

Non-extremal Island in de Sitter Gravity

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YITP $\rightarrow$ Vrije Universiteit Brussel

arXiv:2306.07575 with Taishi Kawamoto, Yu-ki Suzuki, Tadashi Takayanagi

arXiv:2407.21617 with Peng-Xiang Hao, Taishi Kawamoto, Tadashi Takayanagi

Gauge Gravity Duality 2024

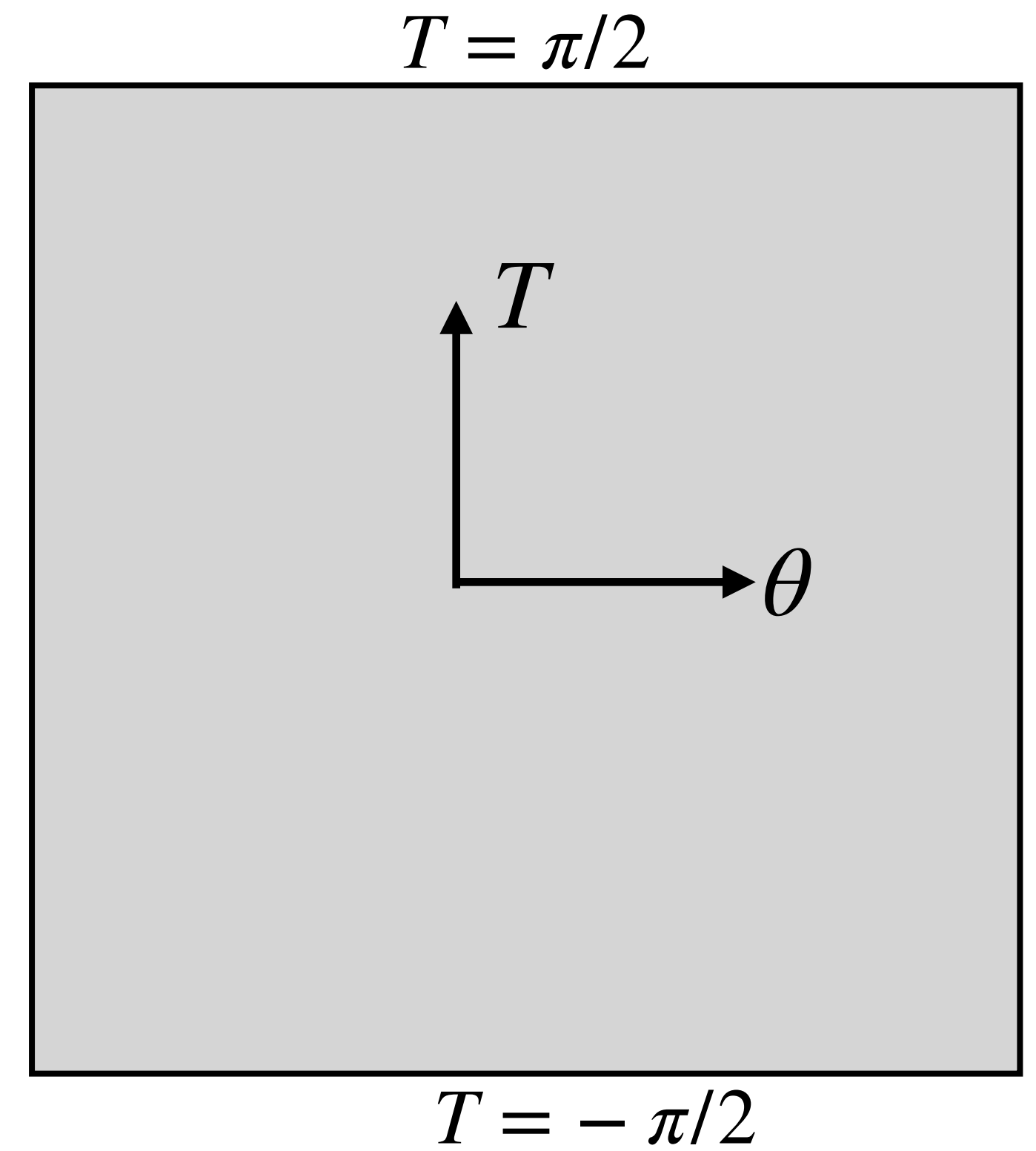
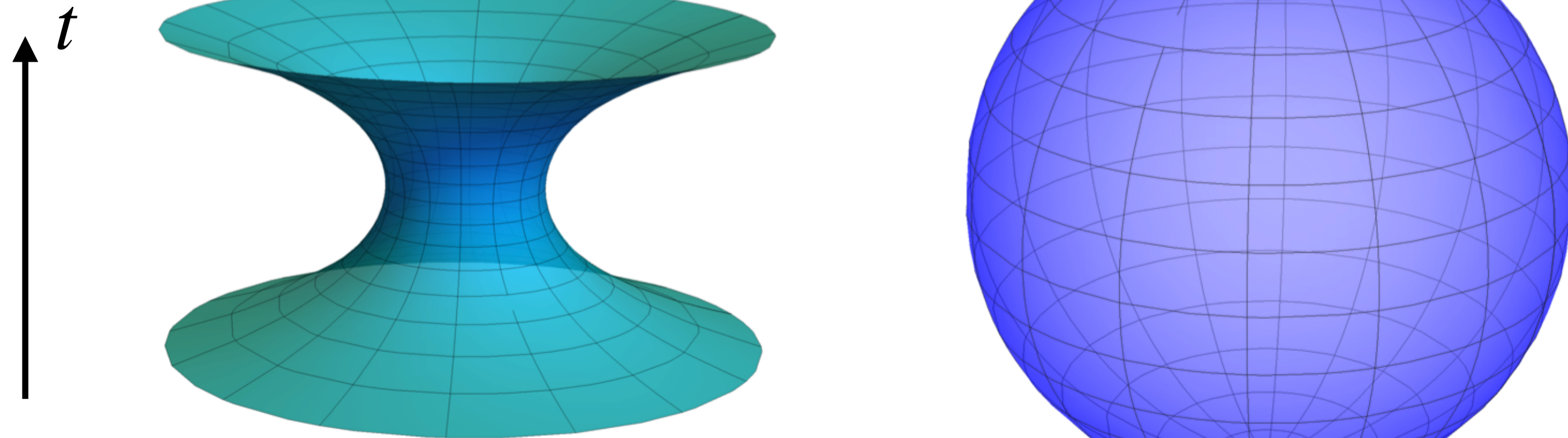
2024-12-03@Tsinghua Sanya International Mathematics Forum

- 01. Introduction
- 02. de Sitter Holography
- 03. Problematic Extremal Islands in dS gravity
- 04. Resolution from Double Holography

01. Introduction- de Sitter Space

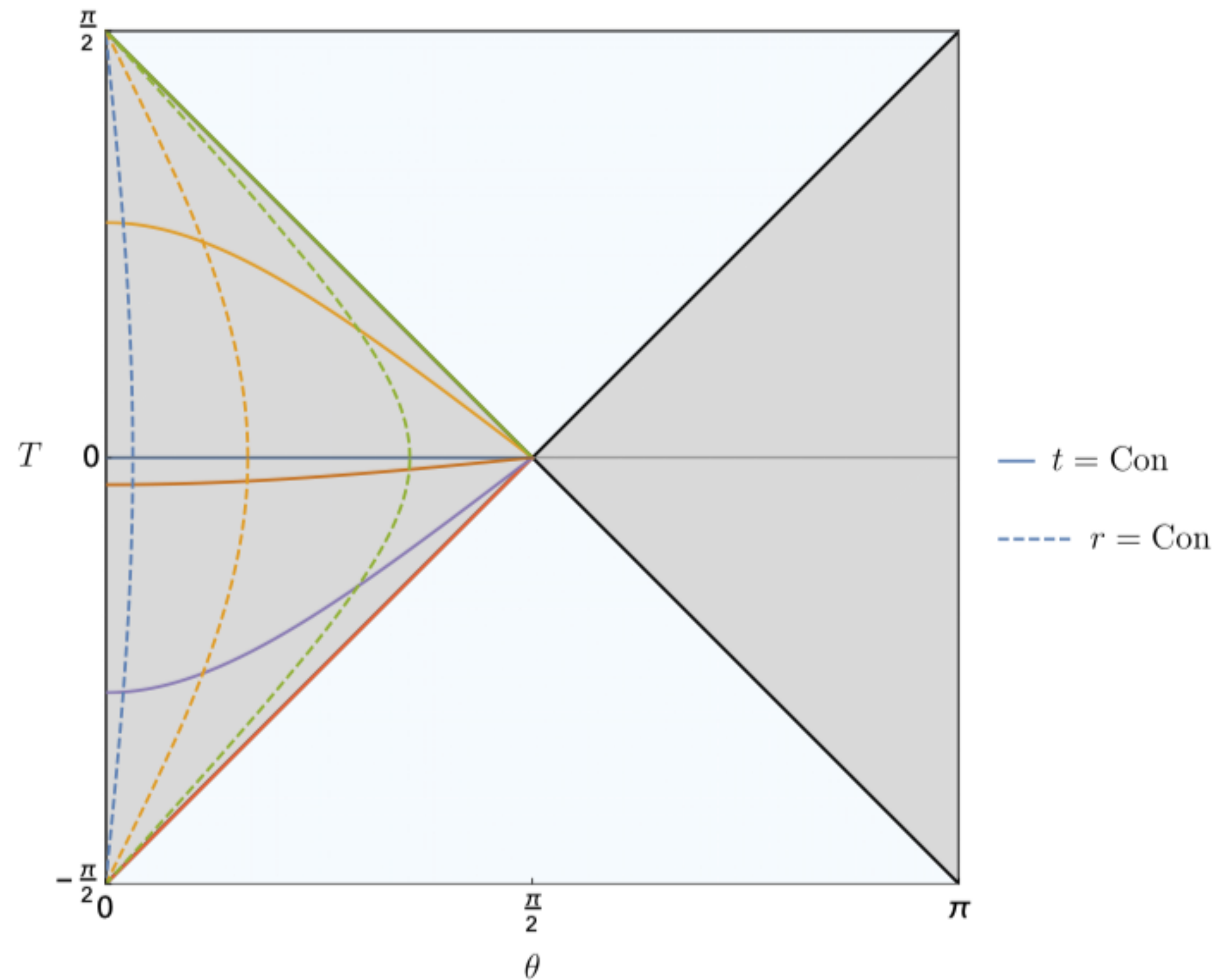
01. Introduction: de Sitter Spacetime

$$ds^2 = L^2 \left(-dt^2 + \cosh^2 t (d\theta^2 + \sin^2 \theta d\Omega_{d-1}^2) \right)$$



$$ds^2 = \frac{L^2}{\cos^2 T} \left(-dT^2 + d\theta^2 + \sin^2 \theta d\Omega_{d-1}^2 \right)$$

Static Patch



$$ds^2 = - \left(1 - \frac{r^2}{L^2} \right) dt^2 + \frac{dr^2}{1 - \frac{r^2}{L^2}} + r^2 d\Omega_{d-1}^2.$$

Hawking temperature of cosmological horizon

$$T = \frac{1}{2\pi L}$$

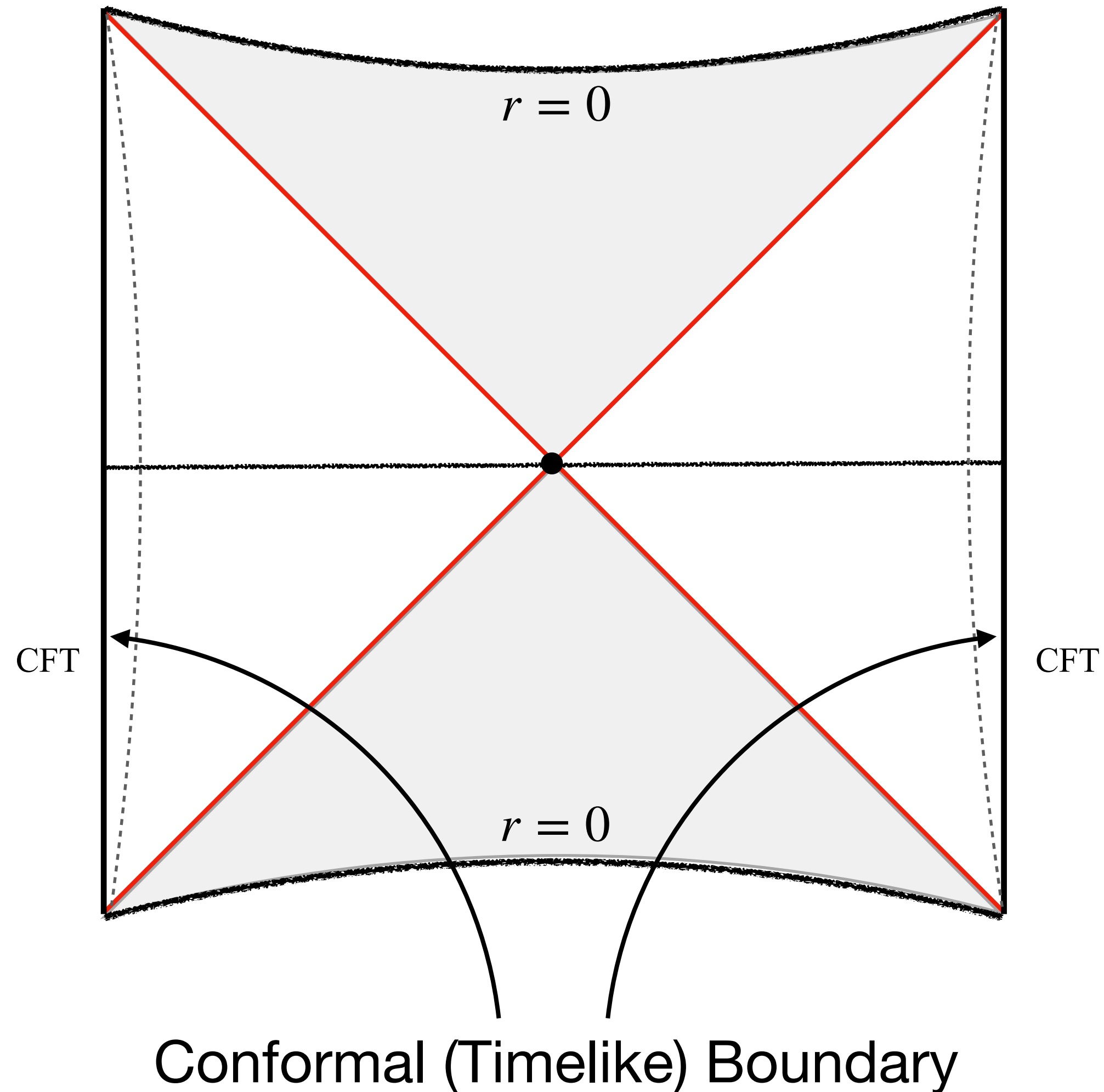
Entropy $\rightarrow$ Bekenstein-Hawking formula

$$S_{\text{dS}} = \frac{\text{Area}}{4G_{\text{N}}} = \frac{L^{d-1} \Omega_{d-1}}{4G_{\text{N}}} \equiv N$$

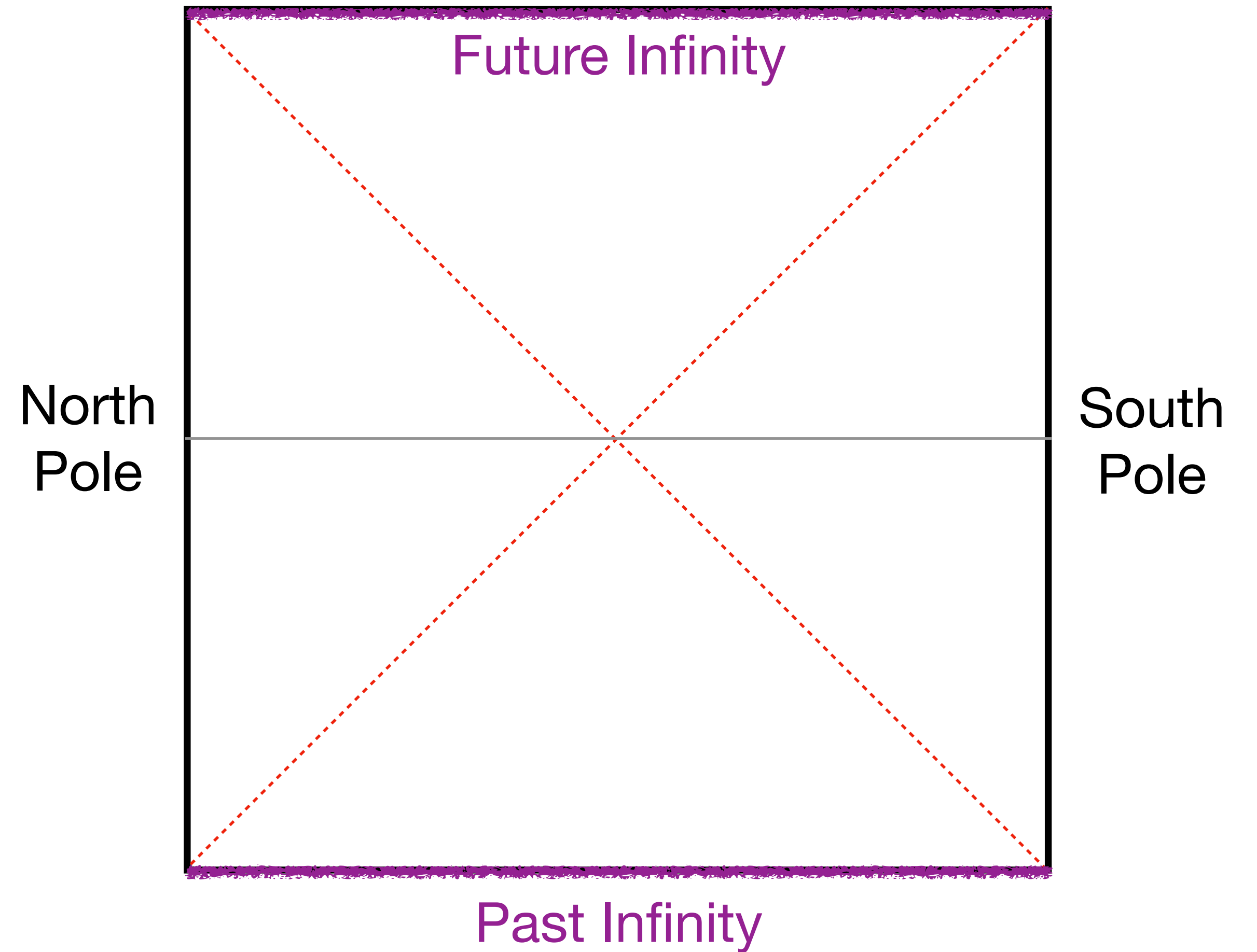
Maximal entropy? Finite Dimension of Hilbert Space?

