

Gauged Global Strings

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Outline

- Introduction
global strings and gauge strings
- Gauge $U(1)_Z \times$ global $U(1)_{PQ}$
and **string solutions**
- Cosmological implication
 - 1) **rich string structure / dynamics**
 - 2) **opening up QCD axion window**
 - 3) **gauge string radiating axions?**
- Conclusion

Global strings

Global strings

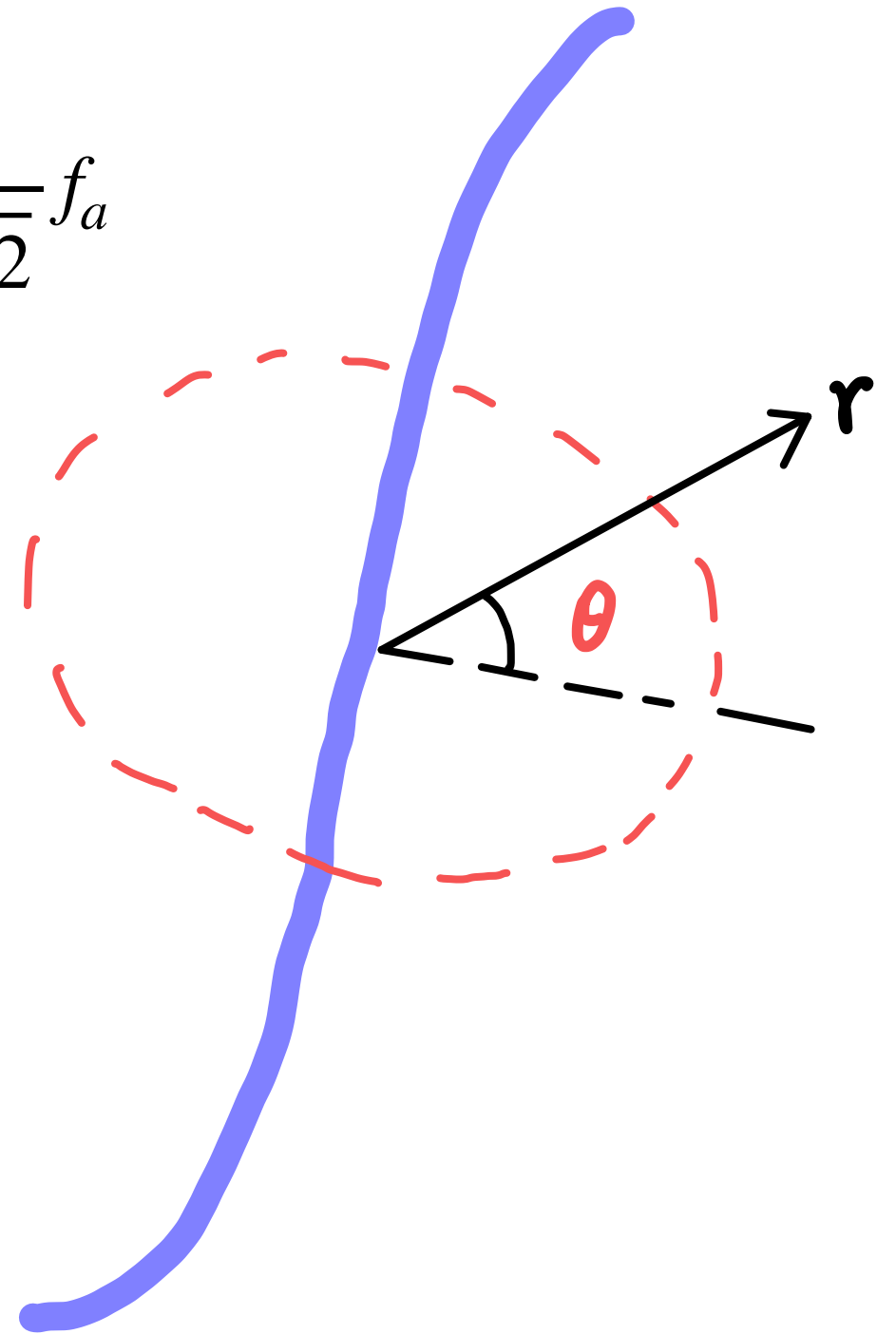
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- global string solution

$$\Phi(r, \theta) = \frac{1}{\sqrt{2}} f_a e^{i\theta}, \quad r \rightarrow \infty$$



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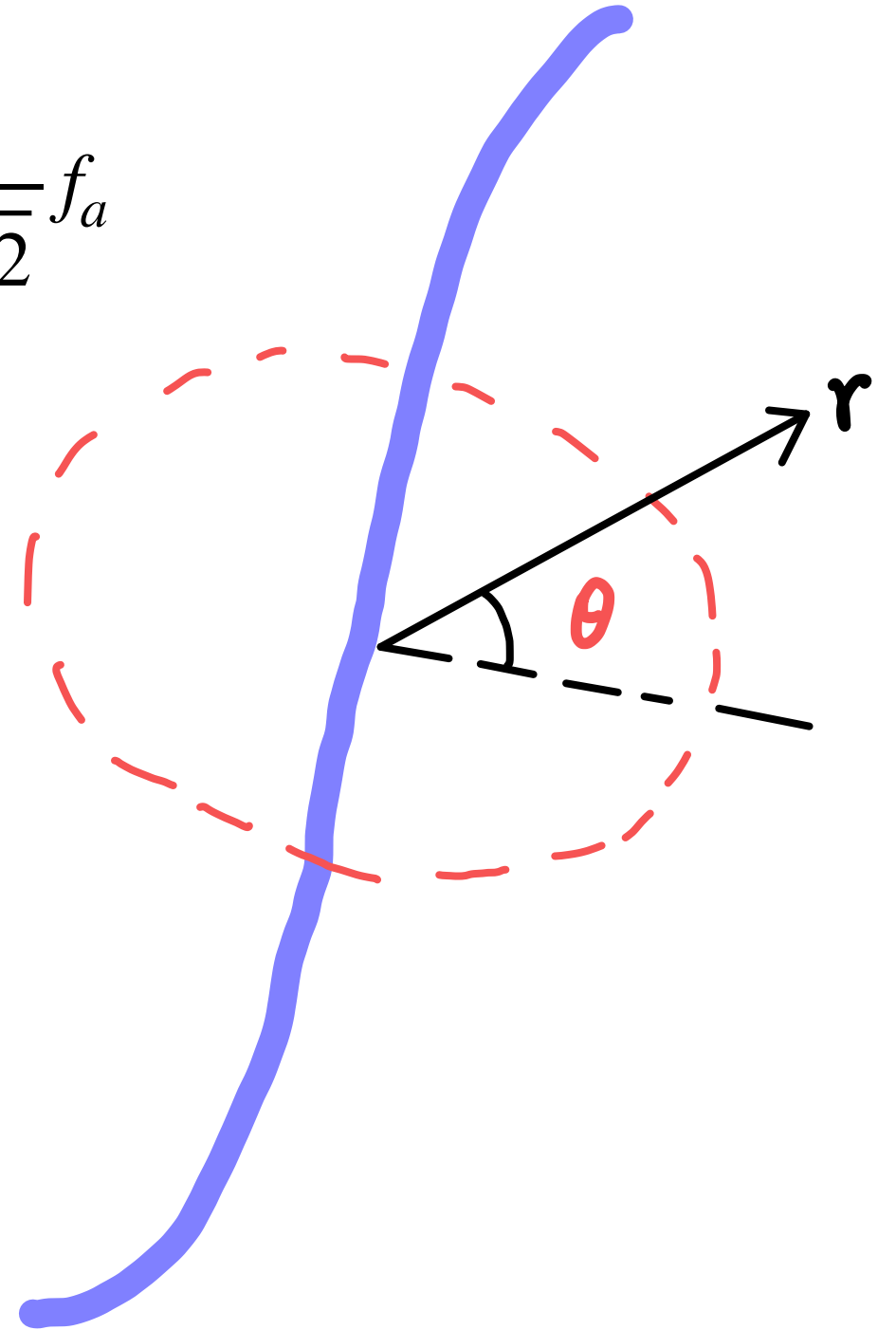
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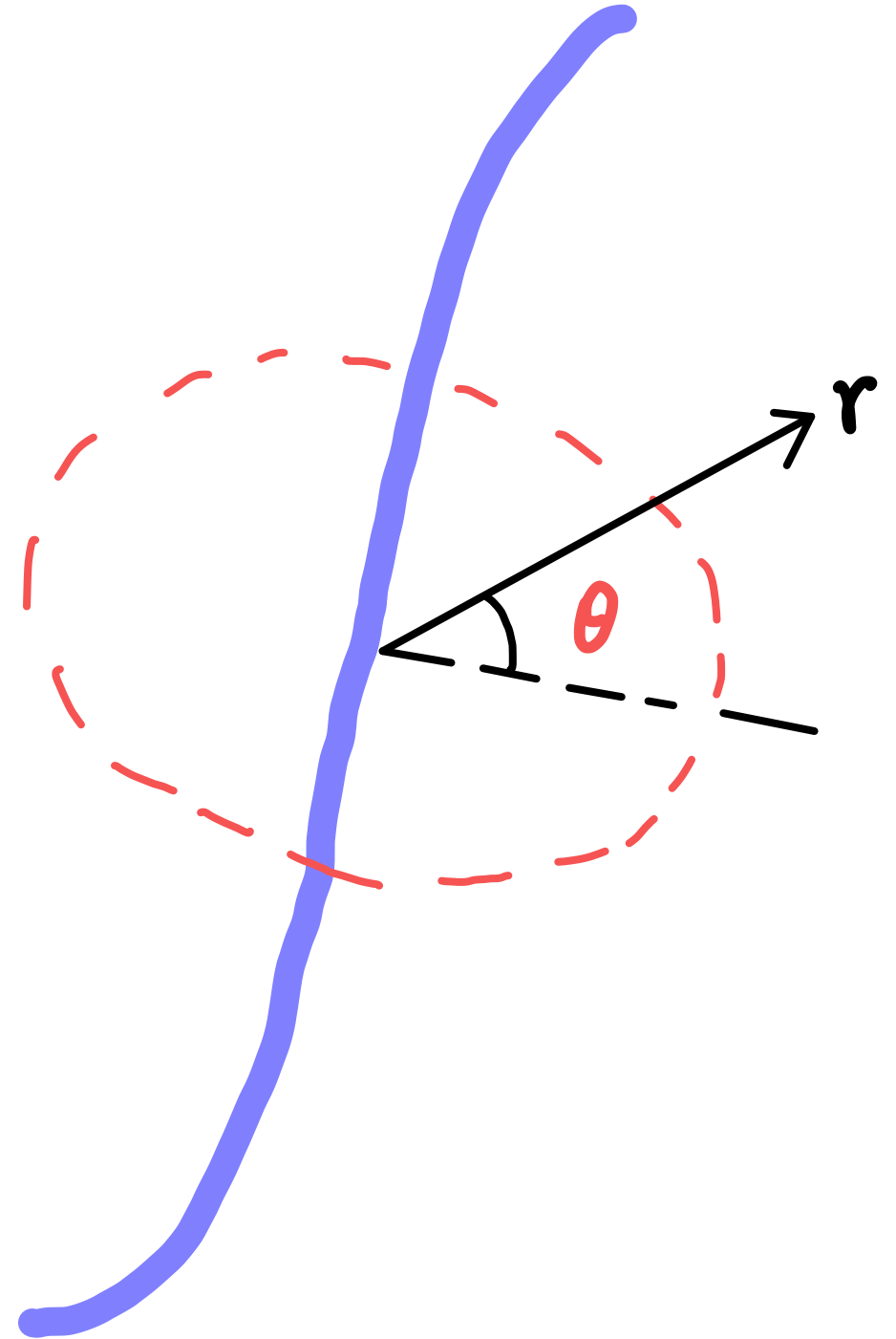
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- tension

