

Gauge Gravity Duality 2024



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Holographic spin alignment for vector mesons

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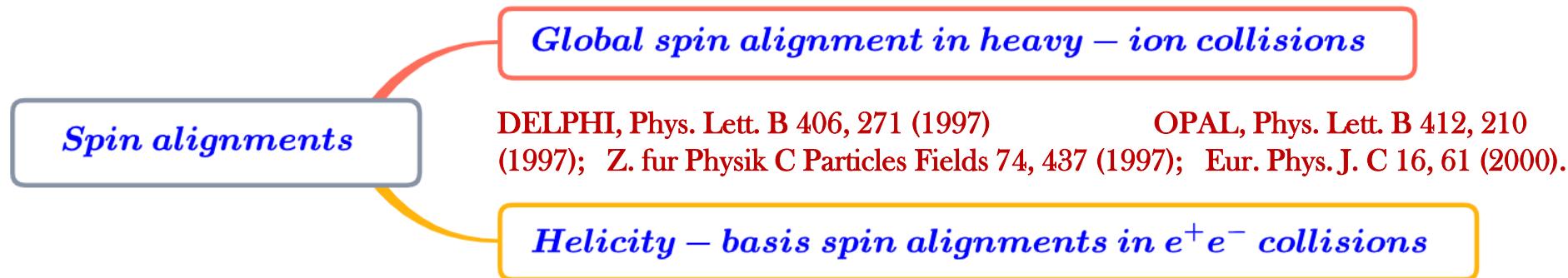
Based on: JHEP 08 (2024) 070 ; PRD 110 (2024) 5, 056047

What spin alignment ?

- The spin state of particles exhibits a certain non-random arrangement in specific directions.

Why spin alignment ?

- Different hadronization mechanisms: STAR, Nature 614, 244 (2023).



Measurements of the global spin alignment in this case provide a good opportunity to study the polarization correlation in the sQGP produced in heavy ion collisions.

How spin alignment ?

- Spin alignment is described by 00-element of meson's normalized spin density matrix.

Theoretical approaches

➤ Coalescence model with spin

- Quark/antiquark polarized by external field
- Non-equilibrium process described by kinetic theory
- Z.-T. Liang, X.-N. Wang, PLB 629, 20 (2005).
- XLS, Q. Wang, X.-N. Wang PRD 102, 056013 (2020).
- A. Kumar, B. Mueller, D.-L. Yang, PRD 108, 016020 (2023).
- XLS, L.Oliva, Z.-T.Liang, Q.Wang, X.-N.Wang, PRL 131, 042304 (2023);

➤ Spin kinetic equation

- D. Wagner, N. Weickgenannt, E. Speranza, PRR 5, 013187 (2023)
- S. Fang, S. Pu, D.-L. Yang, PRD 109, 034034 (2024)
- Y.-L. Yin, W.-B. Dong, J.-Y. Pang, S. Pu, Q. Wang, arXiv:2402.03672

➤ Linear response theory

- F. Li, S. Liu, arXiv: 2206.11890
- W.-B. Dong, Y.-L. Yin, XLS, S.-Z. Yang, Q. Wang,
- arXiv:2311.18400.

➤ Spectral function method

In this talk !!!

- Splitting between spectral functions of longitudinal and transverse modes due to external fields or motion relative to a thermal background, calculated by QFT, NJL model, holographic model...
- Meson at thermodynamical equilibrium
- XLS, S.-Y. Yang, Y.-L. Zou, D. Hou, arXiv: 2209.01872.
- A. Kumar, B. Mueller, D.-L. Yang, PRD 108, 016020 (2023). M. Wei, M. Huang, CPC 47, 104105 (2023).
- W.-B. Dong, Y.-L. Yin, XLS, S.-Z. Yang, Q. Wang, arXiv:2311.18400.
- XLS, Y.-Q. Zhao, S.-W. Li, F. Becattini, D. Hou, arXiv:2403.07522
- Y.-Q. Zhao, XLS, S.-W. Li, D. Hou, arXiv:2403.07468

A talk by Shu Lin in first day !!!

Theoretical approaches

1. Z.-T. Liang, X.-N. Wang, PLB 629, 20 (2005)
2. F. Becattini, L. Csernai, D.-J. Wang, PRC 88, 034905 (2013)
3. Y.-G. Yang, R.-H. Fang, Q. Wang, X.-N. Wang, PRC 97, 034917 (2018)
4. X.-L. S, L. Oliva, Q. Wang, PRD 101, 096005 (2020)
5. F. Li, S. Liu, arXiv: 2206.11890
6. M. Wei, M. Huang, CPC 47, 104105 (2023)
7. S. Fang, S. Pu, D.-L. Yang, arXiv:2311.15197.
8. P. H. D. Moura, K. J. Goncalves, G. Torrieri, PRD 108, 034032 (2023)
9. X.-L. Xia, H. Li, X.-G. Huang, H.-Z. Huang, PLB 817, 136325 (2021)
10. A. Kumar, P. Gubler, D.-L. Yang, arXiv:2312.16900
11. D. Wagner, N. Weickgenannt, E. Speranza, PRR 5, 013187 (2023)
12. W.-B. Dong, Y.-L. Yin, XLS, S.-Z. Yang, Q. Wang, arXiv:2311.18400.
13. F. Sun, J. Shao, R. Wen, K. Xu, M. Huang, arXiv: 2402.16595.
14. X.-L. S, S.-Y. Yang, Y.-L. Zou, D. Hou, arXiv:2209.01872
15. Y.-Q. Zhao, X.-L. S, S.-W. Li, D. Hou, JHEP 08 (2024) 070
16. B. Muller, D.-L. Yang, PRD 105, 1 (2022).
17. J.-H. Gao, PRD 104, 076016 (2021)
18. A. Kumar, B. Muller, D.-L. Yang, PRD 108, 016020 (2023)
19. X.-L. S, S. Pu, Q. Wang, PRC 108, 054902 (2023). 20. X.-L. S, Y.-Q. Zhao, S.-W. Li, F. Becattini, D. Hou, PRD 110 (2024) 5, 056047
21. X.-L. S, L. Oliva, Z.-T. Liang, Q. Wang, X.-N. Wang, PRL 131, 042304 (2023); PRD 109, 036004 (2024)

$\rho_{00} \approx \frac{1}{3} + c_{hadro} + c_{EM} + c_F + c_A + c_h + c_{strong} + \dots$

Smaller deviation from 1/3 Larger deviation from 1/3 ?

