

# The Effective Potential originating from Swampland and the non-trivial Brans-Dicke coupling

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# Outline

- 1) The String Swampland and String Landscape
- 2) The origin and sign of Large Scale Lorentz Violation
- 3) The Cosmology with large scale Lorentz Violation in swampland and string landscape
- 4) Summary

- The possible string vacua compactification choice can be of order  $10^{500}$ . Among them, inequivalent ones constitute the string landscape.
- it is likely that any consistent looking lower dimensional effective field theory (EFT) coupled to gravity can arise in some way from a string theory compactification
- the set of all EFT which do not admit a string theory UV completion as the swampland.
- The generic AdS vacuum is SUSY preserving.
- To account the inflation and accelerating expansion one needs to lift the AdS to dS vacuum.

The uplifting the AdS type of vacua to dS ones comes from  $\bar{D}_3$  branes tension in a sufficiently warped background, in the presence of quantum corrections, by carefully adding  $\bar{D}_3$  branes into the compactification. –KKLT construction Kachru, PRD 03

Shortage: no-go theorems, restrictions on ingredients used in string theory, typically specific combinations of fluxes, D-branes, orientifolds

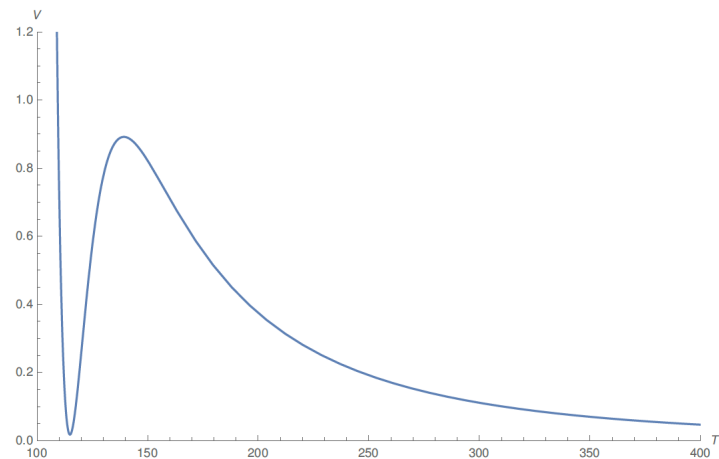
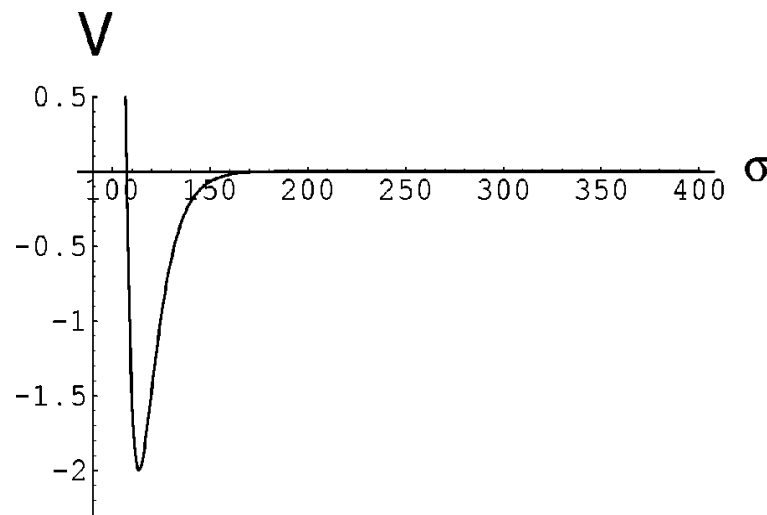
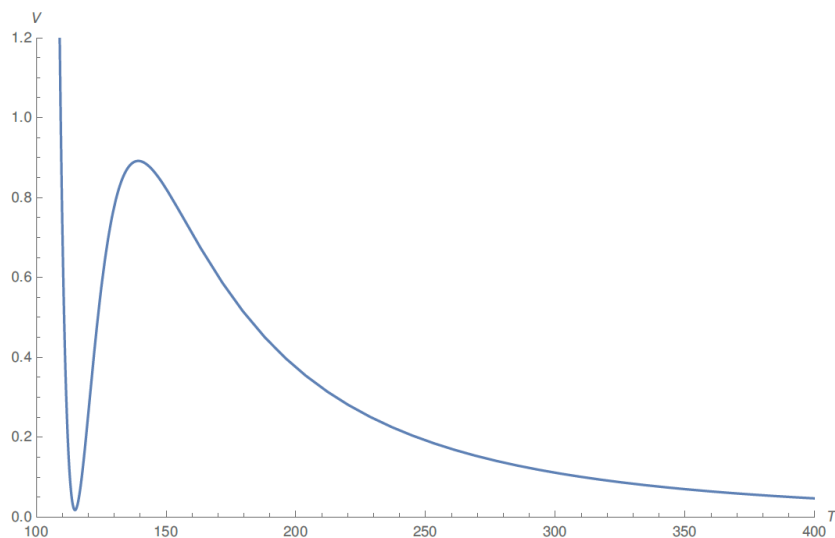


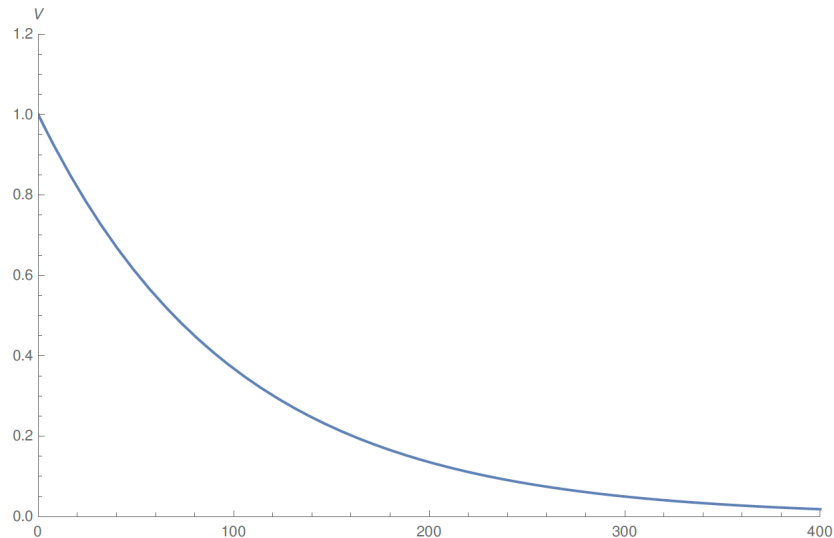
Figure 7. The scalar potential for a metastable dS.

The second criterion of swampland conjecture excludes the EFT with a meta-stable dS vacuum as theory with UV completion. Metastable dS belongs to the swampland.

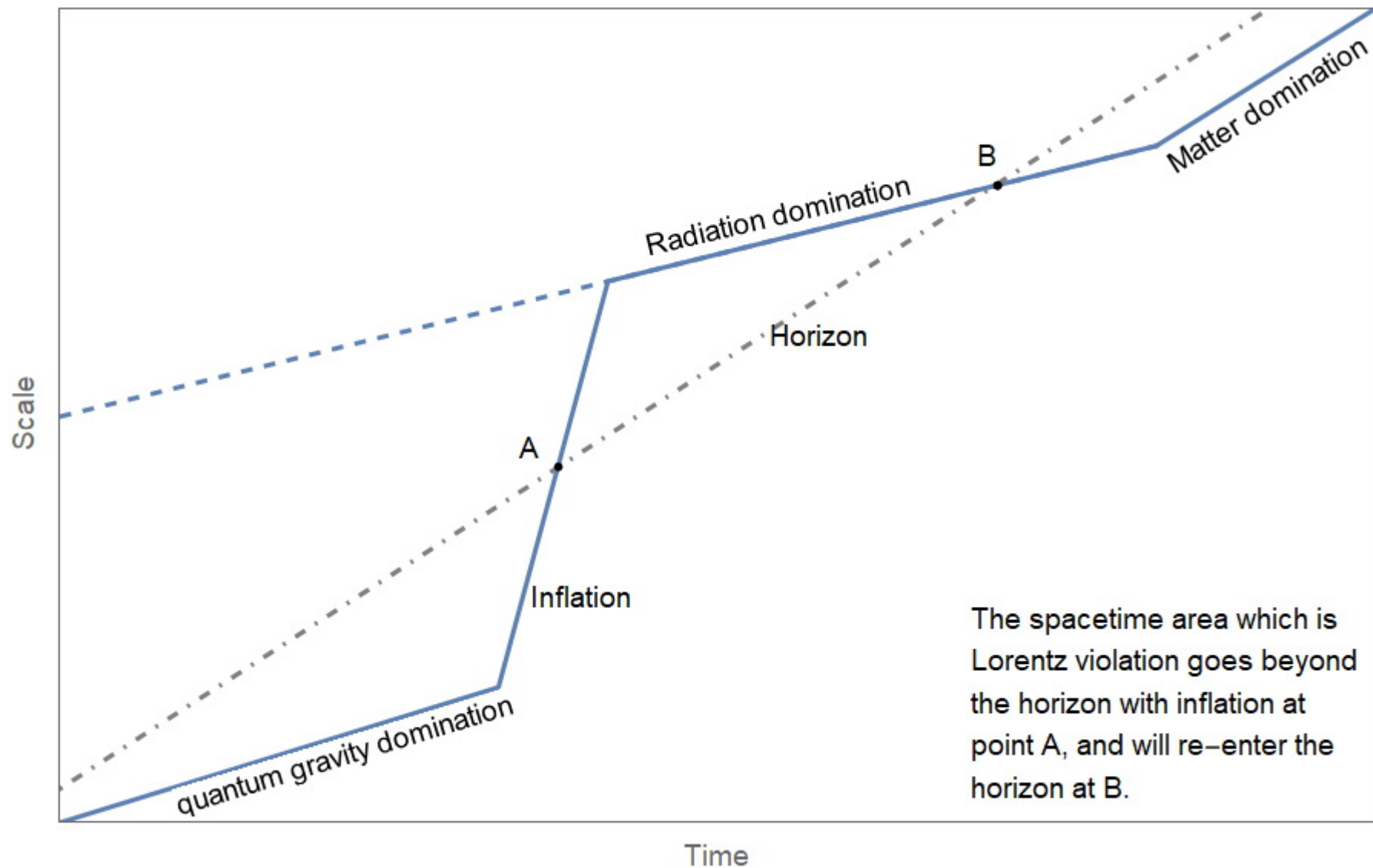
Quintessence model can satisfy the second criterion.