gradient term $\mu \simeq 2\pi \int_{m^{-1}}^L dr \frac{1}{r} |\partial_\theta \Phi(r, \theta)|^2 = \pi f_a^2 \ln(mL)$



Gauge strings



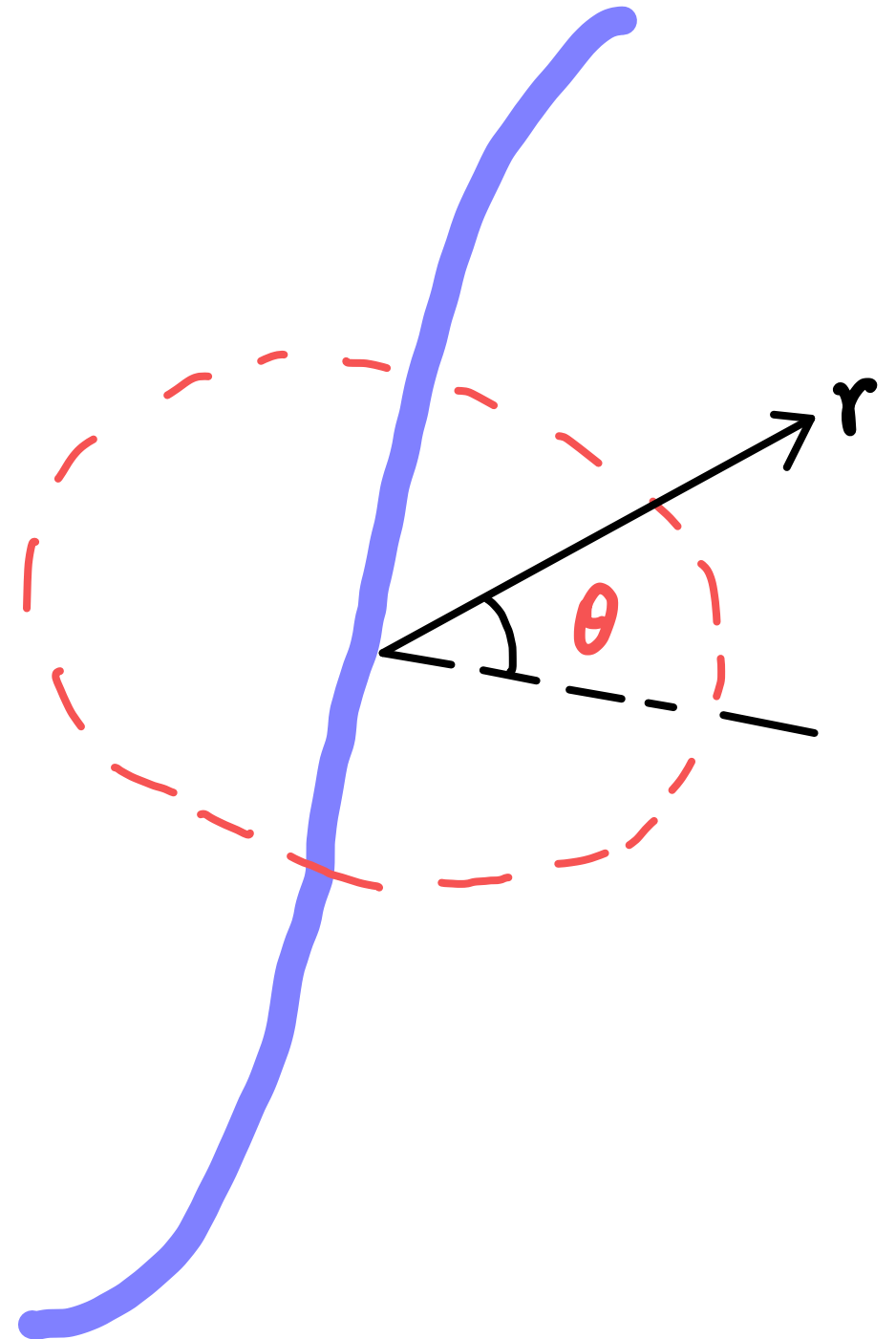
Gauge strings

- gauge string solution

$$\Phi(r, \theta) = \frac{1}{\sqrt{2}} f_a e^{i\theta}$$

$$Z_\mu = \frac{1}{e} \partial_\mu \theta$$

$$r \rightarrow \infty$$

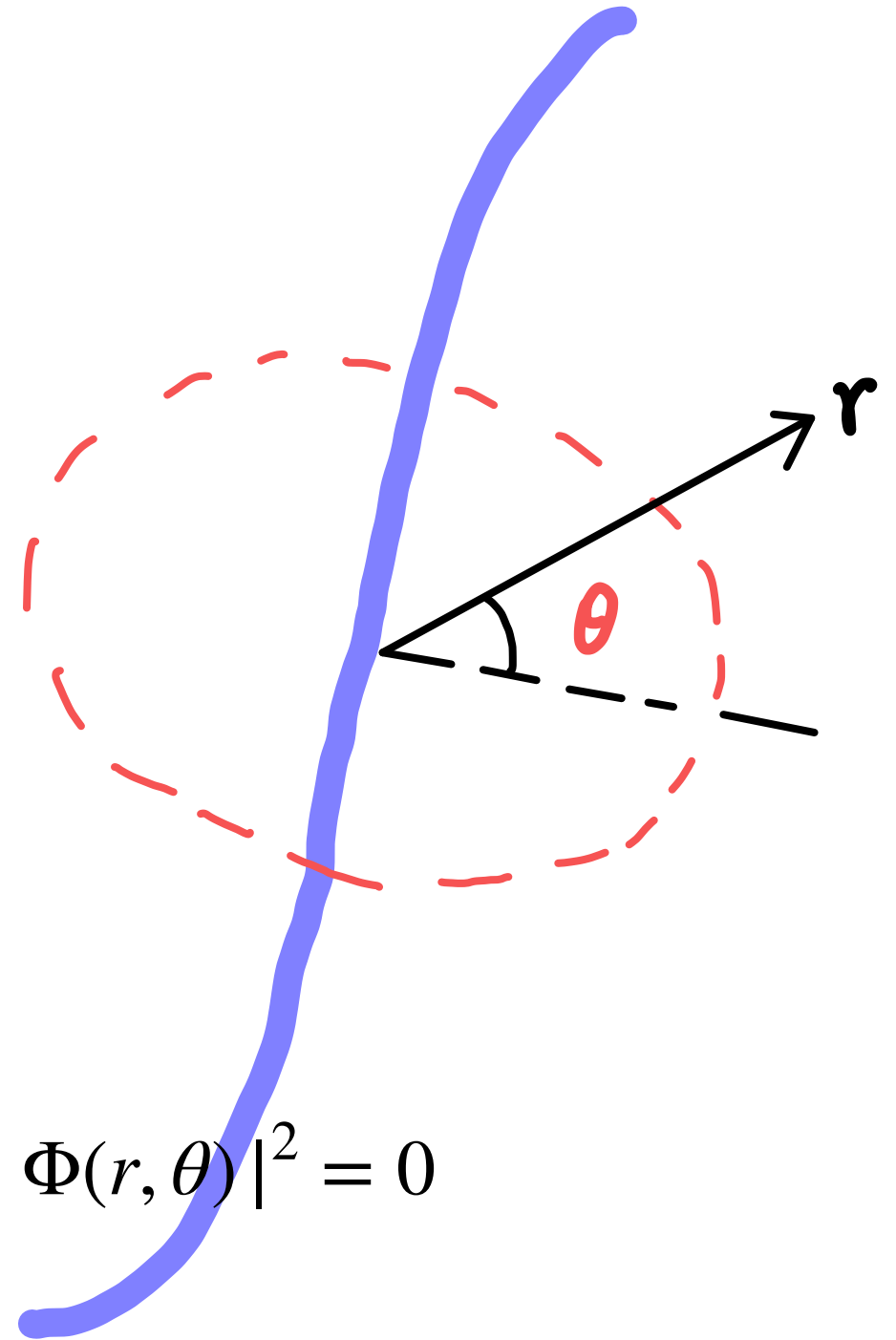


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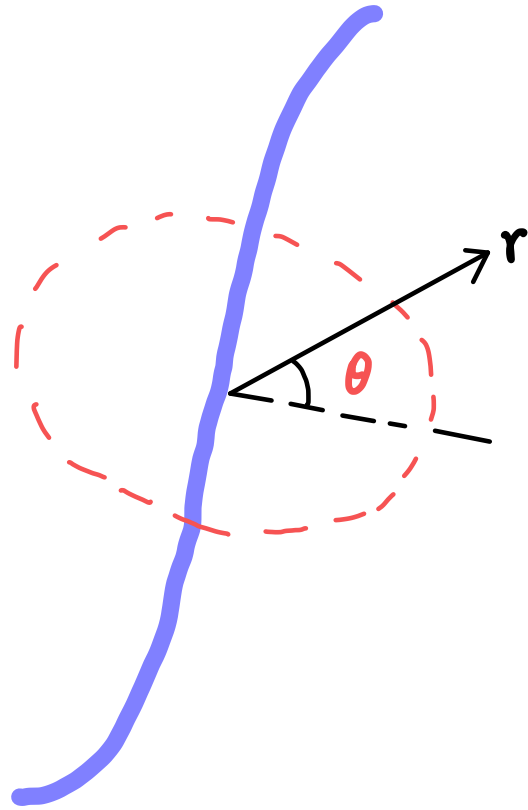
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gradient term $\mu \simeq 2\pi \int_{m^{-1}}^L dr \left| \left(\frac{1}{r} \partial_\theta - ie Z_\mu \right) \Phi(r, \theta) \right|^2 = 0$