$\rho_{00} \approx \frac{1}{3} + c_{hadro} + c_{EM} + c_F + c_A + c_h + c_{strong} + \dots$

Anisotropic strong force [4,15,18-21]

Helicity Polarization[17]

Anomalous color field [16]

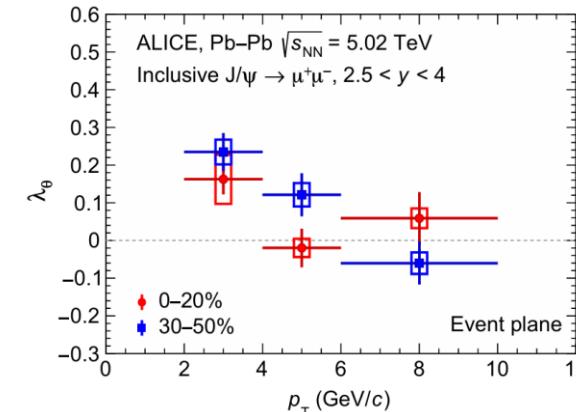
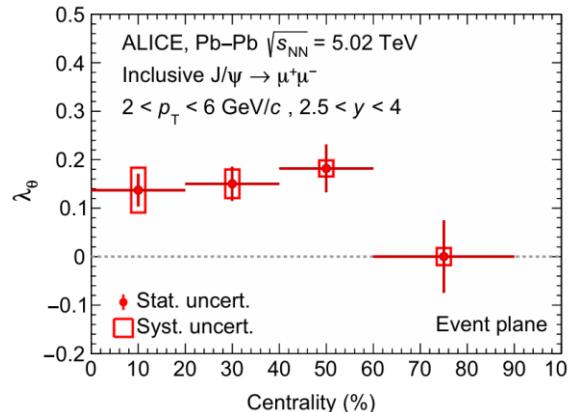
Fragmentation [1]

Electromagnetic fields [3,4,14,15]

Hydrodynamic gradient (vorticity; Acceleraction; shear tensor; second order;) [1-13]

Experiment results

ALICE, Phys. Rev. Lett. 131 (4) (2023) 042303.



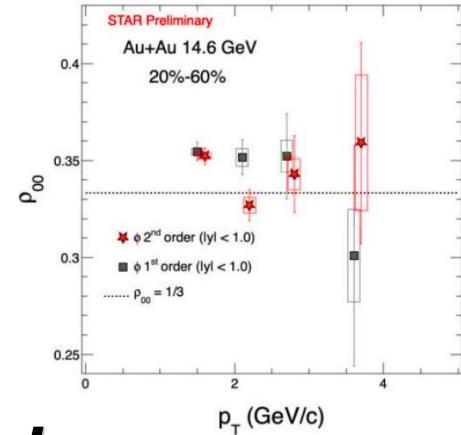
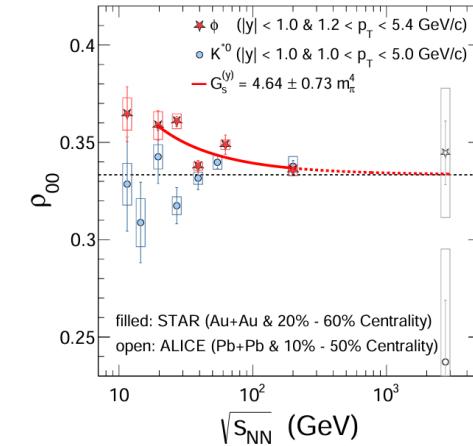
J/ψ

$$\rho_{00} < 1/3$$

$$\lambda_\theta = \frac{1 - 3\rho_{00}}{1 + \rho_{00}},$$

J/ψ are preferably to be transversely polarized.

STAR, Nature 614 (7947) (2023) 244–248.



ϕ

$$\lambda_\theta = \frac{3\rho_{00} - 1}{1 - \rho_{00}}$$

ϕ are preferably to be longitudinally polarized.

$$\rho_{00} = 1/3 \text{ no spin alignment}$$

Outline

- Theory
 - Spin alignment
 - Correlation function
- Momentum induced spin alignment
- Magnetic field induced spin alignment
- Summary and Outlook

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Spin alignment

- S-matrix element($J/\psi \rightarrow l + \bar{l}$ and $\phi \rightarrow l + \bar{l}$):

C. Gale and J. I. Kapusta, Nucl. Phys. B 357 (1991) 65–89

$$S_{fi} = \int d^4x d^4y \langle f, \bar{l}l | J_\mu(y) G_R^{\mu\nu}(x-y) J_\nu^l(x) | i \rangle$$

where

J_μ is the current that couples to V.M. ;

J_ν^l is the leptonic current

$$J_\nu^l(x) = g_M l \bar{\psi}_l(x) \Gamma_\nu \psi_l(x)$$

- Propagators (vacuum):

$$G_{R/A}^{\mu\nu} = -\frac{\eta^{\mu\nu} + p^\mu p^\nu / p^2}{p^2 + m_V^2 \pm im_V\Gamma}$$

Coordinate space

- Retarded current-current correlation:

$$D^{\mu\nu}(x, p) \equiv \int d^4y \theta(y^0) \langle [J^\mu(y), J^\nu(0)] \rangle_{T(x)} e^{-ip \cdot y}$$

- Spectral function:

$$\varrho_{\alpha\beta}(x, p) \equiv -\text{Im}D_{\alpha\beta}(x, p)$$

- Differential production rate:

$$n(x, p) = -\frac{2g_M^2 l \bar{l}}{3(2\pi)^5} \left(1 - \frac{2m_l^2}{p^2}\right) \sqrt{1 + \frac{4m_l^2}{p^2}} p^2 n_B(x, \omega) \times \left(\eta_{\mu\nu} + \frac{p_\mu p_\nu}{p^2}\right) G_A^{\mu\alpha}(p) \varrho_{\alpha\beta}(x, p) G_R^{\beta\nu}(p),$$

Y. Burnier, M. Laine, M. Vepsäläinen,, JHEP 02 (2009) 008.

L. D. McLerran, T. Toimela, Phys. Rev. D 31 (1985) 545.

H. A. Weldon, Phys. Rev. D 42 (1990) 2384–2387.

Spin alignment

➤ Spectral function:

$$Q^{\mu\nu}(x, p) = \sum_{\lambda, \lambda'=0, \pm 1} v^\mu(\lambda, p) v^{*\nu}(\lambda', p) \tilde{Q}_{\lambda\lambda'}(x, p)$$

➤ Polarization vectors:

$$v^\mu(\lambda, p) = \left(\frac{\mathbf{p} \cdot \mathbf{\epsilon}_\lambda}{M}, \mathbf{\epsilon}_\lambda + \frac{\mathbf{p} \cdot \mathbf{\epsilon}_\lambda}{M(\omega + M)} \mathbf{p} \right)$$

Orthonormality conditions:

$$\eta_{\mu\nu} v^\mu(\lambda, p) v^{*\nu}(\lambda', p) = \delta_{\lambda\lambda'}$$

Completeness condition:

$$\sum_\lambda v^\mu(\lambda, p) v^{*\nu}(\lambda, p) = (\eta^{\mu\nu} + p^\mu p^\nu / p^2)$$

➤ Dilepton production rate:

$$n_\lambda(x, p) = -\frac{2g_M^2 l \bar{l}}{3(2\pi)^5} \left(1 - \frac{2m_l^2}{p^2} \right)$$

$$\times \sqrt{1 + \frac{4m_l^2}{p^2}} \frac{p^2 n_B(x, \omega) \tilde{Q}_{\lambda\lambda}(x, p)}{(p^2 + m_V^2)^2 + m_V^2 \Gamma^2}.$$