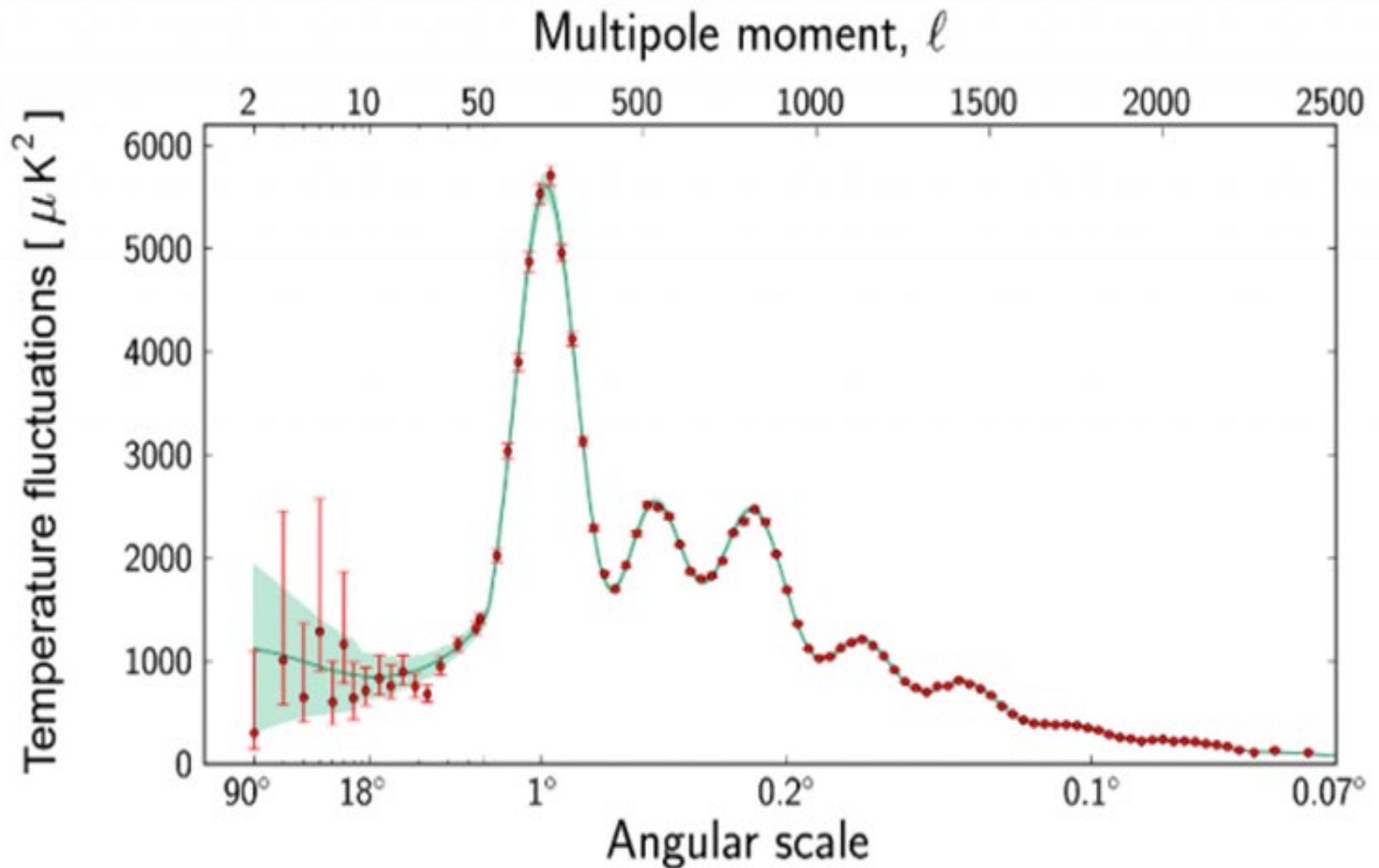
**Figure 7.** The scalar potential for a metastable  $dS$ .



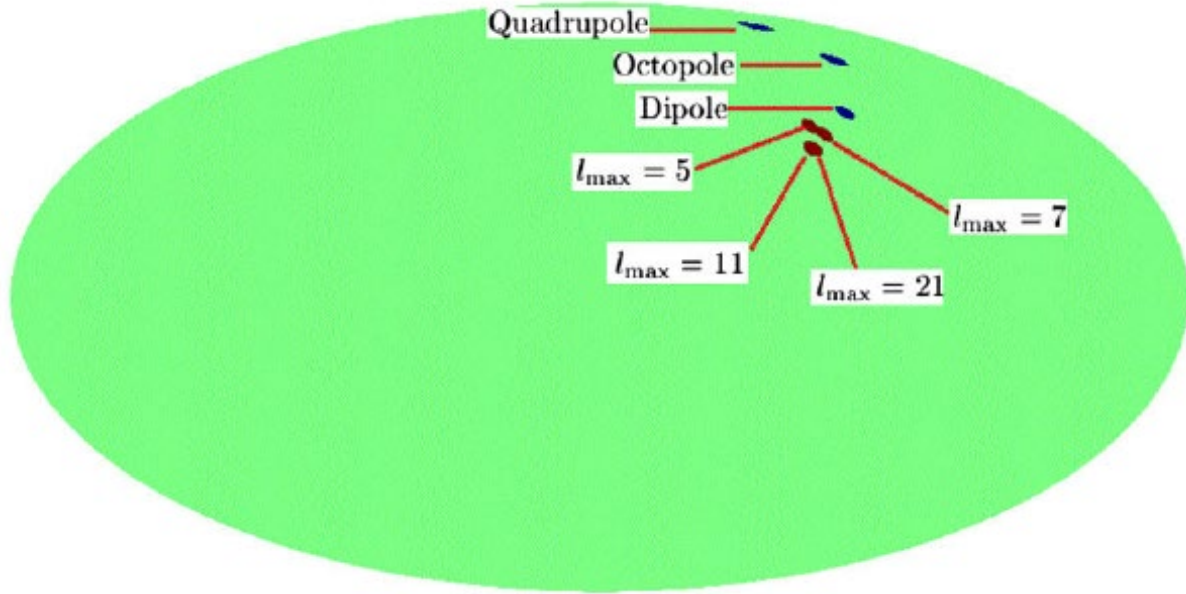
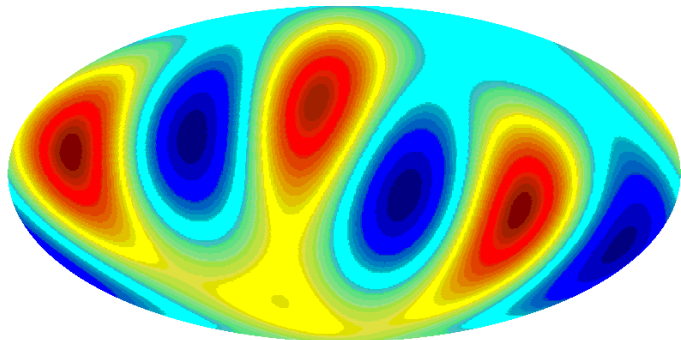
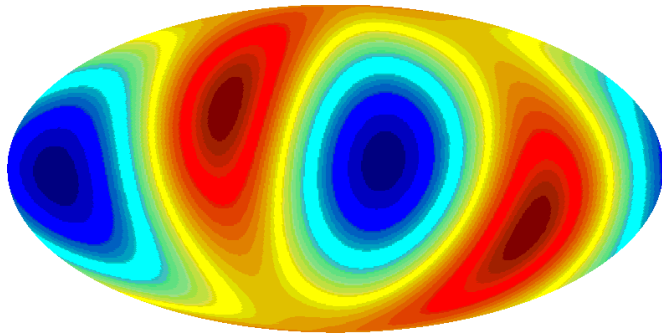
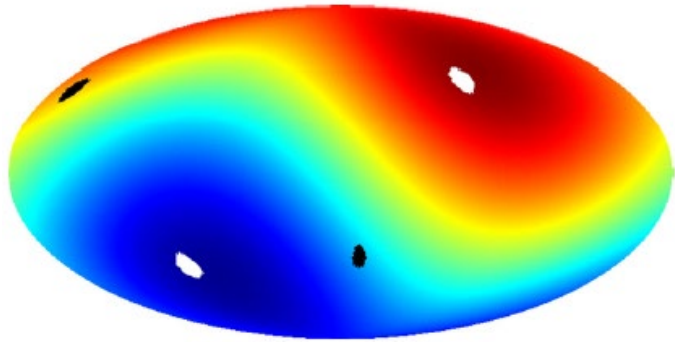
**Figure 8.** The scalar potential for an unstable  $dS$ .



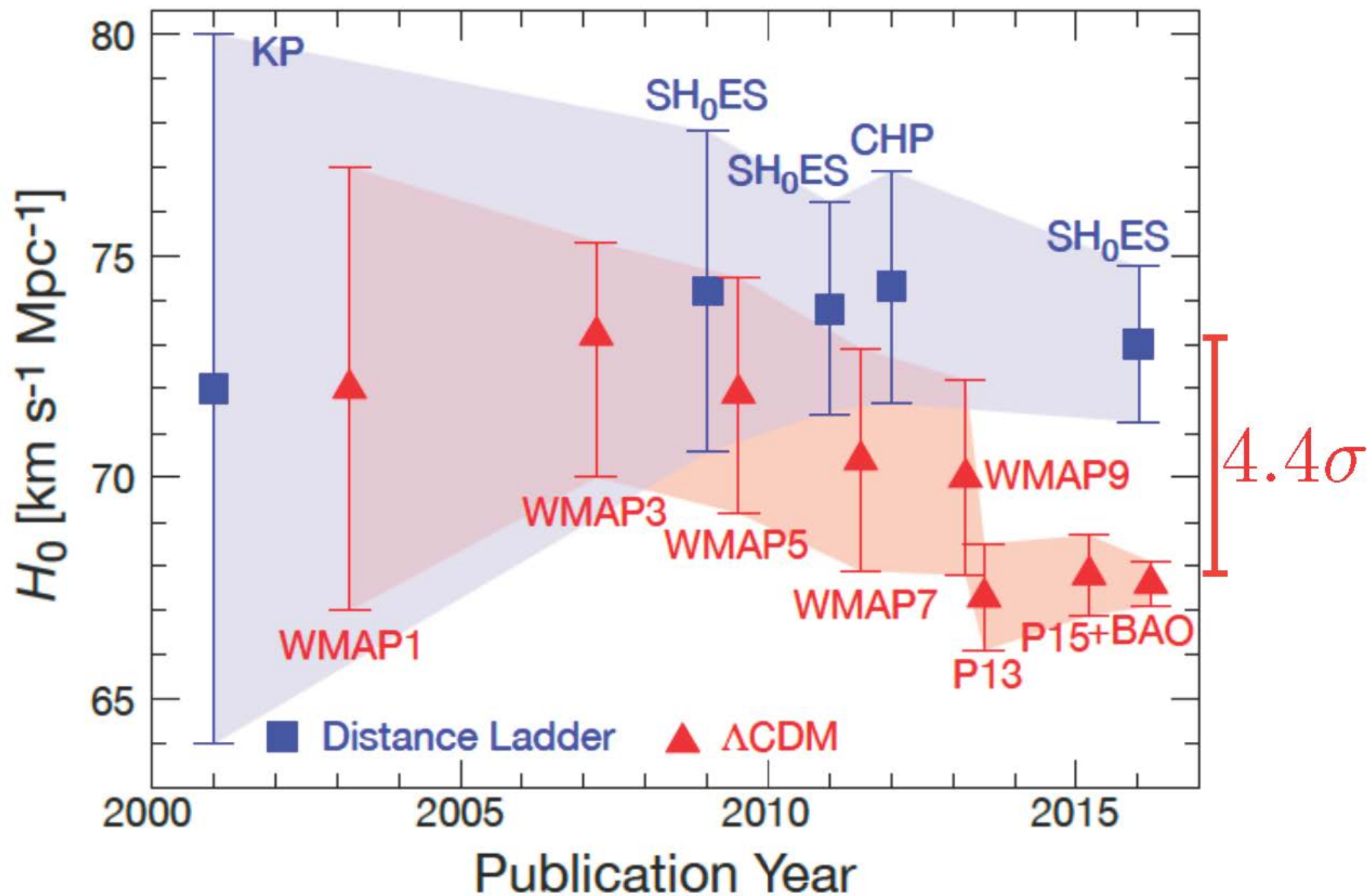
# Anisotropies of CMB



# Comparing with preferred directions in CMB dipole, quadrupole and octopole







# The Lorentz Violated EFT of Gravity

- The action for a sim(2) gravity

$$S_E = \frac{1}{16\pi G} \int d^4x h \left( R^{ab}{}_{ab} + \lambda_1^\mu \left( A^{10}{}_\mu - A^{31}{}_\mu \right) + \lambda_2^\mu \left( A^{20}{}_\mu + A^{23}{}_\mu \right) \right)$$

- The Lagrange-multipliers term can be regarded as an effective angular momentum distribution  $C_{M eff}$

$$\mathcal{D}_\nu \left( h \left( h_a{}^\nu h_b{}^\mu - h_a{}^\mu h_b{}^\nu \right) \right) = 16\pi G \left( C_M + C_{M eff} \right)_{ab}{}^\mu$$

- Lorentz violation leads to non-trivial distribution of contortion
- The non-trivial effective contribution to the energy-momentum distribution by contortion is expected to be responsible for the dark partner of the matter.

$$\tilde{R}_c{}^a - \frac{1}{2} \delta_c{}^a \tilde{R} = 8\pi G \left( T_{eff} + T_M \right)_c{}^a$$

- The Bianchi Identities imply the conservation of  $T_{eff}$

## The Modified Constraint for SO(3)

- For SO(3)  $\Lambda_0^j(x) = 0$

$$\begin{aligned} A'^i_{0\mu} &= \Lambda^i_j(x) A^j_{0\mu} \Lambda_0^0(x) + \Lambda^i_j(x) \partial_\mu \Lambda_0^j(x) \\ &= \Lambda^i_j(x) A^j_{0\mu} \end{aligned}$$

- The Modified Constraint for SO(3) can be

$$S_E = \frac{c^4}{16\pi G} \int d^4x h \left( R - 2\Lambda_0 + \lambda^u \left( \left( A^0_{1u} \right)^2 + \left( A^0_{2u} \right)^2 + \left( A^0_{3u} \right)^2 - f_u^2 \right) \right)$$

- Where  $f_\mu$  can be regarded as the measurement of Lorentz violation.