core $\mu \simeq \mathcal{O}(1) \pi f_a^2$

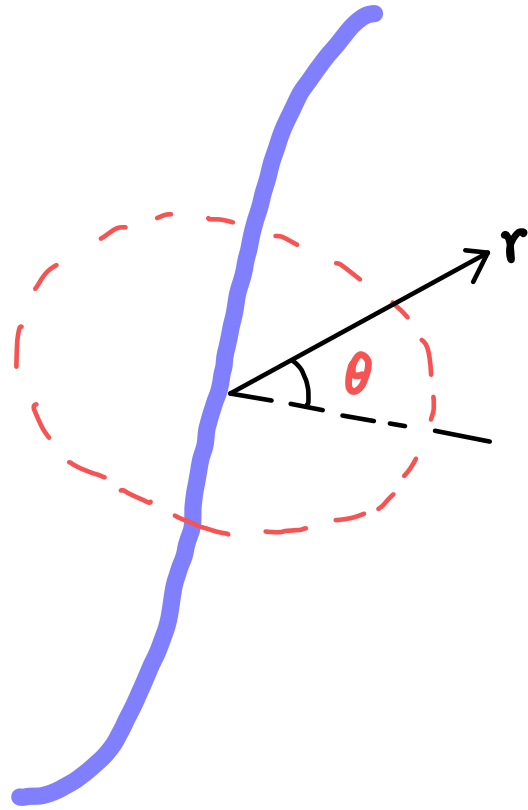
Motivation of cosmic strings

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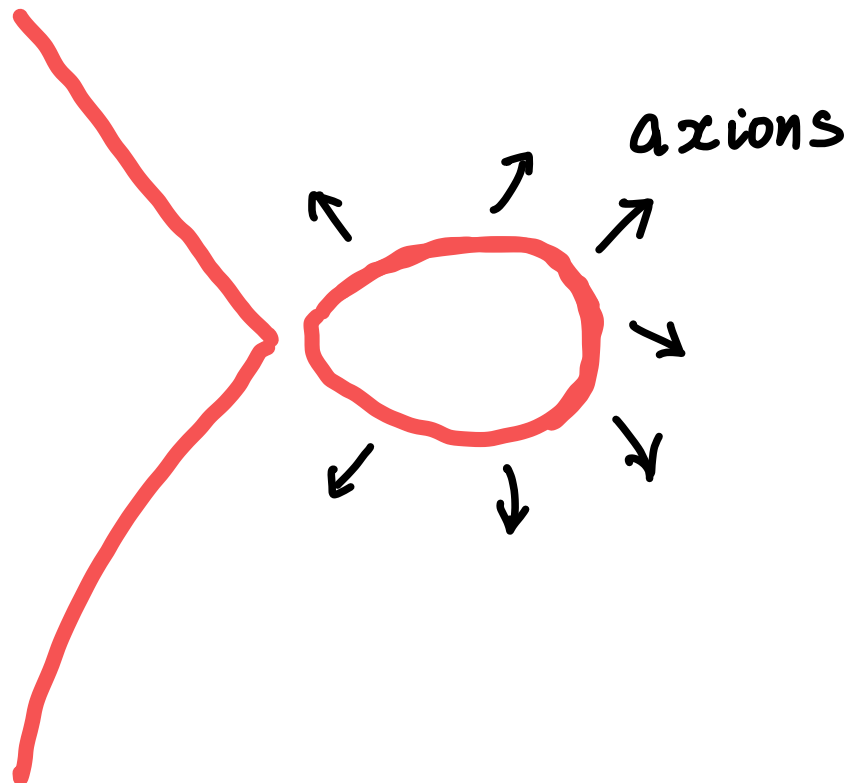


- theoretically interesting classical field solutions

Motivation of cosmic strings



- theoretically interesting classical field solutions



- phenomenological rich cosmology (Kibble mechanism)
axion dark matter abundance
new observables (CMB, ...)

$$U(1)_Z \times U(1)_{PQ}$$

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- Lagrangian

$$\mathcal{L} = -\frac{1}{4}Z_{\mu\nu}Z^{\mu\nu} + D_\mu\Phi_1^\dagger D^\mu\Phi_1 - \frac{\lambda_1}{4}\left(|\Phi_1|^2 - \frac{v_1^2}{2}\right)^2 + D_\mu\Phi_2^\dagger D^\mu\Phi_2 - \frac{\lambda_2}{4}\left(|\Phi_2|^2 - \frac{v_2^2}{2}\right)^2$$

$$D_\mu = \partial_\mu - ieZ_\mu$$

assume that $v_1 > v_2$

$U(1)_Z \times U(1)_{PQ}$

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- $\Phi_1 \rightarrow \Phi_1 e^{i\alpha_Z + i\alpha_{PQ}}$

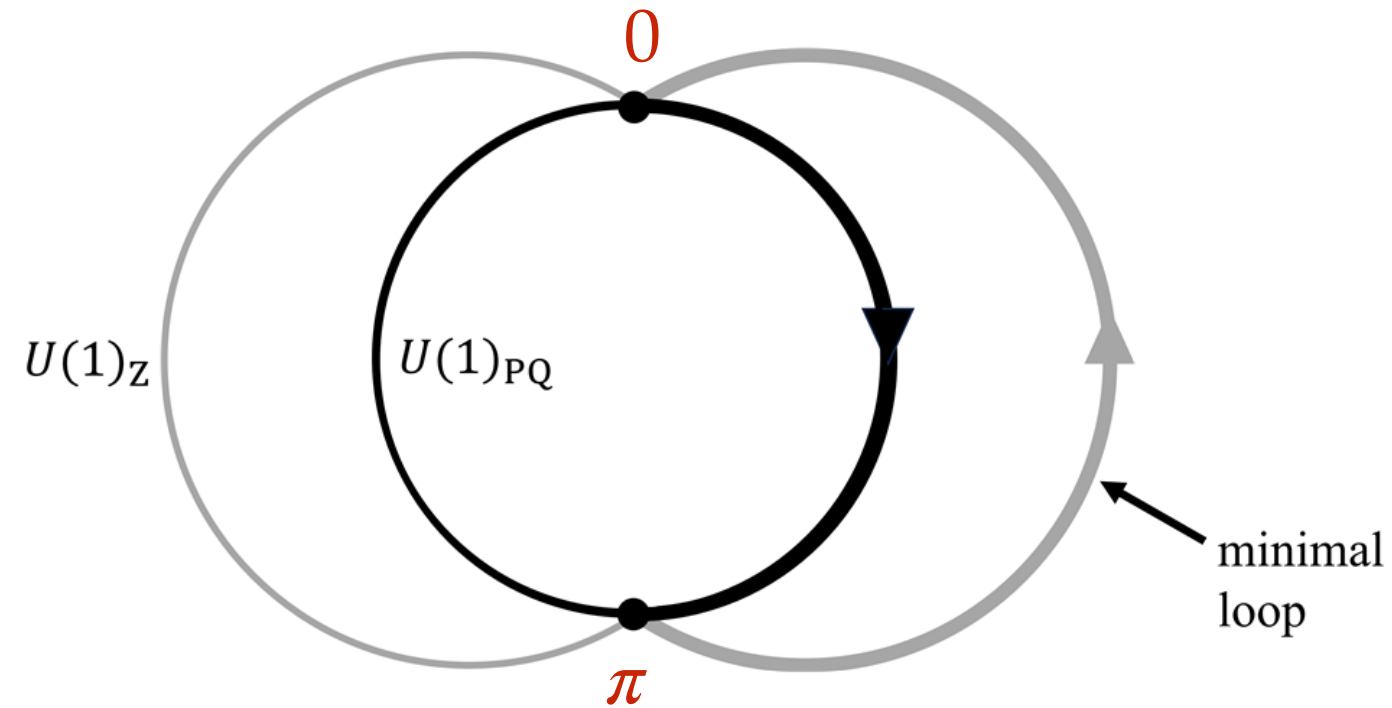
$$\Phi_2 \rightarrow \Phi_2 e^{i\alpha_Z - i\alpha_{PQ}}$$