➤ Spin alignment:

$$\rho_{00}(x, \mathbf{p}) \equiv \frac{\int d\omega n_0(x, p)}{\sum_{\lambda=0, \pm 1} \int d\omega n_\lambda(x, p)}$$

Spin space

- Theory
 - Spin alignment
 - Correlation function
- Momentum induced spin alignment
- Magnetic field induced spin alignment
- Summary and Outlook

Correlation function

- AdS/CFT dictionary:

$$Z_{\text{QFT}}[A_\mu^{(0)}] = Z_{\text{gravity}}[A_\mu]$$

where

$$Z_{\text{QFT}}[A_\mu^{(0)}] = \left\langle \exp \left\{ \int_{\partial\mathcal{M}} J^\mu A_\mu^{(0)} d^4x \right\} \right\rangle,$$

$$Z_{\text{gravity}}[A_\mu] = \exp \{-S_{\text{bulk}}[A_\mu]\}$$

- Bulk geometry: $0 \leq \zeta \leq \zeta_h$

$$ds^2 = g_{\mu\nu} dx^\mu dx^\nu + g_{\zeta\zeta} d\zeta^2$$

- Two-point correlators

$$\langle [J^\mu, J^\nu] \rangle = D^{\mu\nu} \propto \frac{\delta^2 Z_{\text{QFT}}}{\delta A_\mu^{(0)} \delta A_\nu^{(0)}} = \frac{\delta^2 Z_{\text{gravity}}}{\delta A_\mu^{(0)} \delta A_\nu^{(0)}}$$

- Bulk mesonic action:

T. Sakai, S. Sugimoto, Prog. Theor. Phys. 113 (2005) 843882.

T. Sakai, S. Sugimoto, Prog. Theor. Phys. 114 (2005) 1083-1118.

A. Karch, E. Katz, D. T. Son, M. A. Stephanov, Phys. Rev. D 74(2006) 015005.

$$S_{\text{bulk}} = - \int d^4x d\zeta Q(\zeta) F_{MN} F^{MN}$$

- Equation of motion: $\partial_M [Q(\zeta) F^{MN}] = 0$ Radial gauge condition: $A_\zeta = 0$

- Fourier transformation: $A_\mu(x, \zeta) = \int \frac{d^4p}{(2\pi)^4} e^{-ip\cdot x} A_\mu(p, \zeta)$

- Electric fields: $E_i(p, \zeta) \equiv -p_0 A_i(p, \zeta) + p_i A_0(p, \zeta)$

- Boundary condition: $\lim_{\zeta \rightarrow 0} \tilde{E}_i(j, p, \zeta) = \delta_{ij}$

- Equation of motion :
$$\begin{aligned} \partial_\zeta^2 E_i(p, \zeta) + \frac{[\partial_\zeta Q(\zeta) g^{\zeta\zeta}]}{Q(\zeta) g^{\zeta\zeta}} [\partial_\zeta E_i(p, \zeta)] - \frac{p^2}{g^{\zeta\zeta}} E_i(p, \zeta) \\ + (-p_0 g_{i\mu} + p_i g_{0\mu})(\partial_\zeta g^{\mu\nu}) [\partial_\zeta A_\nu(p, \zeta)] = 0, \end{aligned}$$

Correlation function

➤ Correlation function:

D. T. Son and A. O. Starinets, JHEP 09 (2002) 042

$$D^{\mu\nu}(p) = \lim_{\zeta \rightarrow 0} g^{\zeta\zeta} g^{\mu\alpha} Q(\zeta) \frac{\delta[\partial_\zeta A_\alpha(p, \zeta)]}{\delta A_\nu(p, \zeta)} \Big|_{A_\mu(p, 0)=0}$$

$$D^{\mu 0}(p) = -\lim_{\zeta \rightarrow 0} \frac{p_j}{p_0} \left(g^{\mu k} - \frac{p^\mu p^k}{p^2} \right) \frac{1}{\zeta} \partial_\zeta \tilde{E}_k(j, p, \zeta),$$

$$D^{\mu i}(p) = \lim_{\zeta \rightarrow 0} \left(g^{\mu k} - \frac{p^\mu p^k}{p^2} \right) \frac{1}{\zeta} \partial_\zeta \tilde{E}_k(i, p, \zeta).$$

➤ Ward identity: $p_\mu D^{\mu\nu} = p_\mu D^{\nu\mu} = 0$

➤ Spectral function:

$$\tilde{\varrho}_{\lambda\lambda}(p) = v_\mu^*(\lambda, p) v_\nu(\lambda, p) \text{Im}D^{\mu\nu}(p)$$

➤ Theory

➤ Spin alignment

➤ Correlation function

➤ Momentum induced spin alignment

➤ Magnetic field induced spin alignment

➤ Summary and Outlook

Momentum induced spin alignment

➤ Holographic background metric:

$$ds^2 = \frac{L^2}{\zeta^2} \left(-f(\zeta)dt^2 + dx^2 + dy^2 + dz^2 + \frac{d\zeta^2}{f(\zeta)} \right)$$

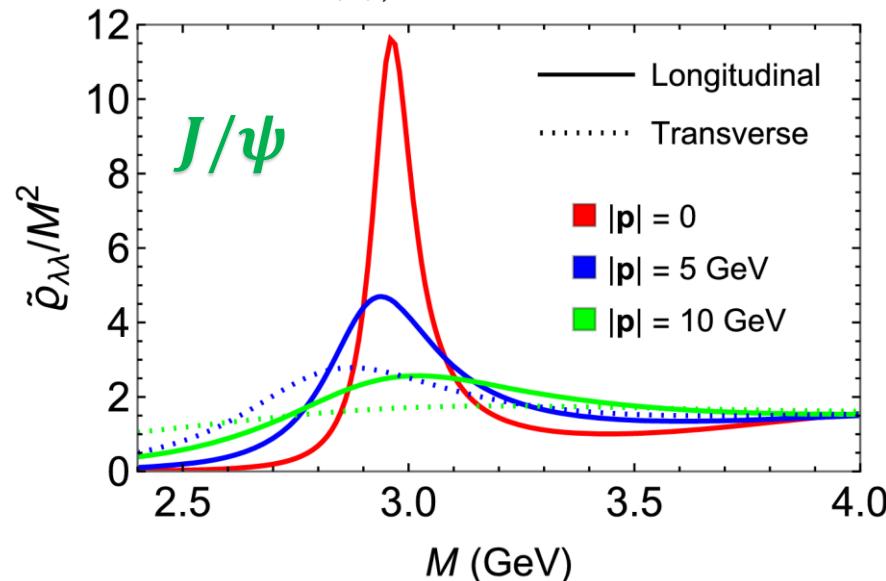
$$f(\zeta) = 1 - \zeta^4/\zeta_h^4$$

➤ Function:

$$\mathcal{Q}(\zeta) = e^{-\Phi(\zeta)} \sqrt{-g}/(4g_5^2)$$

➤ Dilaton field:

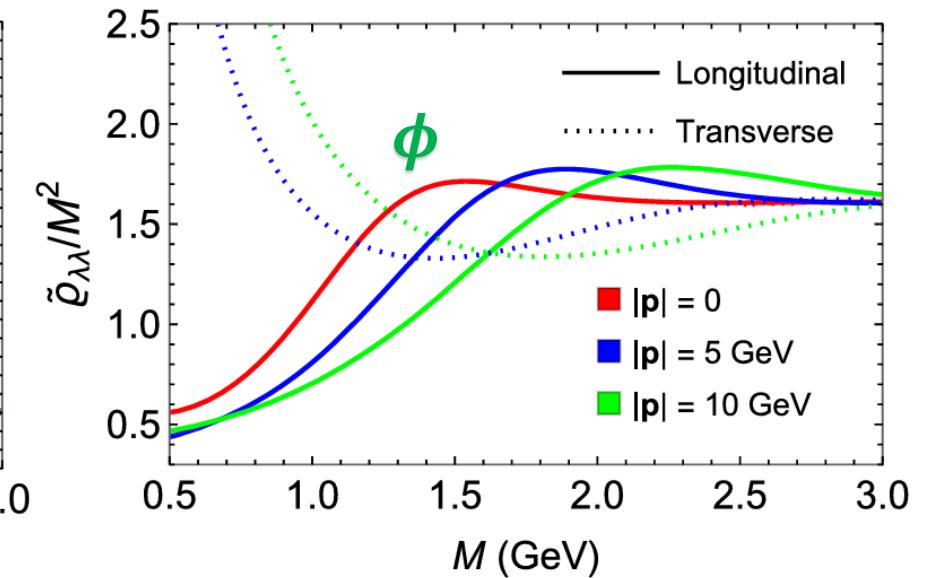
$$\Phi(\zeta) = c\zeta^2$$



- At zero momentum, spectral functions for all spin states are degenerate because of the rotation symmetry.
- High-momentum resonances are harder to be produced than low-momentum resonances.

Spectral function

$T = 150 \text{ MeV}$

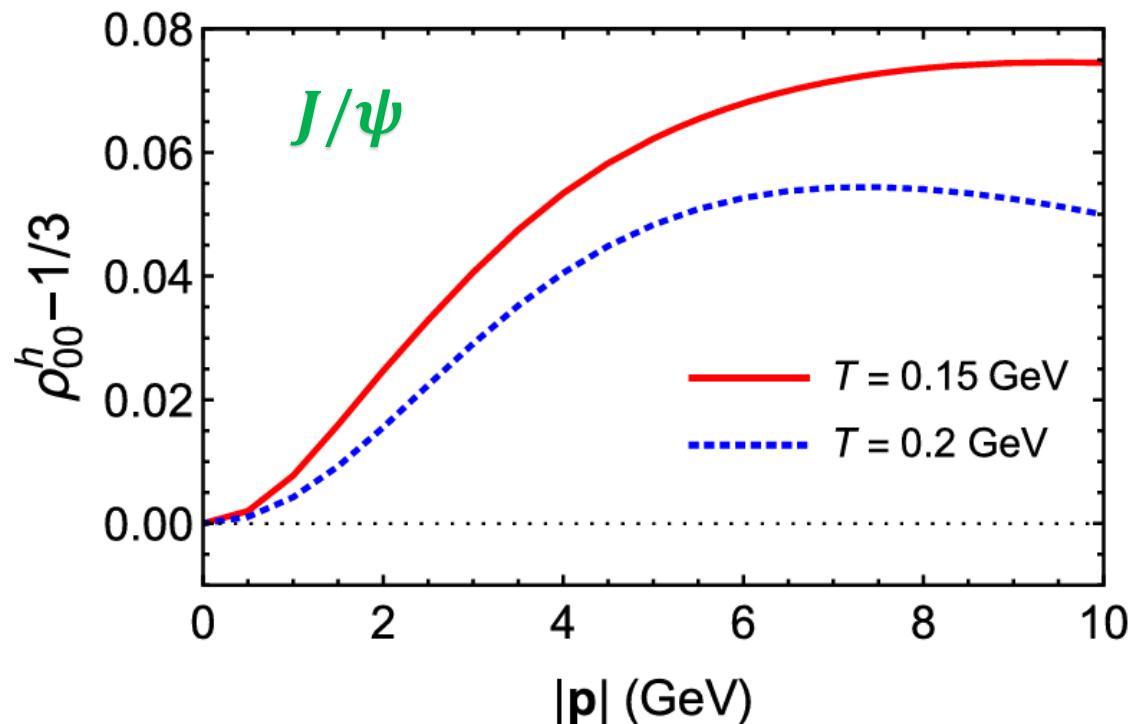


- $s - \bar{s}$ pairs cannot exist as quasi-stable particles.
- Spectral functions should be interpreted as probabilities of $s - \bar{s}$ pairs.