# Accelerating Expansion of the Universe

- To construct the FRW like solution of the model

$$ds^2 = dt^2 - a(t)^2 \left( \frac{dr^2}{1 - kr^2} + r^2 d\theta^2 + r^2 \sin^2 \theta d\varphi^2 \right)$$

- The naïve comoving tetrad can be chosen as

$$h^0 = dt, h^1 = \frac{a(t)}{\sqrt{1 - kr^2}} dr, h^2 = ra(t) d\theta, h^3 = r \sin \theta a(t) d\varphi$$

- And  $h_0 = \frac{\partial}{\partial t}, h_1 = \frac{\sqrt{1 - kr^2}}{a(t)} \frac{\partial}{\partial r}, h_2 = \frac{1}{ra(t)} \frac{\partial}{\partial \theta}, h_3 = \frac{1}{r \sin \theta a(t)} \frac{\partial}{\partial \varphi}$

# Cosmic solution of contortion

- The perfect fluid of cosmic media demands

$$G^1_1 = G^2_2 = G^3_3$$

- With decomposition of connections,  $A^a_{b\mu} = \Gamma^a_{b\mu} + K^a_{b\mu}$  a simple solution can be chosen as

$$K^0_{11} = K^0_{22} = K^0_{33} = \mathcal{K}(t)$$

- With other contortion components vanish.
- Expand the field eq. with quantities with “ $\sim$ ” in Levi-Civita connection

$$\tilde{R}^a_c - \frac{1}{2} \tilde{R} \delta^a_c = 8\pi G (T + T_\Lambda)^a_c, \quad T_\Lambda^a_c = \frac{1}{8\pi G} \Lambda^a_c = \frac{1}{8\pi G} (\tilde{G}^a_c - G^a_c)$$

$$[T_\Lambda]^a_c = \text{Diag}(\rho_\Lambda, -p_\Lambda, -p_\Lambda, -p_\Lambda), \quad \rho_\Lambda = -\frac{c^4}{8\pi G} \left( 3\mathcal{K}^2 + 6\mathcal{K} \frac{\dot{a}}{a} - \Lambda_0 \right)$$

$$p_\Lambda = \frac{c^4}{8\pi G} \left( \mathcal{K}^2 + 4\mathcal{K} \frac{\dot{a}}{a} + 2\dot{\mathcal{K}} - \Lambda_0 \right)$$

- Denote  $\Lambda_0$  as the **bare** cosmological constant in our Lorentz violating model from vacuum energy density,  $\Lambda$  as the **observed one** and take the geometrical unit  $\frac{8\pi G}{c^4} = 1$  and  $x = \frac{\Lambda_0}{\Lambda}$
- **the modified Friedmann Equation**

$$\left( \mathcal{K} + \frac{\dot{a}}{a} \right)^2 = \frac{1}{3}(\rho + \Lambda_0)$$

$$\ddot{a} = -\frac{a}{2} \left( p + \frac{\rho}{3} \right) + \frac{1}{3} \left( a\Lambda_0 - 3 \frac{d}{dt} (a\mathcal{K}) \right)$$

- **The Friedmann Eqns in  $\Lambda$ CDM**

$$\left( \frac{\dot{a}}{a} \right)^2 - \frac{\Lambda}{3} = \frac{\rho}{3}$$

$$\ddot{a} = -\frac{a}{2} \left( p + \frac{\rho}{3} \right) + \frac{1}{3} a \Lambda$$

- Accelerating expansion condition:

$$\frac{a}{2} \left( p + \frac{\rho}{3} - \frac{2}{3} \Lambda_0 \right) + \frac{d}{dt} (a\mathcal{K}) < 0$$

- the modified Friedmann Equation with the Eq of States for cosmic media  $p=w\rho$

$$\dot{H}(t) + \dot{\mathcal{K}}(t) + H(t)(H(t) + \mathcal{K}(t)) + \frac{3w+1}{2} (H(t) + \mathcal{K}(t))^2 - \frac{(w+1)}{2} \Lambda_0 = 0$$

- And  $w \approx 0$  for matter dominated period

- **Initial conditions:**  $\mathcal{K}(t_0)^2 + 2\mathcal{K}(t_0)\frac{\dot{a}(t_0)}{a(t_0)} = \frac{\Lambda_0}{3} - \frac{\Lambda}{3}$

$$\mathcal{K}(t_0) = H_0 \left( \pm \sqrt{1 - \frac{\Lambda - \Lambda_0}{3H_0}} - 1 \right) \rightarrow \Lambda_0 \geq -(3H_0 - \Lambda) \approx -\frac{2}{5}\Lambda$$

- **Three cases of approximation**

- **Case A:**  $\frac{d}{dt}(a\mathcal{K}) = -\frac{1}{3}a(\Lambda - \Lambda_0)$

- **Or**

$$H(t)\mathcal{K}(t) + \dot{\mathcal{K}}(t) = \frac{1}{3}(\Lambda_0 - \Lambda)$$



## ■ Case B:

$$\dot{\mathcal{K}}(t) + (3w + 2)H(t)\mathcal{K}(t) + \frac{3w+1}{2}\mathcal{K}^2(t) = \frac{w+1}{2}(\Lambda_0 - \Lambda)$$

## ■ Case C:

$$[T_\Lambda]^a_c = \text{Diag}(\rho_\Lambda, -p_\Lambda, -p_\Lambda, -p_\Lambda)$$

$$p_\Lambda = w_0 \rho_\Lambda$$

$$(3w_0 + 1)\mathcal{K}^2 + (6w_0 + 4)H\mathcal{K} + 2\dot{\mathcal{K}} - (w_0 + 1)\Lambda_0 = 0$$

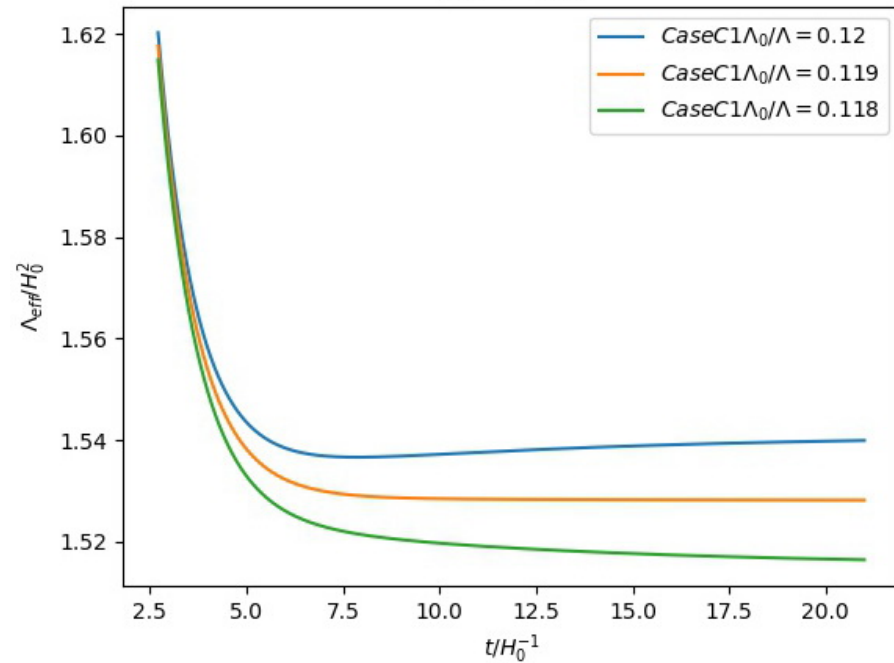
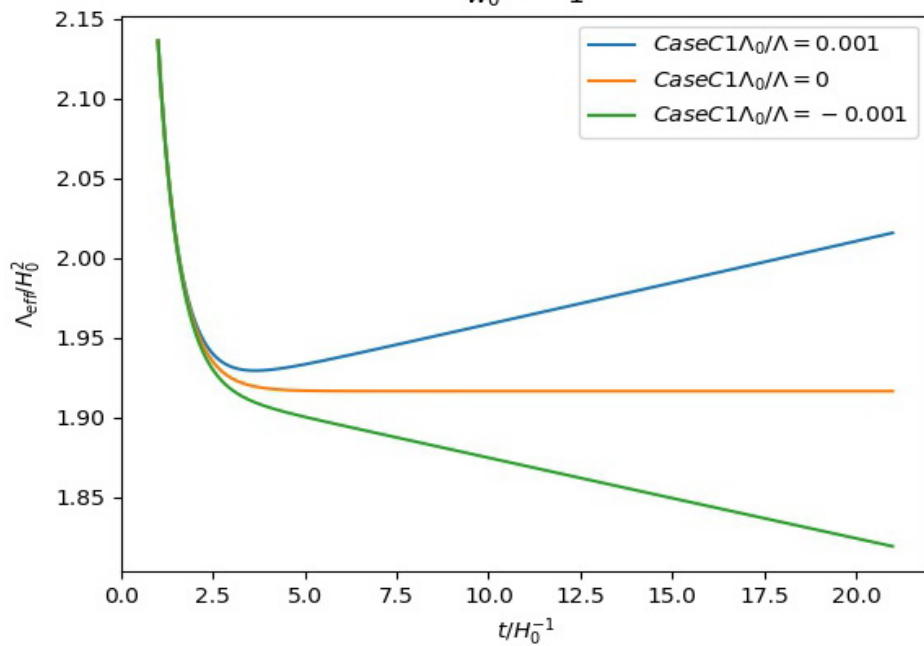
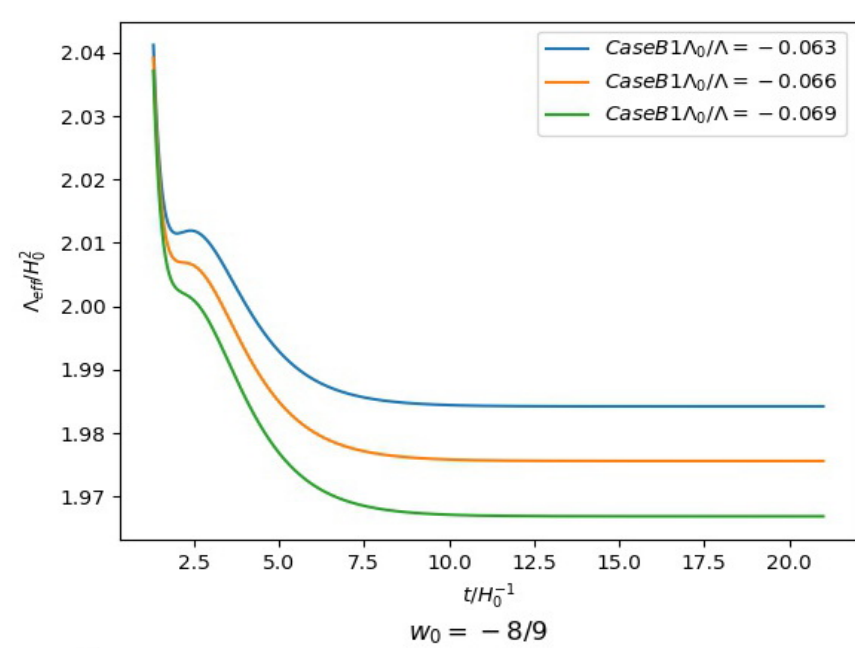
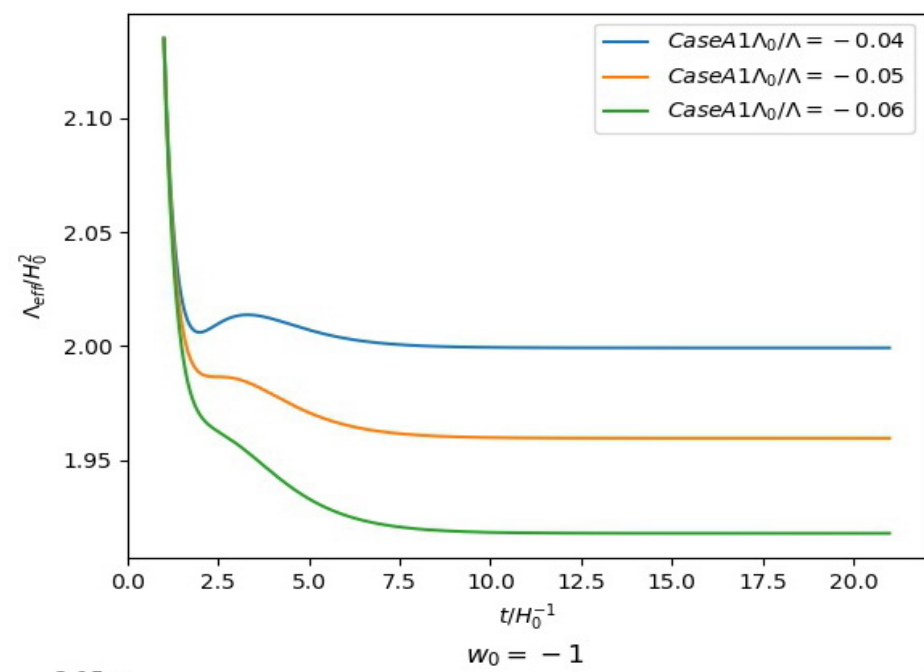
$$\dot{H} + \dot{\mathcal{K}} + H(H + \mathcal{K}) + \frac{3w+1}{2}(H + \mathcal{K})^2 - \frac{(w+1)}{2}\Lambda_0 = 0$$

$$\mathcal{K}(t_0) = H_0 \left( \pm \sqrt{1 - \frac{\Lambda - \Lambda_0}{3H_0^2}} - 1 \right), \quad H_0 = H(t_0)$$