	Φ_1	Φ_2
$U(1)_Z$	1	1
$U(1)_{PQ}$	1	-1

Vacuum and fluctuations

Vacuum and fluctuations

- cross section of vacuum manifold

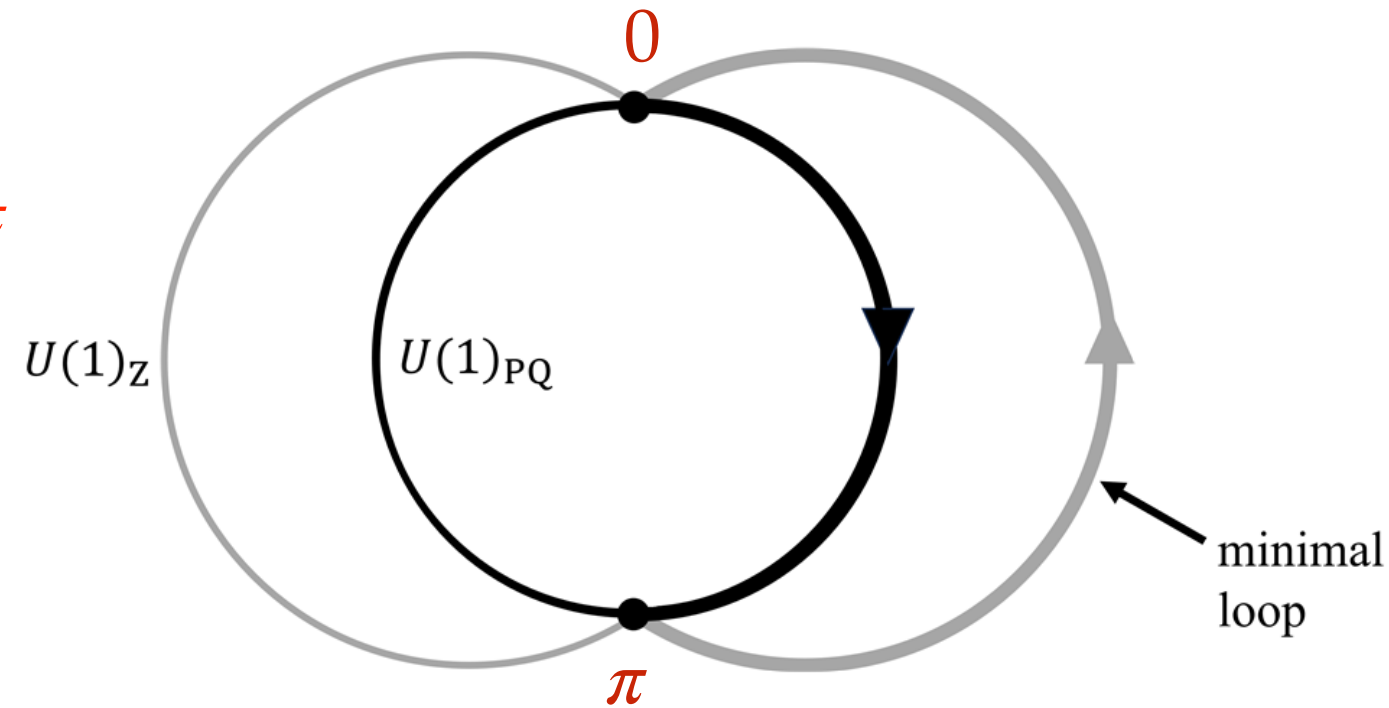


Vacuum and fluctuations

- cross section of vacuum manifold

$$\alpha_Z = \pi$$

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Vacuum and fluctuations

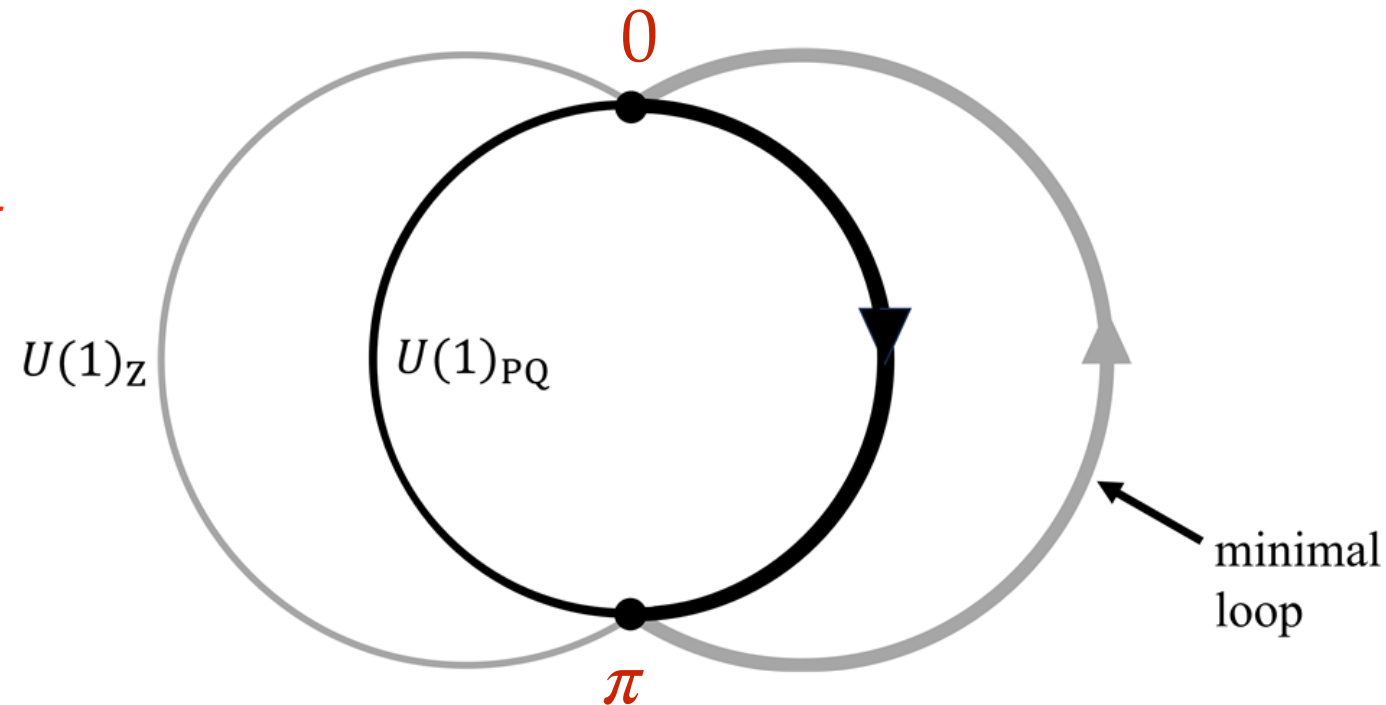
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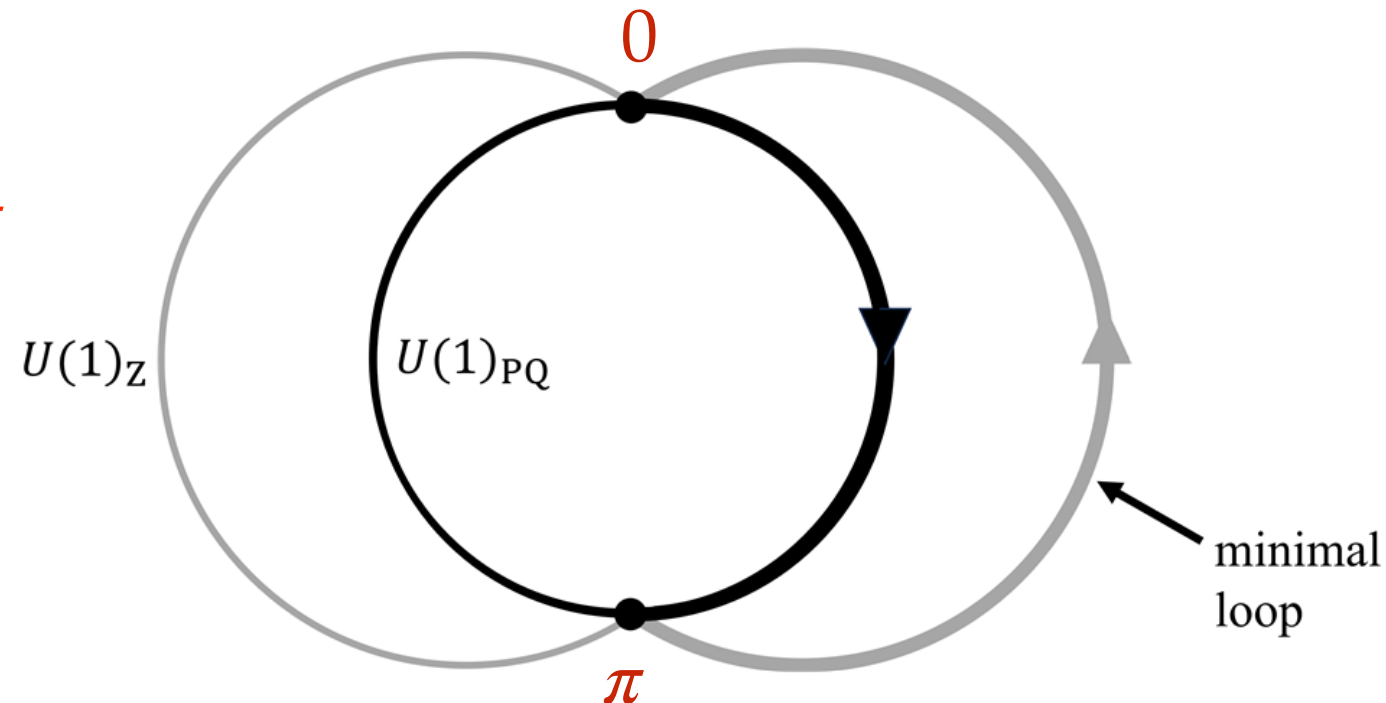
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- axion direction is orthogonal to the longitudinal mode of Z^μ

$$a(x) = v_a \alpha_{PQ}, \quad v_a = \frac{2v_1 v_2}{\sqrt{v_1^2 + v_2^2}} \sim 2v_2$$

Integrating with QCD axion model

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- KSVZ-like model
introduce Q_L and Q_R with color charge and $U(1)_{PQ}$ charge

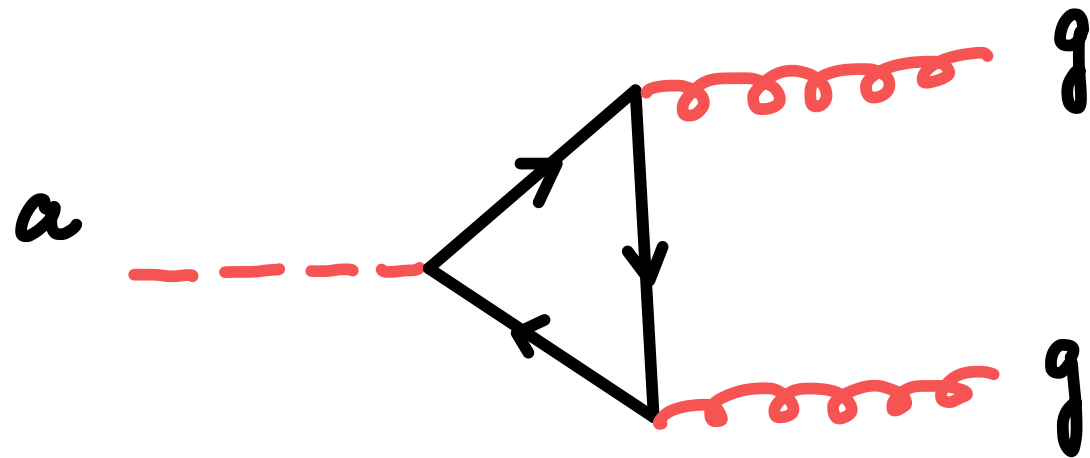
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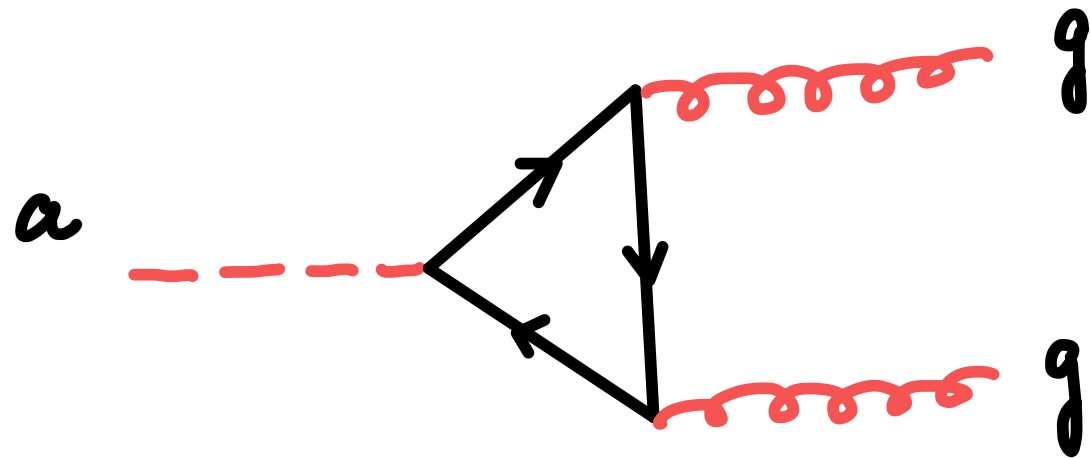