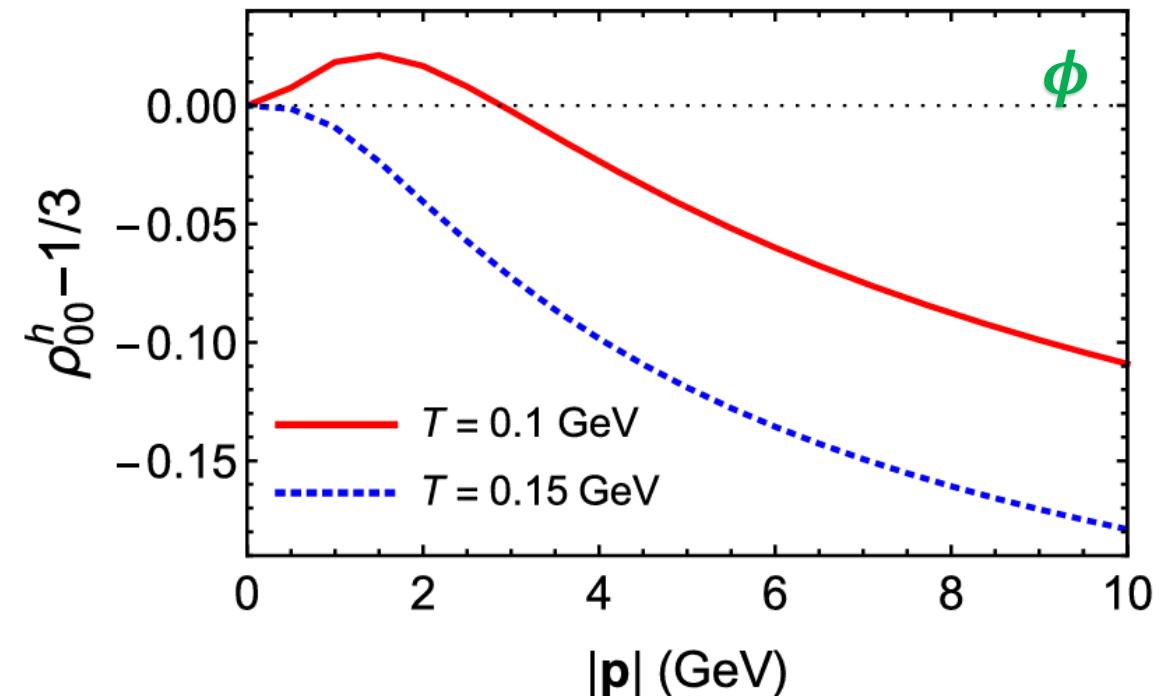
Momentum induced spin alignment

Spin alignment in the helicity frame

$$\lambda_\theta = \frac{1 - 3\rho_{00}}{1 + \rho_{00}},$$



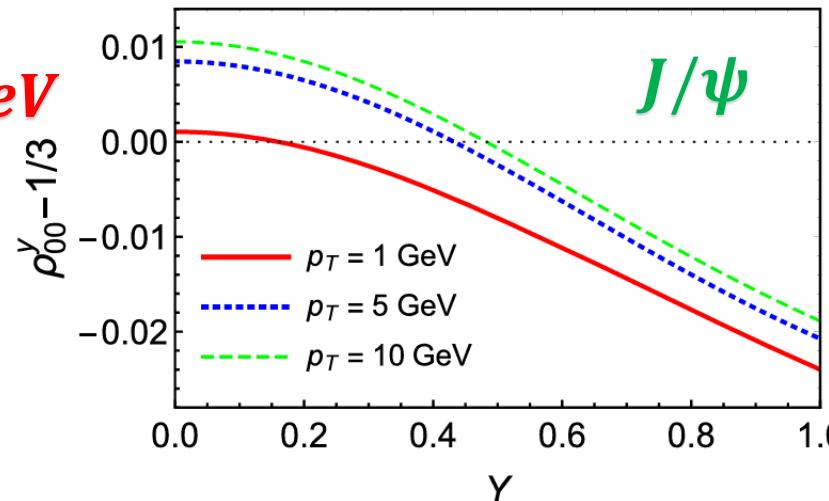
$$\lambda_\theta = \frac{3\rho_{00} - 1}{1 - \rho_{00}}$$



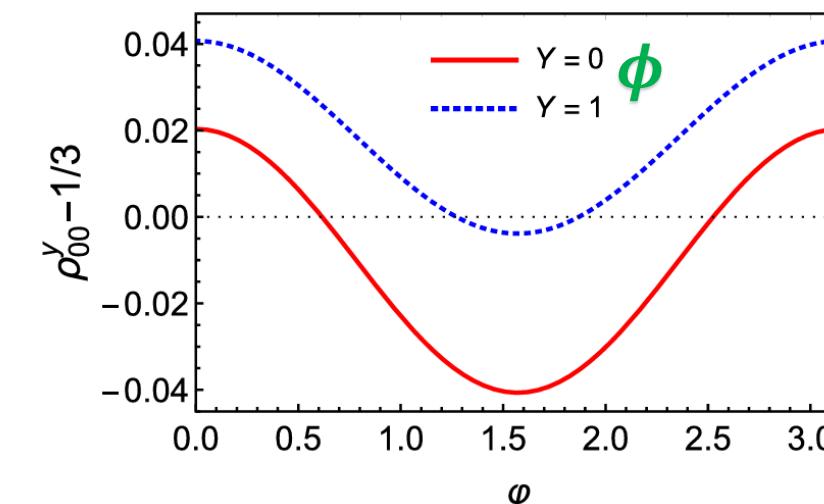
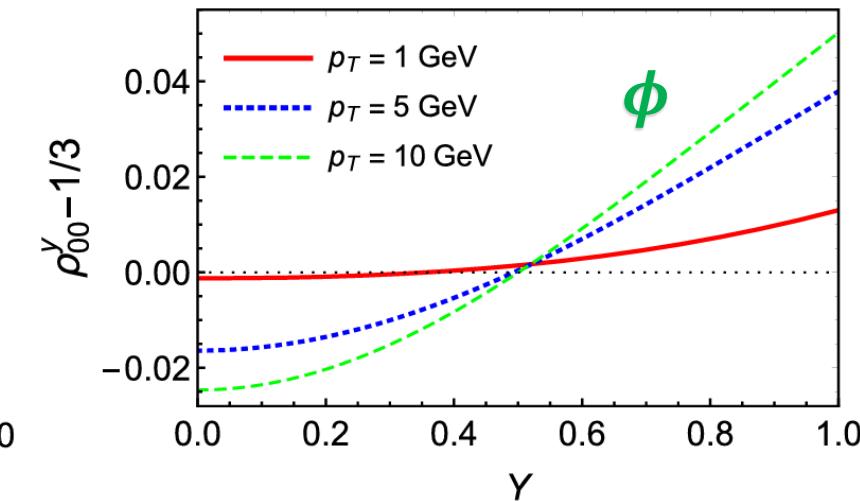
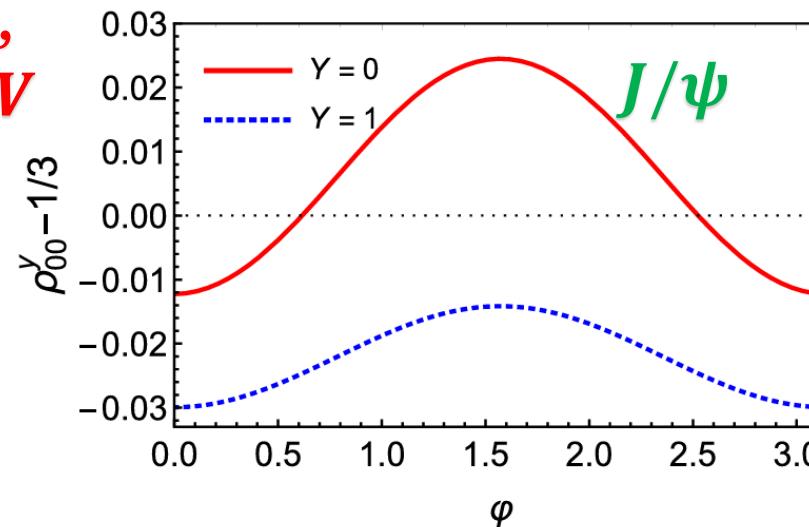
Momentum induced spin alignment

Global Spin alignment

$T = 150 \text{ MeV}$



**$p_T = 2 \text{ GeV},$
 $T = 150 \text{ MeV}$**



- Theory
 - Spin alignment
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- **Magnetic field induced spin alignment**
- Summary and Outlook

Magnetic field induced spin alignment

➤ Action:

J/ψ

➤ Magnetic field constraint conditions

$$S = \frac{1}{16\pi G_5} \int d^5x \sqrt{-g} \left(R + \frac{12}{L^2} \right) + \frac{1}{8\pi G_5} \int d^4x \sqrt{-\gamma} \left(K - \frac{3}{L} \right) + S_f,$$

$$eB < \sqrt{\frac{3}{2}} \frac{1.6}{\zeta_h^2} \approx \frac{1.96}{\zeta_h^2}$$

$$S_f = -\frac{N_c}{16\pi^2} \int d^4x \int_0^{\zeta_h} d\zeta Q(\zeta) \text{Tr} \left(F_L^2 + F_R^2 \right),$$