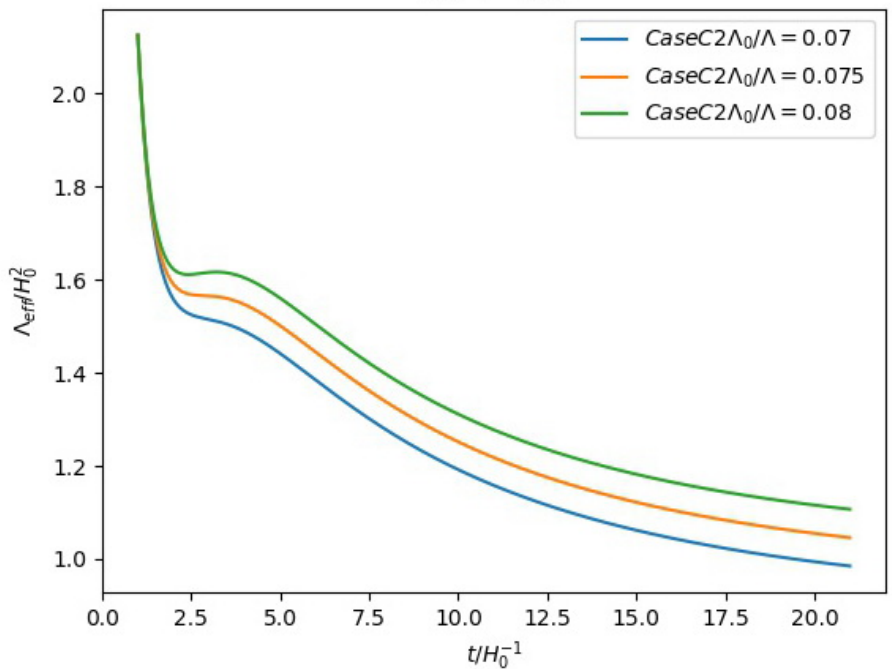
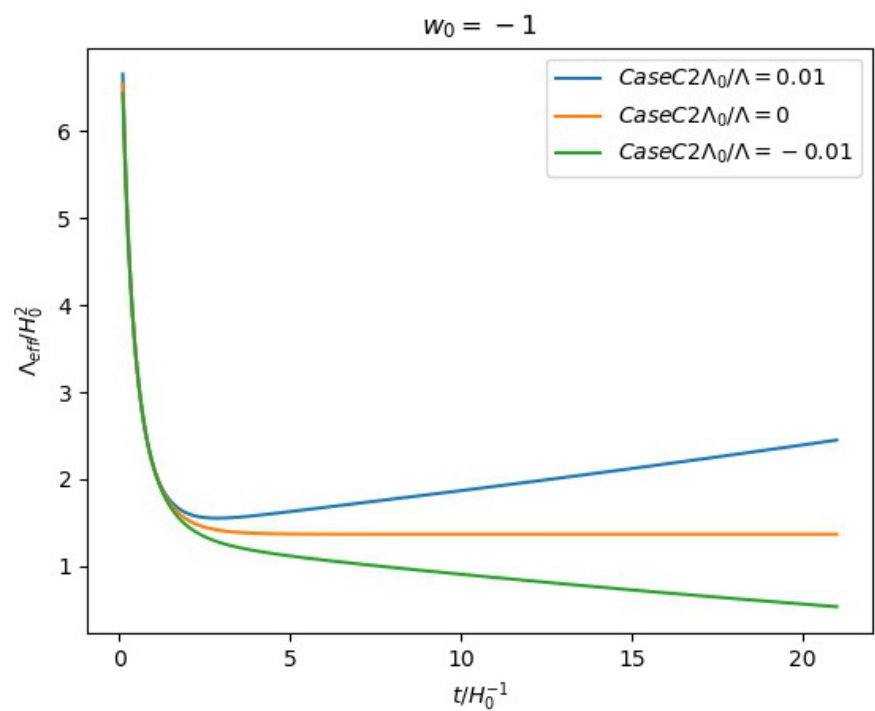
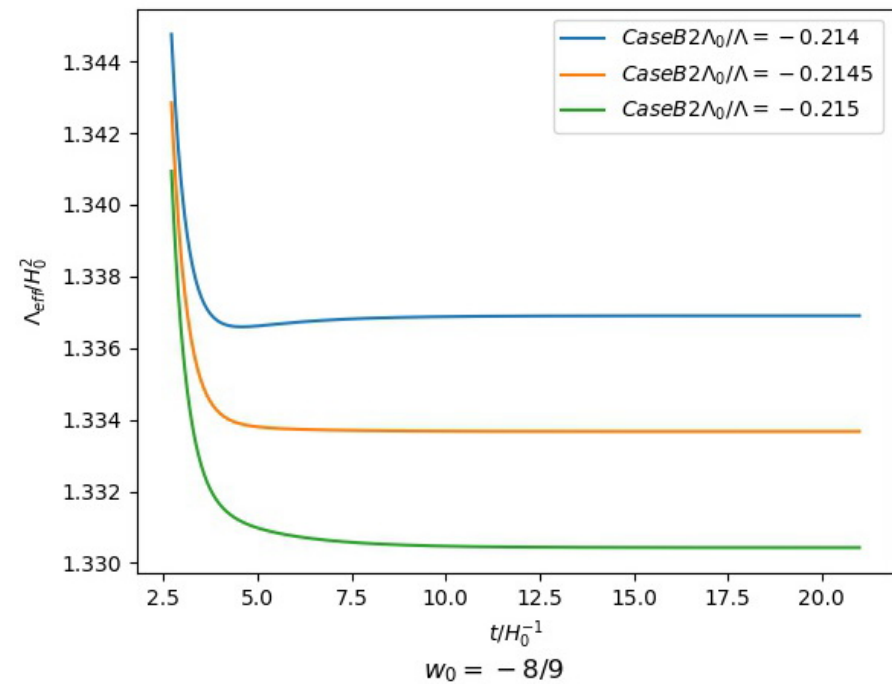
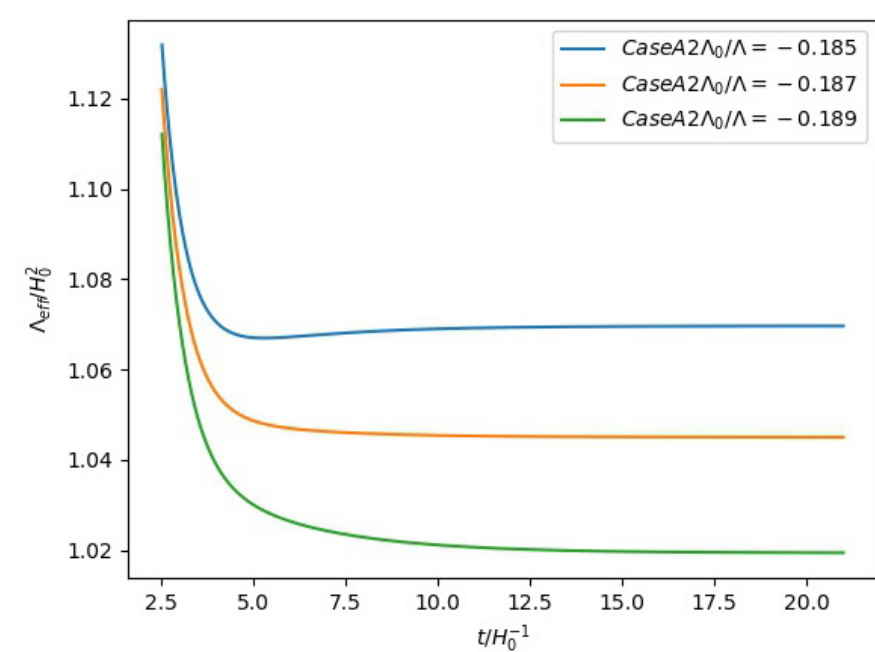
- Define the Effective Cosmological Constant which really responsible to the accelerating expansion

$$\Lambda_{eff}(t) = \Lambda_0 - 3 \left( \mathcal{H}(t)^2 + 2\mathcal{H}(t) \frac{\dot{a}(t)}{a(t)} \right)$$

- the critical value for  $\Lambda_0$  which symbolizes the transformation from a monotonically quintessence like  $\Lambda_{eff}(t)$  to the metastable dS potential can be solved for all of the case of approximations. The critical value  $\Lambda_{0-crit}$  centers around  $\Lambda_{0-crit} = 0$ . It can be conjectured the deviation of  $\Lambda_{0-crit}$  from 0 is caused by the approximations. In a more elaborated model, it should have  $\Lambda_{0-crit} = 0$ .



• The transformation from quintessence to dS



- Phenomenological, the  $\Lambda_{eff}$  can be regarded as the energy density produced by some auxiliary fields which responsible for the accelerating expansion such as quintessence field etc. e.g.

$$S_q = \int d^4x \sqrt{-g} \left[ \frac{1}{2} M_{pl}^2 R (1 + \xi \phi) - \frac{1}{2} g^{\mu\nu} \partial_\mu \phi \partial_\nu \phi - V(\phi) \right]$$

$$R_{\mu\nu} - \frac{1}{2} g_{\mu\nu} R = 8\pi G \frac{T_{\mu\nu}}{1 + \xi \phi}, \quad \ddot{\phi} + 3H\dot{\phi} + V_{,\phi} - \frac{1}{2} M_{pl}^2 R \xi = 0$$

$$\Lambda_{eff} = \frac{\frac{\dot{\phi}^2}{2} + V(\phi(t))}{1 + \xi \phi}, \quad \dot{\Lambda}_{eff} = \dot{\phi} \frac{\left( \frac{1}{2} M_{pl}^2 R - \Lambda_{eff} \right) \xi - 3H\dot{\phi}}{1 + \xi \phi}$$

$$V(\phi(t)) = \Lambda_{eff} (1 + \xi \phi) - \frac{\dot{\phi}^2}{2}$$

- the critical value for  $\Lambda_0$  which symbolizes the transformation from a monotonically quintessence like  $V(\phi(t))$  to the metastable dS potential is almost the same as one for  $\Lambda_{eff}$  in most cases esp in the case of landscape when