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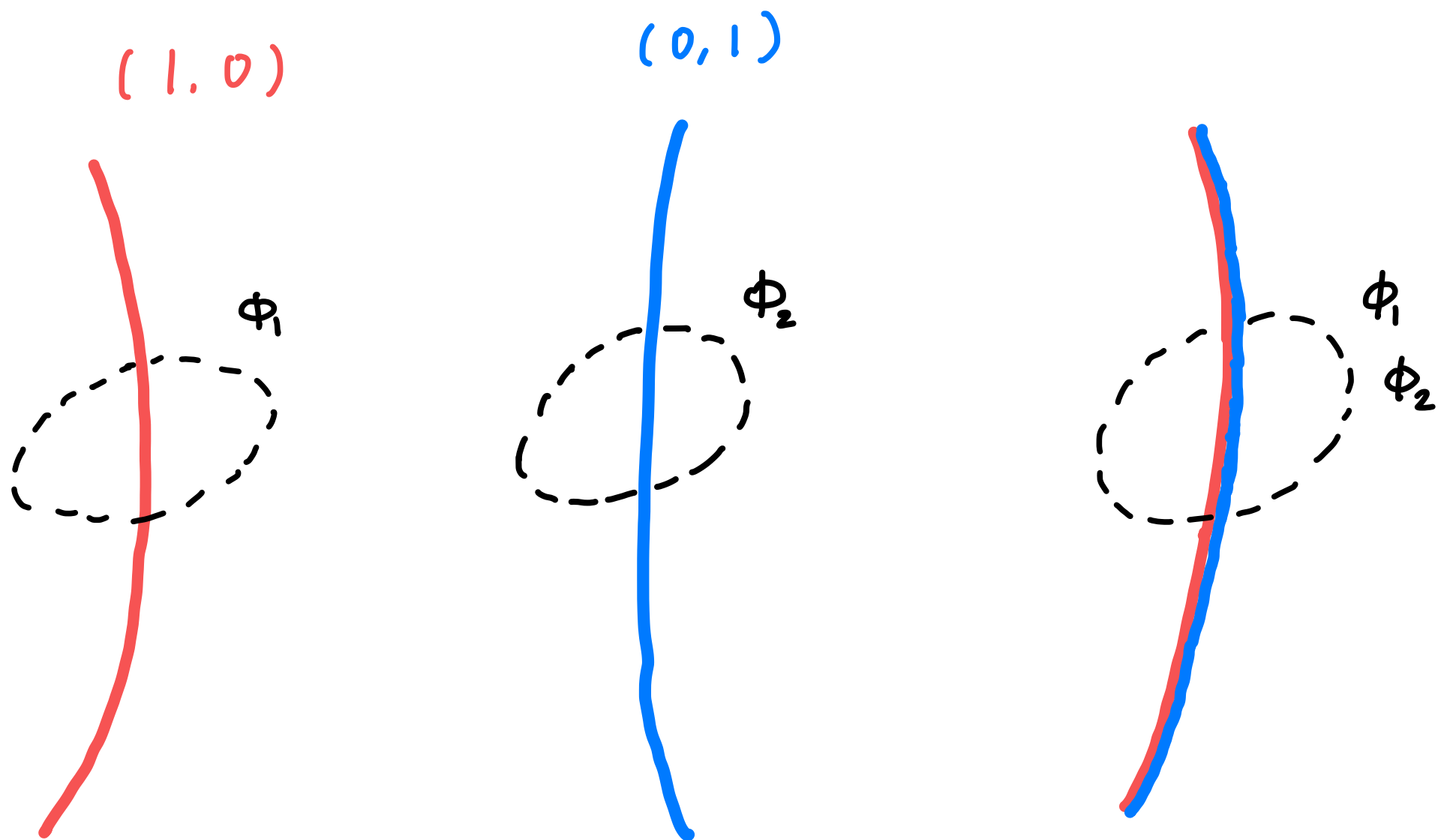
$$\mathcal{L} \supset \frac{g_s^2}{32\pi^2} \frac{a}{f_a} G_{\mu\nu}^a \tilde{G}_{\mu\nu}^a$$



- Barr and Seckel's model
 $Q_{1L} Q_{1R} Q_{2L} Q_{2R}$
color, $U(1)_Z$ and $U(1)_{PQ}$ charges

$$\mathcal{L} = \Phi_1 \bar{Q}_{1L} Q_{1R} + \Phi_2 \bar{Q}_{2L} Q_{2R} + h.c.$$

String Solutions

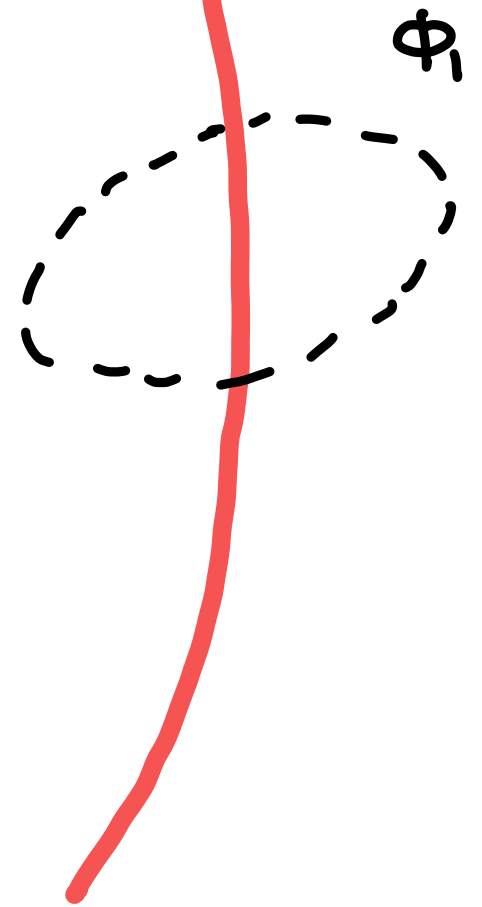


$(1,0)$ strings

(1,0) strings

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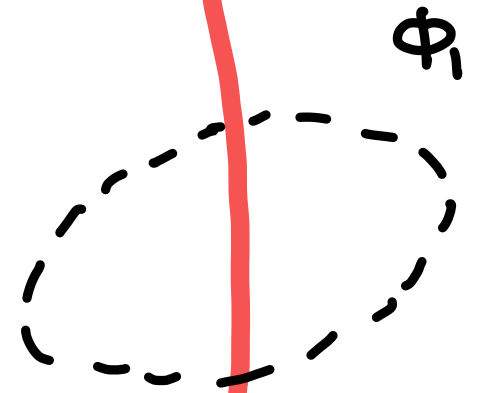
$$\Phi_1 = \frac{1}{\sqrt{2}} v_1 e^{i\theta}, \quad \Phi_2 = \frac{1}{\sqrt{2}} v_2, \quad Z_\mu = c \partial_\mu \theta, \quad r \rightarrow \infty$$



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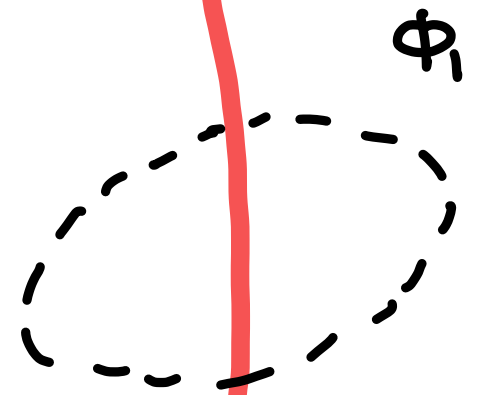
- gradient energy

$$\begin{aligned} \mu_{k,(1,0)} &= \int_0^{2\pi} d\theta \int_\delta^L dr \, r \left(\left| \left(\frac{1}{r} \partial_\theta - ieZ_\theta \right) \Phi_1 \right|^2 + \left| (-ieZ_\theta) \Phi_2 \right|^2 \right) \\ &= \pi \ln\left(\frac{L}{\delta}\right) \left[v_1^2 (1 - ec)^2 + v_2^2 (ec)^2 \right] \end{aligned}$$

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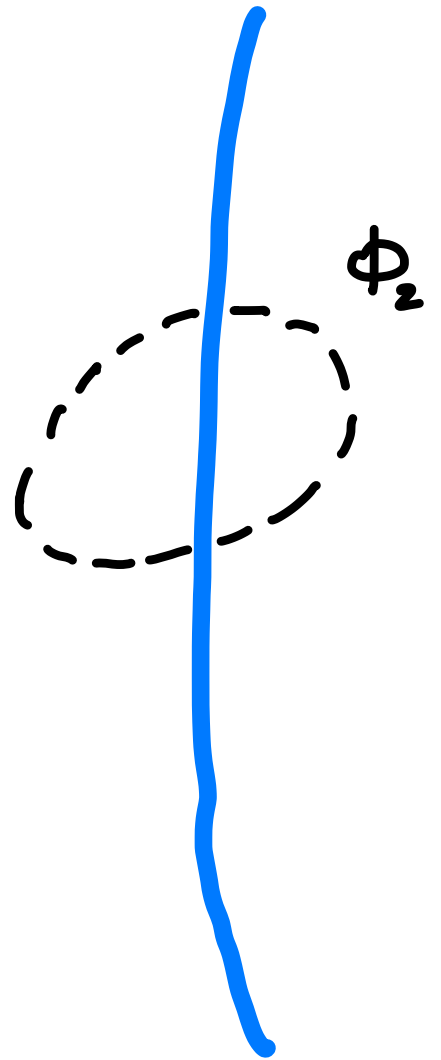
- outside core (minimize it by varying c)

$$\mu_{k,(1,0)} = \pi \frac{v_1^2 v_2^2}{v_1^2 + v_2^2} \ln\left(\frac{L}{\delta}\right) = \pi f_a^2 \ln\left(\frac{L}{\delta}\right)$$

