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Alexander von
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Probing Heavy Axion-Like Particles from Massive Stars with X-rays and Gamma rays

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based on

[1] arXiv: 2412.21163 with James Buckley, Francesc Ferrer and Takuya Okawa

[2] *In preparation* with Chris Cappiello, Takuya Okawa and Soebur Razzaque

The Fourth International Conference on Axion Physics and Experiment (Axion 2025)

Nanjing, China

July 28, 2025

ALP Motivation

- Axions were originally proposed to solve the strong CP problem. [Peccei, Quinn (PRL '77); Weinberg (PRL '78); Wilczek (PRL '78)]
- Axion-like particles (ALPs) are ubiquitous in many BSM constructions.
- String theory provides a strong motivation. [Svrcek, Witten (hep-th/0605206); Arvanitaki, Dimopoulos, Dubovsky, Kaloper, March-Russell (0905.4720)]
- In this talk, we only focus on ALP couplings to photons: $\mathcal{L}_{a\gamma} = -\frac{g_{a\gamma}}{4} a F_{\mu\nu} \widetilde{F}^{\mu\nu}$
- Consider heavy ALPs in the keV-MeV range. Not a dark matter candidate.

$$\tau_a = \frac{64\pi}{g_{a\gamma}^2 m_a^3} \simeq 10^{17} \text{ s} \left(\frac{10^{-12} \text{ GeV}^{-1}}{g_{a\gamma}} \right)^2 \left(\frac{10 \text{ keV}}{m_a} \right)^3$$

- Astrophysical environments provide an ideal testing ground. [Raffelt (Phys. Rep. '90)]



Stellar Plasma as ALP Laboratory

Main sequence stars (e.g. Sun)

$$R \sim 1R_{\odot}, T_{\text{core}} \sim \mathcal{O}(\text{keV})$$

[Nguyen, Tanin, Kamionkowski, 2307.11216]

Red giants, horizontal branch stars

$$R \sim \mathcal{O}(10^{2-3} R_{\odot}), T_{\text{core}} \sim \mathcal{O}(10 \text{ keV})$$

[Capozzi, Raffelt, 2007.03694]

White dwarfs

$$R \sim \mathcal{O}(10^{3-4} \text{ km})$$

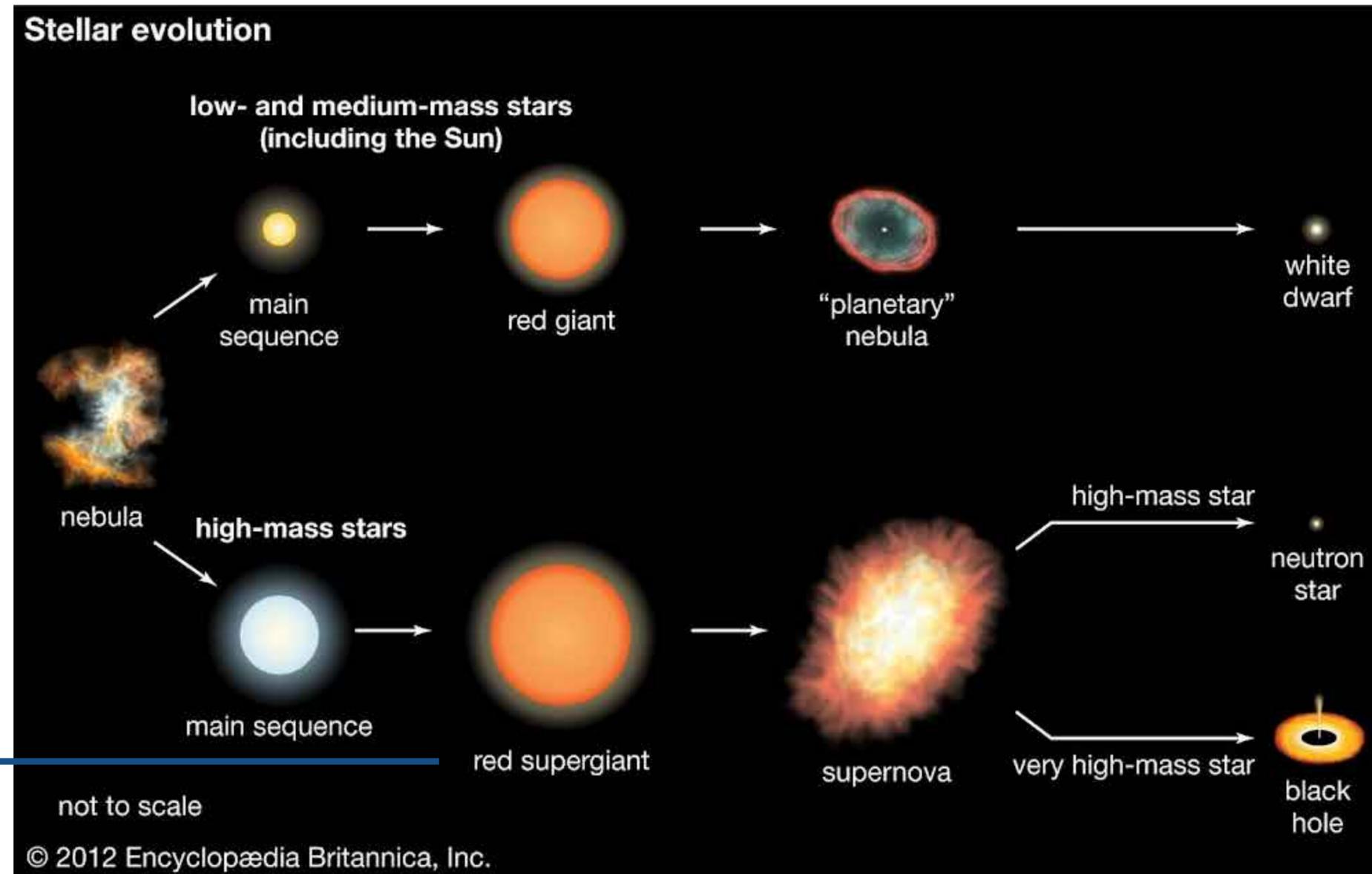
$$T_{\text{core}} \lesssim \mathcal{O}(\text{keV})$$

[Dolan, Hiskens, Volkas, 2102.00379]

Wolf-Rayet stars

$$T_{\text{core}} \simeq 20 \text{ keV}, R \sim 1 - 2R_{\odot}$$

(2412.21163, this talk)



Supernovae,
Neutron star mergers

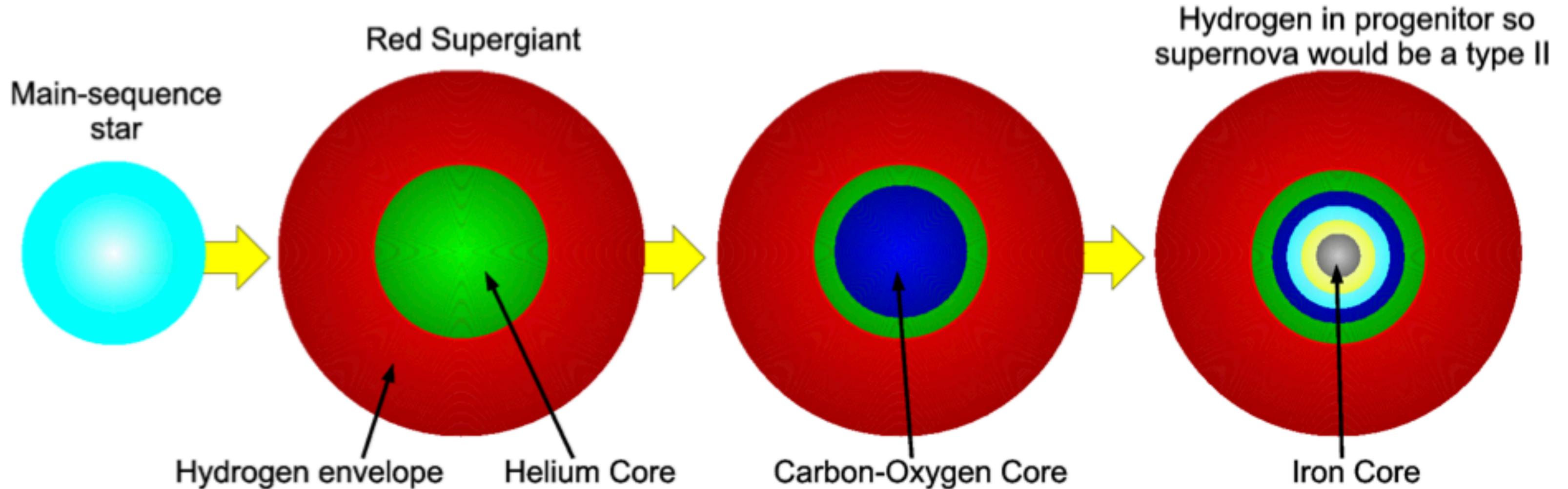
$$R \sim \mathcal{O}(10 \text{ km}),$$

$$T_{\text{core}} \sim \mathcal{O}(10 \text{ MeV})$$

[Muller et al, 2304.01060;
BD et al, 2305.01002;
Diamond et al, 2305.10327]