01. Introduction: de Sitter Spacetime

AdS Black Hole



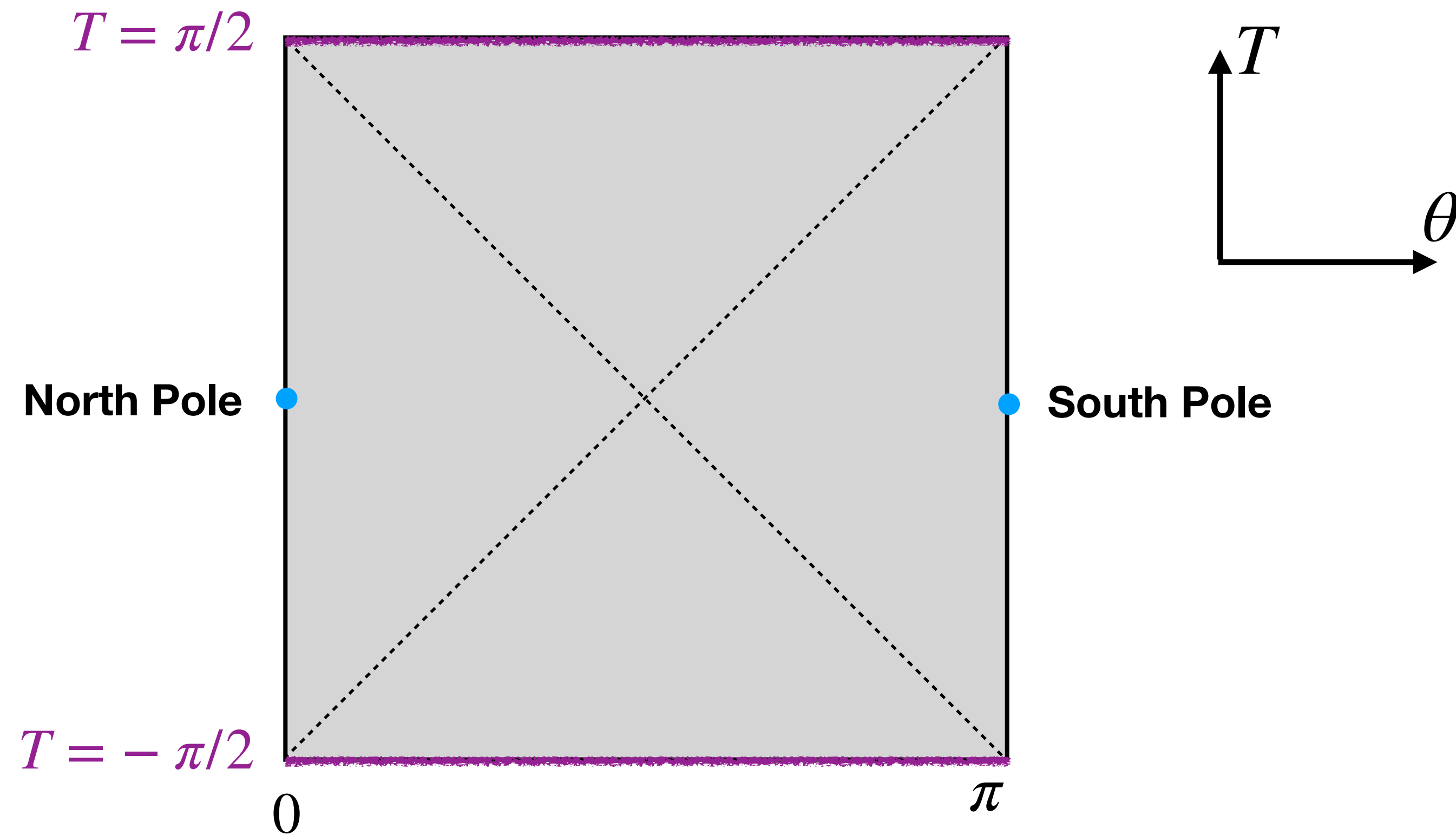
dS_{d+1}



02. de Sitter Holography

dS/CFT Correspondence

[Strominger 2001]



“Analytical continuation” of AdS/CFT correspondence

dS/CFT correspondence: gravity in asymptotically dS_{d+1} spacetime
a d -dimensional CFT living on the **spacelike boundary** at timelike future infinity of dS.

Static Patch Holography

[Parikh, Verlinde 2004]

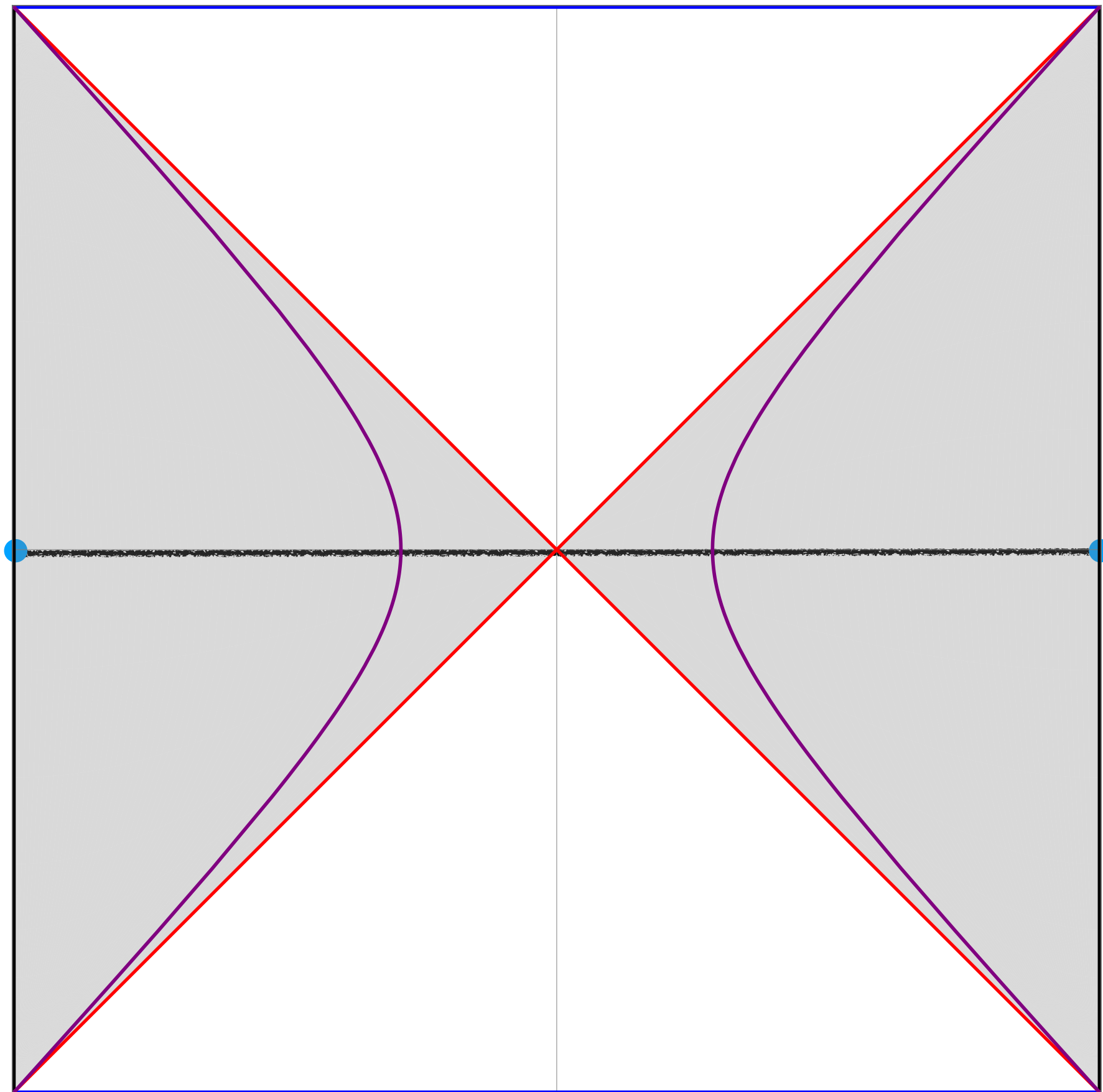
[Alishahiha, Karch, Silverstein, Tong 2004]

[Banks, Fiol and Morisse 2006]

[Susskind 2021]

[Shaghoulian 2021]

[Narovlansky, Verlinde 2023]

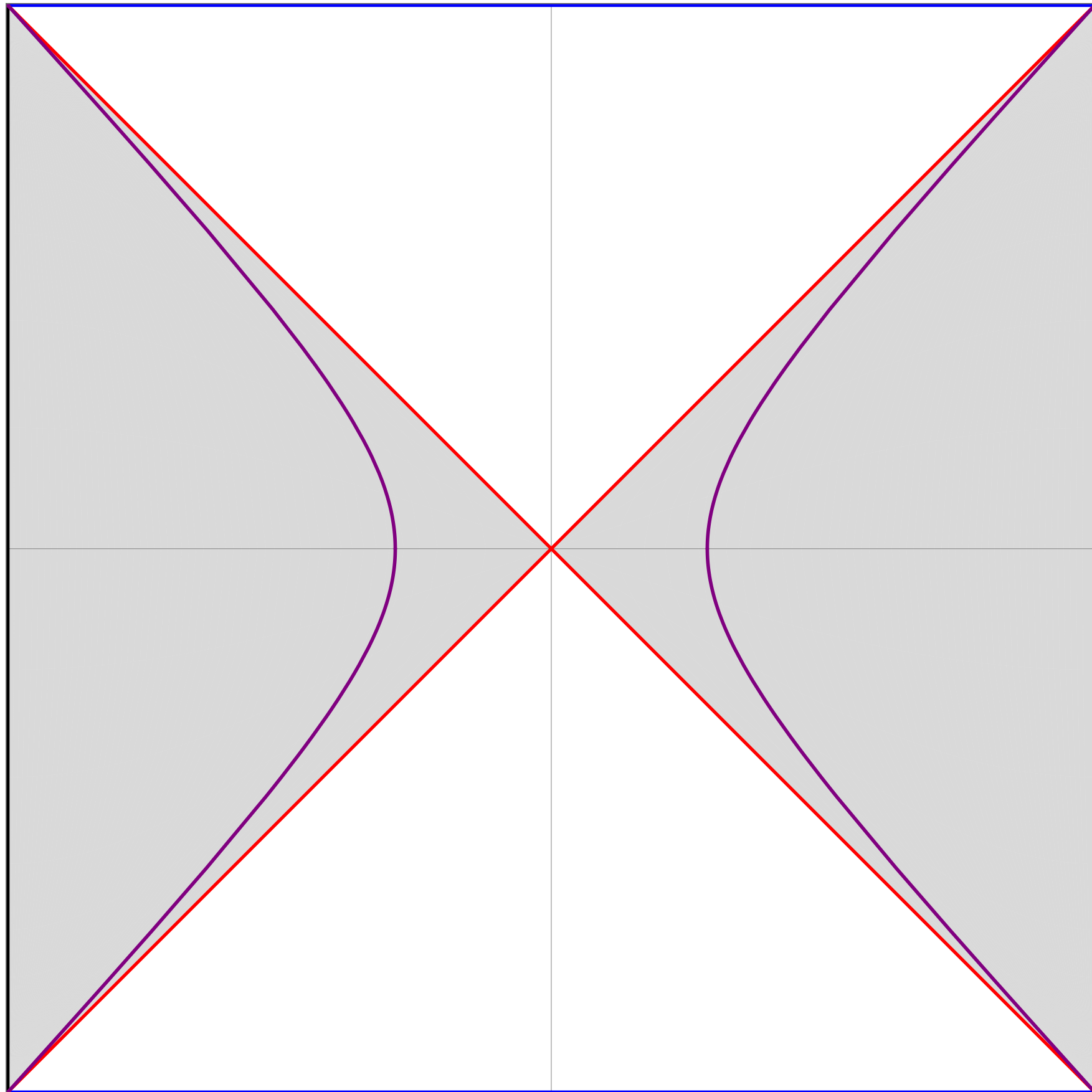


holographic degrees of freedom of de Sitter space



Stretched Horizon

Static Patch Holography



De Sitter version of (H)RT formula

[Susskind 2021]

$$S_{\text{EE}} = \min \max \left(\frac{\text{Area}}{4G_N} \right)$$

?

entanglement entropy between the pole/antipode static patch

$$S_{\text{EE}} = \frac{\text{Horizon Area}}{4G_N} = S_{\text{dS}}$$

Different results: Replica trick + Holographic two-point function
[To appear, with Yu-ki Suzuki]

A Half de Sitter Holography

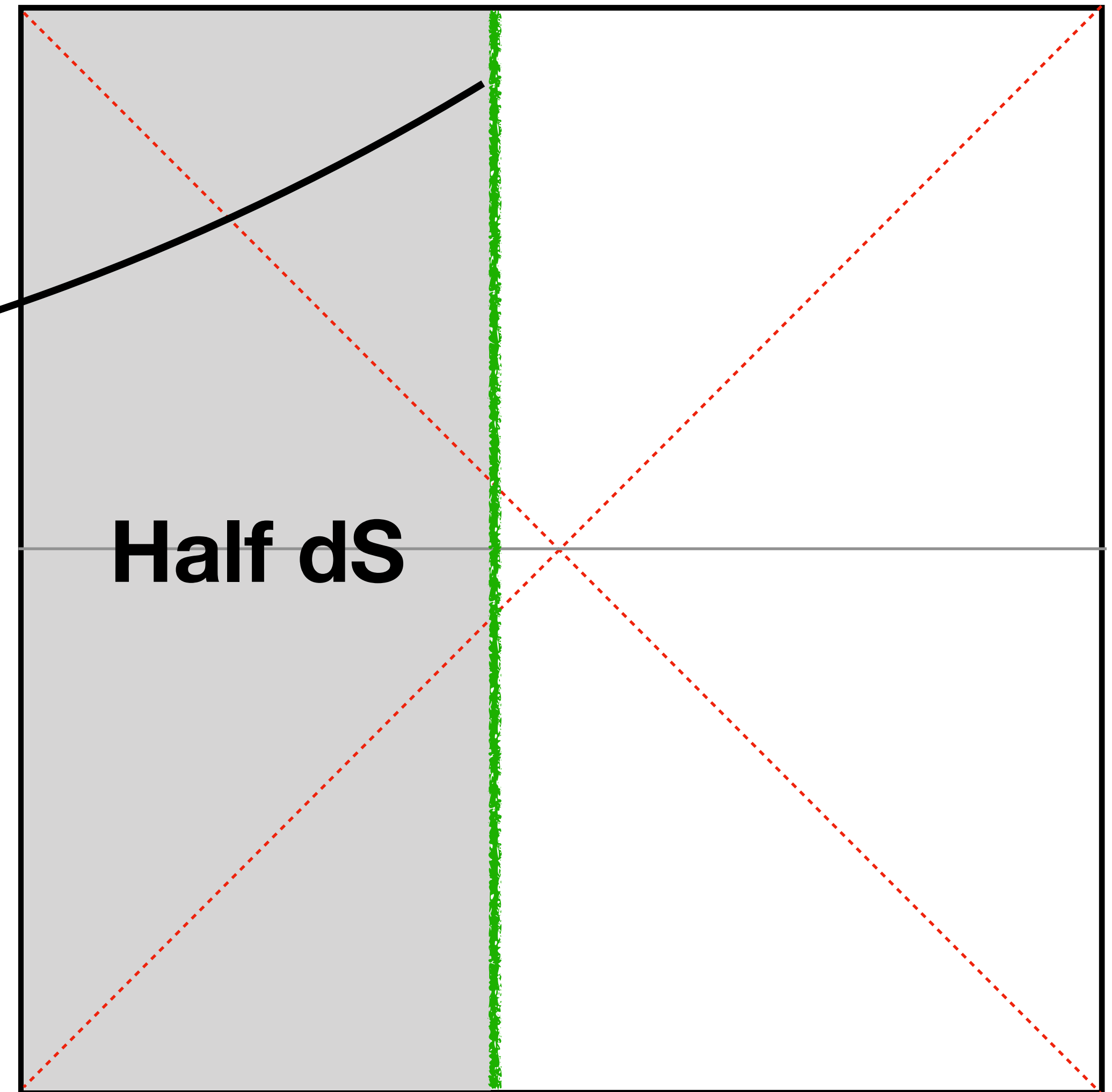
(H)RT formula in a half dS spacetime

Violation of (strong) subadditivity!