(0,1) strings

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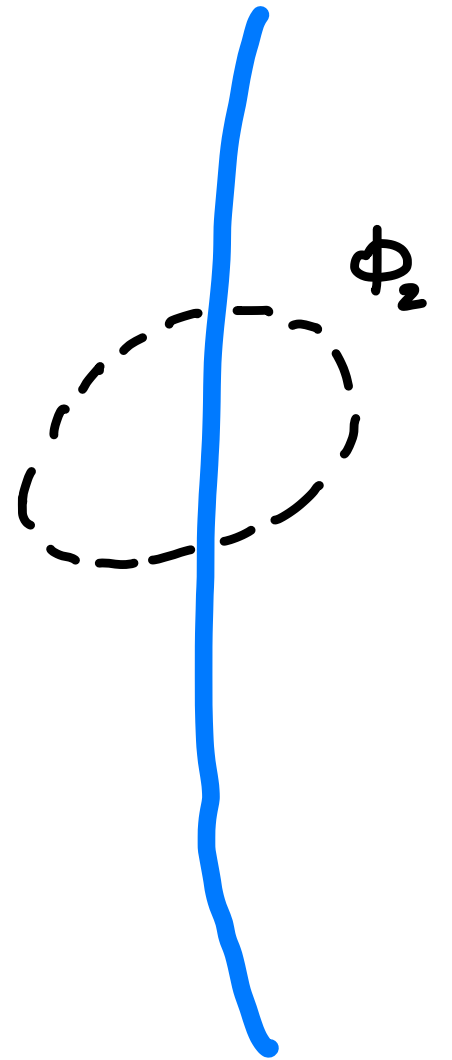
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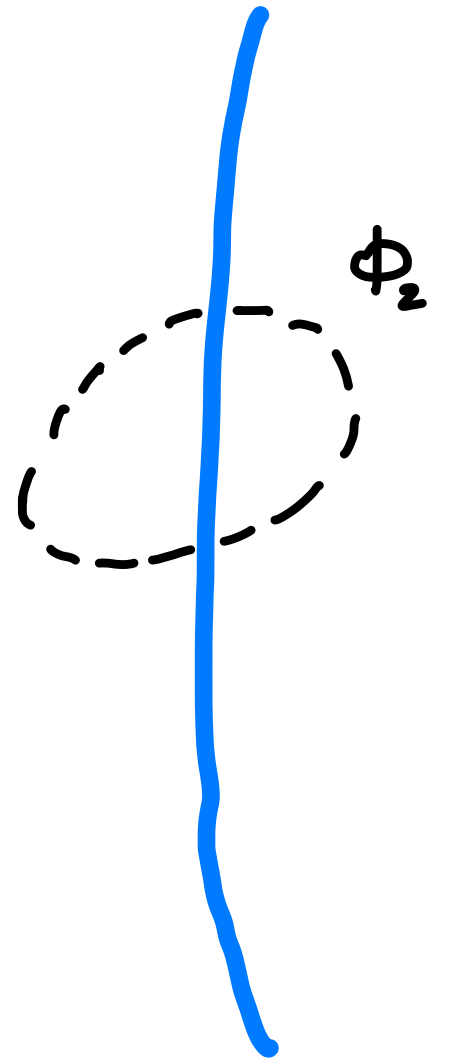
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$$\mu_{k,(0,1)} = \mu_{k,(1,0)}$$

- outside core region

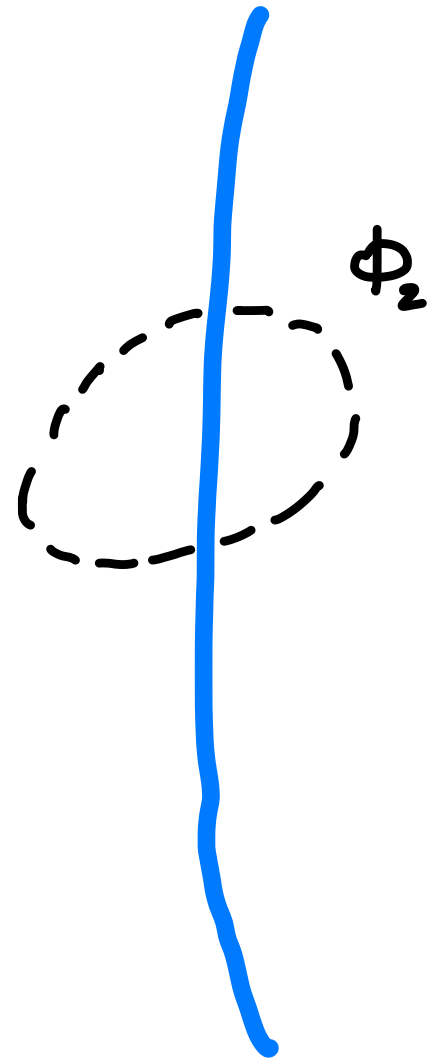
(1,0) string is equivalent to (0,-1) string through a gauge transformation

$$\left(\Phi_1 = \frac{1}{\sqrt{2}}v_1 e^{i\theta}, \Phi_2 = \frac{1}{\sqrt{2}}v_2 \right) \xrightarrow{\alpha_Z \rightarrow \alpha_Z - \theta} \left(\Phi_1 = \frac{1}{\sqrt{2}}v_1, \Phi_2 = \frac{1}{\sqrt{2}}v_2 e^{-i\theta} \right)$$

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$$\mu_{k,(0,1)} = \mu_{k,(1,0)}$$

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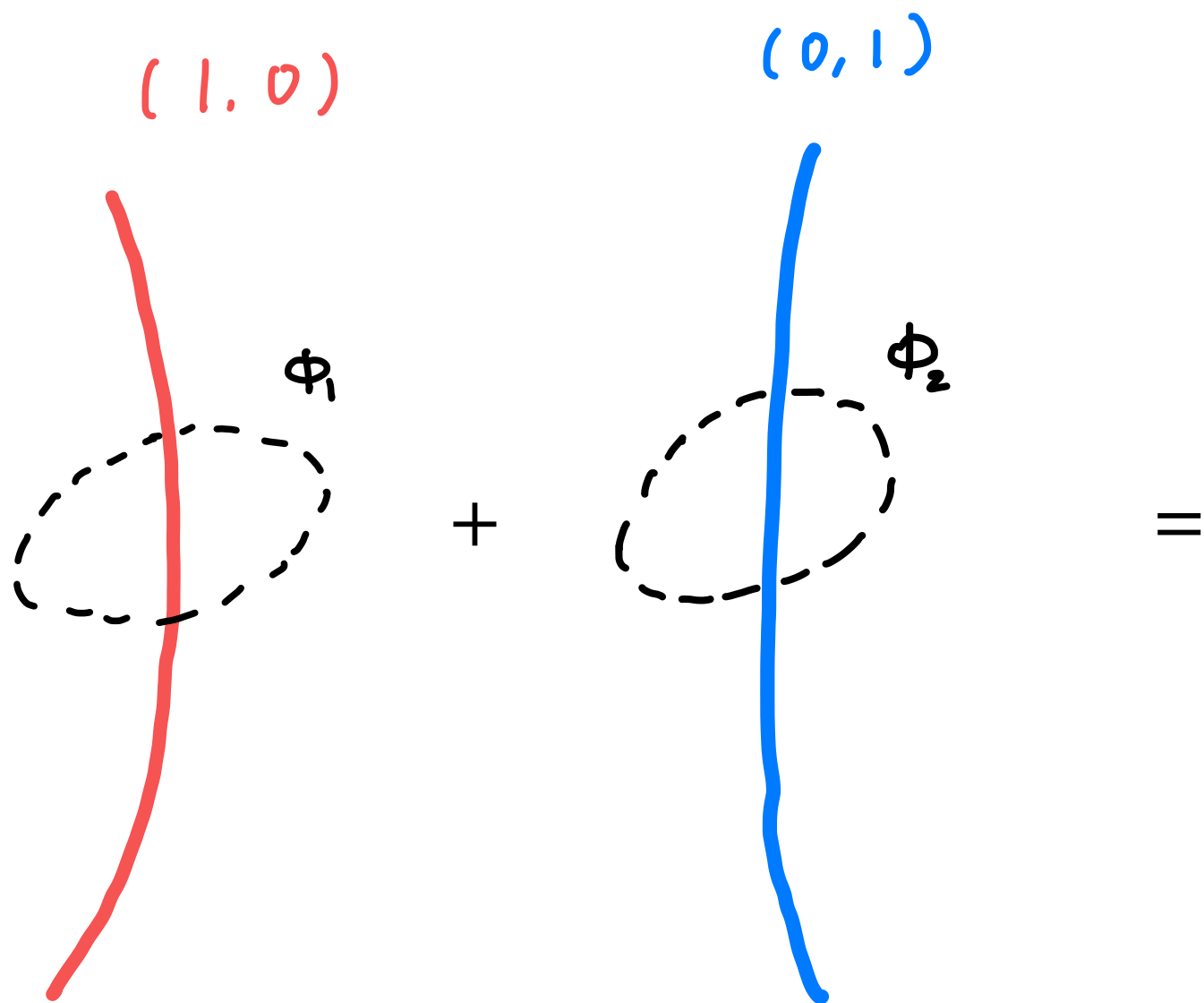
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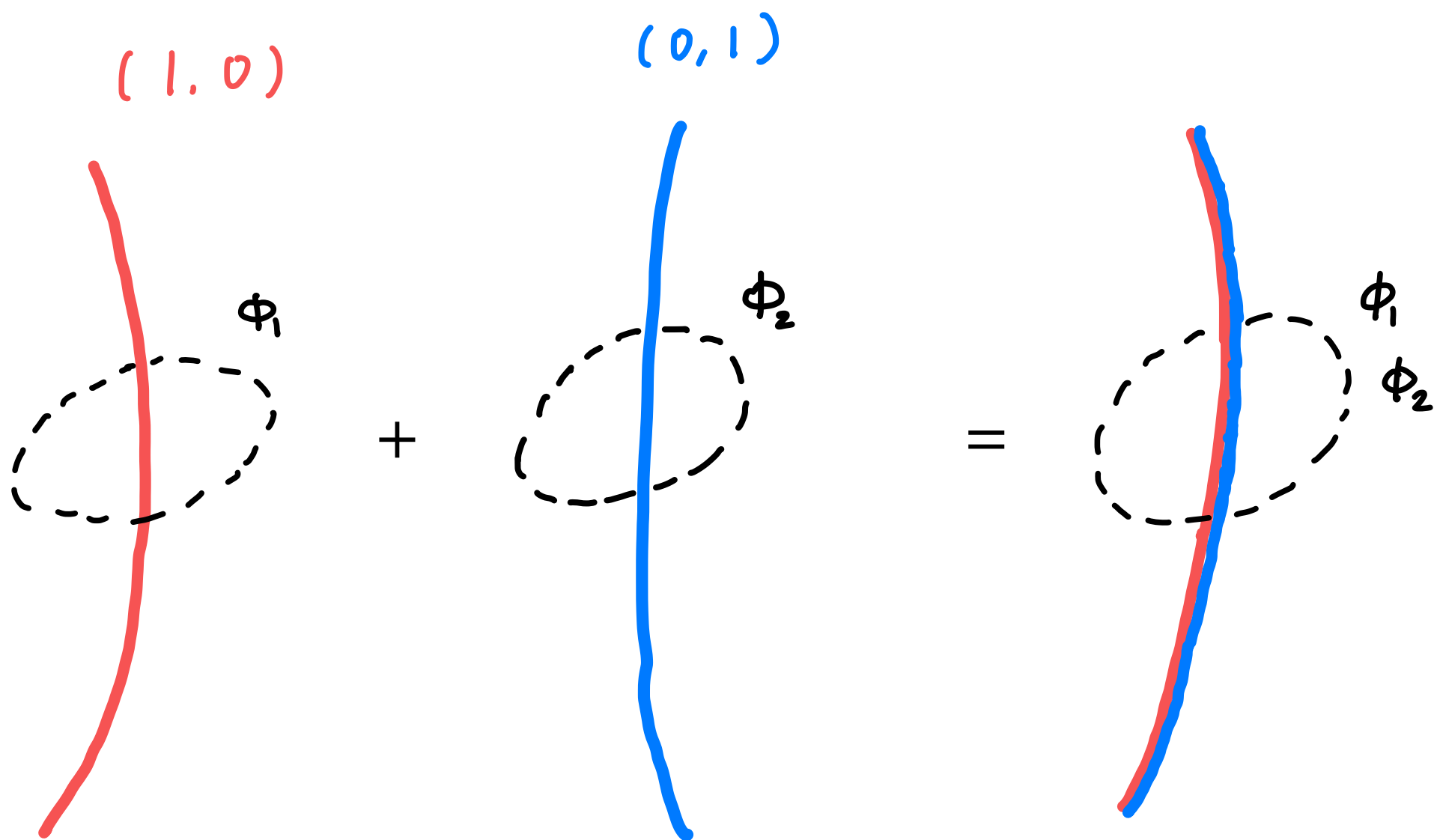
- outside core region