Wolf-Rayet Stars

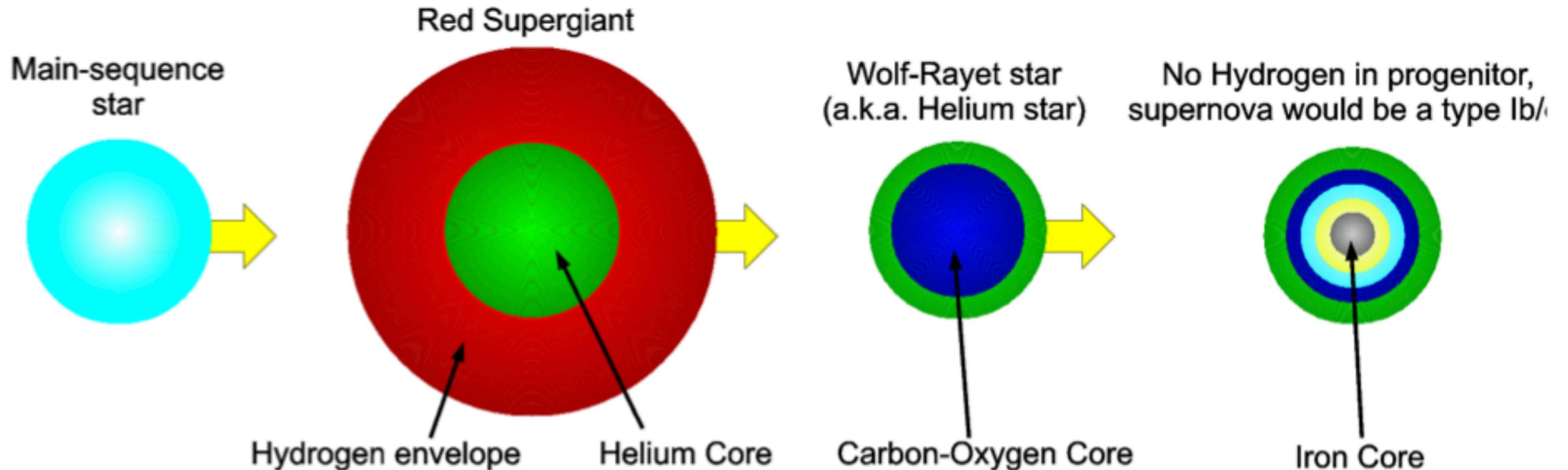
[Eldridge, 0809.1245]



Typical fate of a massive star ($M \gtrsim 8M_{\odot}$). Ends up as type-II supernova.

Wolf-Rayet Stars

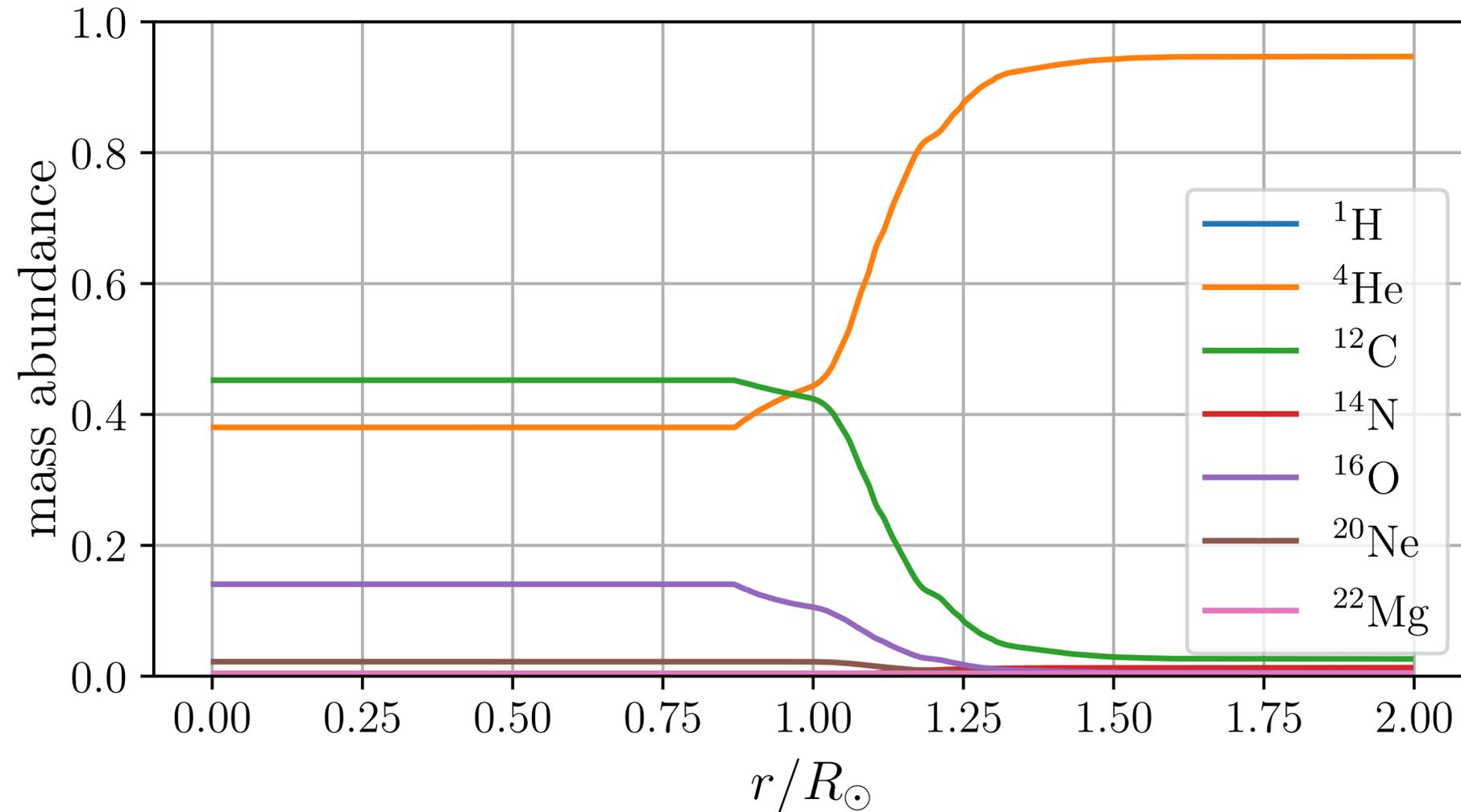
[Eldridge, 0809.1245]



Exception: High-mass stars $> 25M_{\odot}$ with large rotational velocity.

Strong stellar wind (or binary interaction) removes the Hydrogen envelope.