Non-local field theory on the timelike boundary?

[arXiv:2306.07575](https://arxiv.org/abs/2306.07575)

with Taishi Kawamoto, Yu-ki Suzuki, Tadashi Takayanagi



Goal: To Probe a Half de Sitter Space by a Bath System

[arXiv:2407.21617](https://arxiv.org/abs/2407.21617)

With Peng-Xiang Hao, Taishi Kawamoto, Tadashi Takayanagi

Background- Island Formula

A correct formula for
the fine grained entropy of Hawking radiation

Island Formula:
$$S_{\text{EE}}(\hat{\rho}_A) = \text{Min}_I \left[\text{Ext}_I \left(\frac{\text{Area}[\partial\Sigma_I]}{4G_N} + S_{\text{Bulk}}(\Sigma_A \cup \Sigma_{\text{Island}}) \right) \right]$$



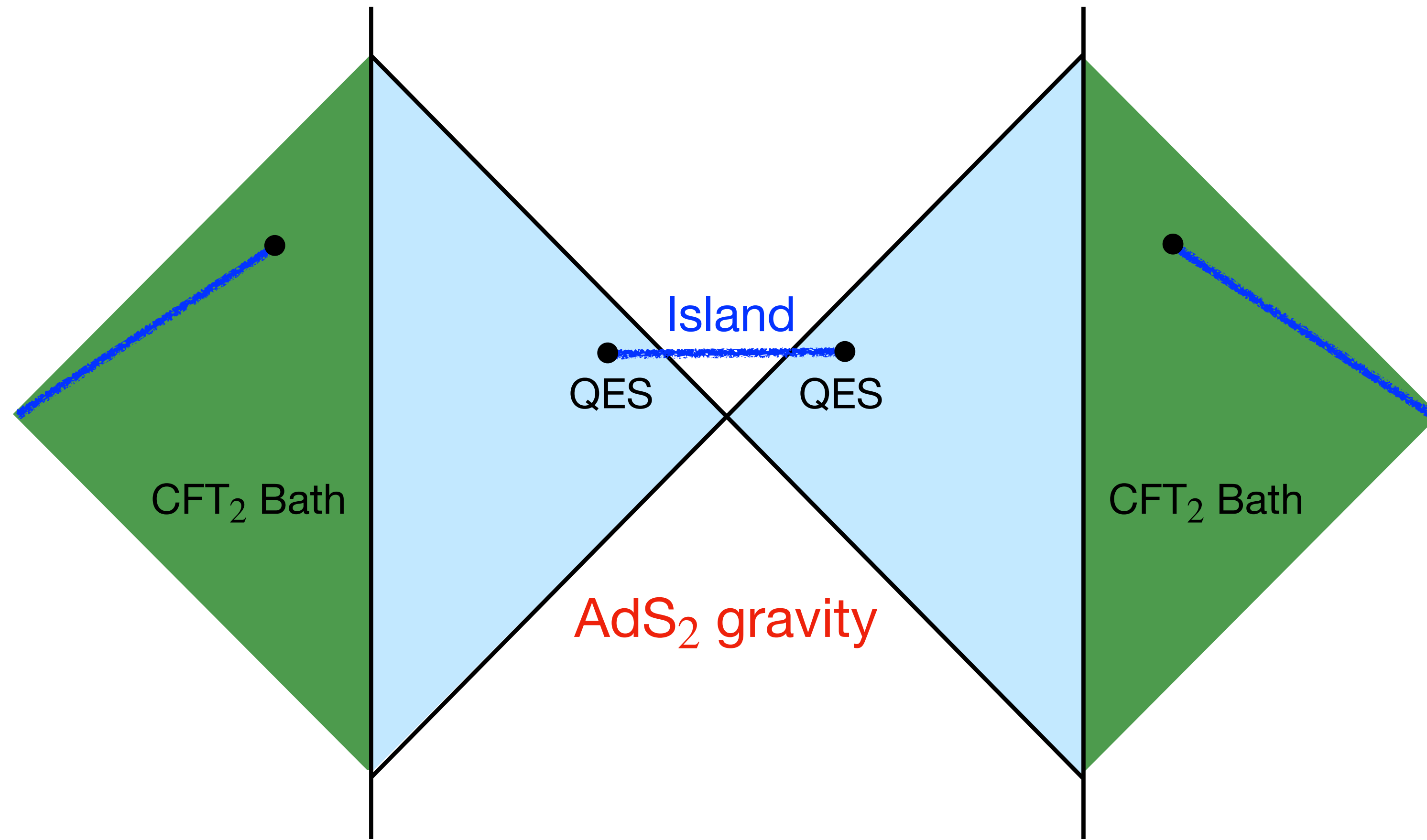
Disconnected
Region

Background- Island Formula

The simplest model with Islands

An information paradox for the eternal black hole [arXiv:1910.11077]

Island Formula:
$$S_{\text{EE}}(\hat{\rho}_A) = \text{Min}_I \left[\text{Ext}_I \left(\frac{\text{Area}[\partial \Sigma_I]}{4G_N} + S_{\text{Bulk}}(\Sigma_A \cup \Sigma_{\text{Island}}) \right) \right]$$



Question: Does Island formula apply to dS gravity?

03: Extremal Islands in dS Gravity

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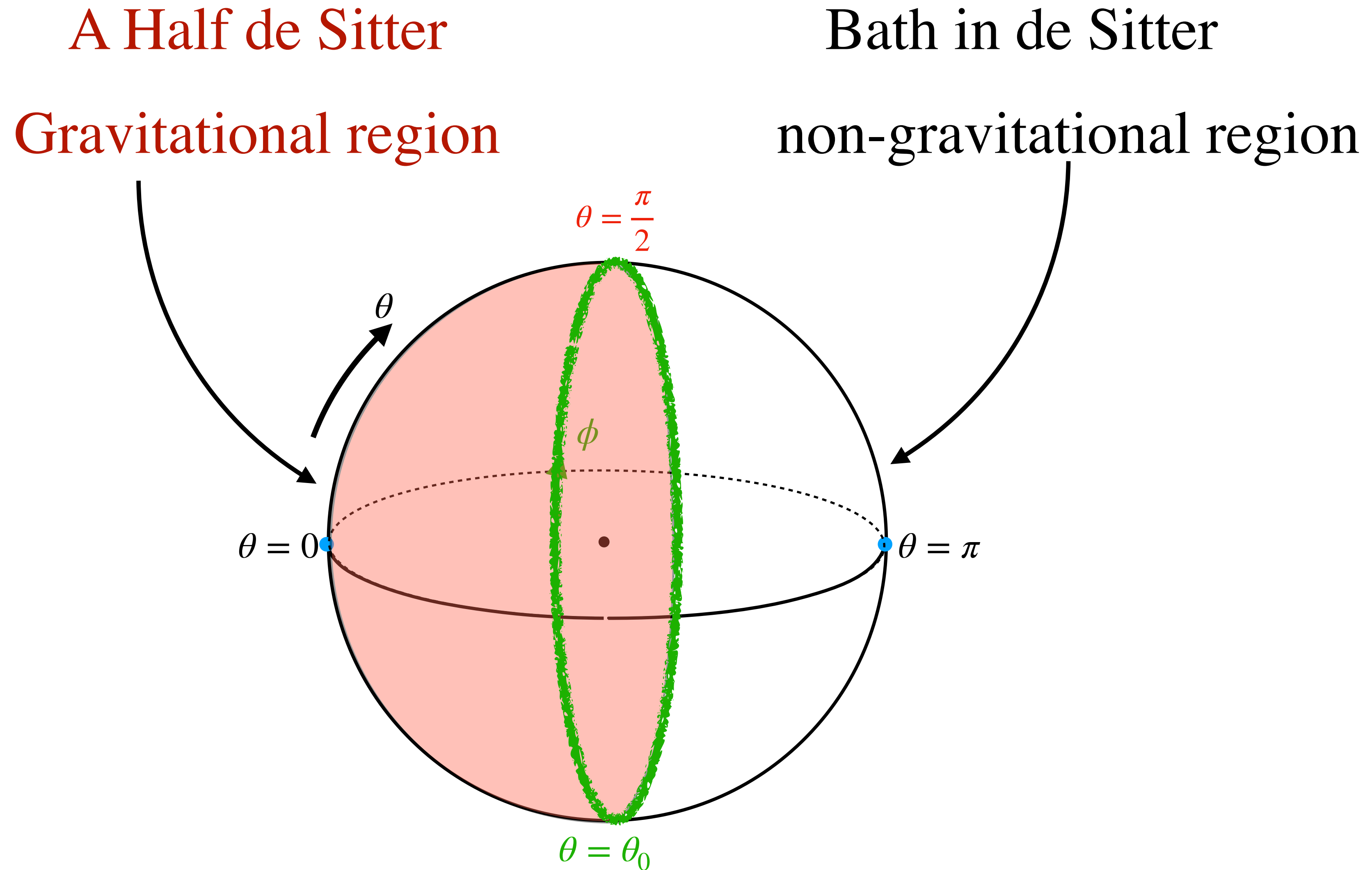
A Half de Sitter

Gravitational region

Bath in de Sitter

non-gravitational region

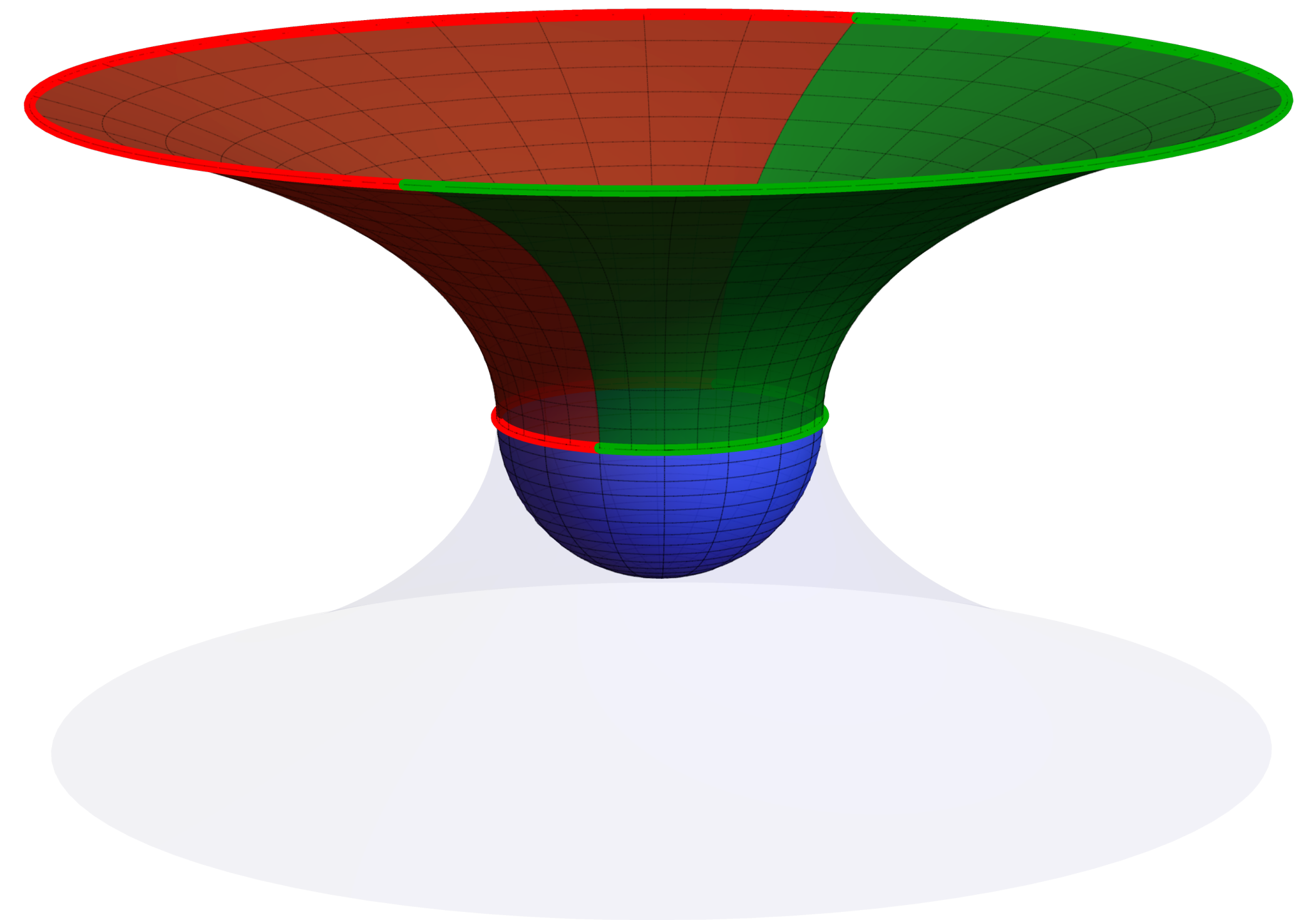
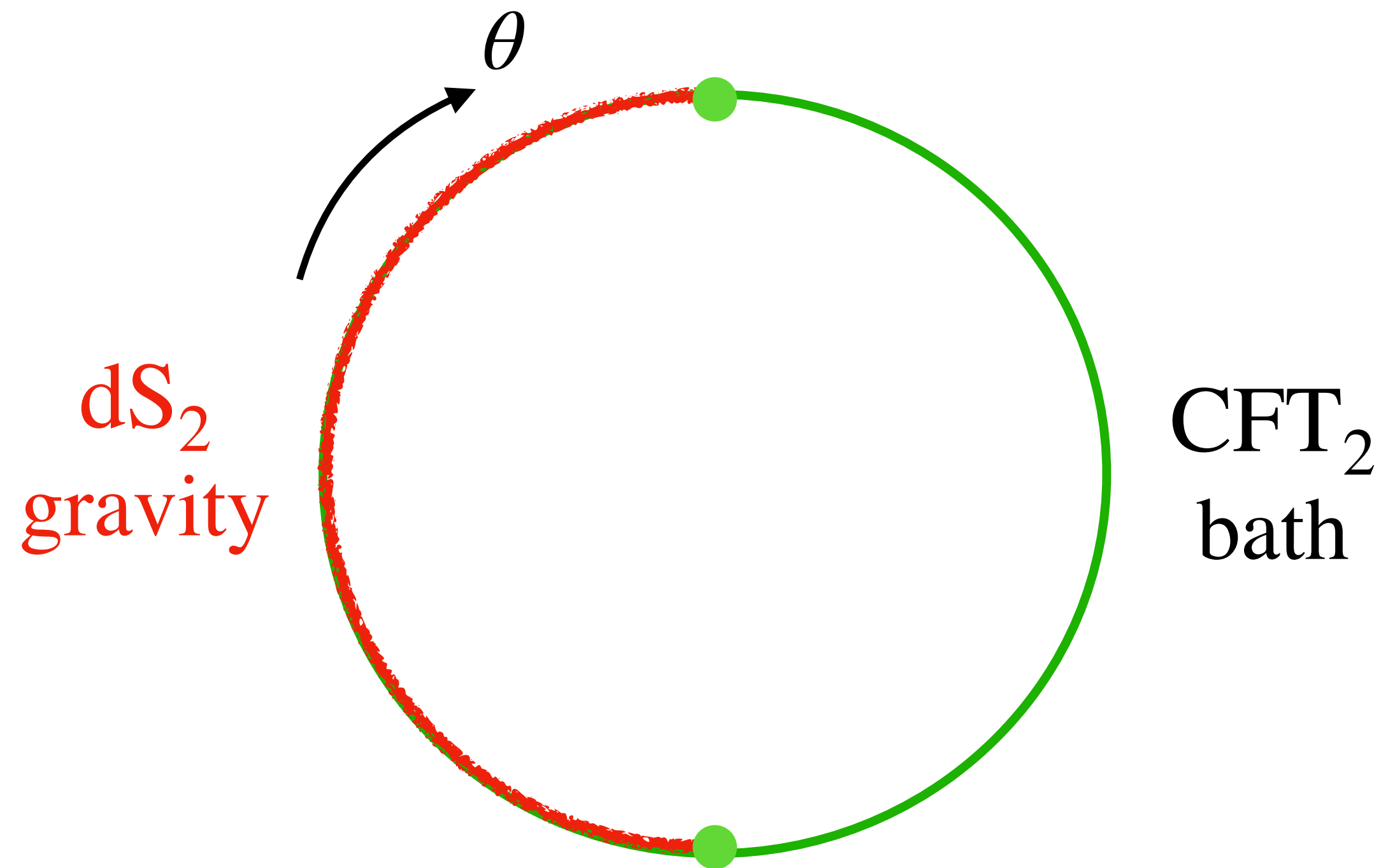
03: Extremal Islands in dS gravity