(1,0) can be viewed as an anti-string of (0,1)

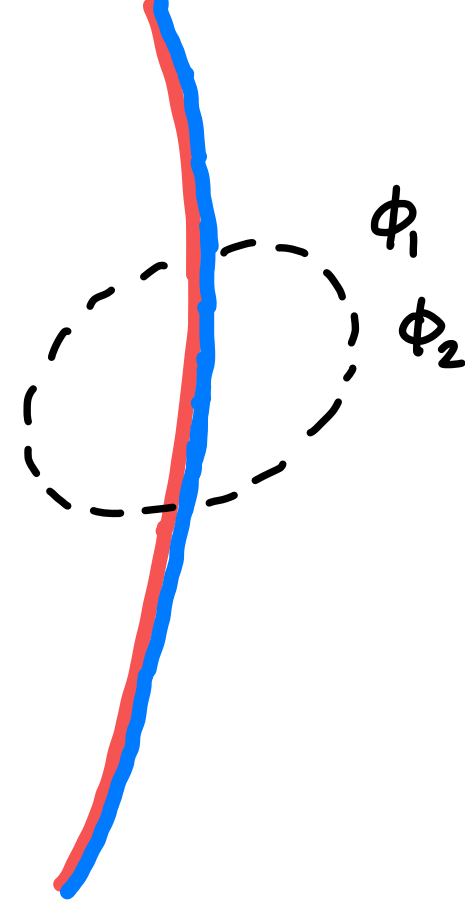
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(1,1) strings

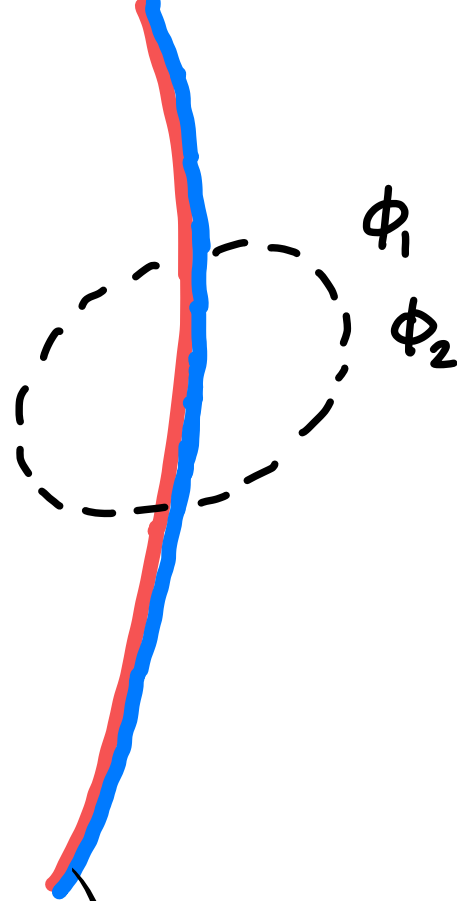


(1,1) strings

- gradient energy of (1,1) string

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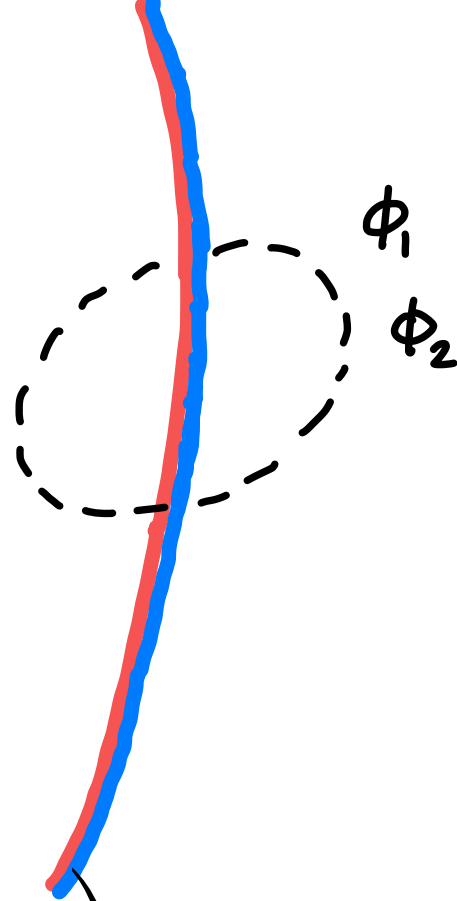
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$$\mu_{k,(1,1)} = 0$$



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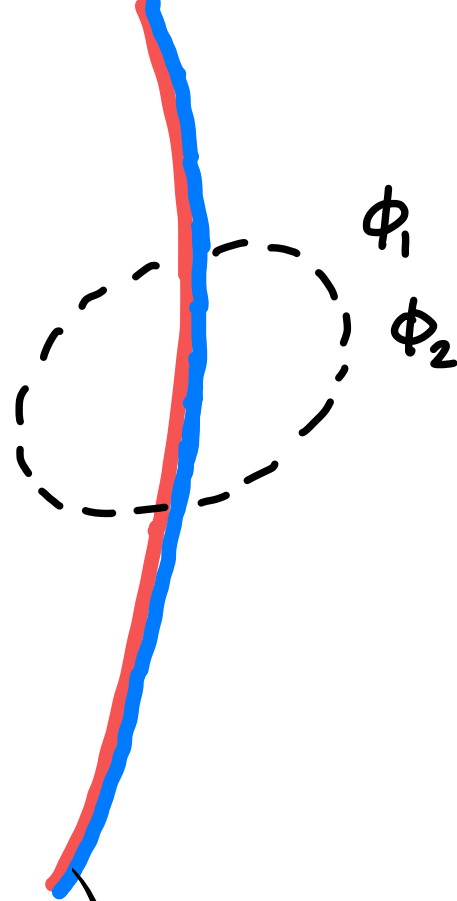
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- (1,1) gauge string

(1,0) and (0,1) global strings



The full tension

The full tension

- magnetic self-energy, scalar potential energy, and gradient energy

The full tension

- magnetic self-energy, scalar potential energy, and gradient energy
- (1,0) string

$$\mu_{(1,0)} \simeq \pi v_1^2 + \pi v_1^2 \ln \left(\frac{m_1}{m_Z} \right) + \pi v_2^2 \ln \left(\frac{m_Z L}{2} \right)$$

(0,1) string

$$\mu_{(0,1)} \simeq \frac{\pi}{2} v_2^2 + \pi v_2^2 \ln \left(\frac{m_2}{m_Z} \right) + \pi v_2^2 \ln \left(\frac{m_Z L}{2} \right)$$