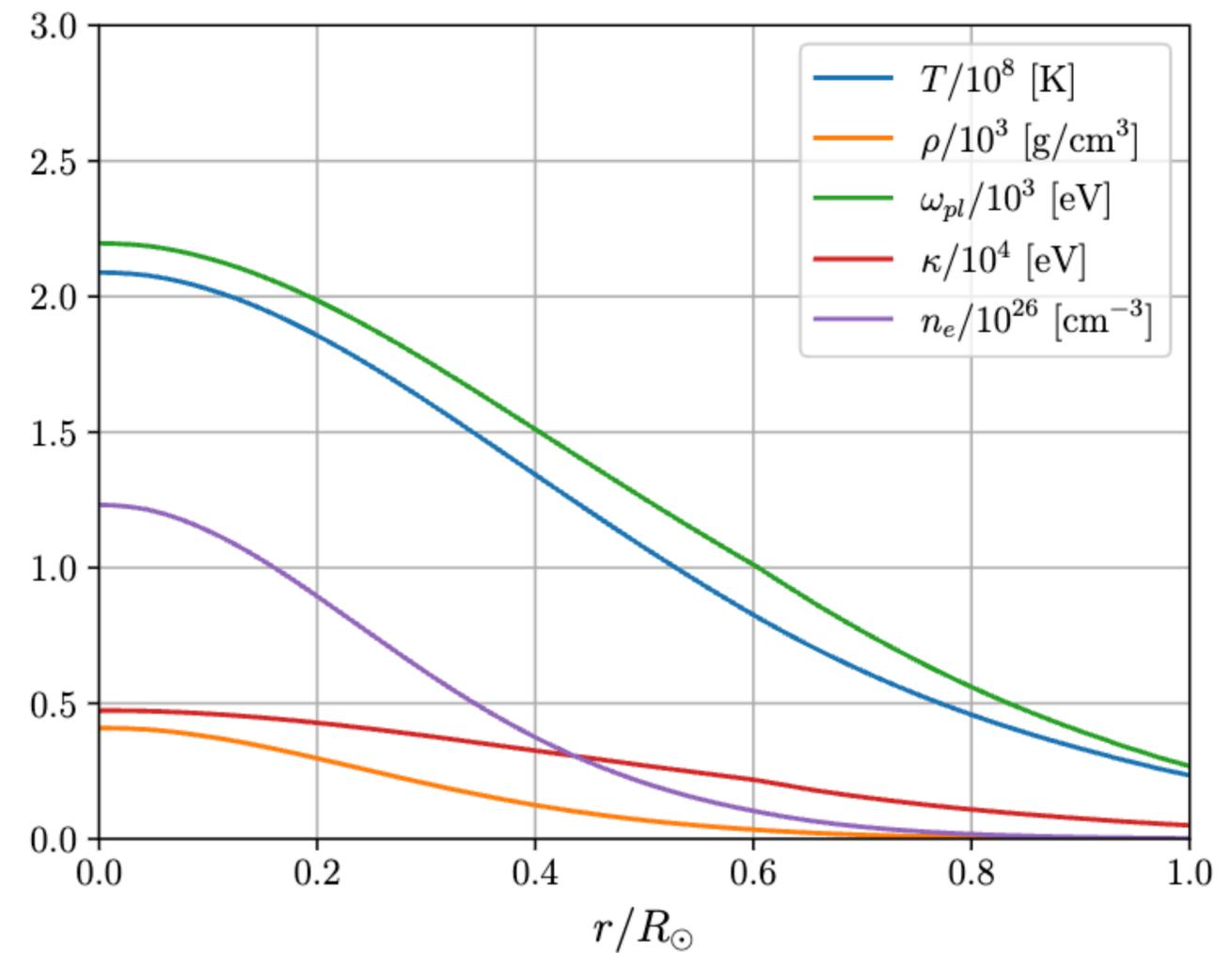
Wolf-Rayet Stars



$$T_{\text{core}} \simeq 20 \text{ keV}, R \sim 1 - 2R_{\odot}$$

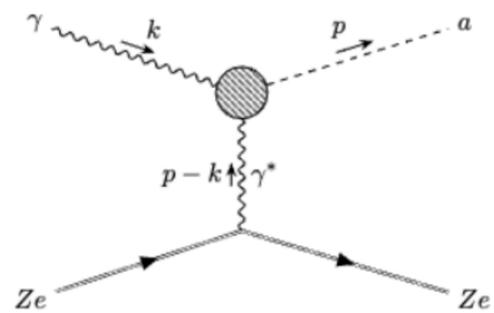
Efficient ALP production and high escape probability

Generated using full stellar evolution implemented in MESA code [1009.1622]

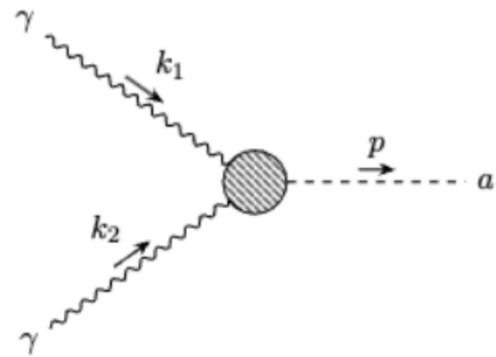


ALP Production in WR Stars

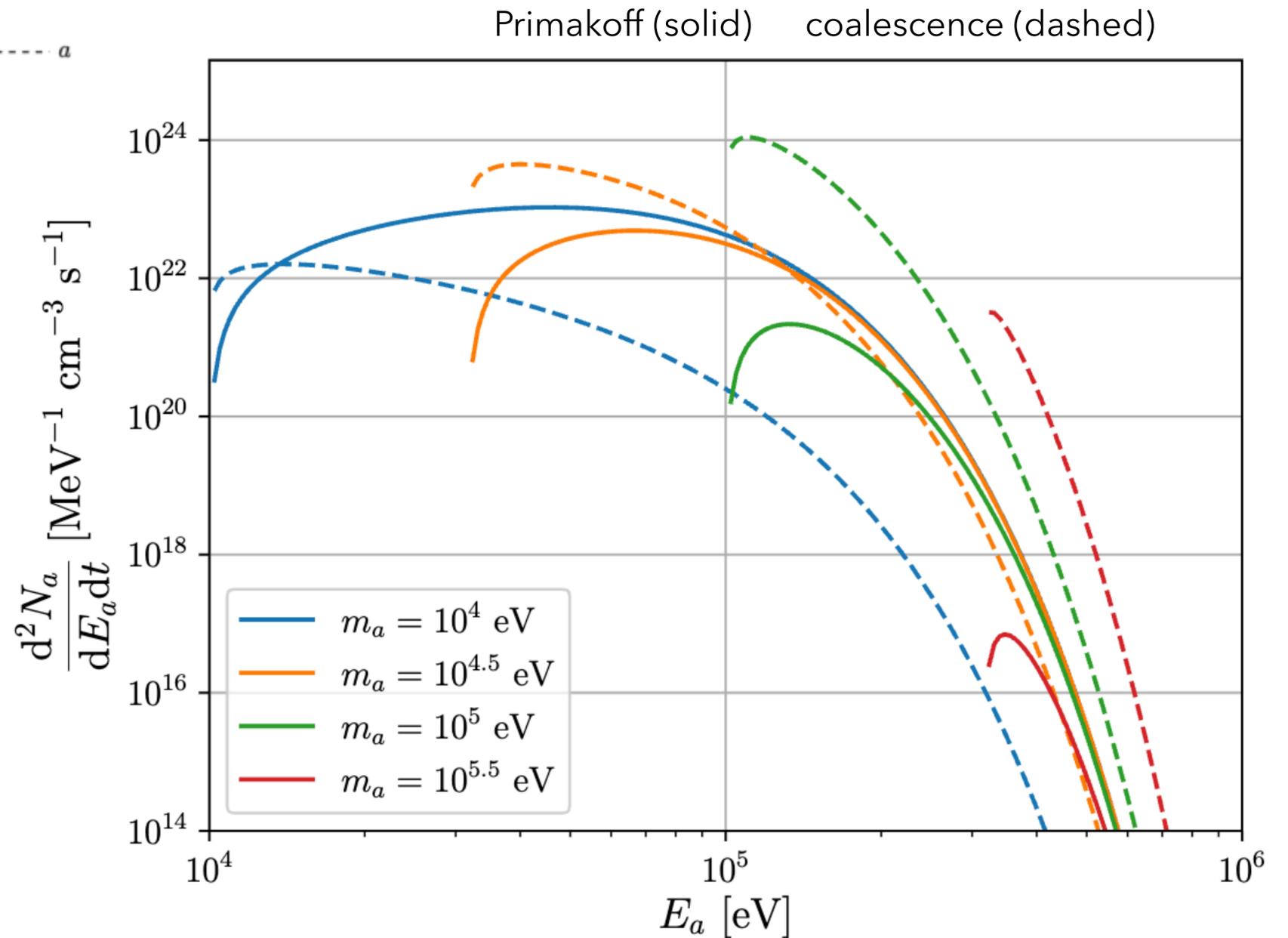
Primakoff process ($\gamma + Ze \rightarrow a + Ze$)