➤ Magnetic field parallel to momentum

➤ Holographic background metric:

$$ds^2 = \frac{L^2}{\zeta^2} \left(-f(\zeta)dt^2 + h_T(\zeta)(dx^2 + dy^2) + h_P(\zeta)dz^2 + \frac{d\zeta^2}{f(\zeta)} \right).$$

$$f(\zeta) = 1 - \frac{\zeta^4}{\zeta_h^4} + \frac{2}{3} \frac{e^2 B^2}{1.6^2} \zeta^4 \ln \frac{\zeta}{\zeta_h} + \mathcal{O}(e^4 B^4),$$

$$h_T(\zeta) = 1 - \frac{4}{3} \frac{e^2 B^2}{1.6^2} \zeta_h^4 \int_0^{\zeta/\zeta_h} \frac{y^3 \ln y}{1-y^4} dy + \mathcal{O}(e^4 B^4),$$

$$h_P(\zeta) = 1 + \frac{8}{3} \frac{e^2 B^2}{1.6^2} \zeta_h^4 \int_0^{\zeta/\zeta_h} \frac{y^3 \ln y}{1-y^4} dy + \mathcal{O}(e^4 B^4),$$

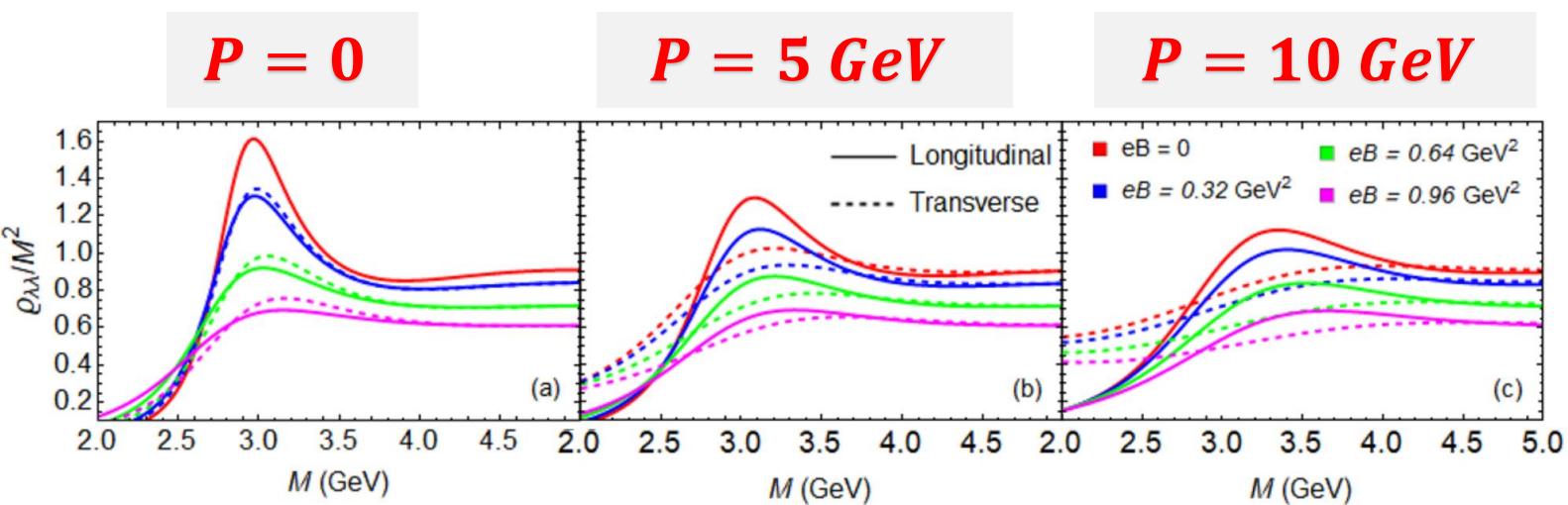
➤ Hawking temperature:

$$T = \frac{1}{4\pi} \left| \frac{4}{\zeta_h} - \frac{2}{3} \frac{e^2 B^2}{1.6^2} \zeta_h^3 \right|.$$

Spectral Function:

$$T = 0.2 \text{ GeV}$$

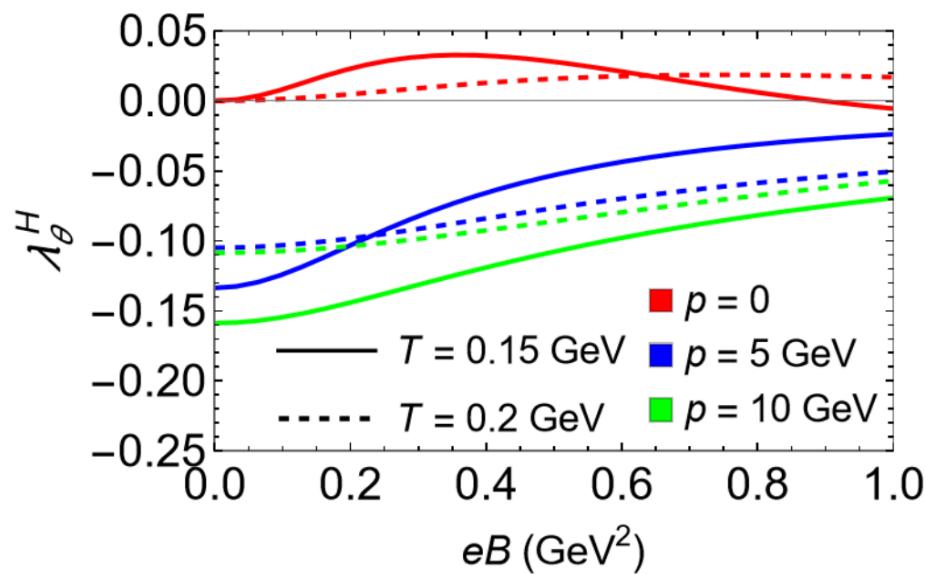
$$\mathbf{p} = (0, 0, p)$$



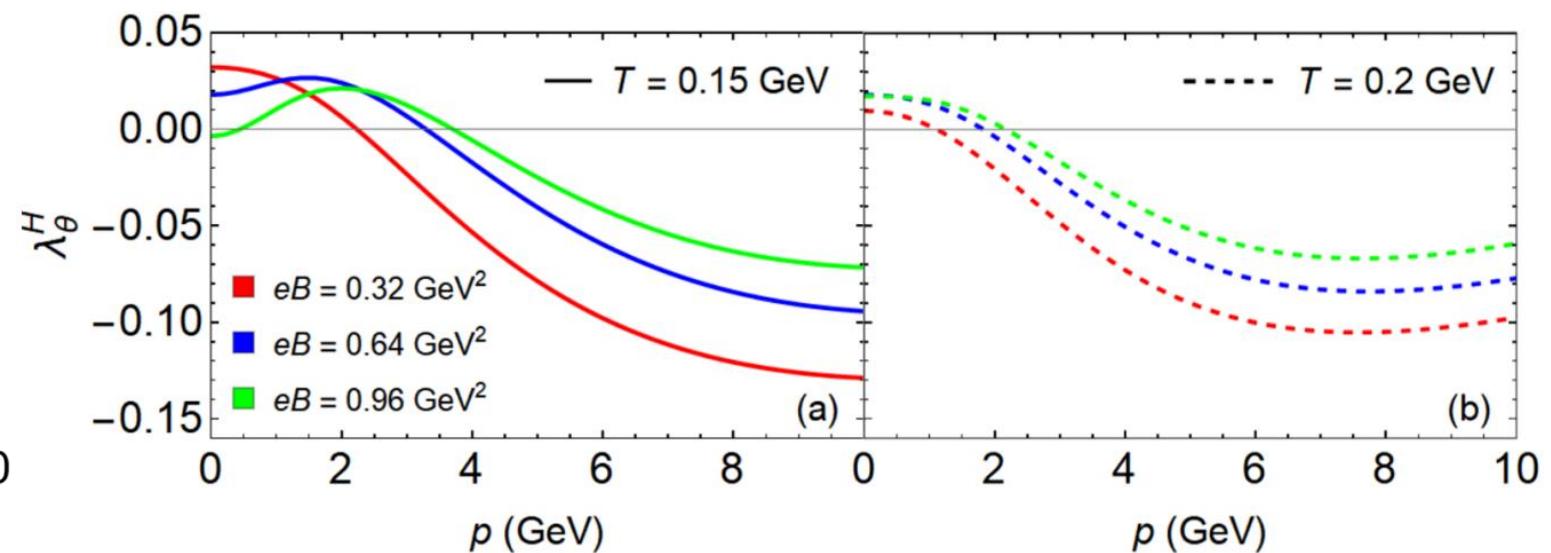
Magnetic field induced spin alignment

➤ Magnetic field parallel to momentum $\mathbf{p} = (0, 0, p)$

Magnetic Field:

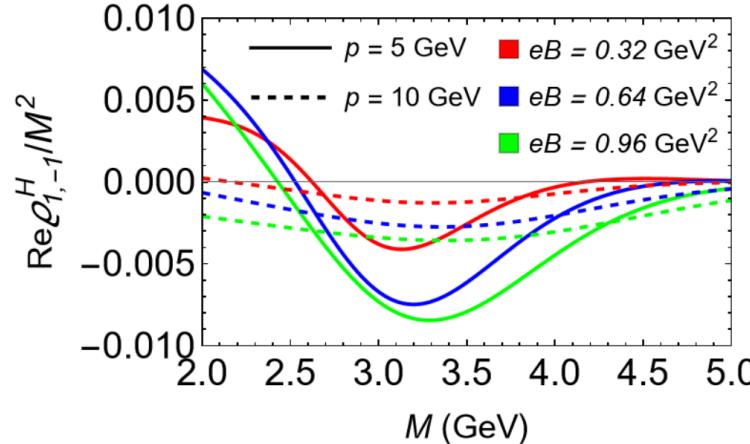


Momentum:

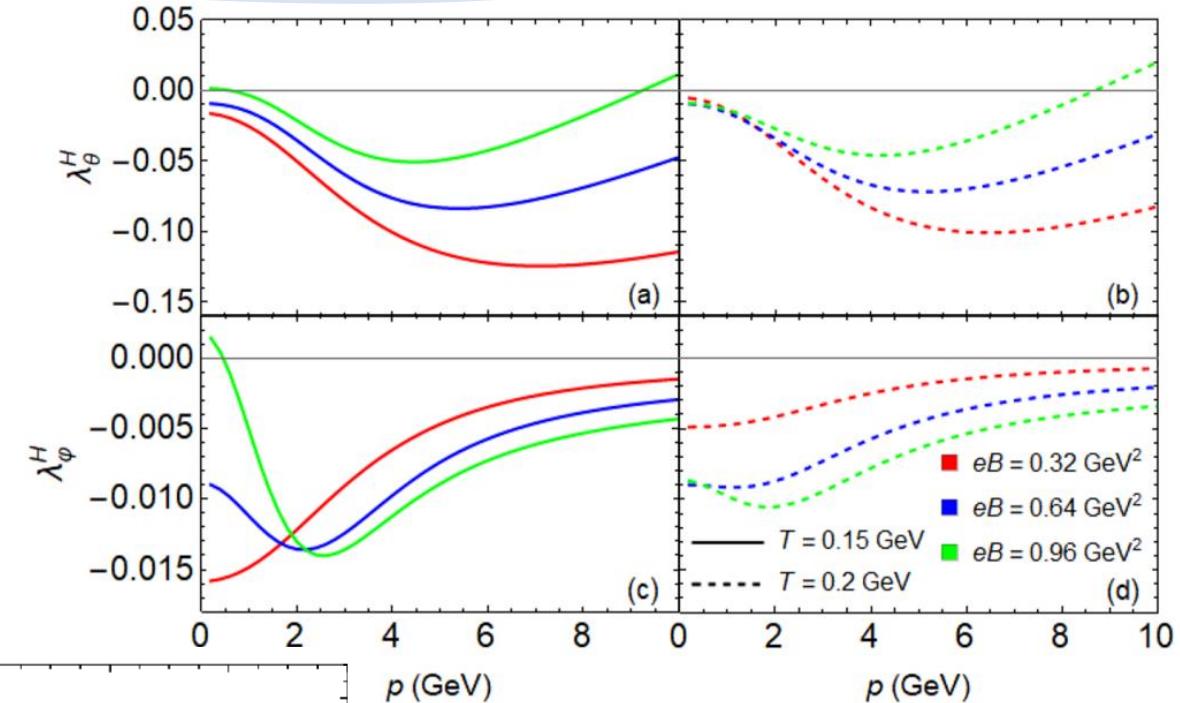
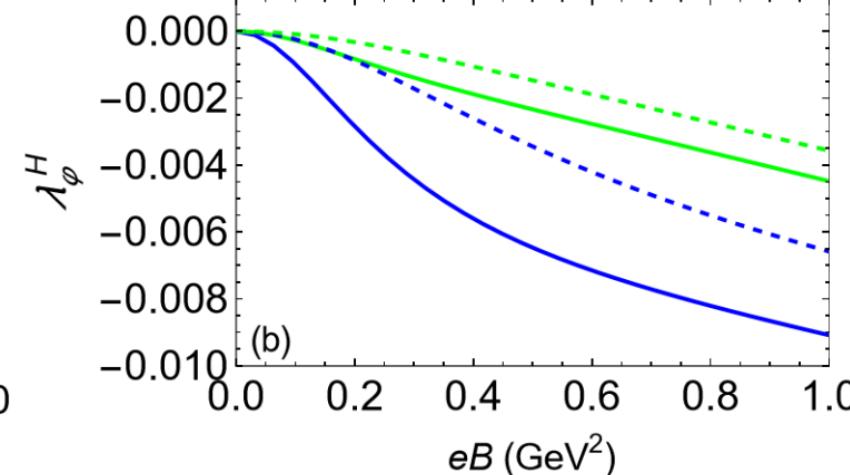
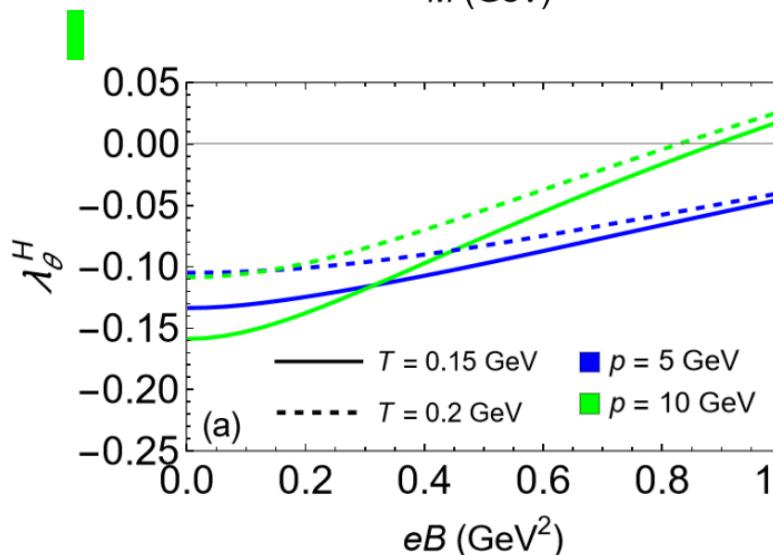


