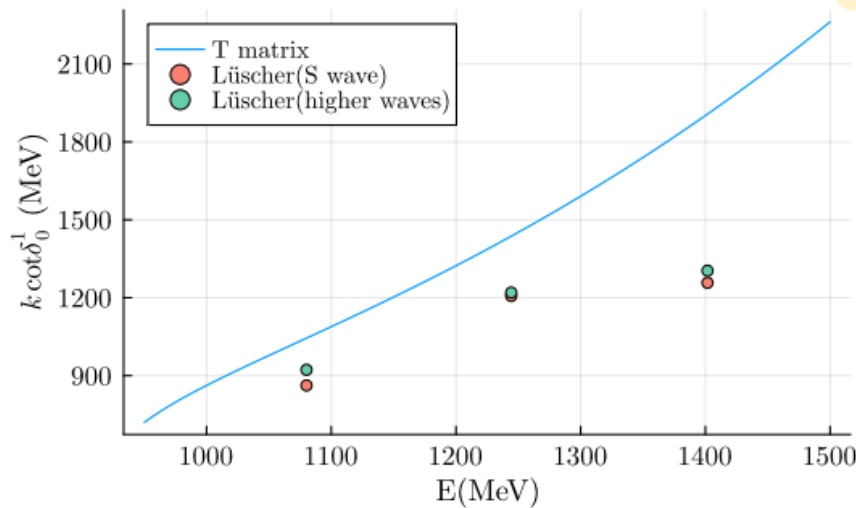
Application: Toy Model 4

pure S-wave short + full long

$$m_{\text{eff}}L = 2$$



(c) repulsive potential: $c_S > 0, g < 0$



(d) attractive potential: $c_S < 0, g > 0$

- Luscher formula (only S-wave)
- Luscher formula (input some higher $T_{l'l}^{J \leq 5}$)

nucl-th/0104027 W. Glockle et al.

wave sums for increasing j_{max} . It should be noted that the convergence of the partial wave sums is especially slow for the backward angle differential cross section. A sum up to $j_{\text{max}}=16$ is needed to obtain convergence within 1 %. In Fig. 4 A_y is displayed, which

Application: Toy Model

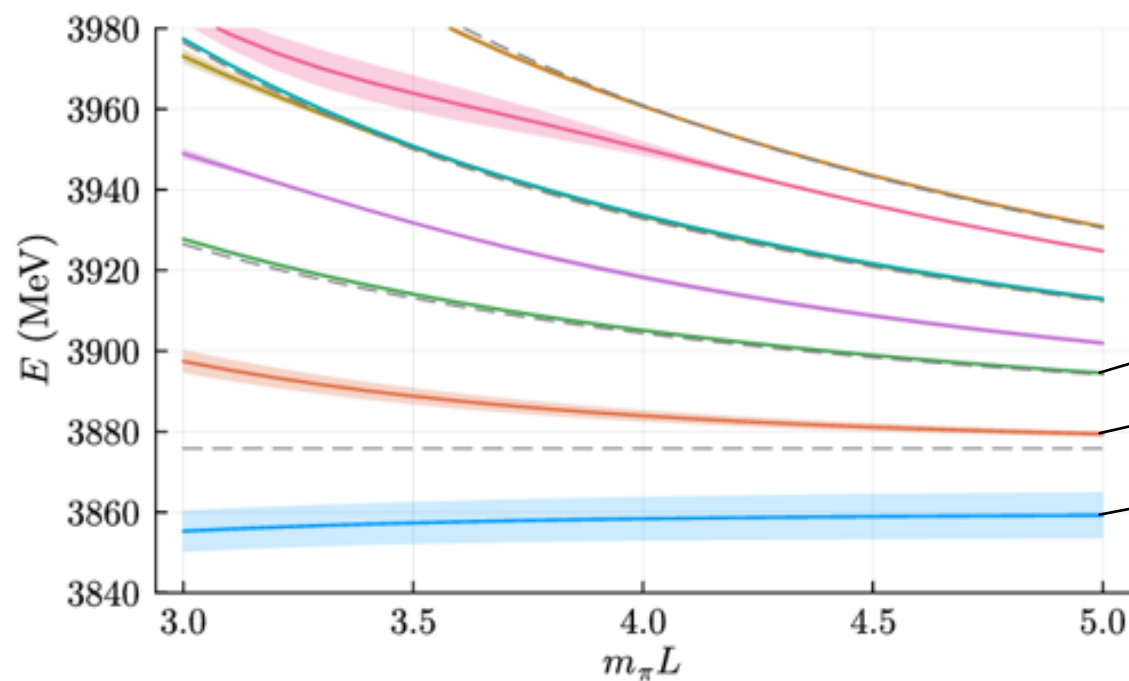
Universal issues

- **For short-ranged system, FVH method is equivalent to Lüscher formula**
- **When long-range potential is important, then**
 - 1' Lüscher formula cannot be applied below left-hand cut**
 - 2' Finite volume correction may be under-estimated**
 - 3' contribution from high partial-wave may be significant**