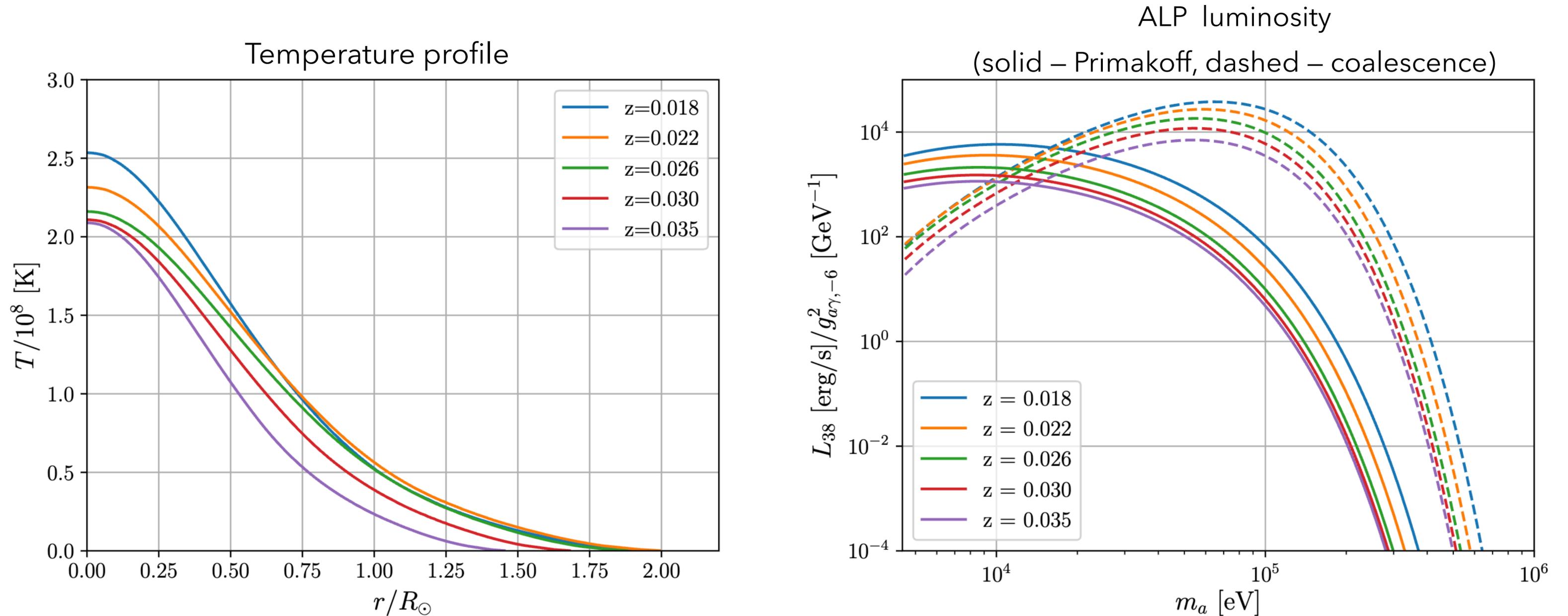
Photon coalescence ($2\gamma \rightarrow a$)



- The ALPs production rate density of each process has different dependence on T, m_a
- Dominant process:
 - $m_a \ll T \rightarrow$ Primakoff process
 - $m_a \gg T \rightarrow$ photon coalescence

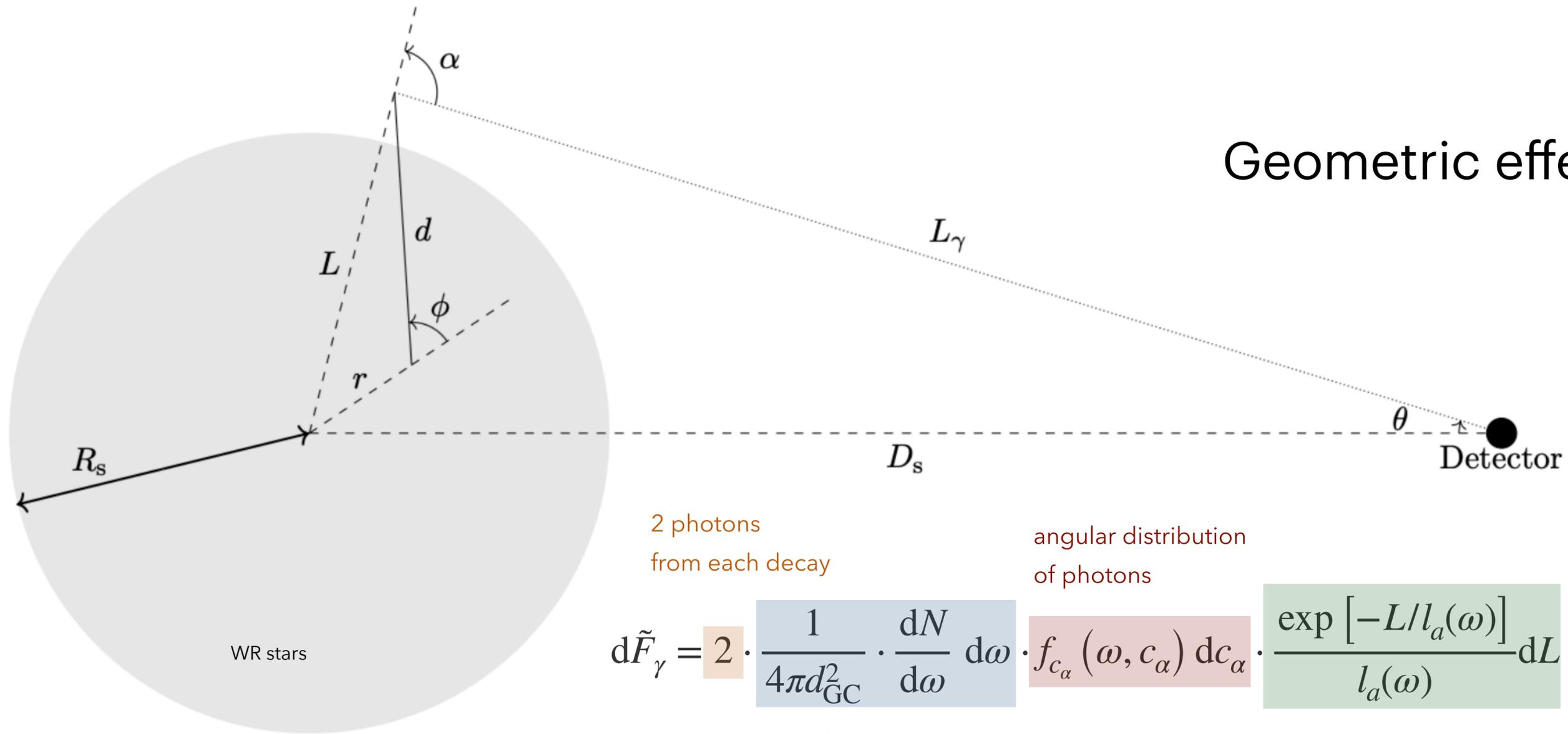


Dependence on Metallicity



- Some uncertainty on the exact metallicity of the WR stars, $z \in [0.018, 0.035]$.
- Lower metallicity results in higher plasma temperature, producing more ALPs.

Photon Flux on Earth



2 photons
from each decay

angular distribution
of photons

constraints on ω, c_α, L

$$d\tilde{F}_\gamma = 2 \cdot \frac{1}{4\pi d_{GC}^2} \cdot \frac{dN}{d\omega} d\omega \cdot f_{c_\alpha}(\omega, c_\alpha) dc_\alpha \cdot \frac{\exp[-L/l_a(\omega)]}{l_a(\omega)} dL \cdot \Theta_{\text{cons.}}(\omega, c_\alpha, L)$$