Magnetic field induced spin alignment

➤ Magnetic field perpendicular to momentum



$\mathbf{p} = (p, 0, 0)$



$$\lambda_\theta = \frac{1 - 3\rho_{00}}{1 + \rho_{00}},$$

$$\lambda_\varphi = \frac{2\text{Re}\rho_{1,-1}}{1 + \rho_{00}},$$

Magnetic field induced spin alignment

➤ Application to heavy-ion collisions (the direction of magnetic field along the y-direction)

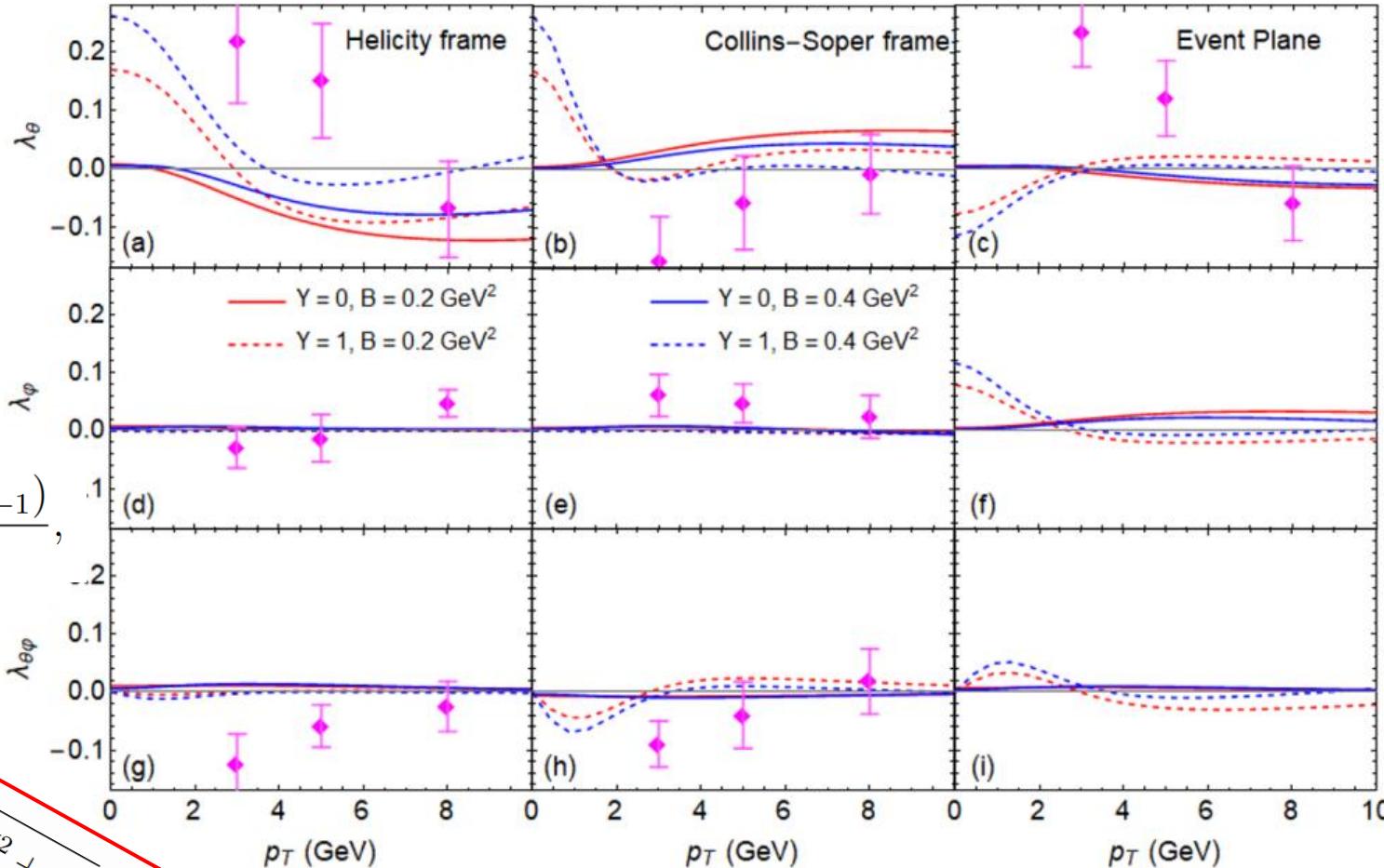
$$T = 0.15 \text{ GeV}$$

$$\lambda_\theta = \frac{1 - 3\rho_{00}}{1 + \rho_{00}},$$

$$\lambda_\varphi = \frac{2\text{Re}\rho_{1,-1}}{1 + \rho_{00}}, \quad \lambda_\varphi$$

$$\lambda_{\theta\varphi} = \frac{\sqrt{2}\text{Re}(\rho_{01} - \rho_{0,-1})}{1 + \rho_{00}},$$

$$\mathbf{p} = \left(p_T \cos \varphi, p_T \sin \varphi, \sqrt{M^2 + p_T^2} \sinh(Y) \right)$$



ALICE Collaboration,
S. Acharya et al., Phys.
Rev. Lett. 131 no.
4,(2023) 042303

ALICE Collaboration,
S. Acharya et al., Phys.
Lett. B 815 (2021)
136146

- Theory
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- **Summary and Outlook**

Summary and Outlook

- The spin alignment can be purely induced by the motion of vector meson relative to the background.
- The holographic prediction shows that J/ψ and ϕ have opposite behaviours. $J/\psi(\phi)$ are preferably to be transversely(longitudinally) polarized.
- The meson's spin alignment is a non-perturbative property in the strongly interacting matter.
- Magnetic field induces $\lambda_\theta^H > 0$ when the meson's momentum p is very small, while $\lambda_\theta^H < 0$ when p is large enough.
- Comparisons with experimental data show qualitative agreement for spin parameters λ_θ and λ_φ in the helicity and Collins-Soper frames.
- We also find significant differences between our results for $\lambda_{\theta\varphi}^H$ and λ_θ^{EP} with experiments .

Thanks!