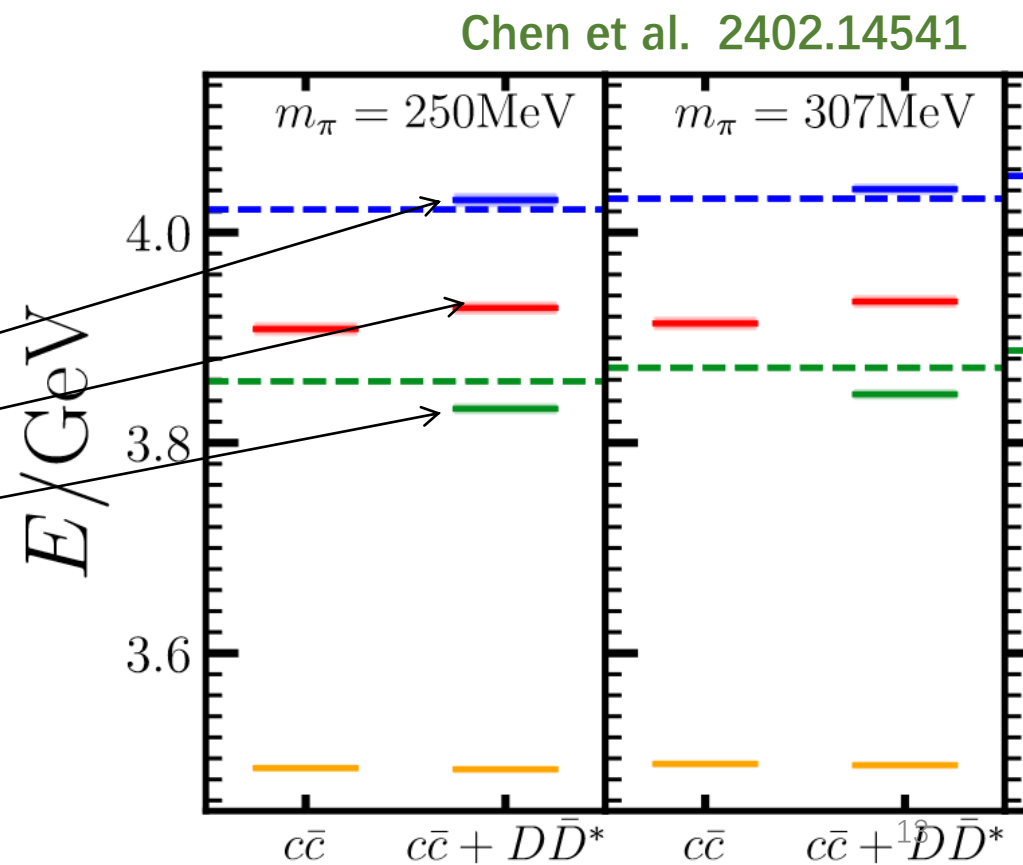
FVH method can resolve them simultaneously.

Application: X(3872)

$$V = V^{D\bar{D}^*} + V^{c\bar{c}} \quad \begin{array}{l} \text{short-range part } V^{c\bar{c}} : 3P0 \text{ model} \\ \text{long-range part } V^{D\bar{D}^*} : \text{OPE at hadronic level} \end{array}$$



(a) full model

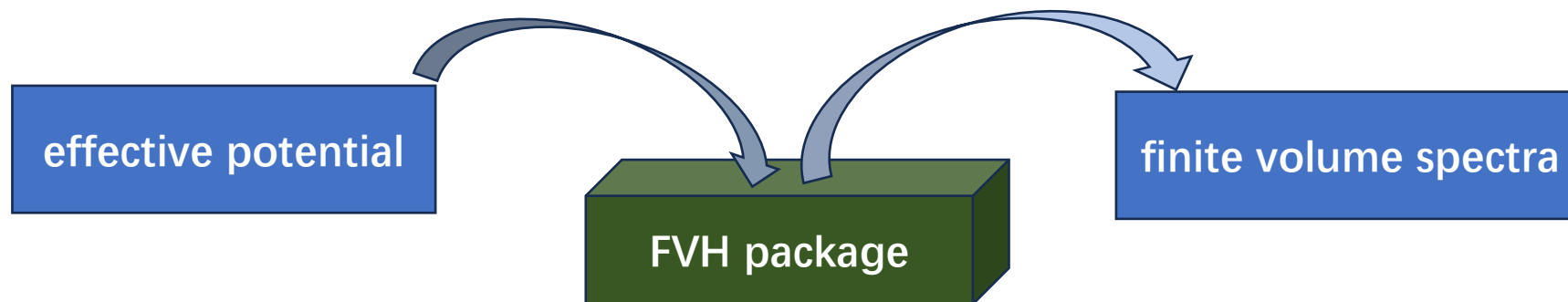


Prospect

1' a practical method to analyze lattice spectra of system whose long-range dynamic may be significant

- DD^* scattering (T_{cc}) see Xie Jia-Jun's report at 14:50 tomorrow
- NN scattering (Almost done...)

2' package



3' Extension to three-particle system

- $\omega \rightarrow 3\pi$ system (Almost done...)
- Roper ($N^*(1440)$) resonance (in preparation...)