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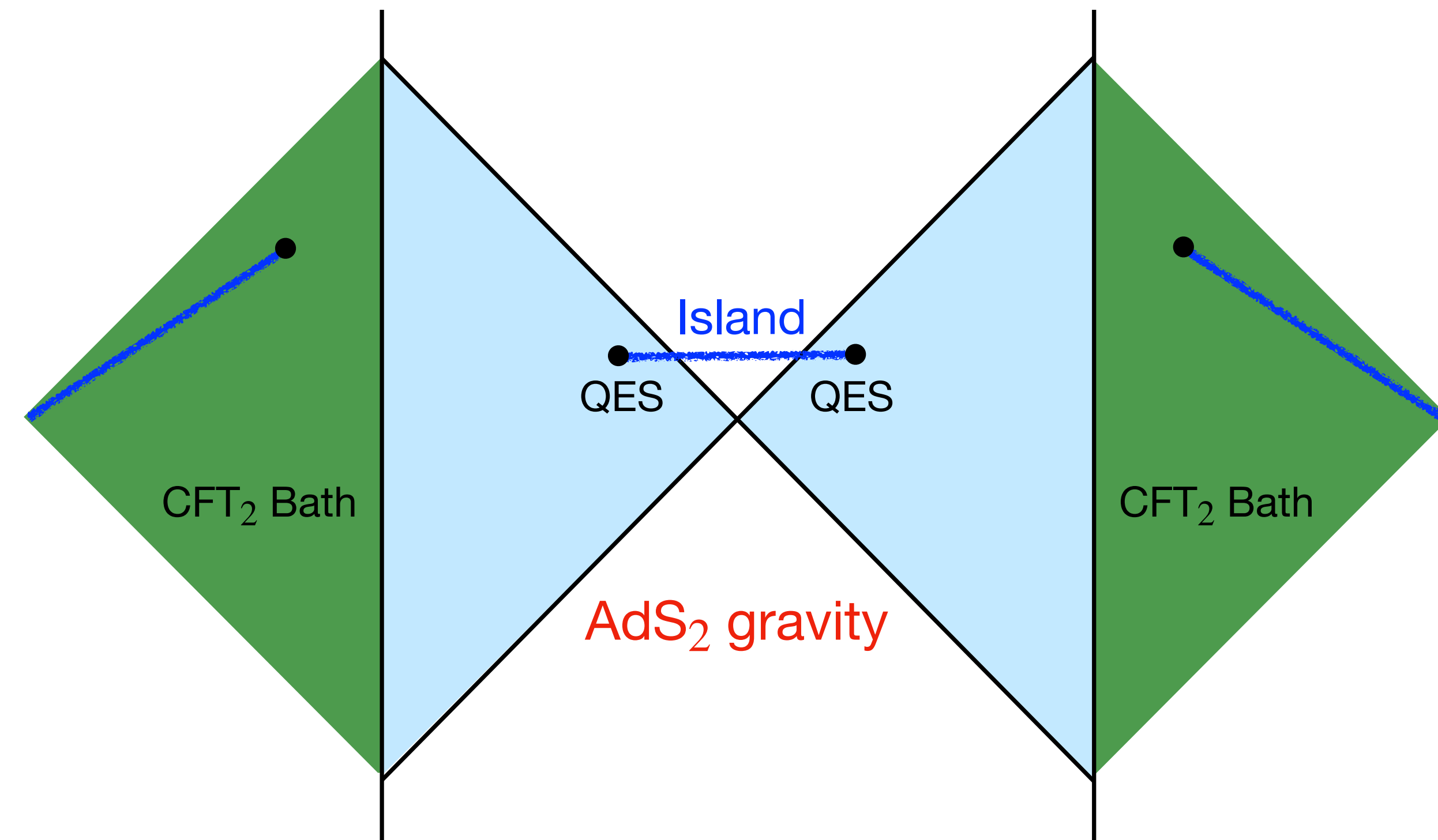
Simplest Model: **Half dS₂ gravity** + Half dS₂ bath

$$ds^2 = L^2 (-dt^2 + \cosh^2 t d\theta^2) = \frac{L^2}{\cos^2 T} (-dT^2 + d\theta^2)$$

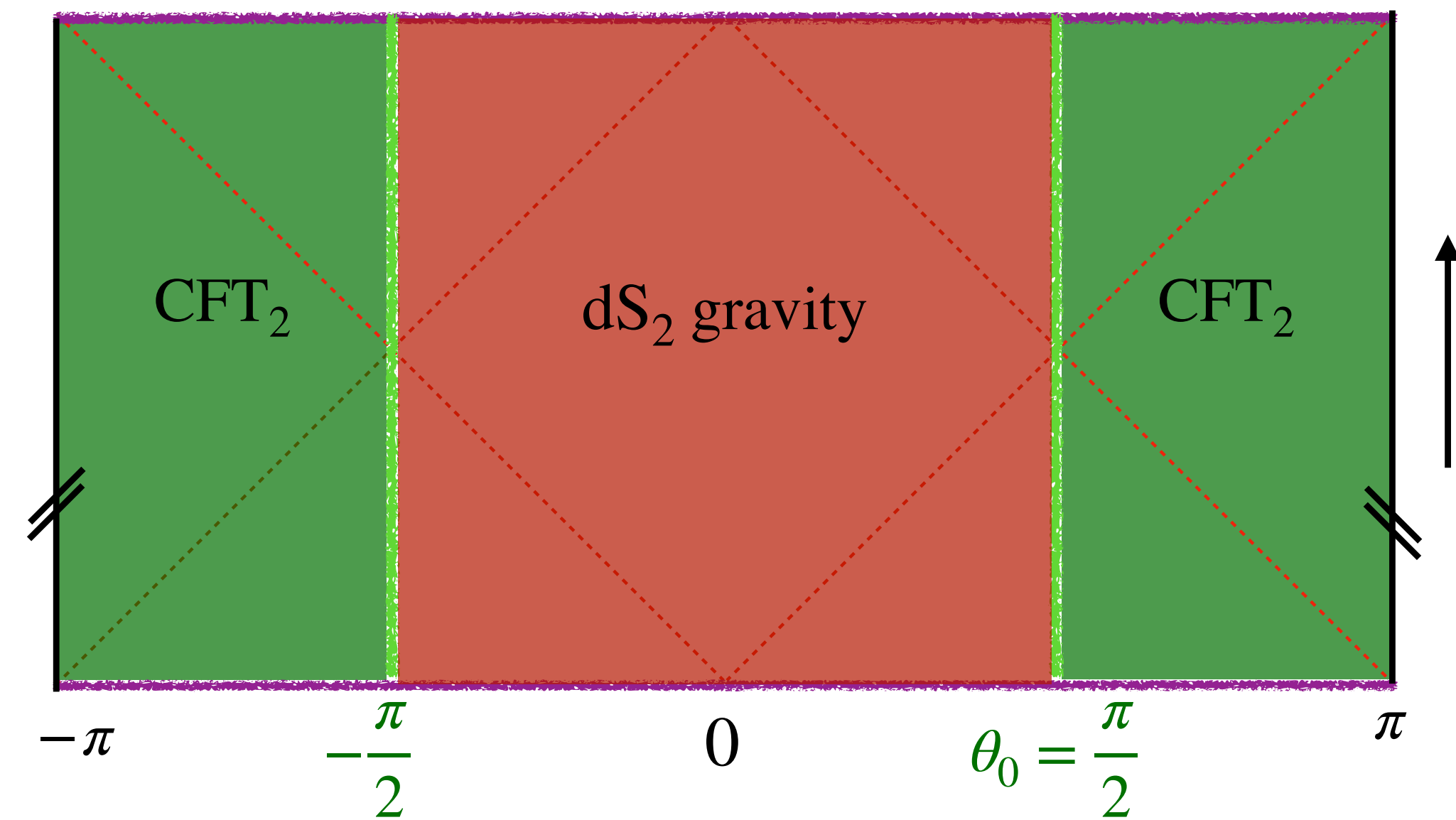


03: Extremal Islands in dS gravity

AdS+Bath



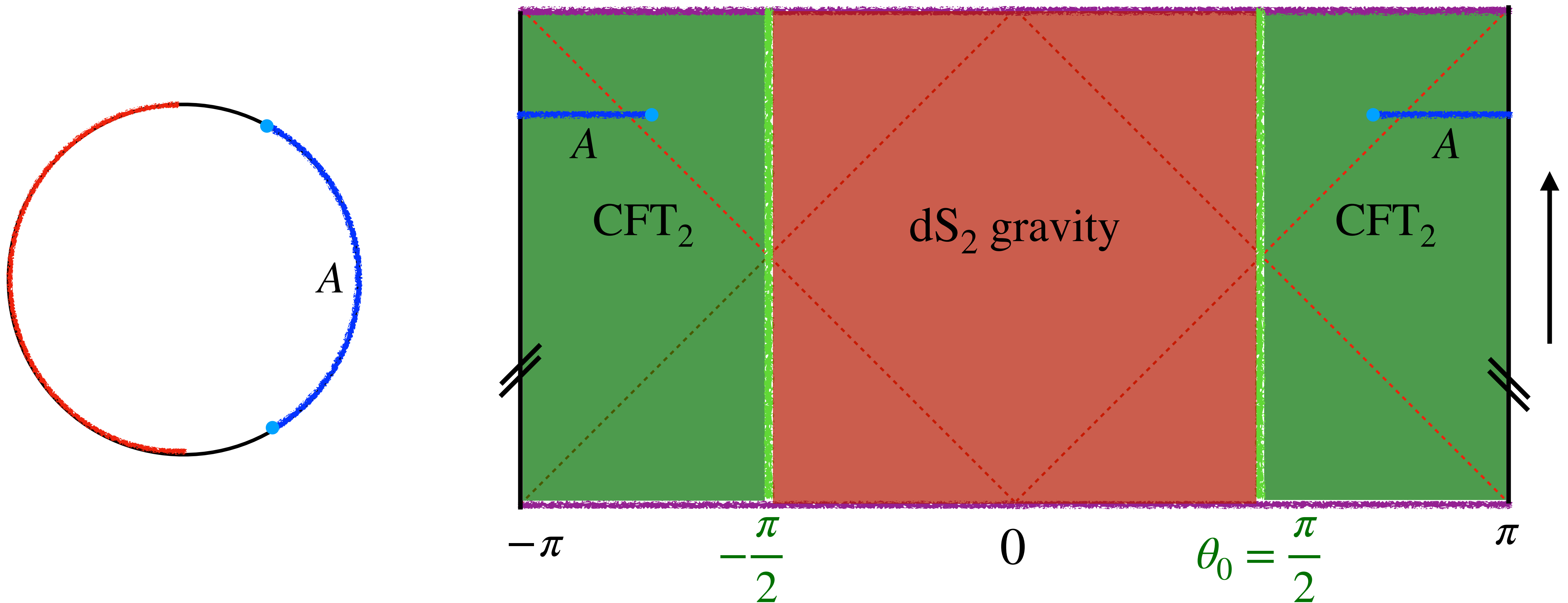
a half dS+Bath



03: Extremal Islands in dS gravity

Simplest Model: **Half dS₂ gravity** + Half dS₂ bath (CFT₂)

Entanglement Entropy of a subsystem A:
$$S_{\text{EE}}(A) = \lim_{n \rightarrow 1} \frac{1}{1-n} \log \left\langle \sigma_n(T_1, \theta_1,) \tilde{\sigma}_n(T_2, \theta_2) \right\rangle$$



03: Extremal Islands in dS gravity

Simplest Model: **Half dS₂ gravity** + Half dS₂ bath (CFT₂)

Entanglement Entropy of a subsystem A: $S_{\text{EE}}(A) = \lim_{n \rightarrow 1} \frac{1}{1-n} \log \left\langle \sigma_n(T_A, \theta_1) \tilde{\sigma}_n(T_A, \theta_2) \right\rangle$

$$ds^2 = \Omega^{-2} dz d\bar{z} = \frac{4}{(1 + z\bar{z})^2} dz d\bar{z},$$

conformal dimension

$$\Delta_n = \frac{c}{12} \left(n - \frac{1}{n} \right)$$

$$S_{\text{Bulk}} = \frac{c}{6} \log \left(\frac{(z_1 - z_2)(\bar{z}_1 - \bar{z}_2)}{\epsilon^2 \Omega_1 \Omega_2} \right) = \frac{c}{6} \log \left(\frac{2 \left(\cos(T_1 - T_2) - \cos(\theta_1 - \theta_2) \right)}{\epsilon^2 \cos T_1 \cos T_2} \right)$$

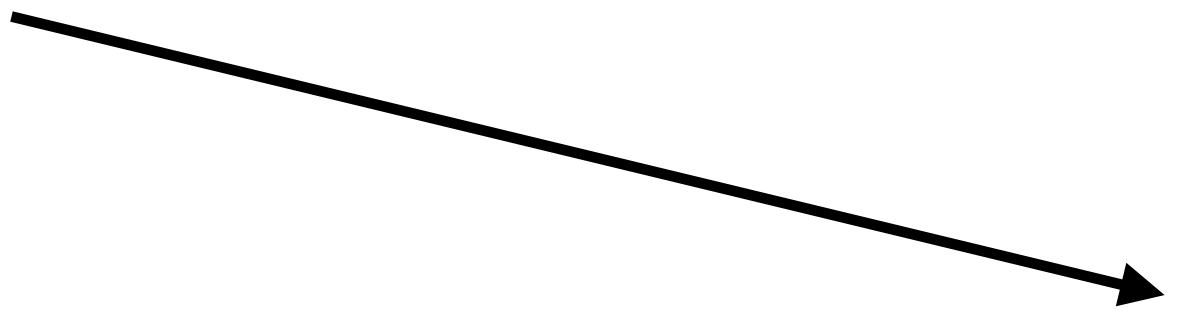
Connected Phase
(No Island Phase)

$$S_{\text{EE}}^{\text{con}}(t_0) = \frac{c}{3} \log \cosh t_0 + \frac{c}{3} \log \left(\frac{2}{\epsilon} \left| \sin \left(\frac{\theta_1 - \theta_2}{2} \right) \right| \right).$$

03: Extremal Islands in dS gravity

$$\text{Island Formula: } S_{\text{EE}}(\hat{\rho}_A) = \text{Min}_I \left[\text{Ext}_I \left(\frac{\text{Area}[\partial \Sigma_I]}{4G_{\text{N}}} + S_{\text{Bulk}}(\Sigma_A \cup \Sigma_{\text{Island}}) \right) \right]$$

Area term



Reduction from AdS3	—>	01. Einstein gravity higher curvature gravity f(R) gravity scalar-tensor theory	constant ~ log g
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Reduction from dS3	—>	02. dS JT gravity	$\Phi = \frac{\cos \theta}{\cos T} > 0$
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Reduction from Extremal dS4 BH (Nariai solution)		03. Topological term + dS JT gravity	
		$S = \frac{\Phi_0}{16\pi G} \int d^2x \sqrt{-g} R + \frac{1}{16\pi G} \int d^2x \sqrt{-g} \Phi (R - 2),$	$\Phi_0 + \Phi > 0$

03: Extremal Islands in dS gravity

Simplest Model: **Half dS₂ gravity** + Half dS₂ bath (CFT₂)

Island Formula: $S_{\text{EE}}(\hat{\rho}_A) = \text{Min}_I \left[\text{Ext}_I \left(\frac{\Phi_0}{4G_N} + S_{\text{Bulk}}(\Sigma_A \cup \Sigma_{\text{Island}}) \right) \right] = \text{constant}$

$$\partial_{\theta_I} S_{\text{gen}} = 0 \longrightarrow \cosh t_I \cosh t_A \sin(\theta_I - \theta_A) = 0,$$

$$\partial_{t_I} S_{\text{gen}} = 0 \longrightarrow \cosh t_I \sinh t_A - \sinh t_I \cosh t_A \cos(\theta_I - \theta_A) = 0$$