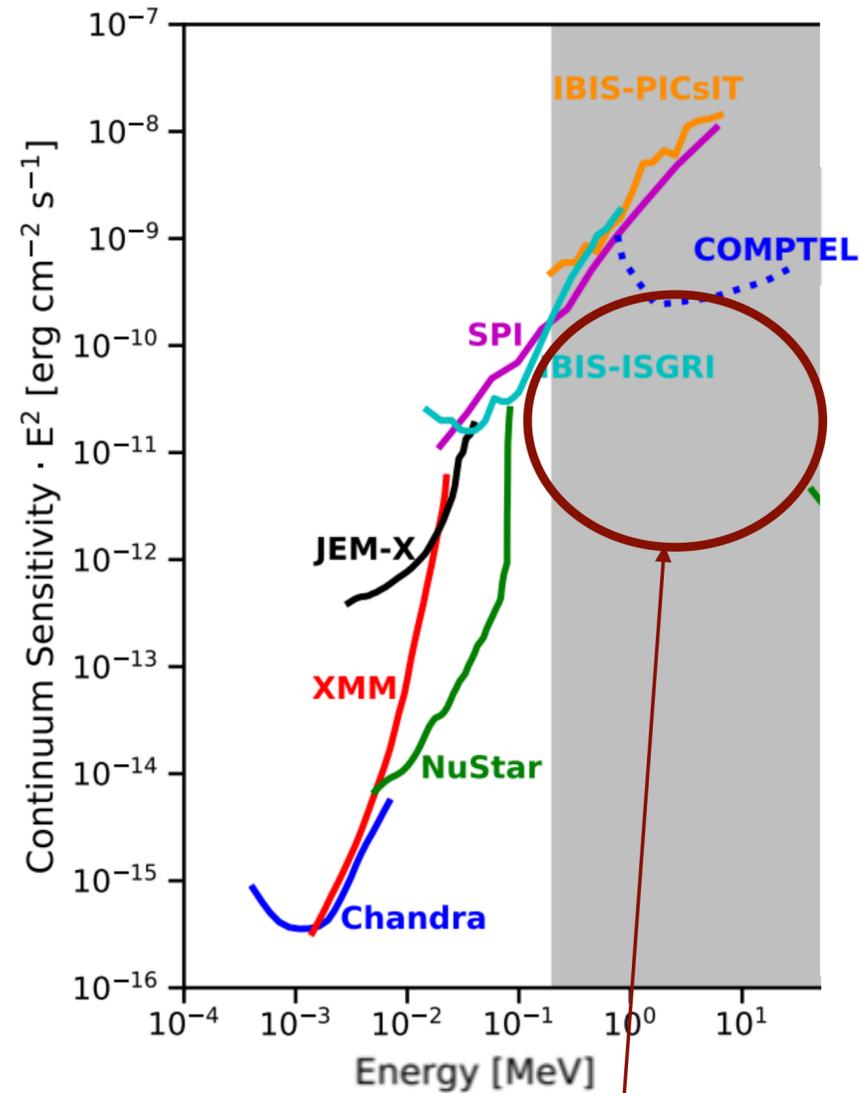
spectral fluence of axions
at the earth

decay probability



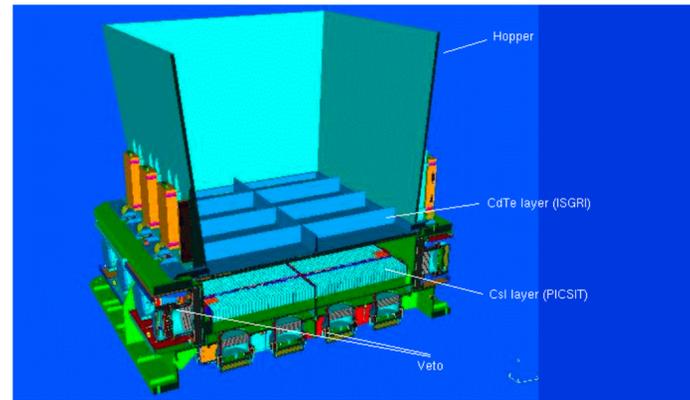
Telescopes Considered

G. Lucchetta et al. 2204.01325

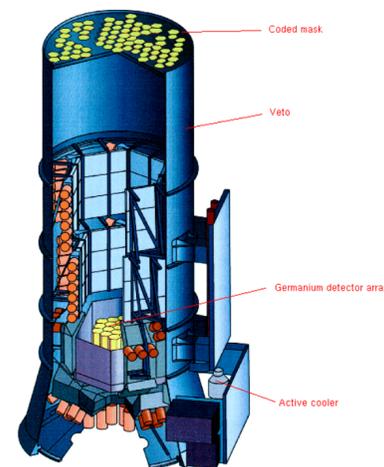


IBIS-ISGRI

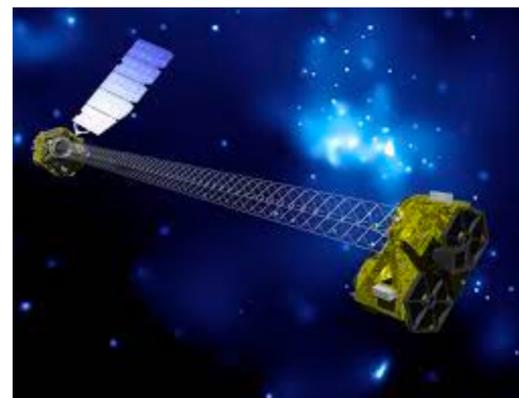
IBIS-PICsIT



SPI



NuSTAR



Mission	Sensitivity Range	Angular Resolution (at Energy)	Mission Status	Effective Area	Observation Time
XMM-Newton [89]	0.1-15 keV	12'' (2-10 keV)	1999-present	~3000 cm ²	10 ⁴ s
NuSTAR [90]	5 keV - 80 keV	18''	2012-present	~1000 cm ²	10 ⁶ s
INTEGRAL IBIS/ISGRI [91]	15 keV - 1 MeV	12'	2002-2023	250 cm ²	10 ⁵ s
INTEGRAL IBIS/PICsIT [92]	170 keV - 10 MeV	12'	2002-2023	~1400 cm ²	10 ⁵ s
INTEGRAL SPI [93]	20 keV - 8 MeV	2.5°	2002-2023	~3000 cm ²	10 ⁶ s
INTEGRAL JEM-X [94]	3 keV - 35 keV	3'	2002-2023	400 cm ²	10 ⁶ s
SWIFT (BAT) [95]	15-150 keV	22'	2004-present	5200 cm ² (15 keV)	0.7×19 yrs
eROSITA [96]	0.2-10 keV	35'' (2-8 keV)	2019-present	1500 cm ²	10 ⁵ s
Insight-HXMT/HE [97]	20 keV - 250 keV	6'	2017-present	4096 cm ²	~10 ⁵ s
COSI [98]	200 keV - 5 MeV	~4° (1 MeV)	2027 (planned)	~300 cm ²	2 yrs
AMEGO (Compton) [99]	200 keV - 10 MeV	~4° (1 MeV)	Concept	~300 cm ²	0.24× 5 yrs
AMEGO-X [100]	100 keV - 1 GeV	~10°	Concept	3500 cm ²	3 yrs
APT (Compton) [101]	200 keV - 10 MeV	~5° (1 MeV)	Concept	10,000 cm ²	0.82× 5 yrs
ASTROGAM [102]	100 keV - 1 GeV	~1.5°	Concept	1000 cm ²	1 yr
e-ASTROGAM [103]	300 keV - 3 GeV	~2°	Concept	9025 cm ²	1 yr
GECCO [104]	100 keV - 10 MeV	2°	Concept	2000 cm ²	10 ⁶ s

Best limit obtained from INTEGRAL-SPI

Future telescopes:

- ASTROGAM
- AMEGO-X
- GECCO

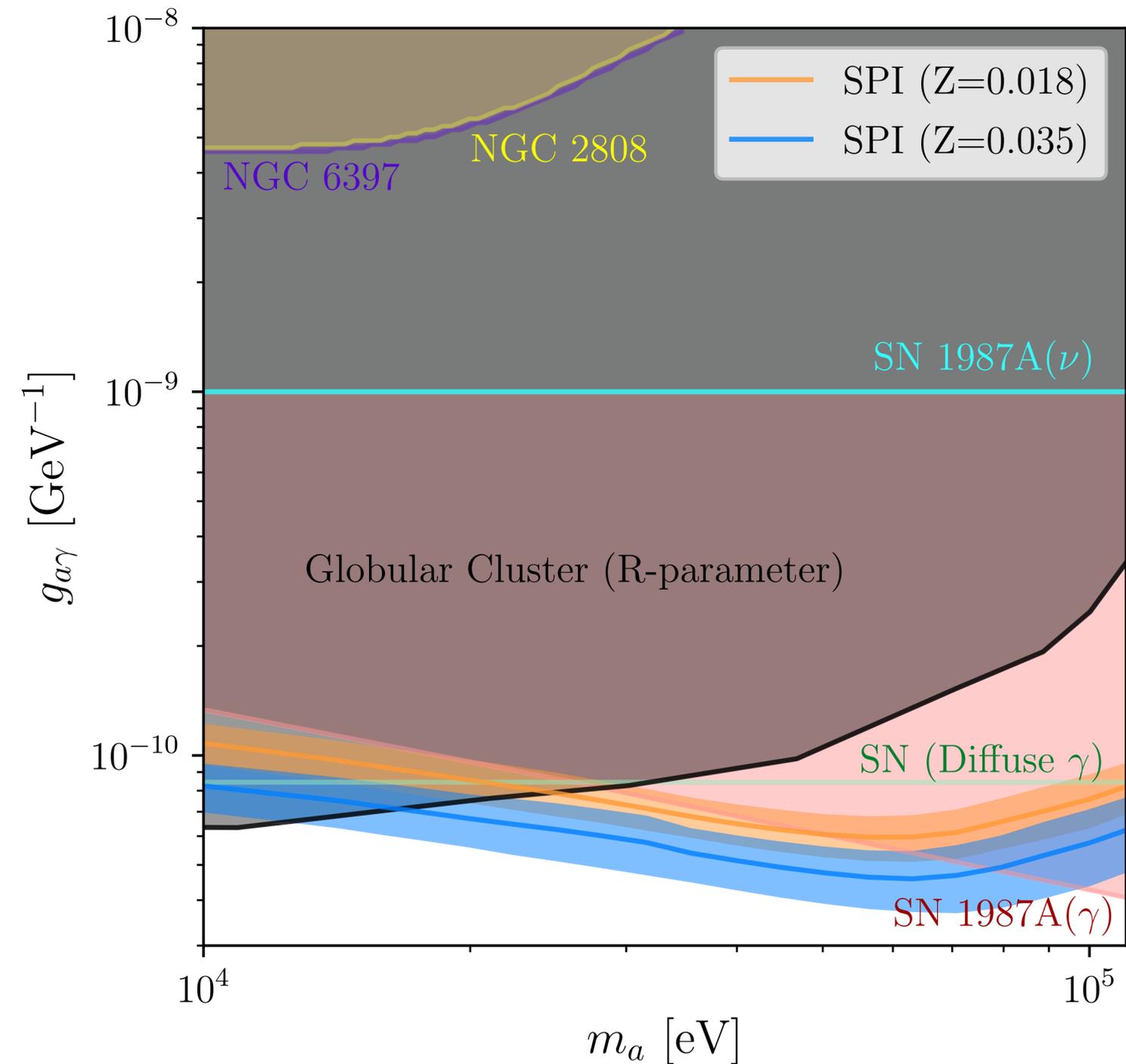
A. De Angelis et al. 2102.02460

R. Caputo et al. 1907.07558

E Orlando et al.2112.07190



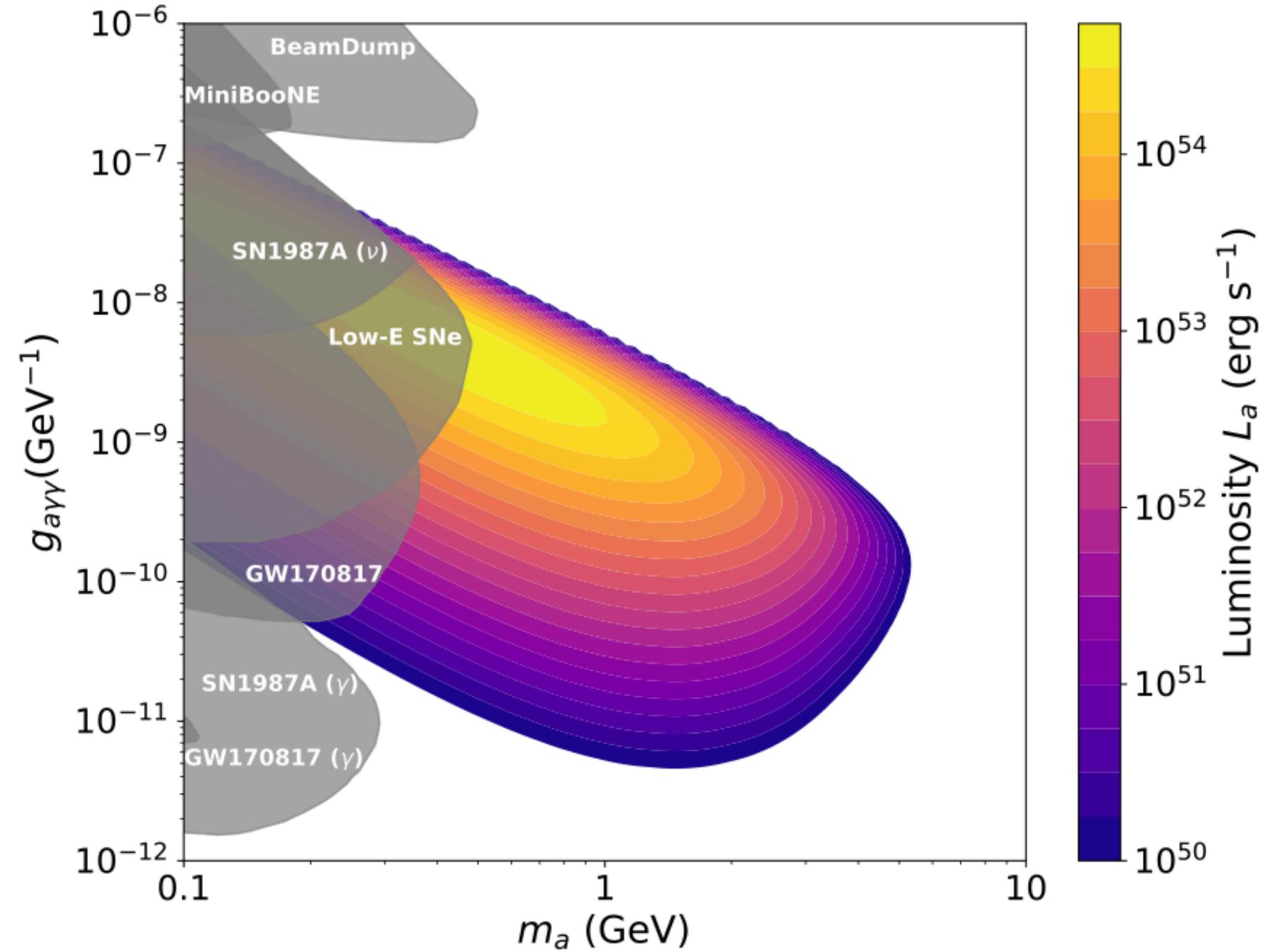
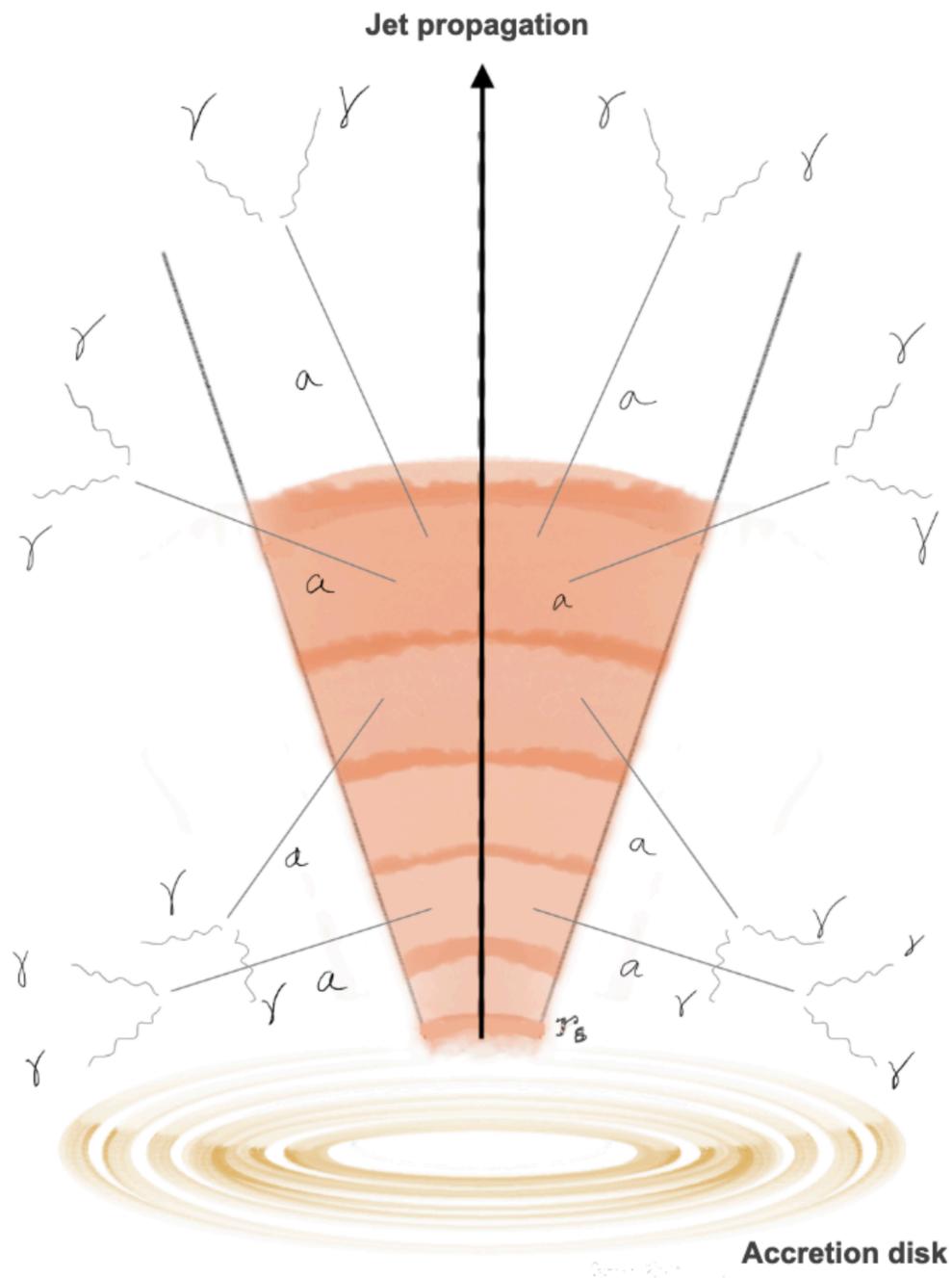
Results for WR Star Analysis