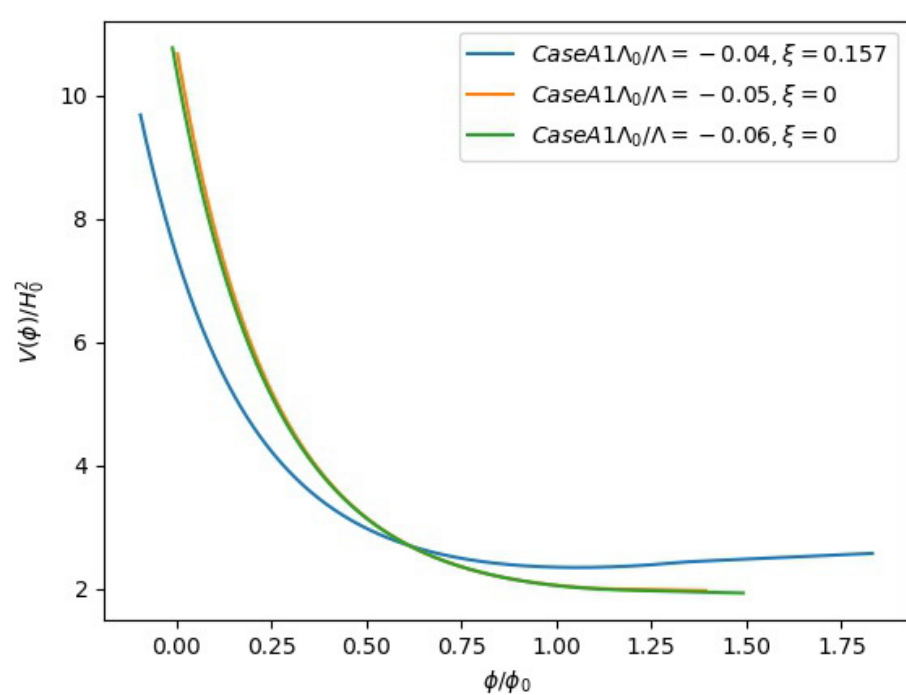
$$\Lambda_0 < \Lambda_{0-crit}$$

- However, when  $\Lambda_0 > \Lambda_{0-crit}$  the non-trivial Brans-Dicke type of coupling  $\xi_{min} > 0$  is required

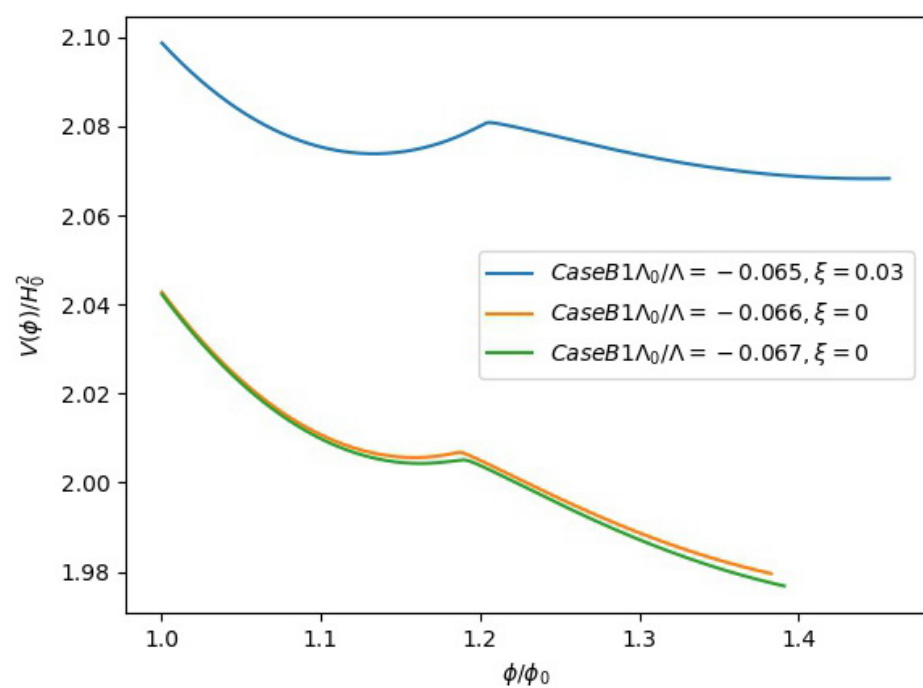
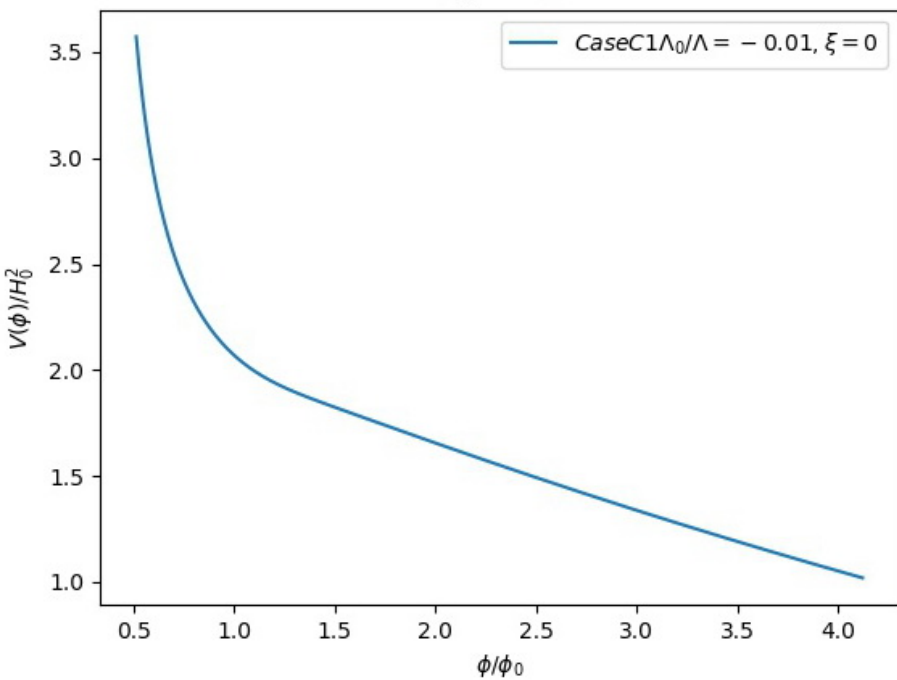
$$\dot{\phi} = \xi \frac{3(\dot{H} + 2H^2) - \Lambda_{eff}}{6H} \pm \sqrt{\left[ \frac{3(\dot{H} + 2H^2) - \Lambda_{eff}}{6H} \right]^2 \xi^2 - \frac{\dot{\Lambda}_{eff}}{3H} (1 + \xi\phi)}$$

- The restriction on  $\xi$

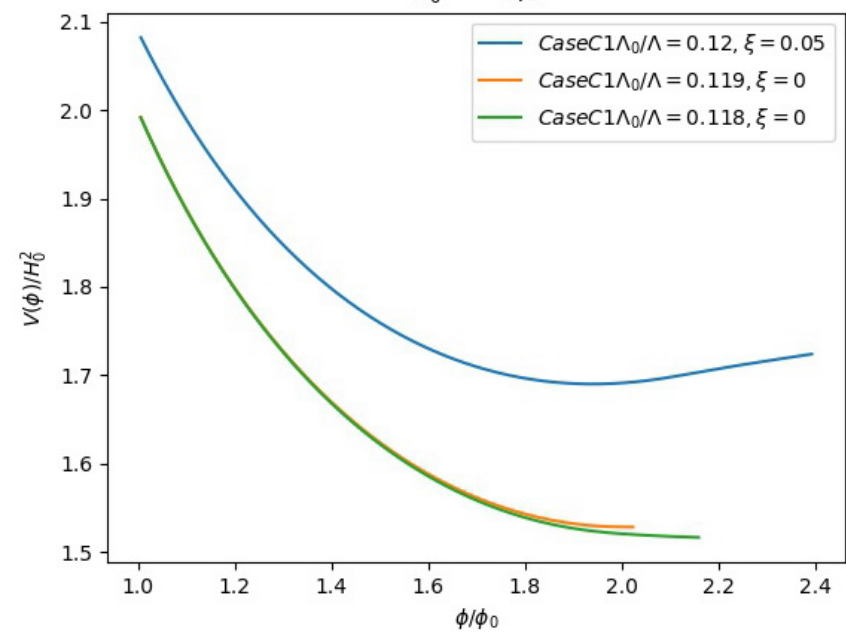
$$\left[ 3(\dot{H} + 2H^2) - \Lambda_{eff} \right]^2 \xi^2 - 12H\dot{\Lambda}_{eff} (1 + \xi\phi) \geq 0$$

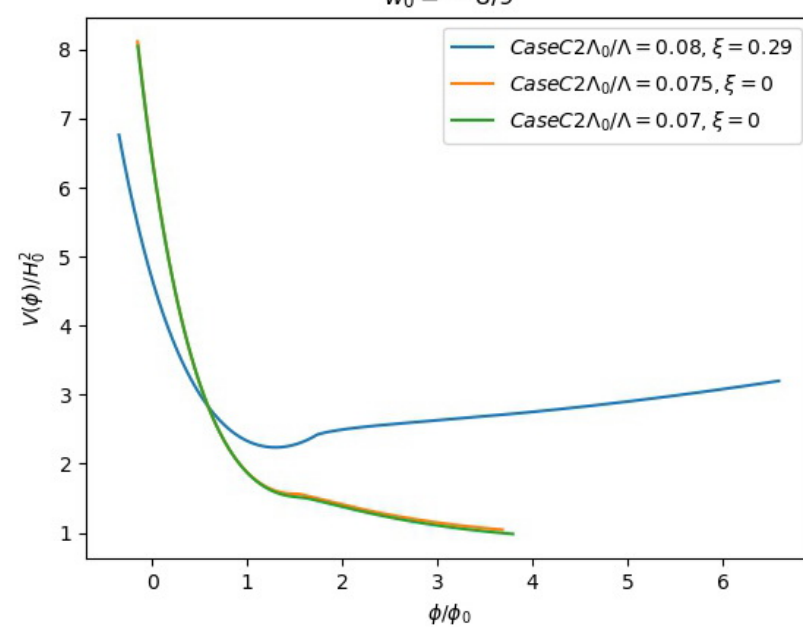
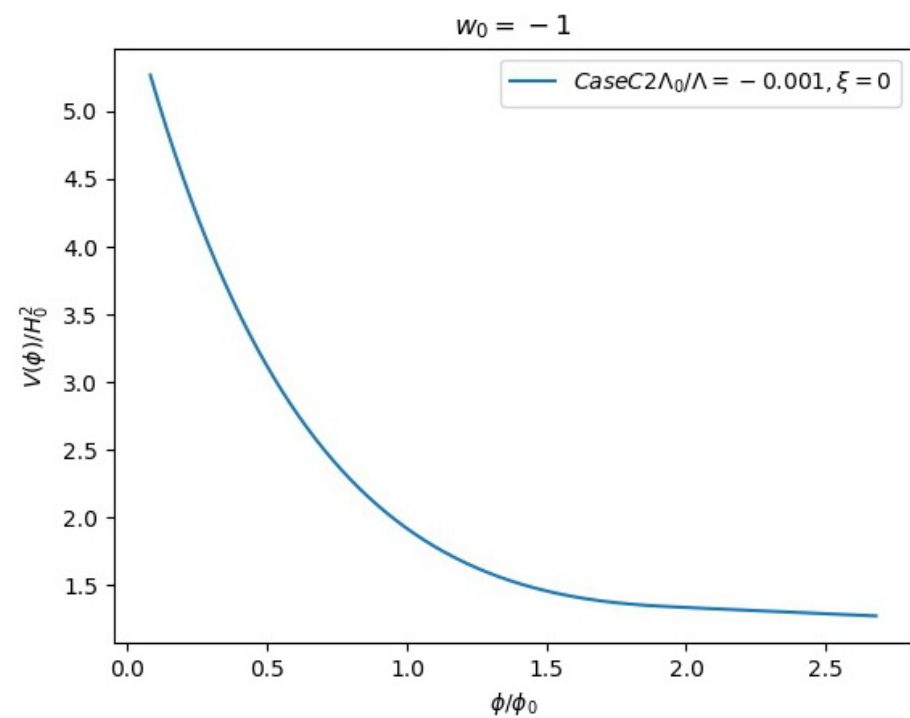
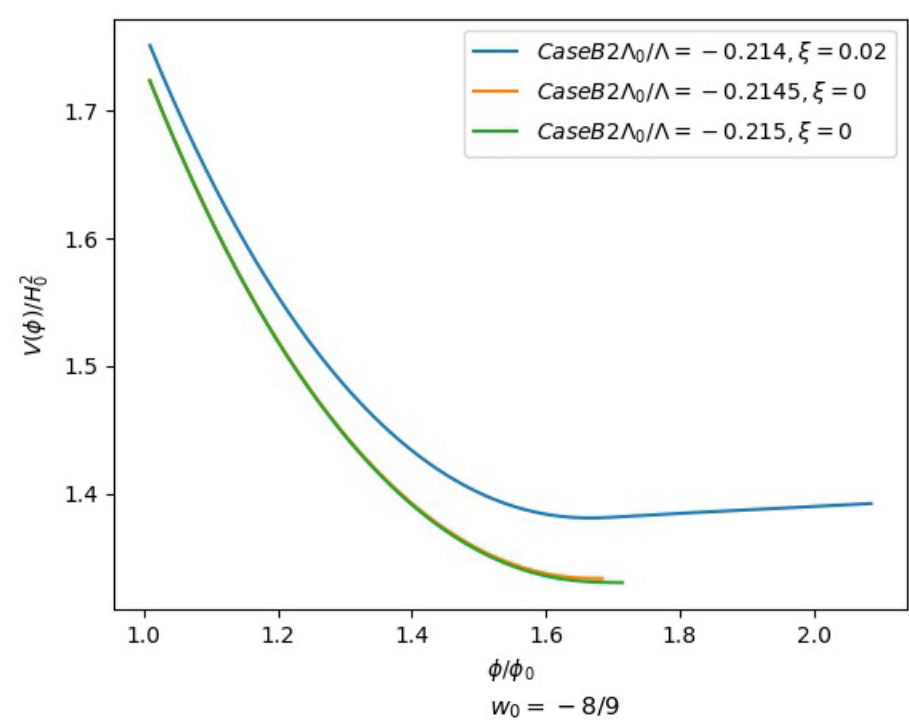
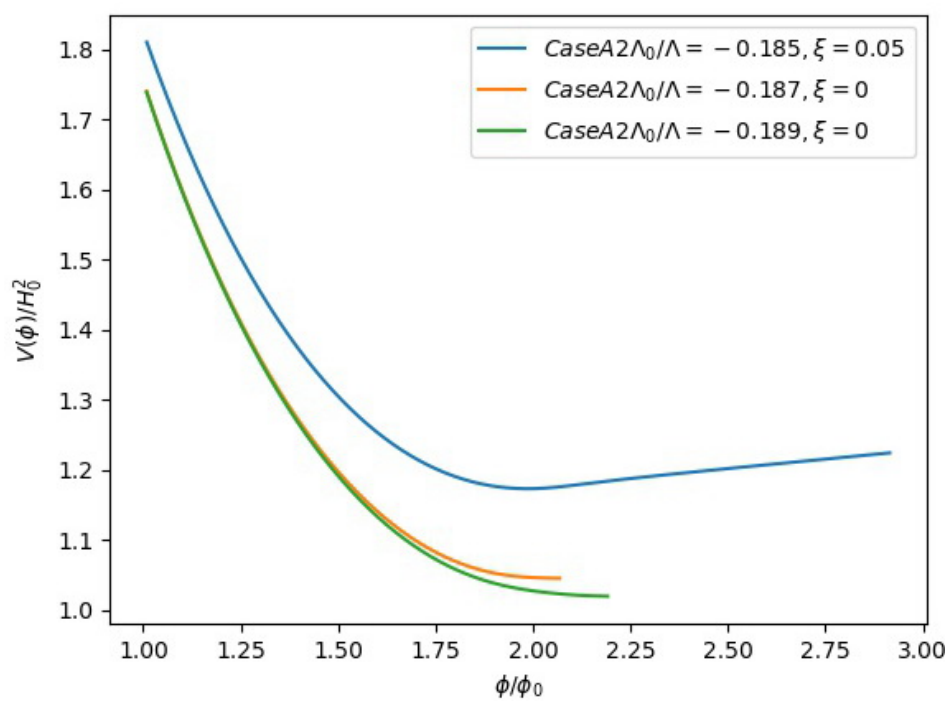