Unique Solution: $\theta_I = -\pi + \theta_A, \quad t_I = -t_A$

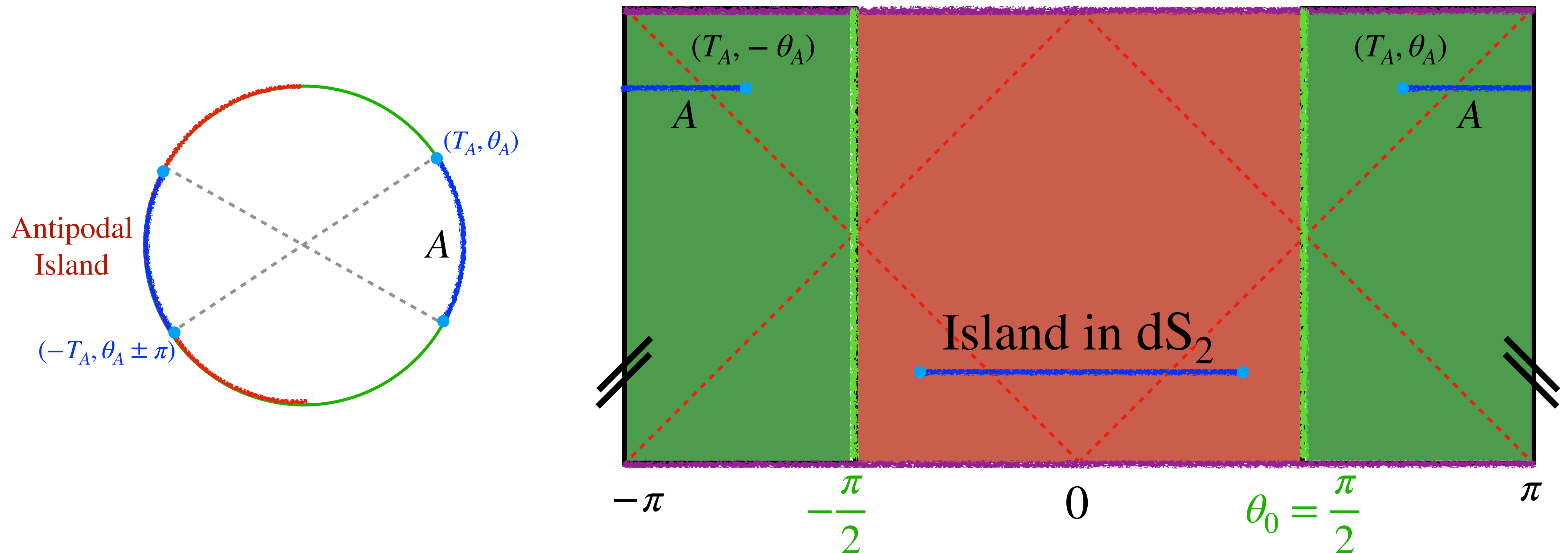
Antipodal Island

03: Extremal Islands in dS gravity

Simplest Model: **Half dS₂ gravity** + Half dS₂ bath (CFT₂)

Island Formula: $S_{\text{EE}}(\hat{\rho}_A) = \text{Min}_I \left[\text{Ext}_I \left(\frac{\Phi_0}{4G_N} + S_{\text{Bulk}}(\Sigma_A \cup \Sigma_{\text{Island}}) \right) \right] = \text{constant}$

Unique Solution $\rightarrow$ Antipodal Island:

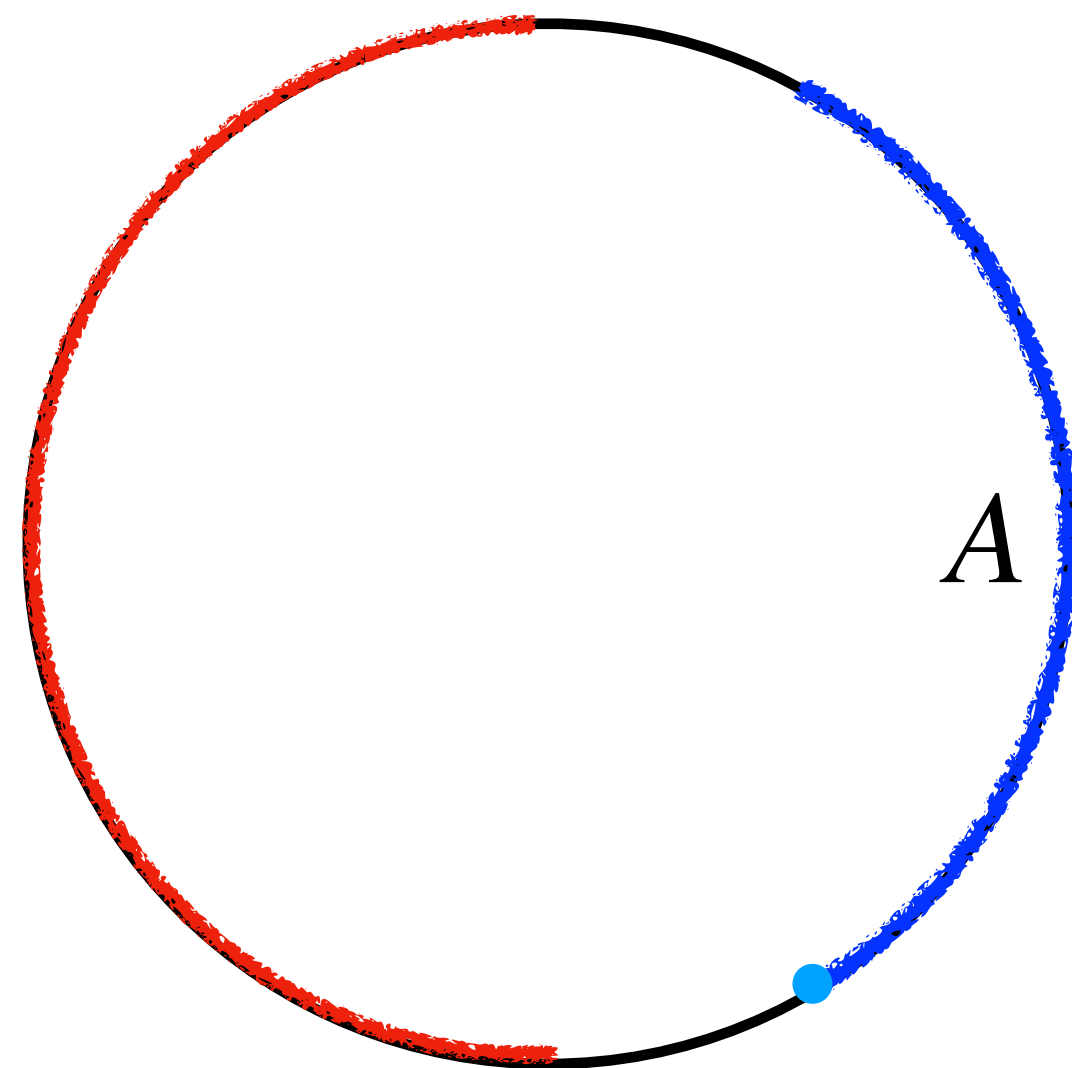


03: Extremal Islands in dS gravity

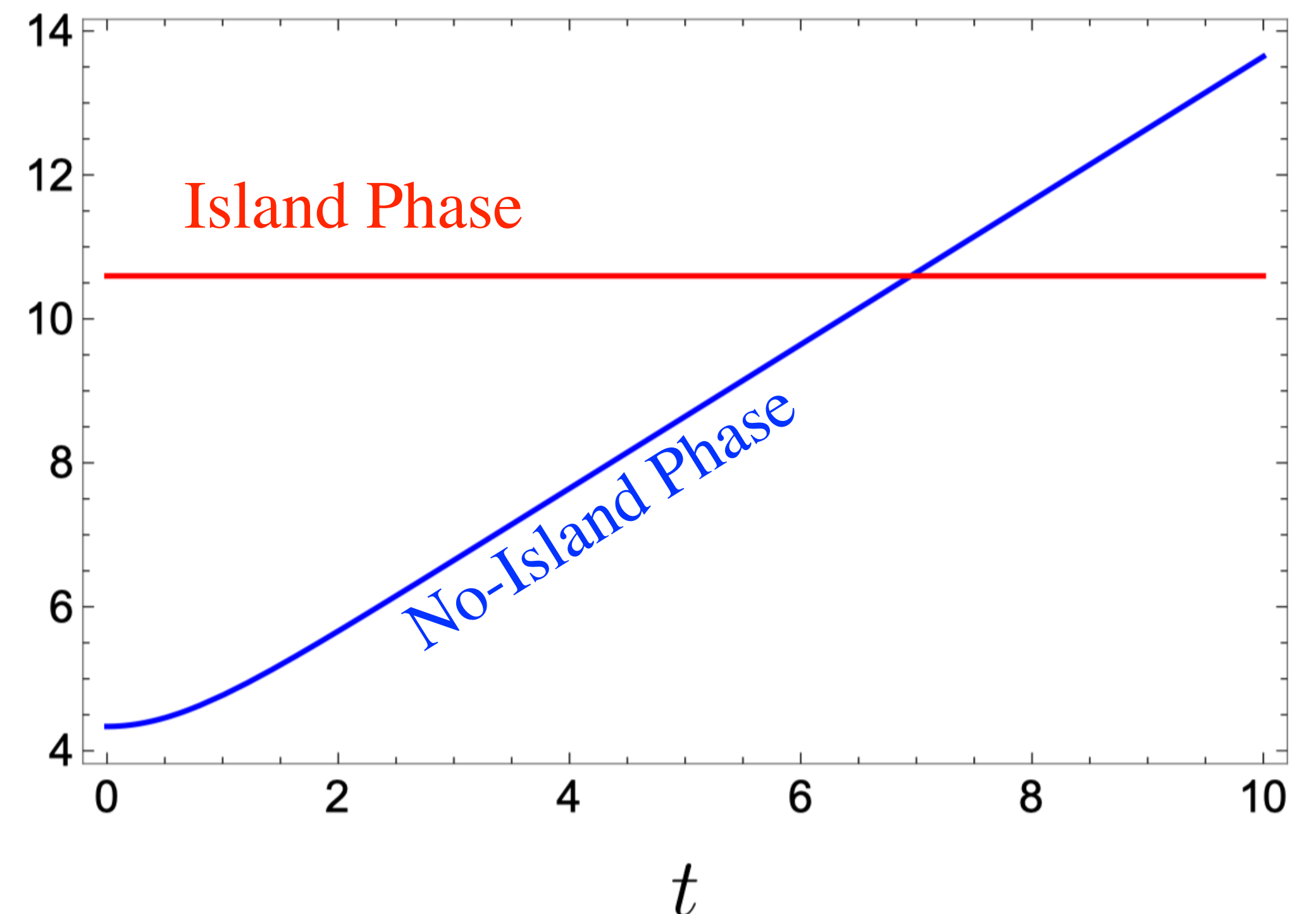
Simplest Model: **Half dS₂ gravity** + Half dS₂ bath (CFT₂)

$$S_{\text{EE}}(\hat{\rho}_A) \Big|_{\text{island}} = S_0 + \frac{2c}{3} \log \frac{2}{\epsilon} = \tilde{S}_0 + \frac{c}{3} \log \left(\frac{4}{\epsilon} \right)$$

dS
Gravity



$$S_{\text{EE}}(A) = ?$$



03: **Problematic** Extremal Islands in dS gravity

However, there are **many reasons** that “Island”/QES in dS is **not physical!**
violate SSA, violate entanglement wedge nesting...

It is a saddle point but it cannot dominate!

$$\frac{\partial^2 S_{\text{bulk}}}{\partial t_I \partial t_I} > 0 \quad \text{Local minimum in the Lorentzian timelike direction}$$

$$\frac{\partial^2 S_{\text{bulk}}}{\partial \theta_I \partial \theta_I} < 0 \quad \text{Local maximum in the spatial direction}$$

(Same conclusion for any **positive dilaton Φ**)

The quantum extremal surface in dS is a minimax surface!

?

04: Resolution from Double Holography

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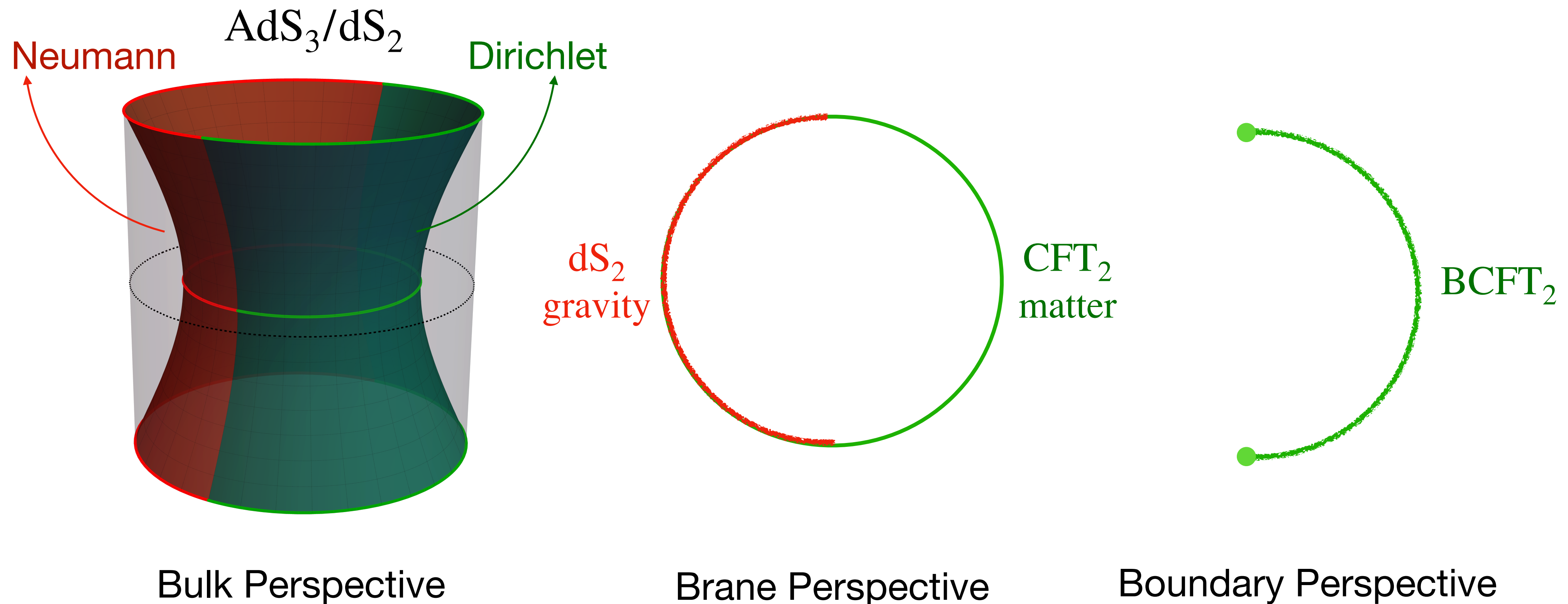
Simplest Model: **Half dS₂ gravity** + Half dS₂ bath

$$ds^2 = L^2 (-dt^2 + \cosh^2 t d\theta^2) = \frac{L^2}{\cos^2 T} (-dT^2 + d\theta^2)$$

AdS₃/dS₂ Slicing: $ds^2 = d\eta^2 + \sinh^2 \eta (-dt^2 + \cosh^2 t d\theta^2)$

tension of EOW brane

$$T = \frac{\coth \eta_b}{L_{\text{AdS}}}$$

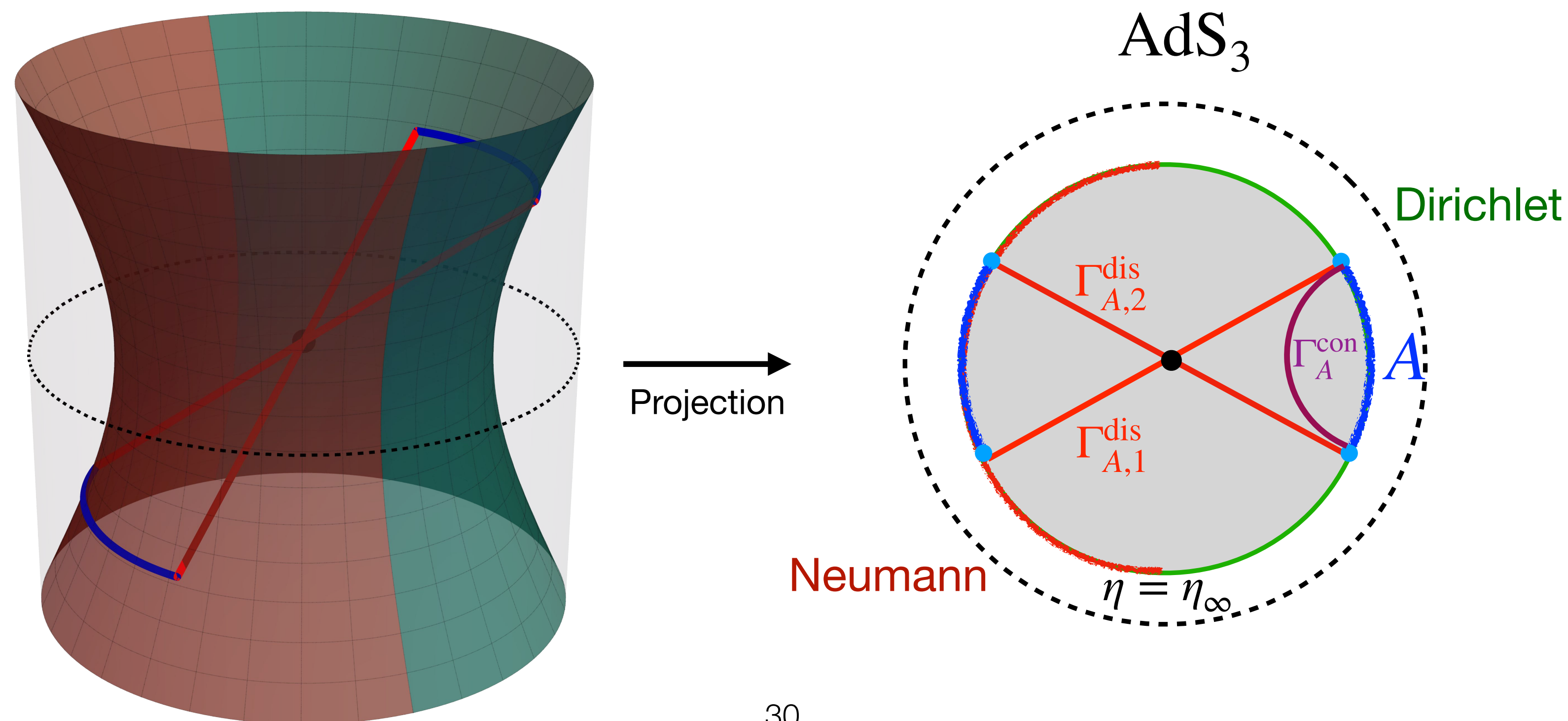


04: Resolution from Double Holography

RT formula provides the same results as the Island formula:

$$S_A = \min \{ S_A^{\text{con}}, S_A^{\text{dis}} \} = \min \left\{ \frac{\text{Area}(\Gamma_A^{\text{con}})}{4G_N}, \frac{\text{Area}(\Gamma_A^{\text{dis}})}{4G_N} \right\},$$

Naive holographic entanglement entropy is problematic!