- Used super star cluster Quintuplet located near the galactic center, which hosts 71 massive stars.
- 13 (WC)+1 (WN) are in the WR phase.
- Stellar evolution parameters taken from a previous study. [Dessert, Foster, Safdi, 2008.03305]
- INTEGRAL-SPI gives the best limit in the parameter space of interest.
- New NuSTAR bound from starburst galaxy M82. [Candon, Fiorillo, Lucente, Vitagliano, Vogel, 2412.03660]

ALPs from GRB Fireball

Ghosh, Jacobsen, Linden, 2501.08978



Fireball ALP-induced Diffuse Gamma-ray Bkg

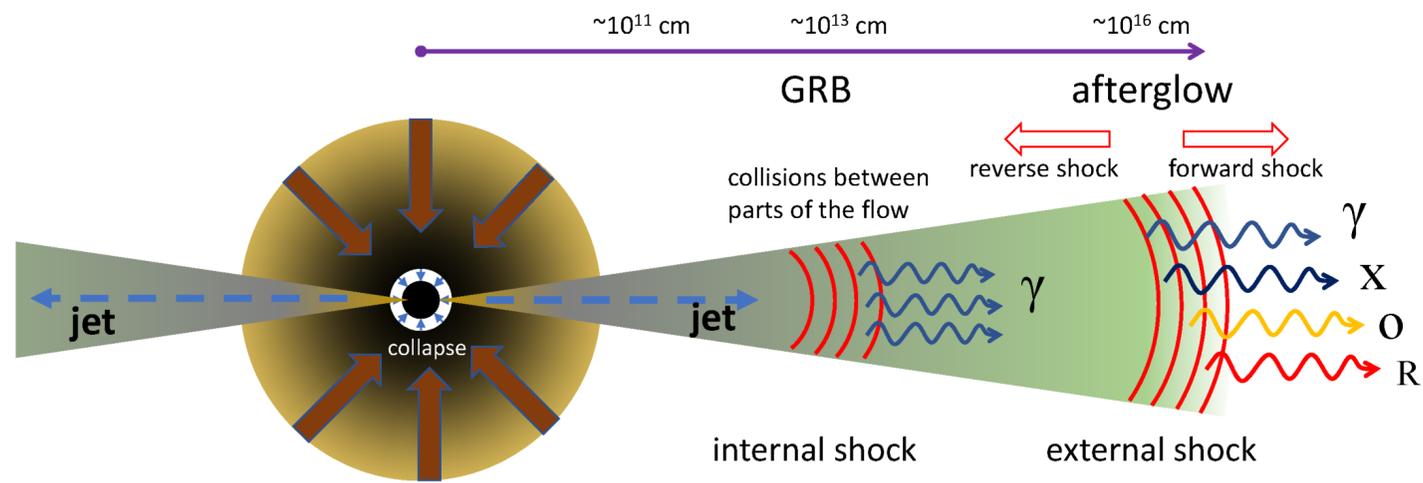


Image credit: S. Dado et al., Universe 2022, 8(7), 350

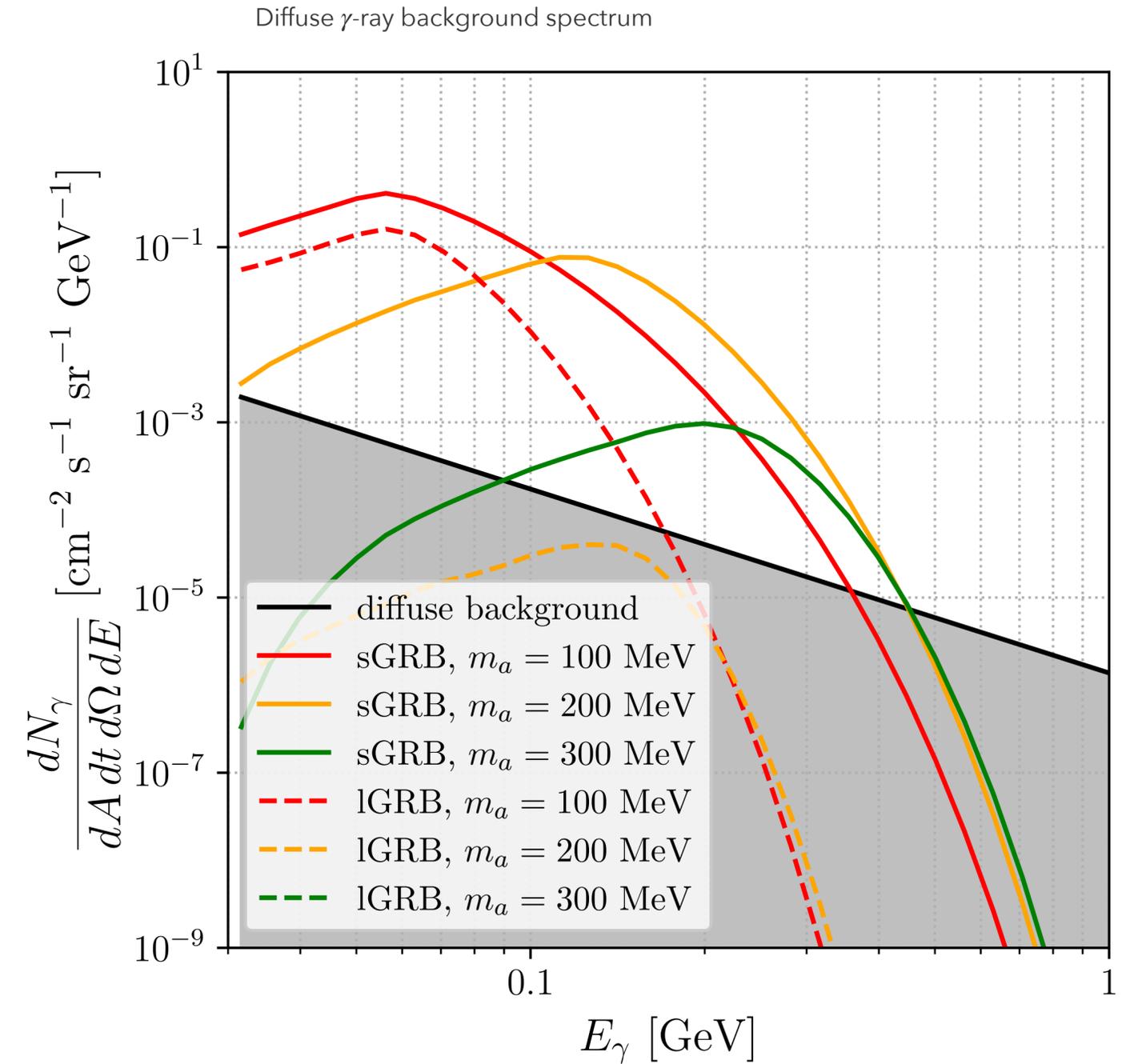
- The innermost (radiation-dominated) region follows

$$T(r) = T_i \frac{r_i}{r}, \quad \gamma(r) = \frac{r}{r_i}$$
- Launched at $r_i \approx (2 - 3) \times r_s$ where $T_i \approx \mathcal{O}(10 \text{ MeV})$
- Extended to $r_m \approx 10^{7-9} \text{ cm}$