$w_0 = -1$



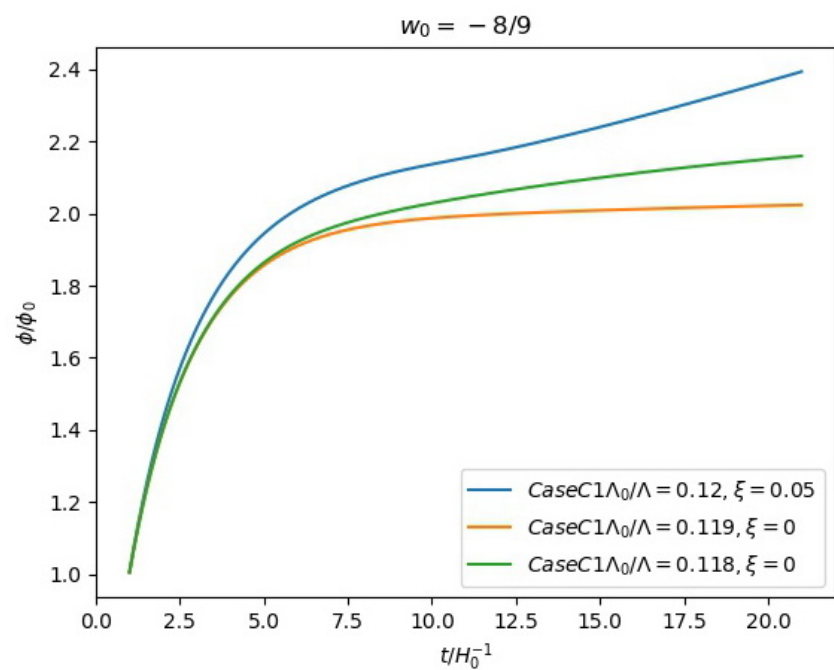
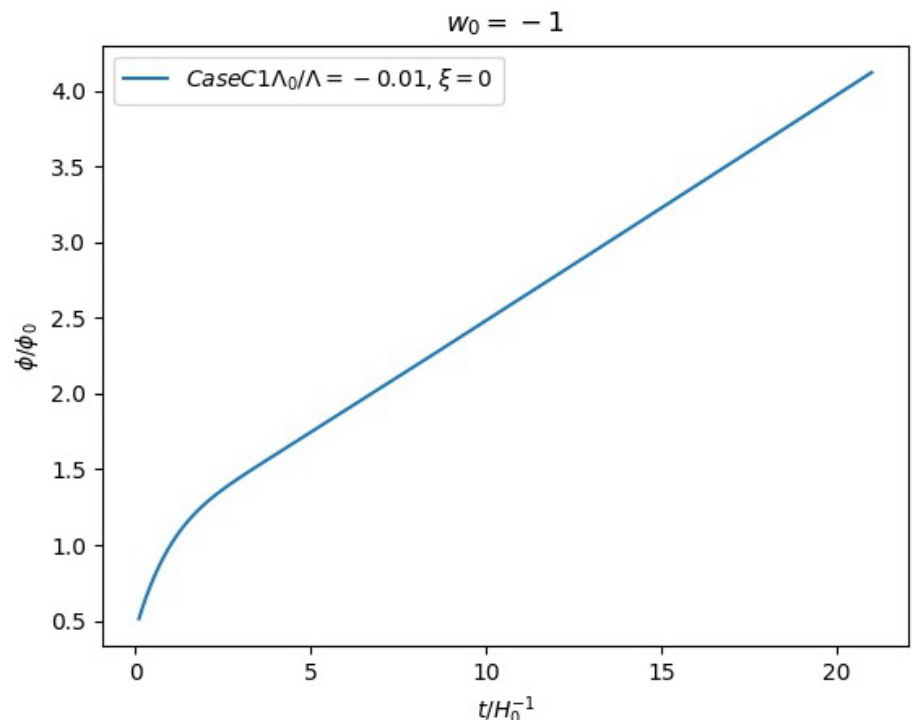
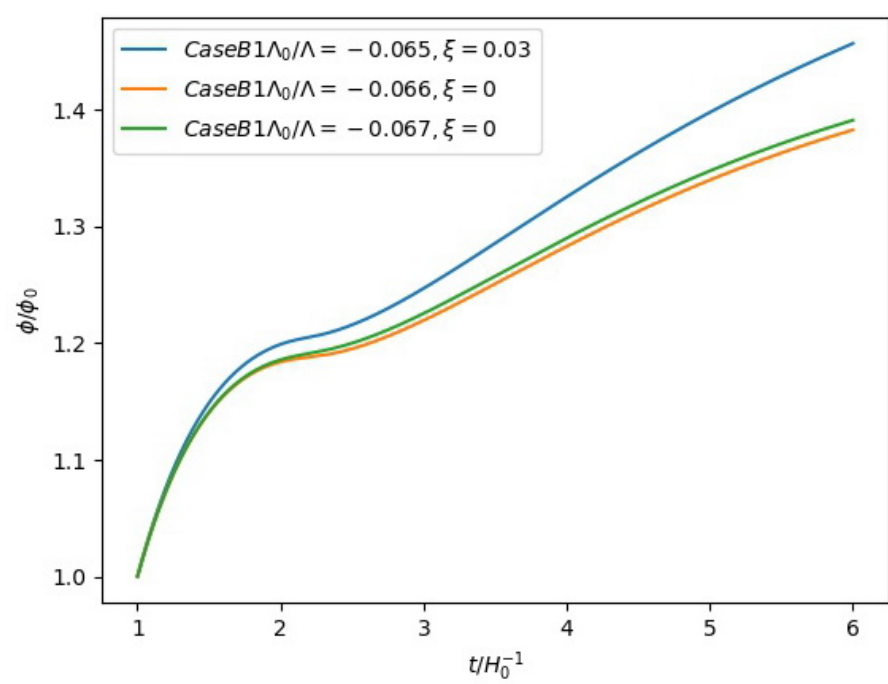
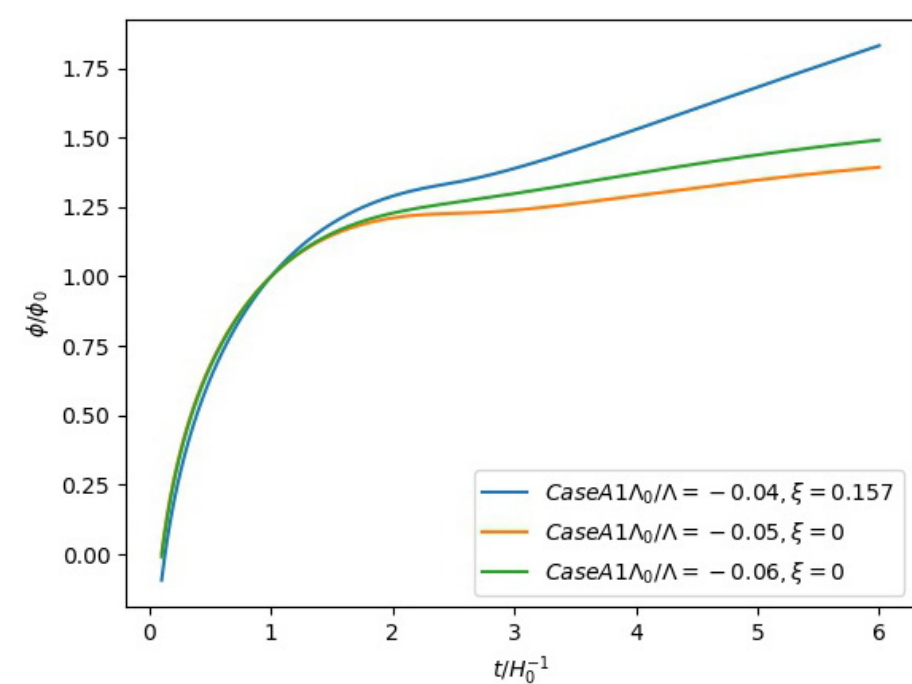
$w_0 = -8/9$