04: Resolution from Double Holography

Entanglement Entropy of a subsystem A: $S_{\text{EE}}(A) = \lim_{n \rightarrow 1} \frac{1}{1-n} \log \left\langle \sigma_n(T_A, \theta_1) \tilde{\sigma}_n(T_A, \theta_2) \right\rangle$

$$\langle \sigma_n(t_{A_1}, \theta_{A_1}) \bar{\sigma}_n(t_{A_2}, \theta_{A_2}) \rangle_{\text{BCFT}} = \max \begin{cases} \langle \sigma_n(t_{A_1}, \theta_{A_1}) \bar{\sigma}_n(t_{A_2}, \theta_{A_2}) \rangle \\ \langle \sigma_n(t_{A_1}, \theta_{A_1}) \rangle \langle \bar{\sigma}_n(t_{A_2}, \theta_{A_2}) \rangle \end{cases}$$

One-Point Function $\left\langle \sigma_n(T_0, \theta_1) \right\rangle?$

Advantage from the double holography:
The one-point function can be derived by using the bulk-to-boundary propagator:

$$\langle \mathcal{O}(\zeta) \rangle = \int_{\text{EOW}} d^d \hat{x} \sqrt{h} K_{\text{Bb}}^\Delta(X, Y(\hat{x}))$$

$$\left(\square_g + m^2 \right) \underline{G_{\text{BB}}^\Delta(X, X')} = \frac{\delta(X - X')}{\sqrt{g}}.$$

Bulk-to-Bulk propagator

$$G_{\text{BB}}^\Delta(X, X') = G_{\text{BB}}^\Delta(\xi) = \frac{C_\Delta}{2^\Delta(2\Delta - d)} \xi^\Delta \cdot {}_2F_1 \left(\frac{\Delta}{2}, \frac{\Delta + 1}{2}; \Delta - \frac{d}{2} + 1; \xi^2 \right)$$

$$\xi = \frac{1}{\cosh d_{\text{geodesic}}(X, X')}$$

04: Resolution from Double Holography

From (Euclidean) AdS3: $ds^2 = d\eta^2 + \sinh^2 \eta (dt_E^2 + \cos^2 t_E d\theta^2)$

$$K_{\text{Bb}}^\Delta (t_{\text{Eb}}, \eta_b, \theta_b; t_{\text{E1}}, \eta_1 = \infty, \theta_1) = C_\Delta \left(\frac{1}{\cosh \eta_b - \sinh \eta_b \left(\cos(t_{\text{E1}}) \cos(t_{\text{Eb}}) \cos(\theta_b - \theta_1) + \sin(t_{\text{E1}}) \sin(t_{\text{Eb}}) \right)} \right)^\Delta$$

$$\langle \mathcal{O}(x) \rangle_{\text{BCFT}} = \alpha \cdot \int_{\text{EOW}} d^d x_b \sqrt{g} K_{\text{Bb}}^\Delta (x; x_b) .$$

Large dimension limit $\Delta_n \rightarrow \infty$

$$\Delta_n = \frac{c}{12} \left(n - \frac{1}{n} \right)$$

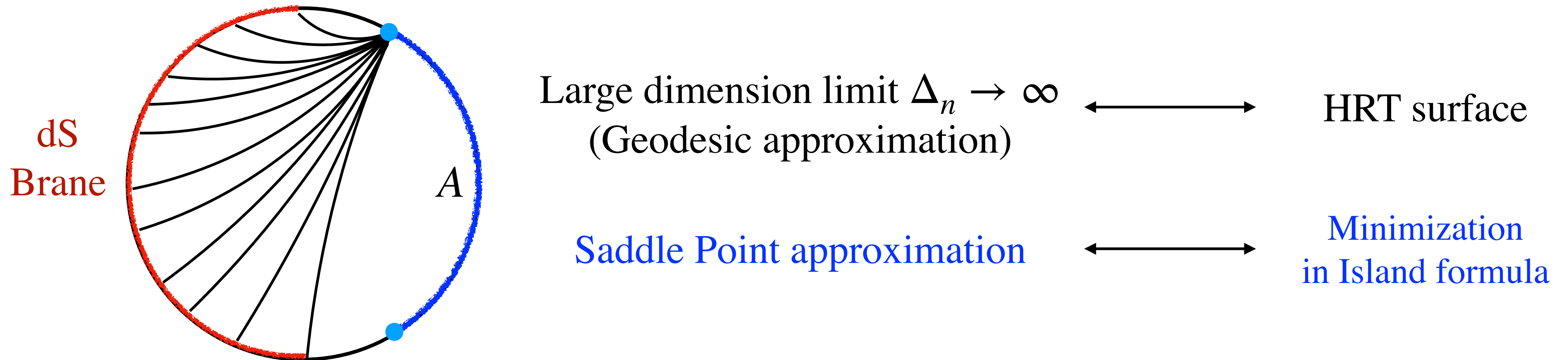
Geodesic approximation $\left\langle \sigma_n(x_A) \right\rangle \approx \text{const.} \times \int_{\text{EOW}} d^2 x_b \sqrt{g} e^{-\Delta_n D(x_A; x_b)}$

04: Resolution from Double Holography

From (Euclidean) AdS3: $ds^2 = d\eta^2 + \sinh^2 \eta (dt_E^2 + \cos^2 t_E d\theta^2)$

$$\left\langle \sigma_n(x_A) \right\rangle \approx \text{const.} \times \int_{\text{EOW}} d^2 x_b \sqrt{g} e^{-\Delta_n D(x_A; x_b)}$$

Geometric/Holographic Interpretation:

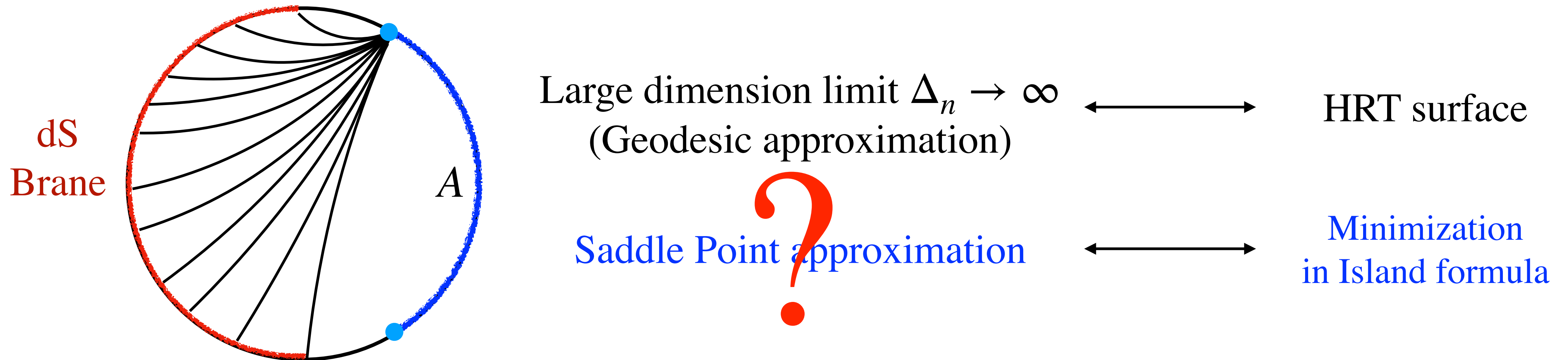


04: Resolution from Double Holography

From (Euclidean) AdS3: $ds^2 = d\eta^2 + \sinh^2 \eta (dt_E^2 + \cos^2 t_E d\theta^2)$

$$\left\langle \sigma_n(x_A) \right\rangle \approx \text{const.} \times \int_{\text{EOW}} d^2 x_b \sqrt{g} e^{-\Delta_n D(x_A; x_b)}$$

Geometric/Holographic Interpretation:



04: Resolution from Double Holography

$$\langle \mathcal{O}(\zeta) \rangle = \int_{\text{dS brane}} d^d \hat{x} \sqrt{h} K_{\text{Bb}}^\Delta(X, Y(\hat{x})) \quad \text{No saddle point approximation!}$$

$$\left\langle \mathcal{O}(t_E, \theta) \right\rangle_{\text{dS}} \approx \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} dt_{\text{Eb}} \int_{-\frac{\pi}{2}}^{\frac{\pi}{2}} d\theta_b \cos(t_{\text{Eb}}) e^{-\Delta D(t_E, \theta; t_{\text{Eb}}, \theta_b)},$$

only one saddle point for the integral: $x_b^* = (\theta_b = \theta - \pi, t_{\text{Eb}} = -t_E)$

This saddle point is a local maximum!

$$\left. \frac{\partial^2 D(t_E, \theta; t_{\text{Eb}}, \theta_b)}{\partial t_{\text{Eb}}^2} \right|_{x_b=x_b^*} < 0, \quad \left. \frac{\partial^2 D(t_E, \theta; t_{\text{Eb}}, \theta_b)}{\partial \theta_b^2} \right|_{x_b=x_b^*} < 0.$$

04: Resolution from Double Holography

$$\langle \mathcal{O}(\zeta) \rangle = \int_{\text{dS brane}} d^d \hat{x} \sqrt{h} K_{\text{Bb}}^\Delta(X, Y(\hat{x})) \quad \text{Dominated by the edge}$$

$$\langle \sigma_n(t_{A_1}, \theta_{A_1}) \bar{\sigma}_n(t_{A_2}, \theta_{A_2}) \rangle_{\text{BCFT}} = \max \begin{cases} \langle \sigma_n(t_{A_1}, \theta_{A_1}) \bar{\sigma}_n(t_{A_2}, \theta_{A_2}) \rangle \\ \langle \sigma_n(t_{A_1}, \theta_{A_1}) \rangle \langle \bar{\sigma}_n(t_{A_2}, \theta_{A_2}) \rangle \end{cases}$$