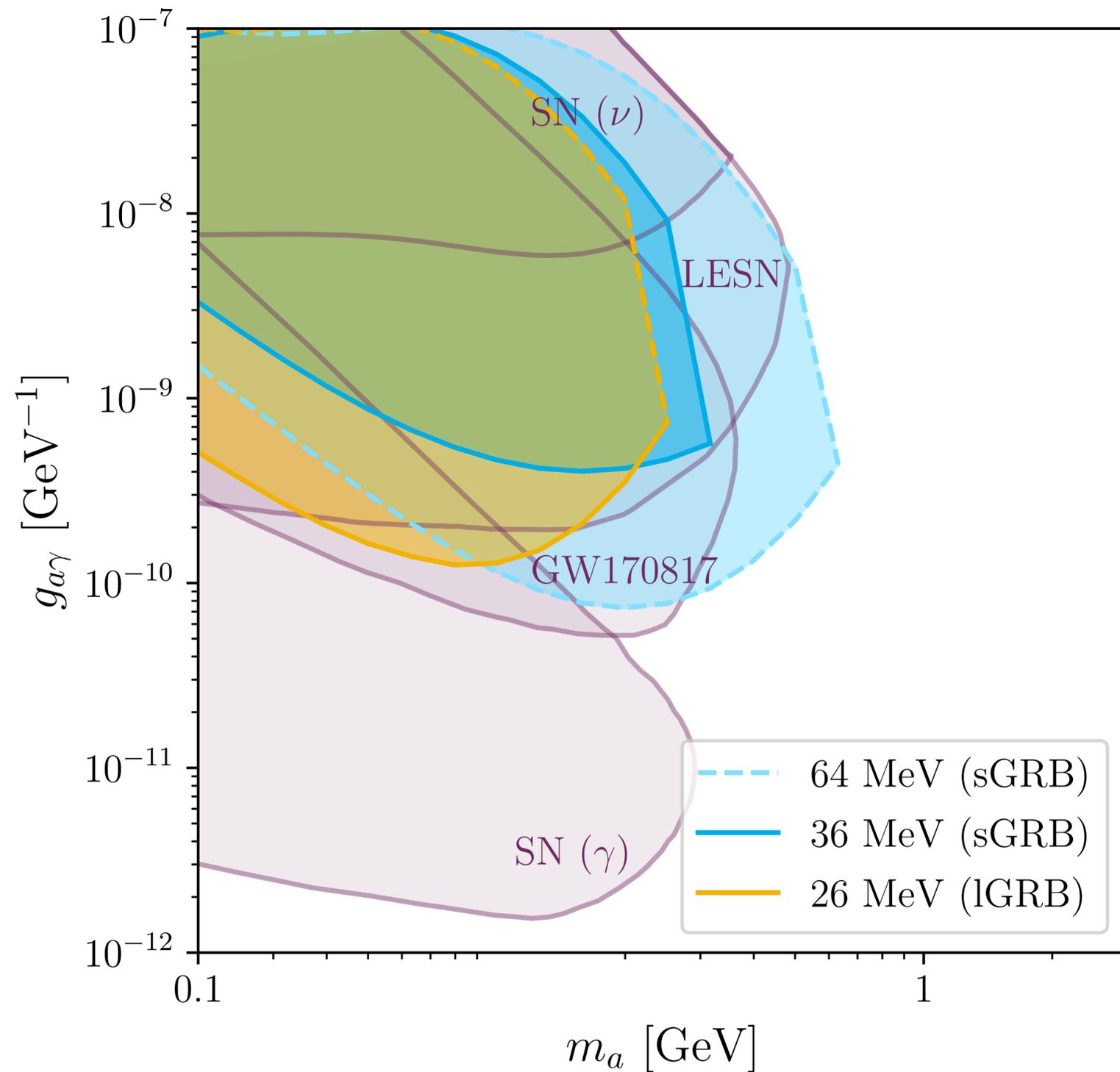
Integrate over a fireball volume energy loss per unit volume

$$L_a = \Delta\Omega \int_{r_i}^{r_m} dr r^2 \int_{m_a}^{\infty} dE_a E_a \frac{d\dot{n}_a}{dE_a} \times \Theta \left(E_a - \frac{m_a}{\sqrt{1 - \frac{2GM_r}{rc^2}}} \right) f_{\text{decay}},$$

condition for ALPs to escape from a gravitational potential escape probability



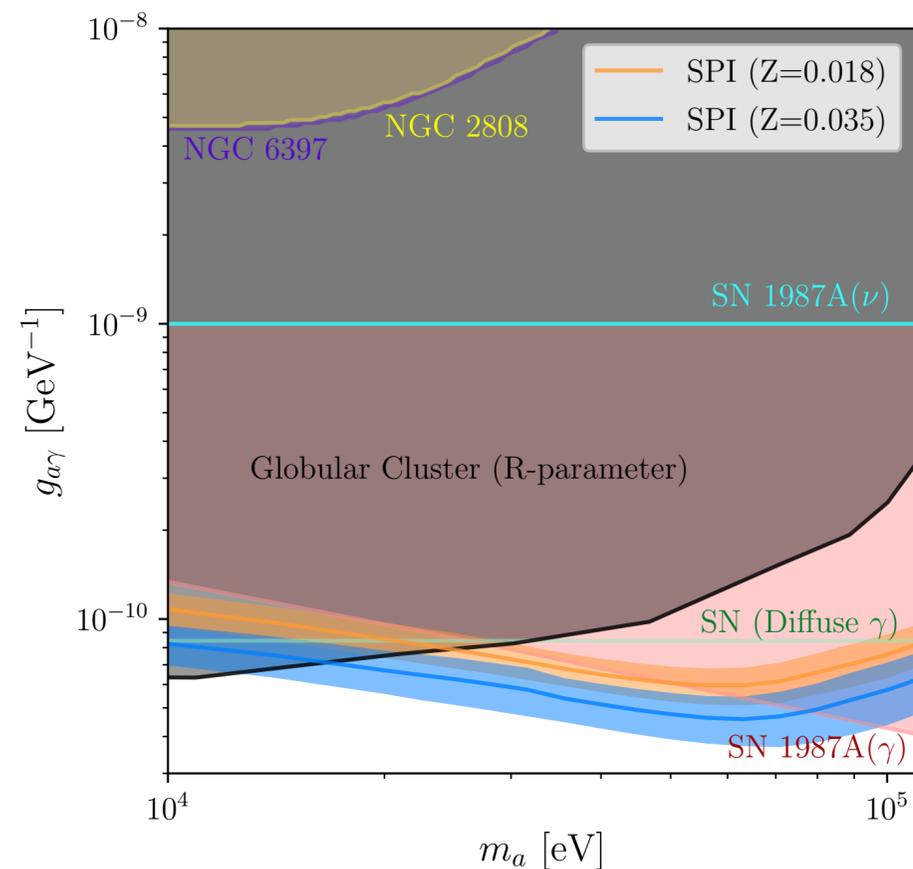
New ALP Bounds



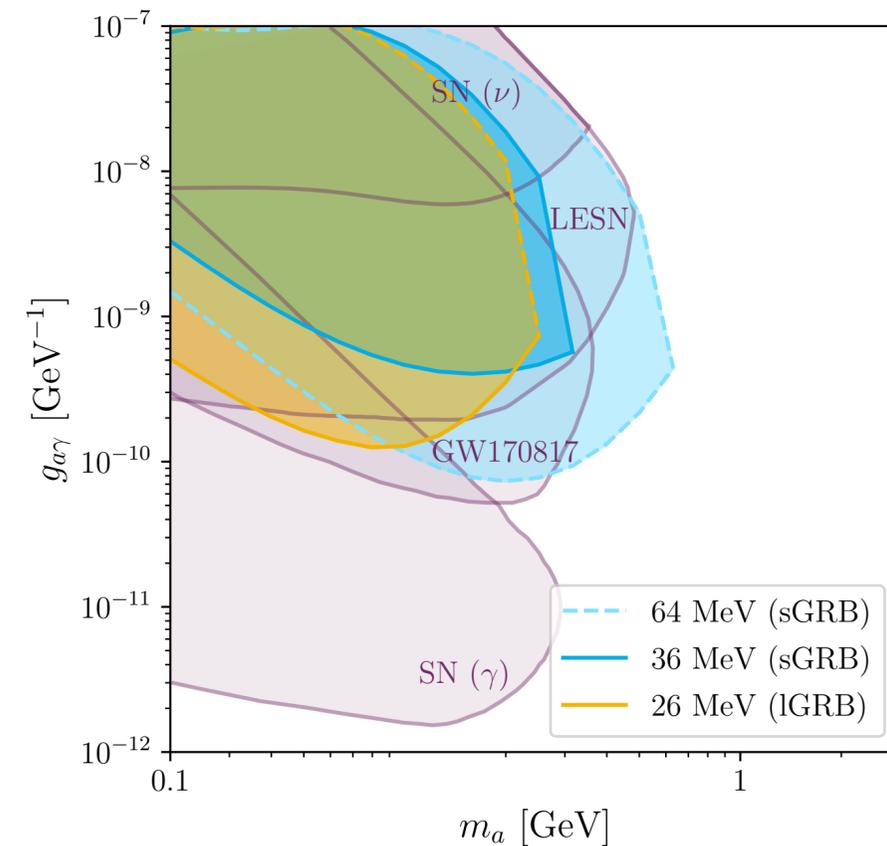
- Strongly depends on initial GRB fireball temperature and launch radius.
- Previous study used an **explosion model** which predicts higher temperatures $T \sim \mathcal{O}(100 \text{ MeV})$.
- Detailed modeling of GRB jet flow suggests a **wind model** with lower temperatures $T \sim \mathcal{O}(10 \text{ MeV})$.
- Results in significantly weaker ALP bounds than those shown in [Ghosh, Jacobsen, Linden, 2501.08978].

Conclusions

- ALPs are well-motivated BSM candidates.
- Astrophysical plasma is an efficient ALP production site.
- Considered massive Wolf-Rayet stars and GRB fireball as two examples.
- Derived new astrophysical bounds on heavy ALPs in the keV-MeV range.



WR Star



GRB Fireball