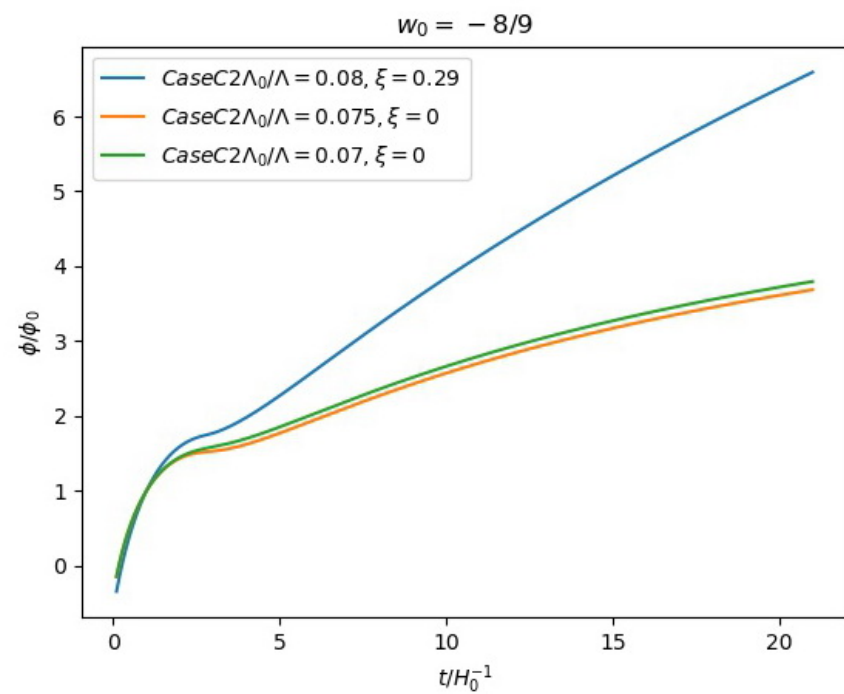
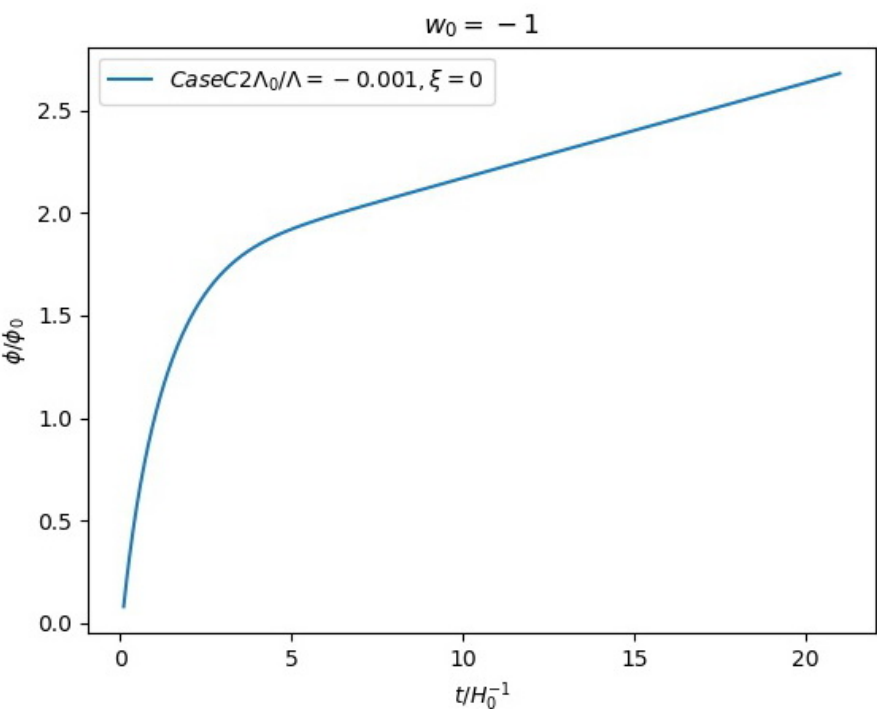
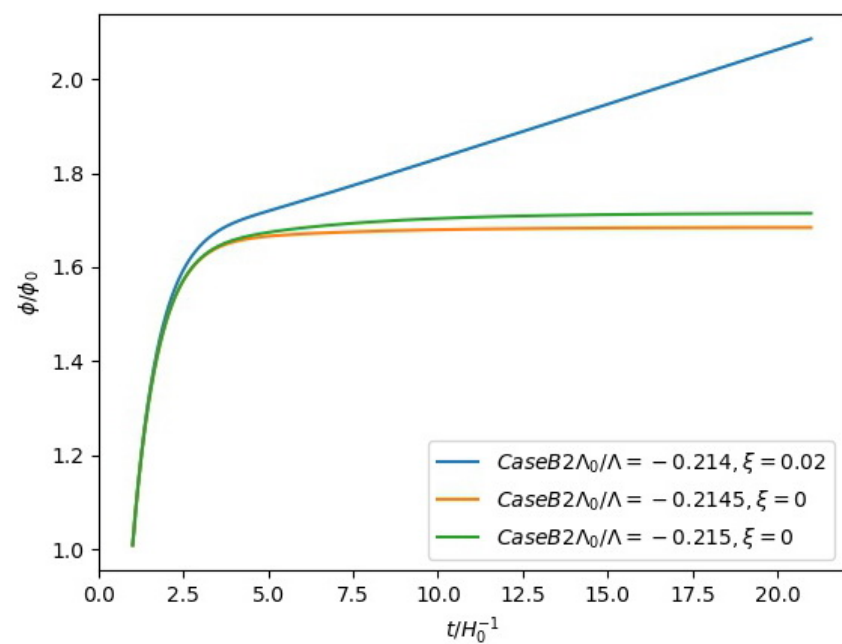
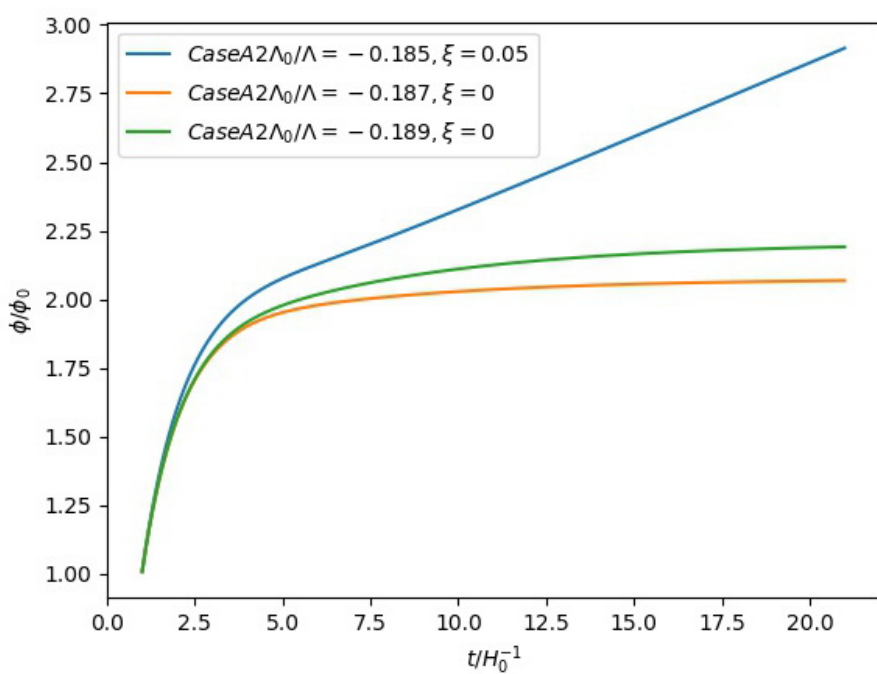


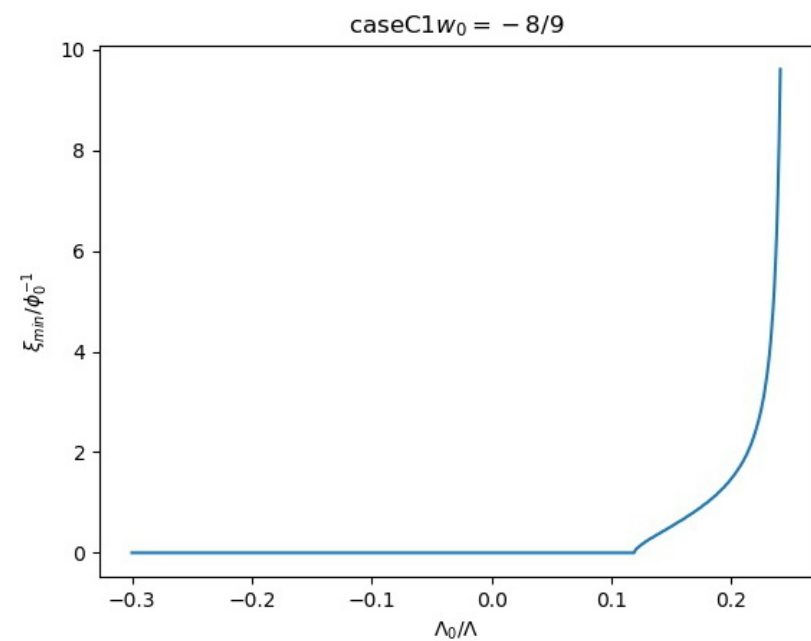
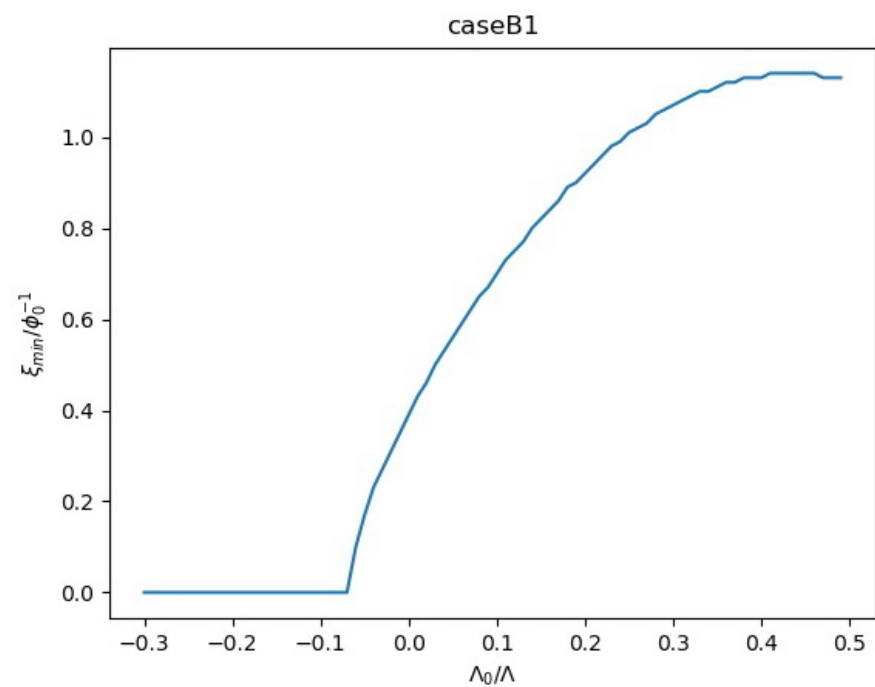
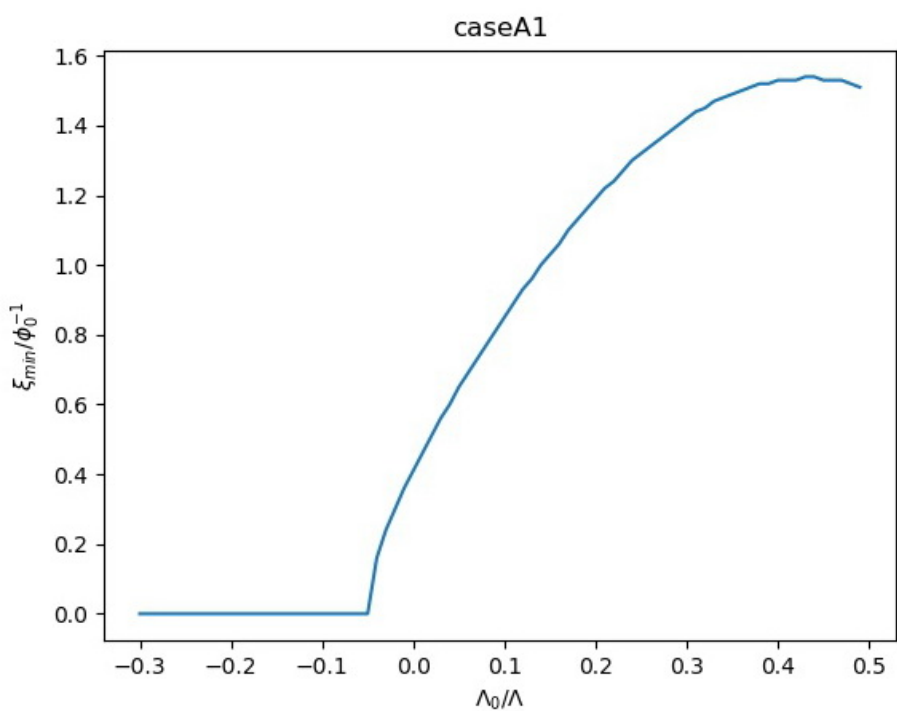