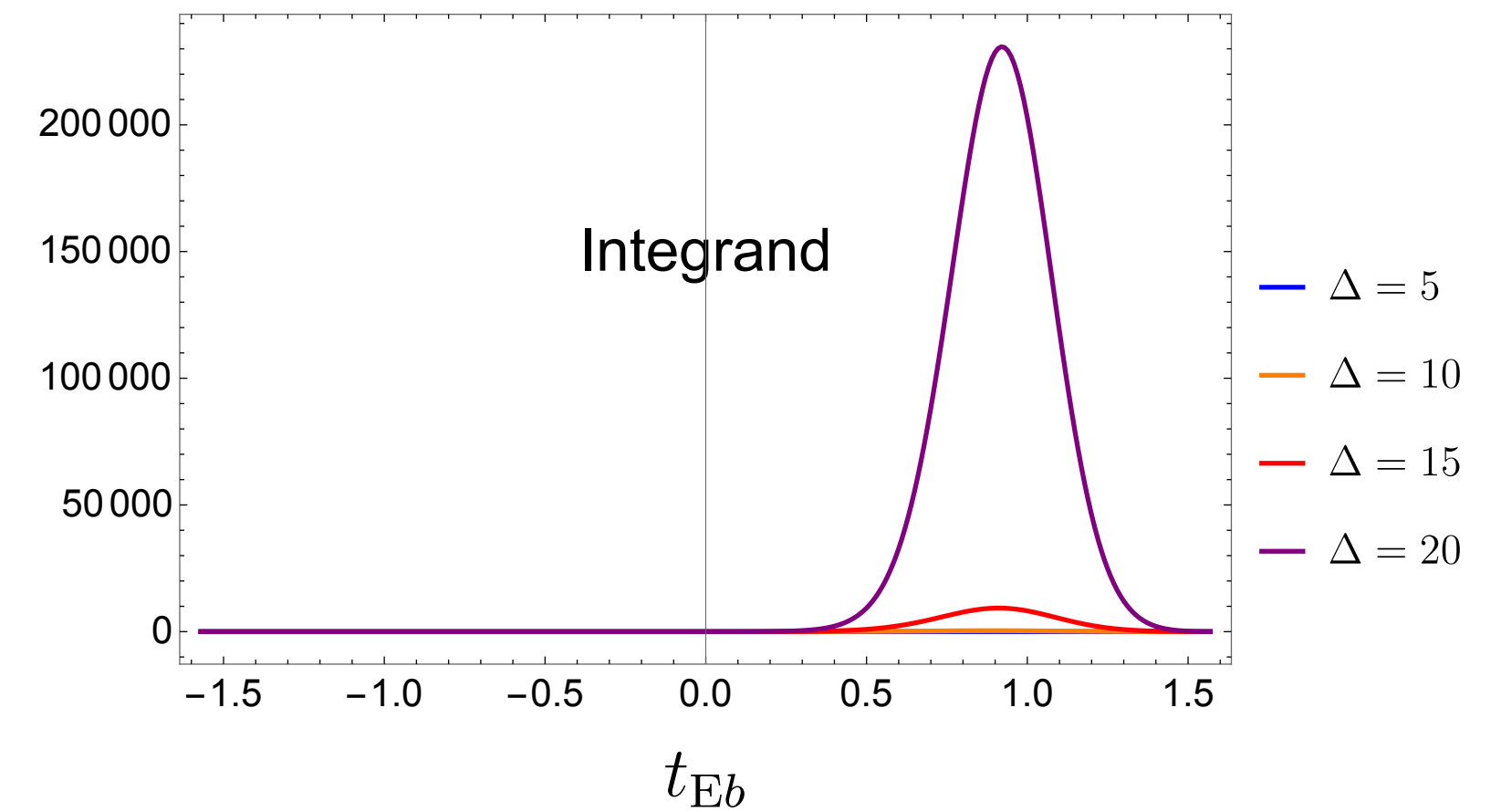
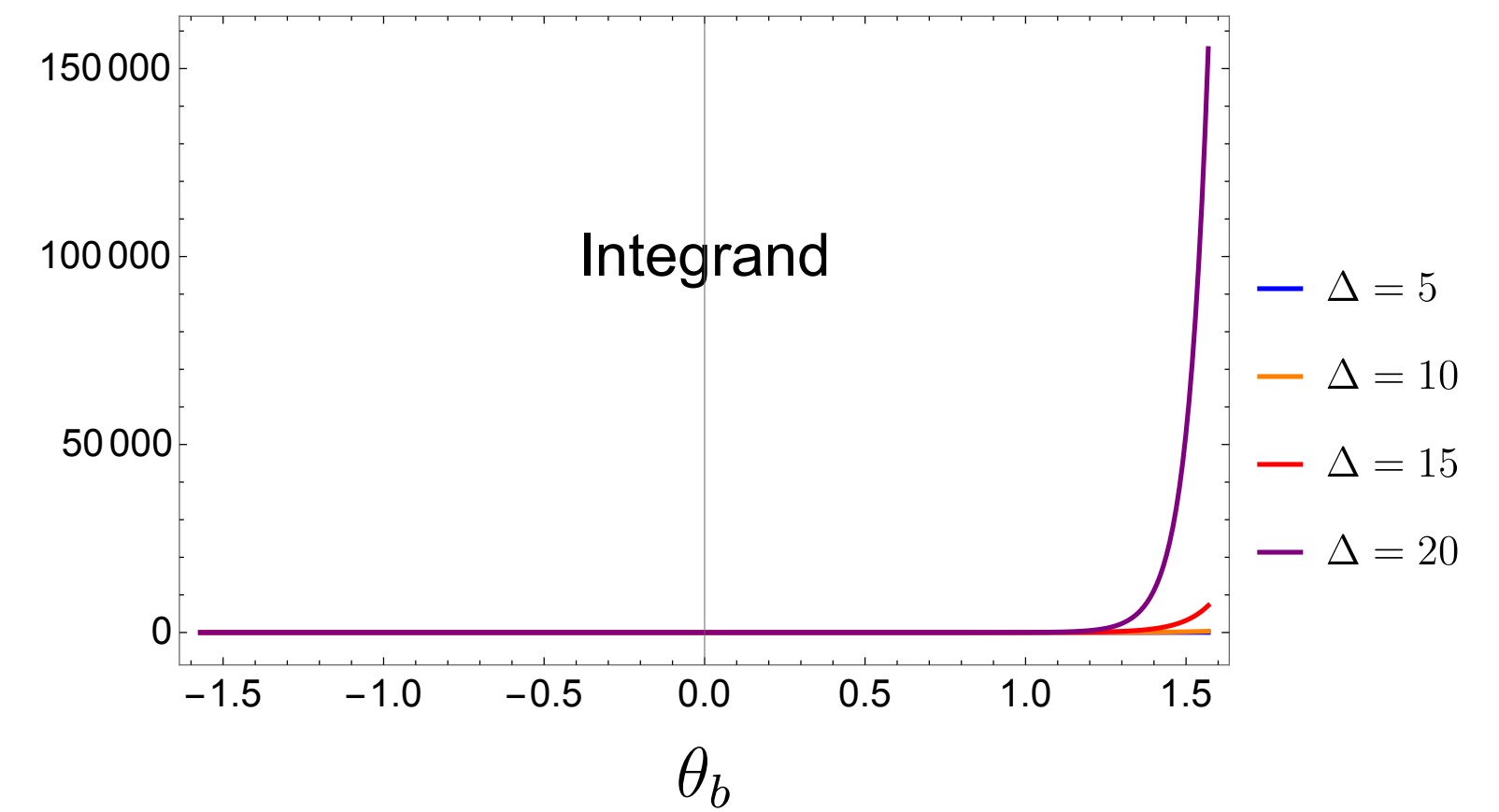
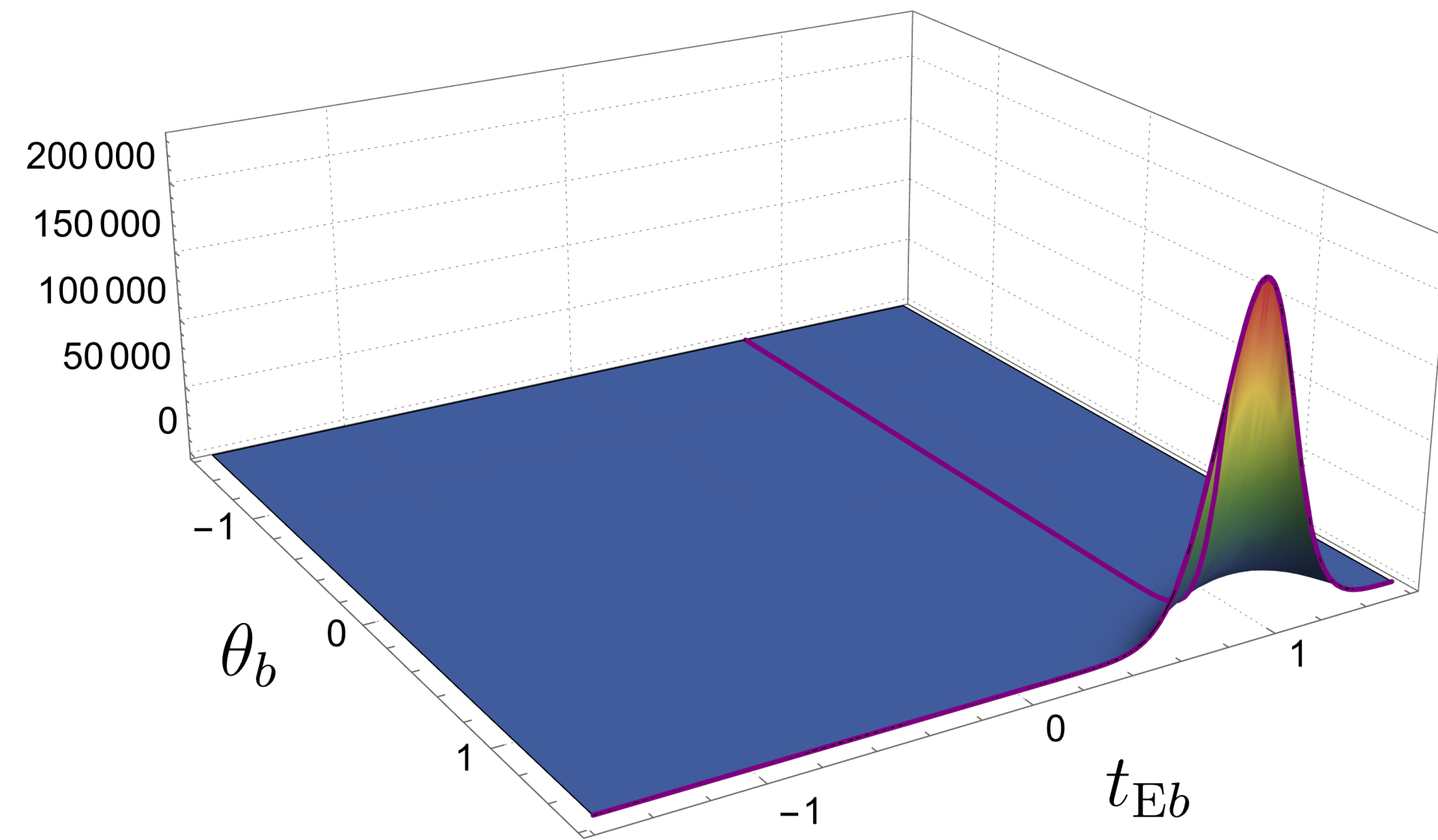
Double Holography

Leading Result from the integral:

$$\text{Entanglement Entropy: } S_{\mathcal{A}} = \min \begin{cases} S_{\mathcal{A}}^{\text{con}} = \frac{c}{3} \log \left(\frac{2 \cosh t_A \cdot \sin \theta_A}{\epsilon} \right) \\ \tilde{S}_{\mathcal{A}}^{\text{dis}} = \frac{c}{3} \log \left(\frac{2}{\epsilon} \left(\cosh \eta_b - \sinh \eta_b \sqrt{\cosh^2 t_A \sin^2 \theta_A - \sinh^2 t_A} \right) \right) \end{cases}$$

04: Resolution from Double Holography

$$K_{\text{Bb}}^{\Delta} (t_{\text{Eb}}, \eta_b, \theta_b; t_{\text{E1}}, \eta_1 = \infty, \theta_1) = C_{\Delta} \left(\frac{1}{\cosh \eta_b - \sinh \eta_b \left(\cos(t_{\text{E1}}) \cos(t_{\text{Eb}}) \cos(\theta_b - \theta_1) + \sin(t_{\text{E1}}) \sin(t_{\text{Eb}}) \right)} \right)^{\Delta}$$



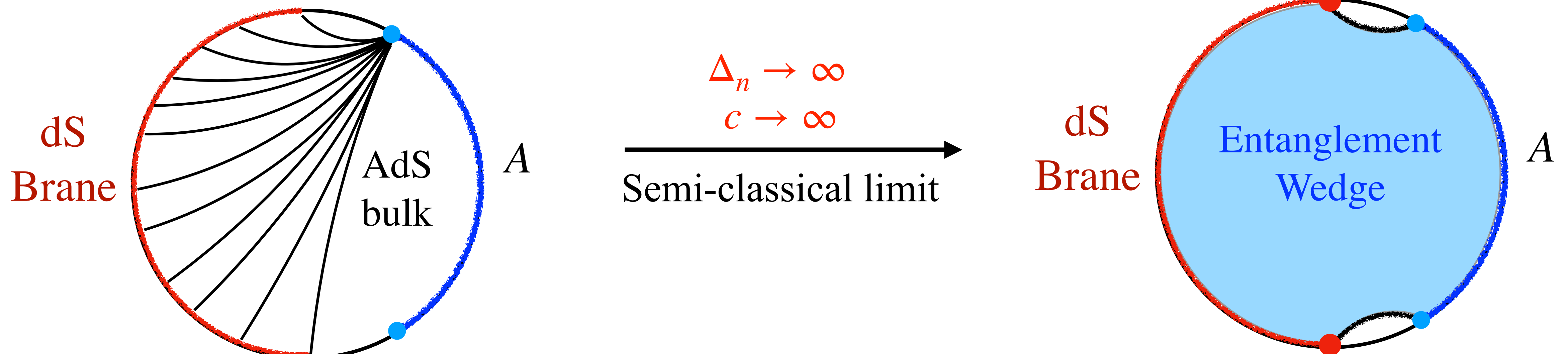
04: Resolution from Double Holography

To get the same results from holography

$$S_A = \min \{ S_A^{\text{con}}, \tilde{S}_A^{\text{dis}} \} = \min \left\{ \frac{\text{Area}(\Gamma_A^{\text{con}})}{4G_N}, \frac{\text{Area}(\tilde{\Gamma}_A^{\text{dis}})}{4G_N} \right\},$$

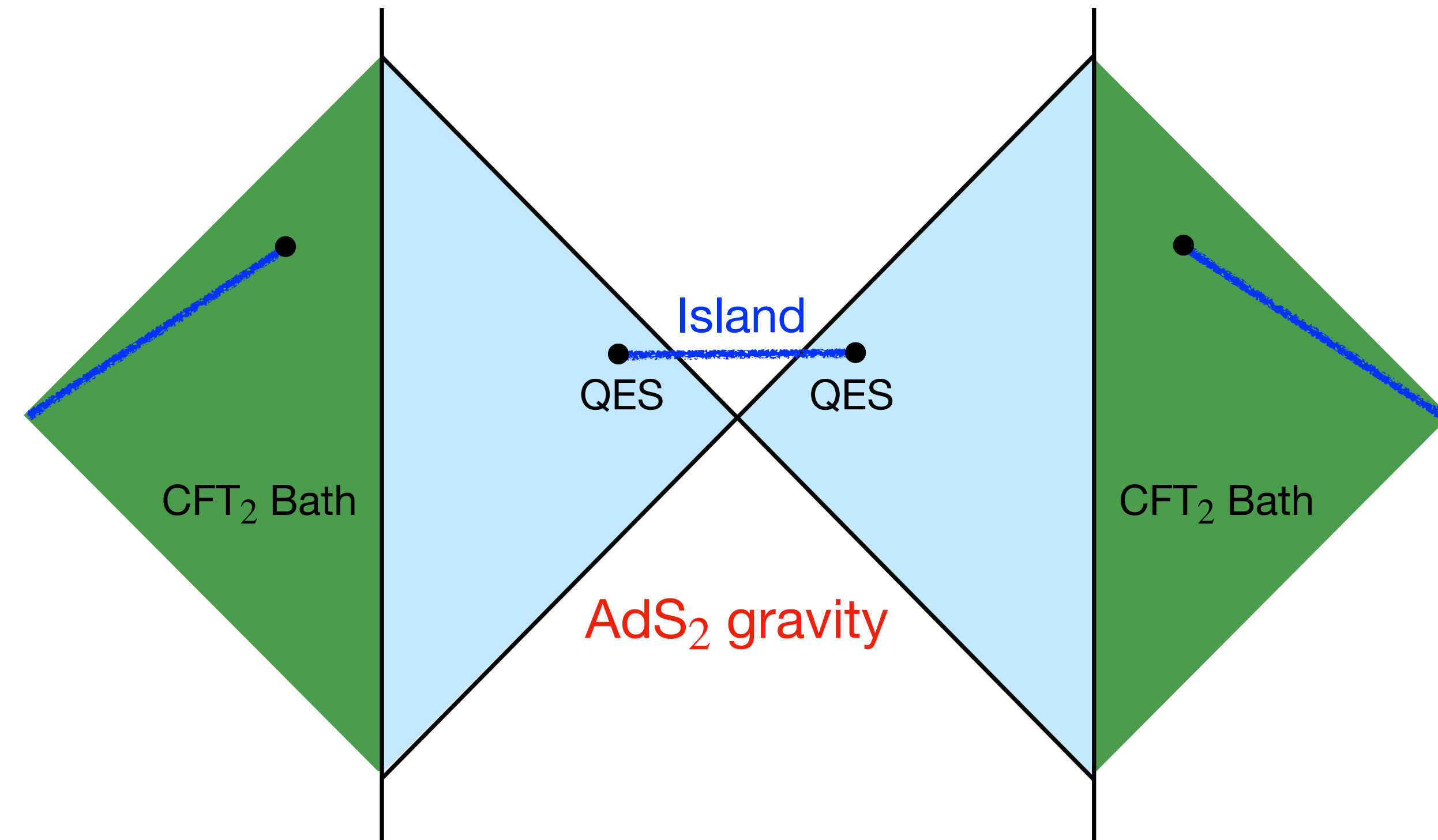
By anchoring the non-extremal surface at the edge of dS gravity region

Holographic dual: Non-extremal Island in dS



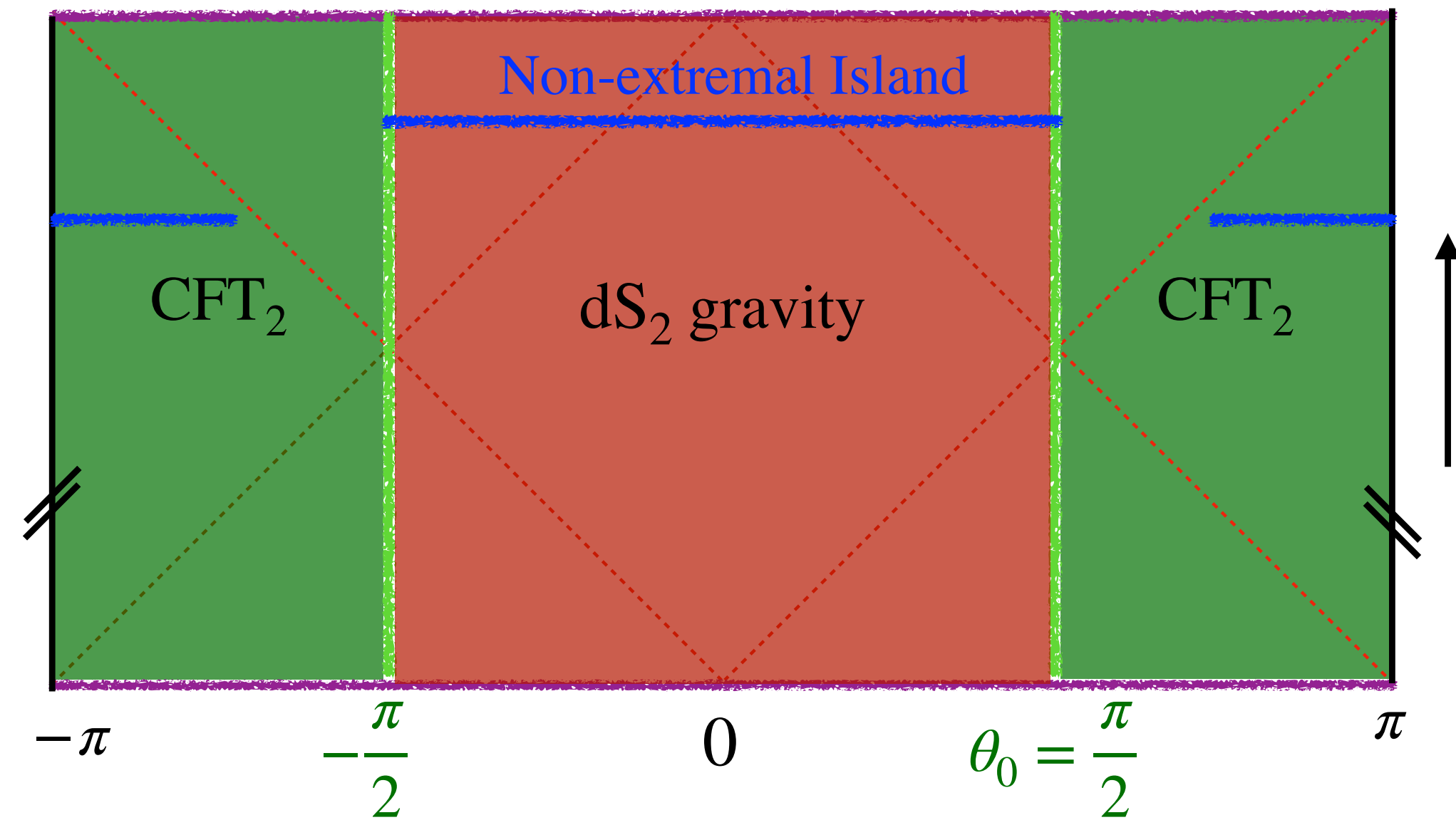
Holographic Suggestion: Non-extremal Island