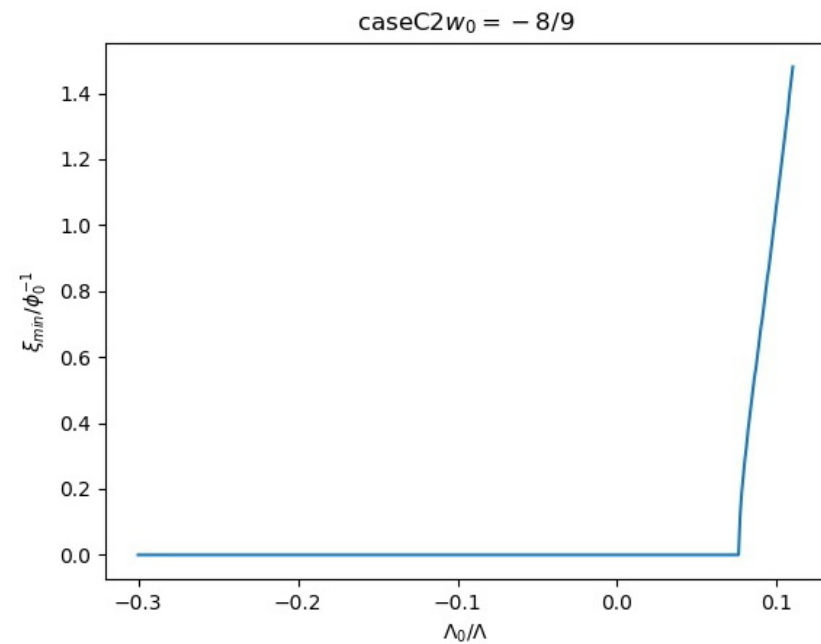
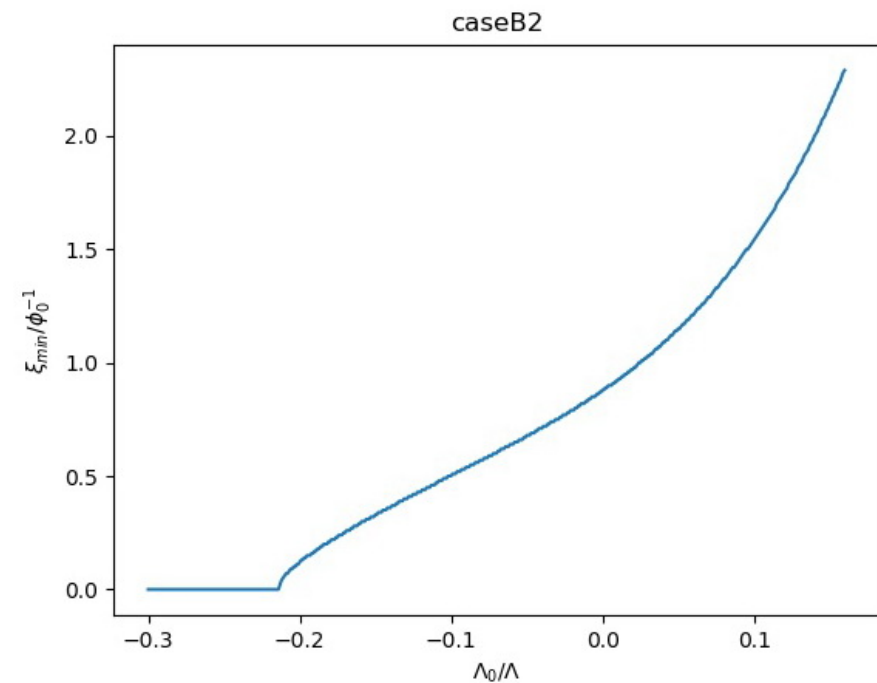
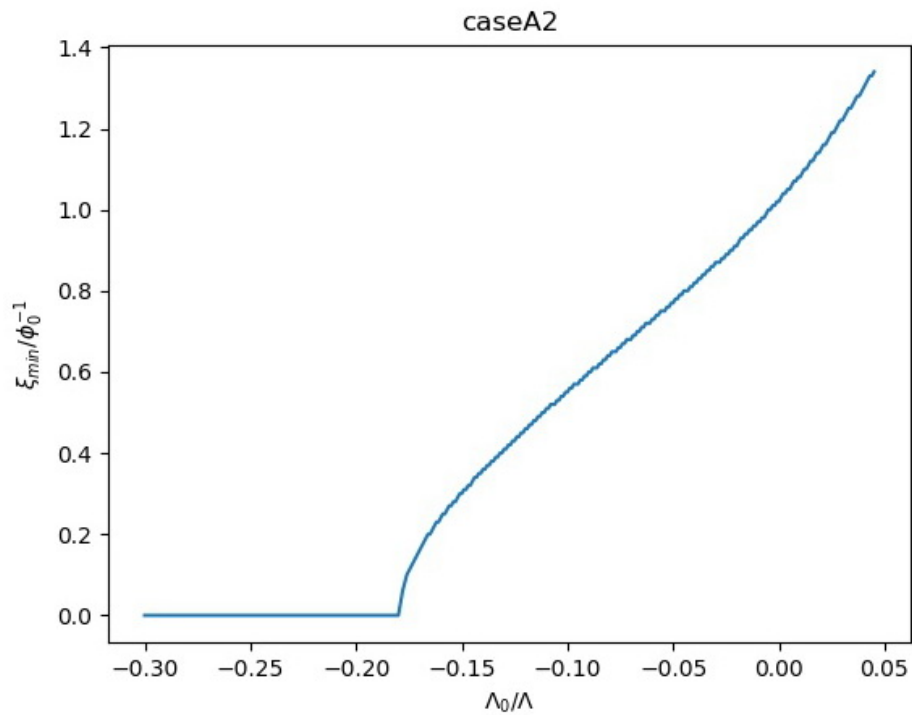
- For Case C, when  $w_0 > -8/9$  there doesn't exist a solution of the critical value for  $\Lambda_0$  which signs the transformation from a monotonically quintessence potential to a dS potential
- The  $\phi$  evolution versus  $t$  is monotonical which may be regarded as an intrinsic time for the evolution of the universe.







- The restrictions on  $\xi$ . The minimal  $\xi$  line.



- The restrictions on  $\xi$ . The minimal  $\xi$  line.

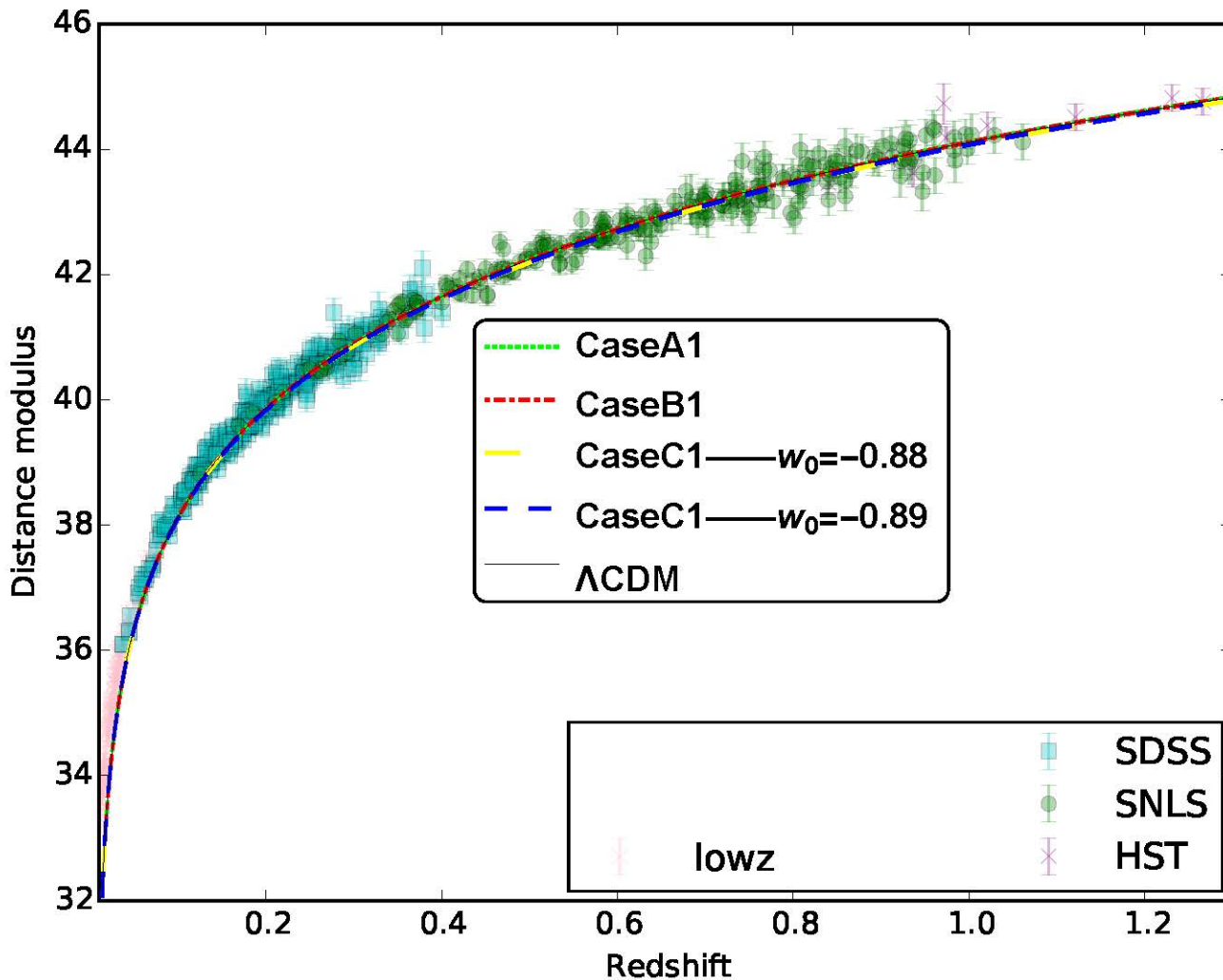
- Actually Case C approximation is not a good one from the comparison between Hubble constant vs  $t$  and the luminosity distance vs redshift  $z$ .
- The reason may be we use a fixed  $w_0$  in the equation of state of dark partner part. Ignore the case  $w_0 > -8/9$  (excluded by observation of luminosity distance with redshift relation), we can make the conclusion:
- Quintessence potential is generated from string landscape AdS vacuum effectively.
- The critical value of cosmological constant separating quintessence from metastable dS is approximately zero
- The metastable dS potential needs non-trivial coupling between quintessence field and gravitation. The  $w_0 = -1$  dS potential is actually stable but metastable and the non-trivial  $\xi$  coupling is divergent in this case.

- The metastable dS potential needs non-trivial Brans-Dicke coupling between quintessence field and gravitation. The  $w_0 = -1$  dS potential is actually stable but metastable and the non-trivial  $\xi$  coupling does not exist in this case.

$$\dot{\Lambda}_{eff} = \dot{\phi} \frac{\left( \frac{1}{2} M_{pl}^2 R - \Lambda_{eff} \right) \xi - 3H\dot{\phi}}{1 + \xi\dot{\phi}}, \quad \left( \frac{1}{2} M_{pl}^2 R - \Lambda_{eff} \right) \xi - 3H\dot{\phi} > 0$$

$$\frac{1}{2} M_{pl}^2 R - \Lambda_{eff} > 0, \quad \xi > \frac{3H\dot{\phi}}{\frac{1}{2} M_{pl}^2 R - \Lambda_{eff}}$$

- If  $\Lambda_{eff}$  grows too large to make  $\frac{1}{2} M_{pl}^2 R - \Lambda_{eff}$  from positive to negative.  $\xi$  would have no solution.
- The effective description of large scale Lorentz violation effect by quintessence is failed in the case.



- Comparison of distance modulus for

$$\Lambda_0 = 0.2\Lambda$$

$$\mathcal{K}(t_0) = H_0 \left( \sqrt{1 - \frac{\Lambda - \Lambda_0}{3H_0^2}} - 1 \right)$$



# Summary

- For string landscape with  $\Lambda_0 > -(3H_0 - \Lambda) \approx -\frac{2}{5}\Lambda$ , the effective cosmological constant naturally give a quintessence like potential which satisfies the dS Swampland conjecture
- The uplifting of AdS to a positive effective cosmological constant by frozen large scale Lorentz violation mechanism avoids the meta-stable dS swampland puzzle and have a quantum gravity origin
- For string swampland with positive cosmological constant for most reasonably approximation, the effective quintessence potential behaves as a metastable dS potential and the non-trivial Brans-Dicke type of coupling is required.
- For the stable dS type behavior, there is no solution for  $\xi$  and the effective potential approach is failed.

**THANKS!**