AdS+Bath



Island Formula
(Quantum Extremal Surface)

dS+Bath



Non-extremal Island
(Non-extremal Surface)

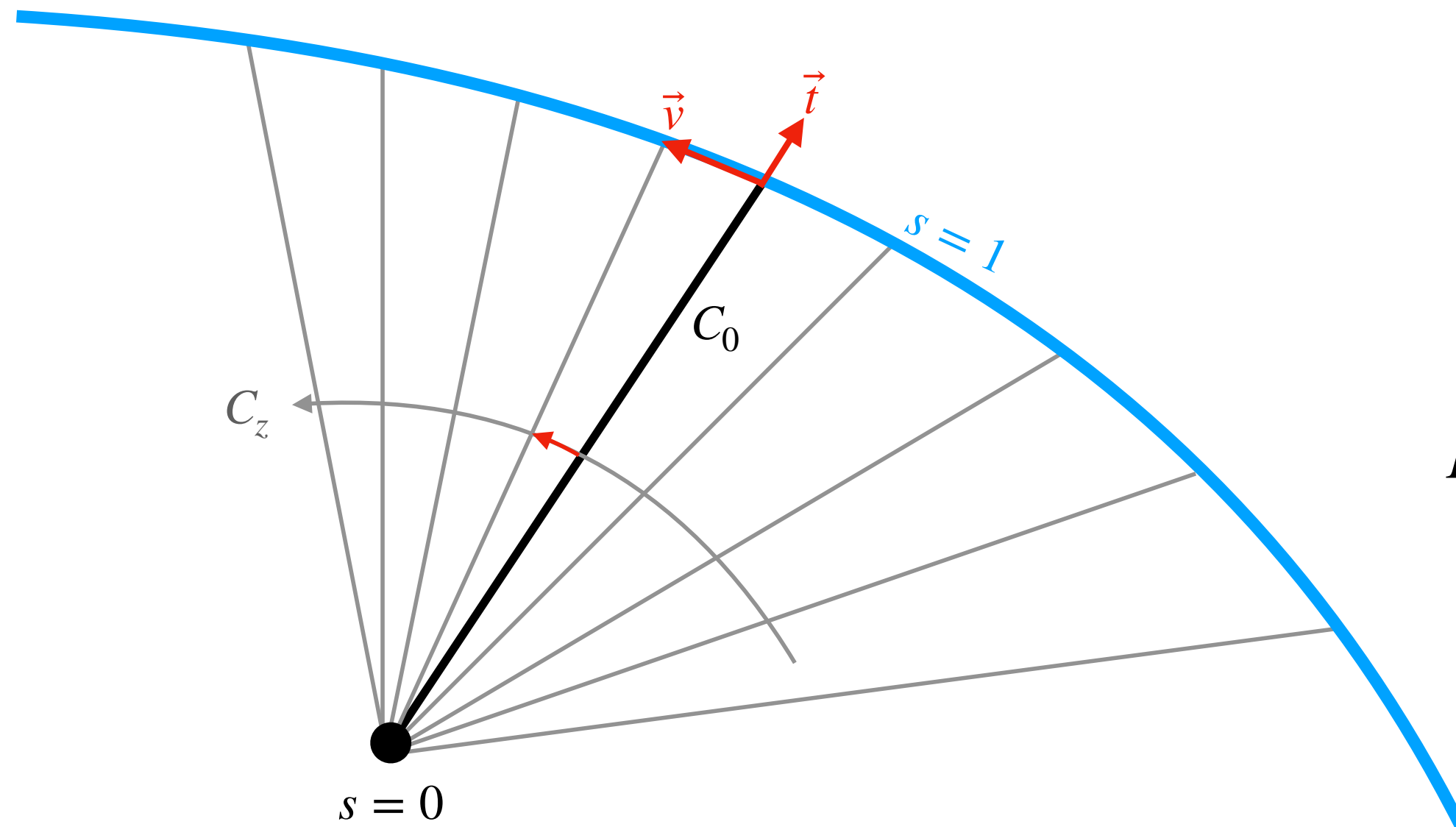
Question: Does Island formula apply to dS gravity?

NO!

Why is there no quantum extremal surface on dS gravity?

04: Resolution from Double Holography

Why is there no quantum extremal surface on dS gravity?



Synge's formula for the second variation of geodesics

$$L''(C_0) = -K(\mathbf{v}, \mathbf{v}) + \frac{1}{L_0} \int_0^1 \left[\left\| (\nabla_t \mathbf{v}) \right\|^2 - K_{\text{set}}(\mathbf{v}, \mathbf{t}) \|\mathbf{v} \wedge \mathbf{t}\|^2 \right] ds$$

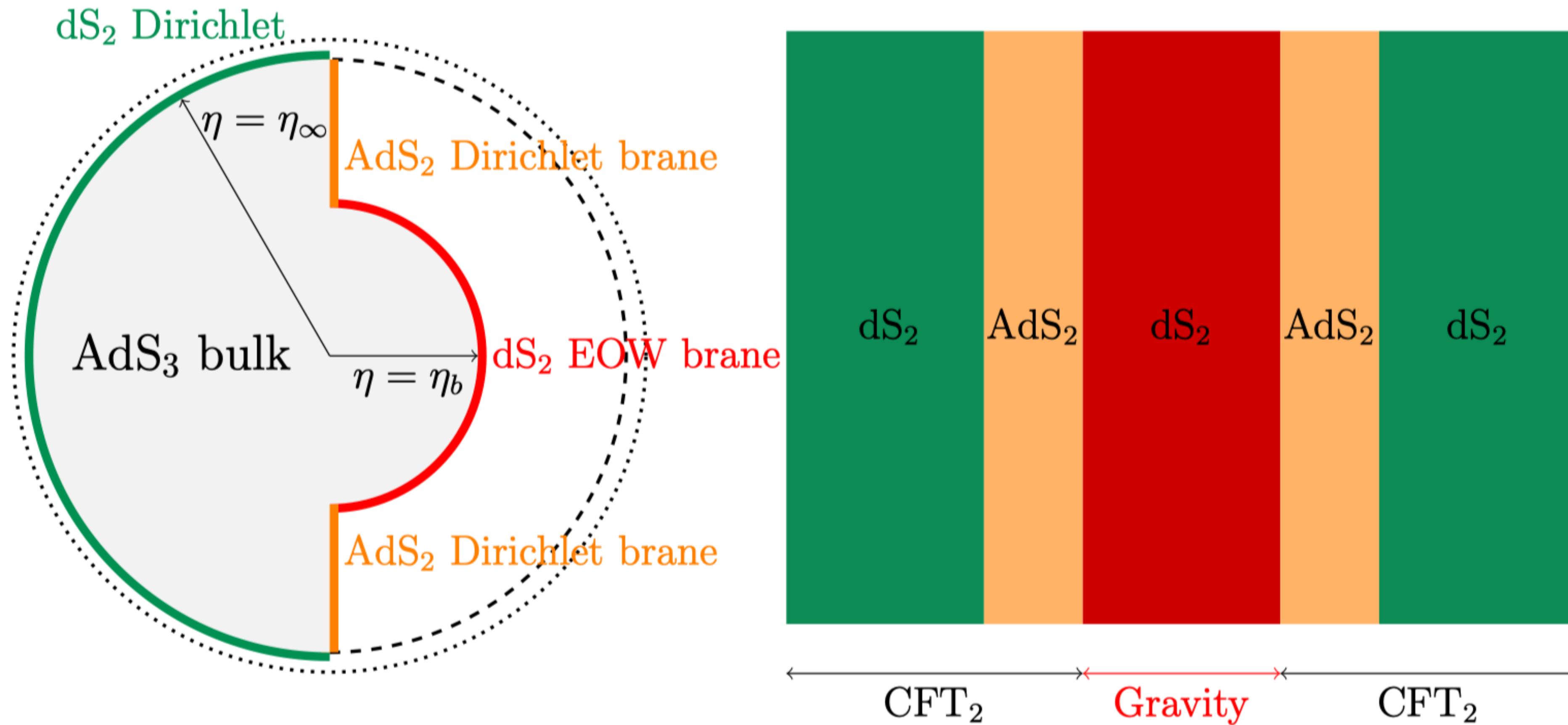
A geometric explanation from the bulk AdS:
the (positive) extrinsic curvature of the dS brane is too large!

$$K(\mathbf{v}, \mathbf{v}) \Big|_{\text{dS brane}} \propto K \Big|_{\text{dS brane}} > \frac{d}{L_{\text{AdS}}}$$

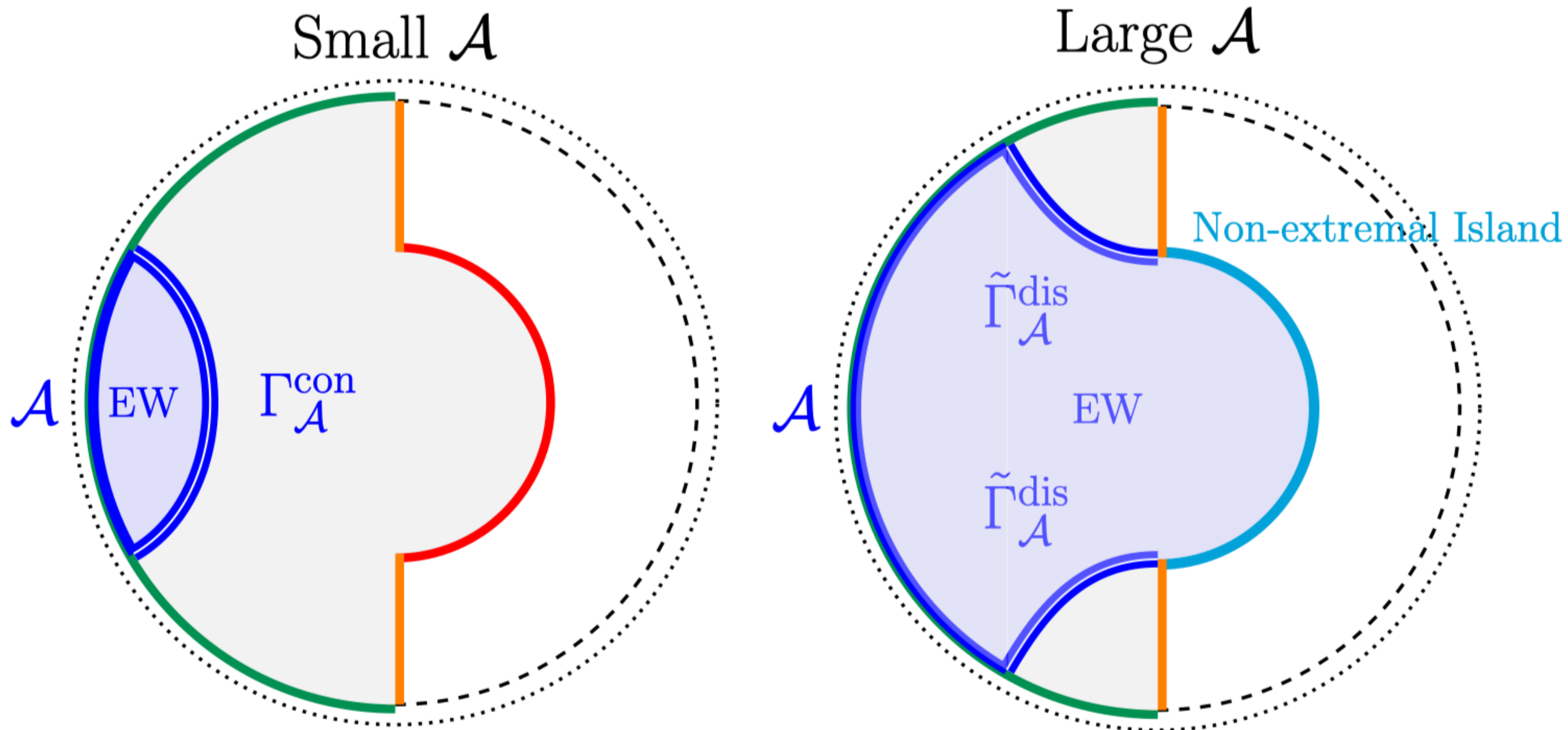
Thanks for your attention!

Extra Slices

Improved doubly holographic model



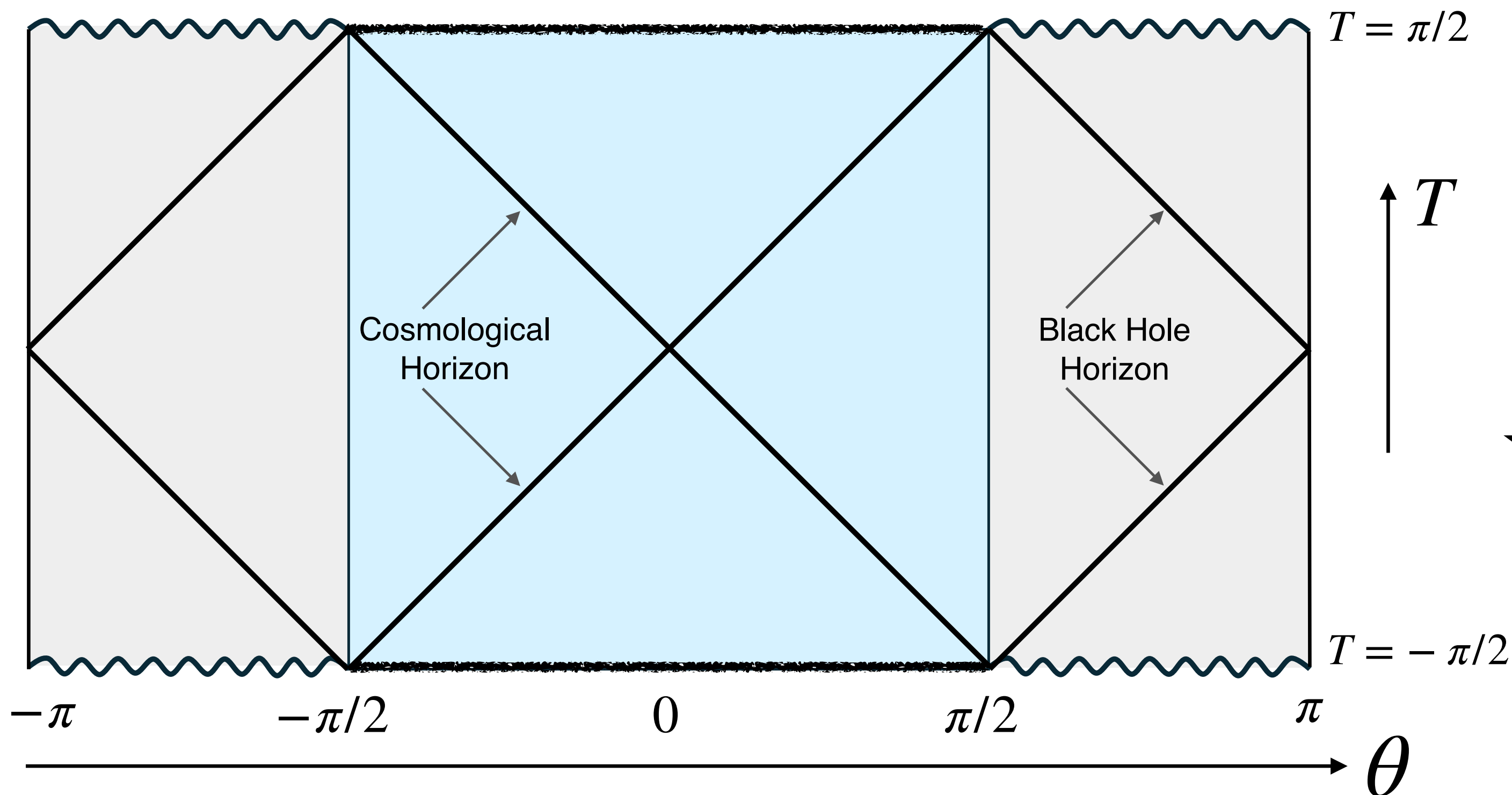
Improved doubly holographic model



Appendix: Extremal Islands in dS JT gravity?

$$S_{\text{JT}'} = \frac{\Phi_0}{16\pi G_N^{(2)}} \int d^2x \sqrt{-g} R + \frac{1}{16\pi G_N^{(2)}} \int d^2x \sqrt{-g} \Phi \left(R - \frac{2}{L^2} \right) + \text{boundary terms},$$

Dilation profile: $\Phi(T, \theta) = \Phi_r \frac{\cos \theta}{\cos T}$



★ QES could be locally minimal in the spatial direction only if

$$\Phi(T, \theta) < 0$$

Behind the black hole horizon

★ QES has to be close to the BH singularity

$$\frac{1}{\cos T_I} > \frac{c}{12\Phi_r} \gg 1$$