(1,1) string

$$\mu_{(1,1)} = \pi v_1^2 + \pi v_1^2 \ln \left(\frac{m_1}{m_Z} \right) + 0$$

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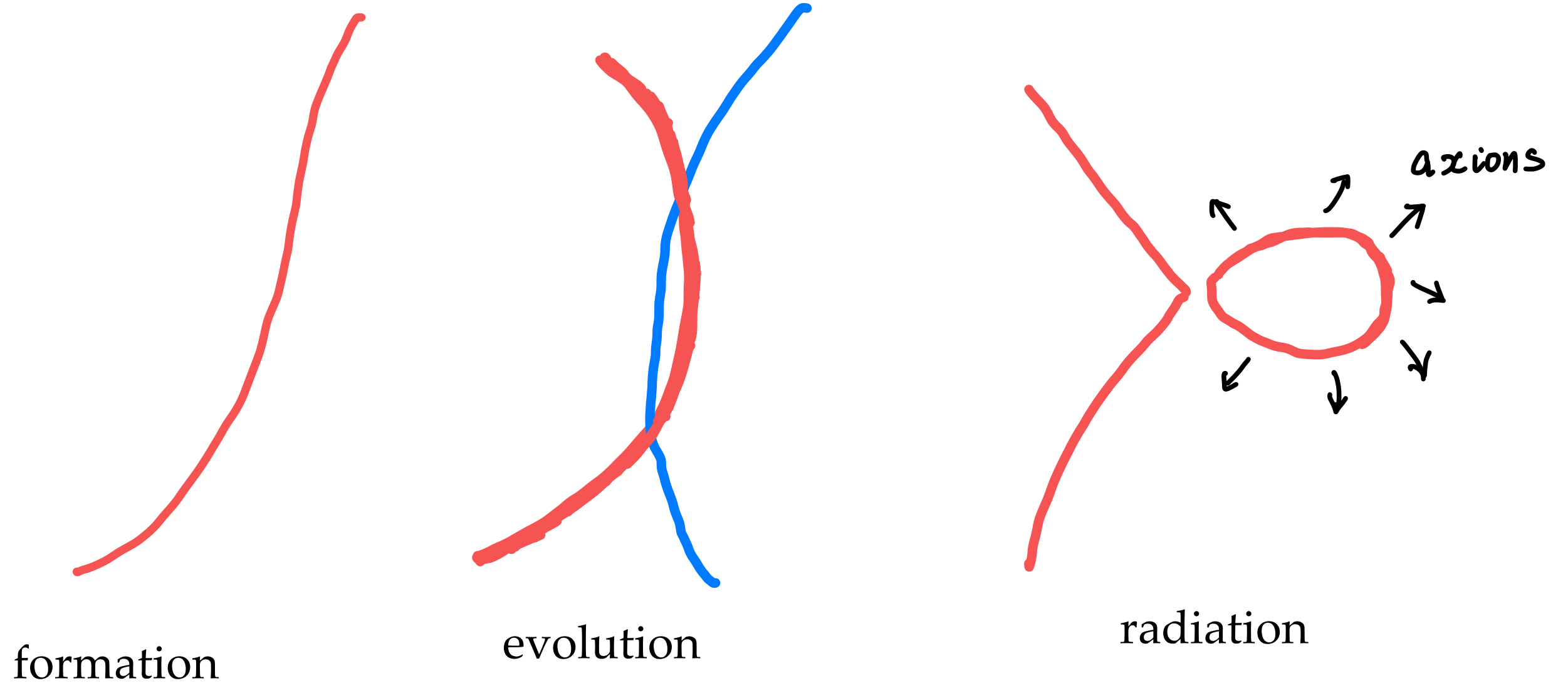
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- heavy core of (1,0) string $\mu_{(1,0)} > \mu_{(0,1)}$

- binding energy of (1,1) string

$$\mu_{(1,0)} + \mu_{(0,1)} - \mu_{(1,1)} = \pi v_2^2 \left[2 \ln \left(\frac{m_Z L}{2} \right) - 1 \right]$$

Cosmological Implication



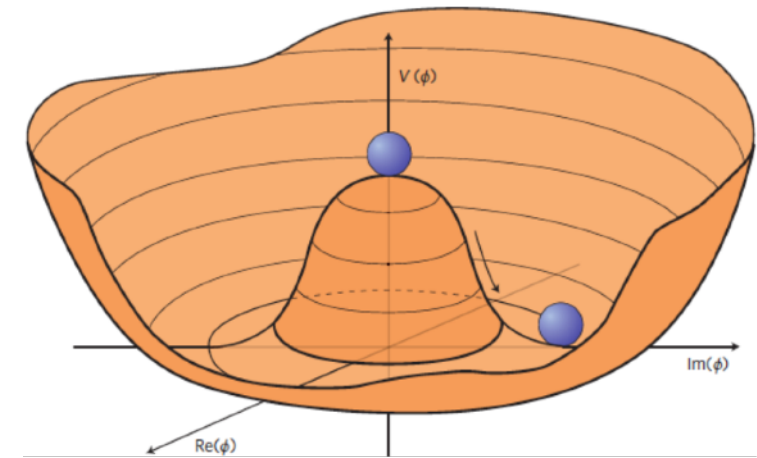
First phase transition and string formation

First phase transition and string formation

- consider $v_1 \gg v_2$

first phase transition

$$\langle \Phi_1(x) \rangle = \frac{v_1}{\sqrt{2}} \text{ and } \langle \Phi_2(x) \rangle = 0$$

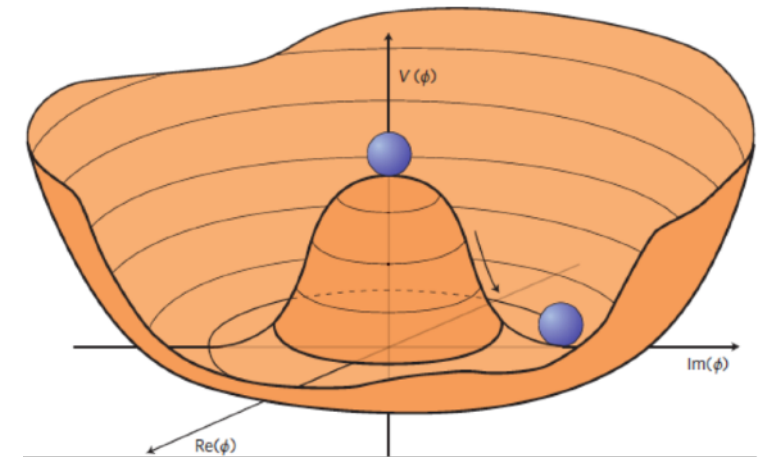


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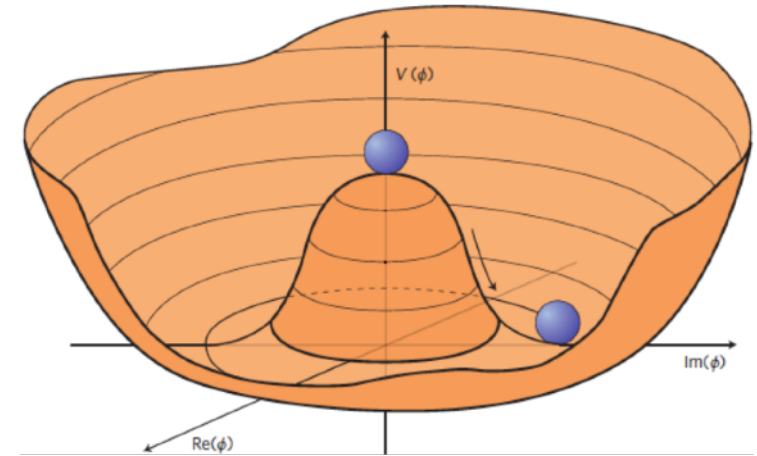
- string formation, the correlation length $\sim 1/v_1$

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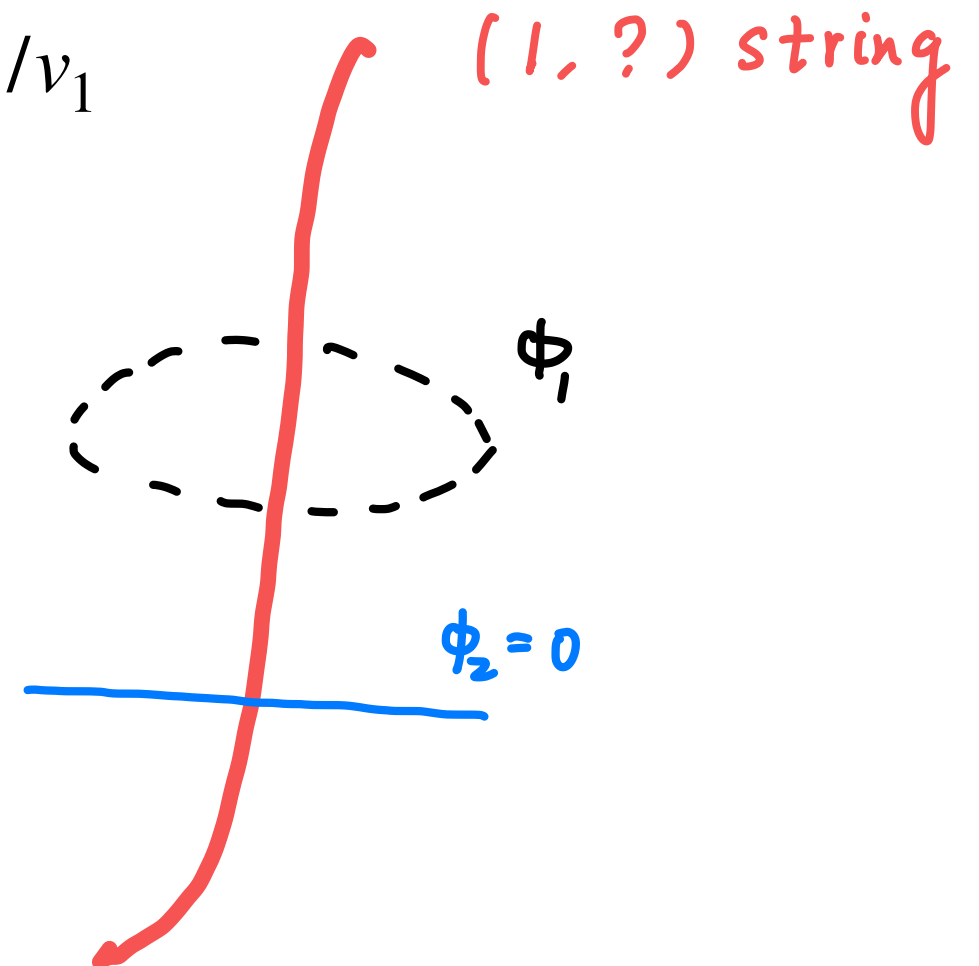
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- string formation, the correlation length $\sim 1/v_1$

- U(1) gauge strings form
(1, **n**) string



Second phase transition

Second phase transition

- second phase transition

$$\langle \Phi_1(x) \rangle = \frac{v_1}{\sqrt{2}} \text{ and } \langle \Phi_2(x) \rangle = \frac{v_2}{\sqrt{2}}$$

Second phase transition

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Second phase transition

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- string formation, the correlation length $\sim 1/v_2$
- $(0,1)$ strings form via Kibble mechanism

$(1,n)$ string in the second phase transition

- $(1, \mathbf{n})$ string $\rightarrow ?$

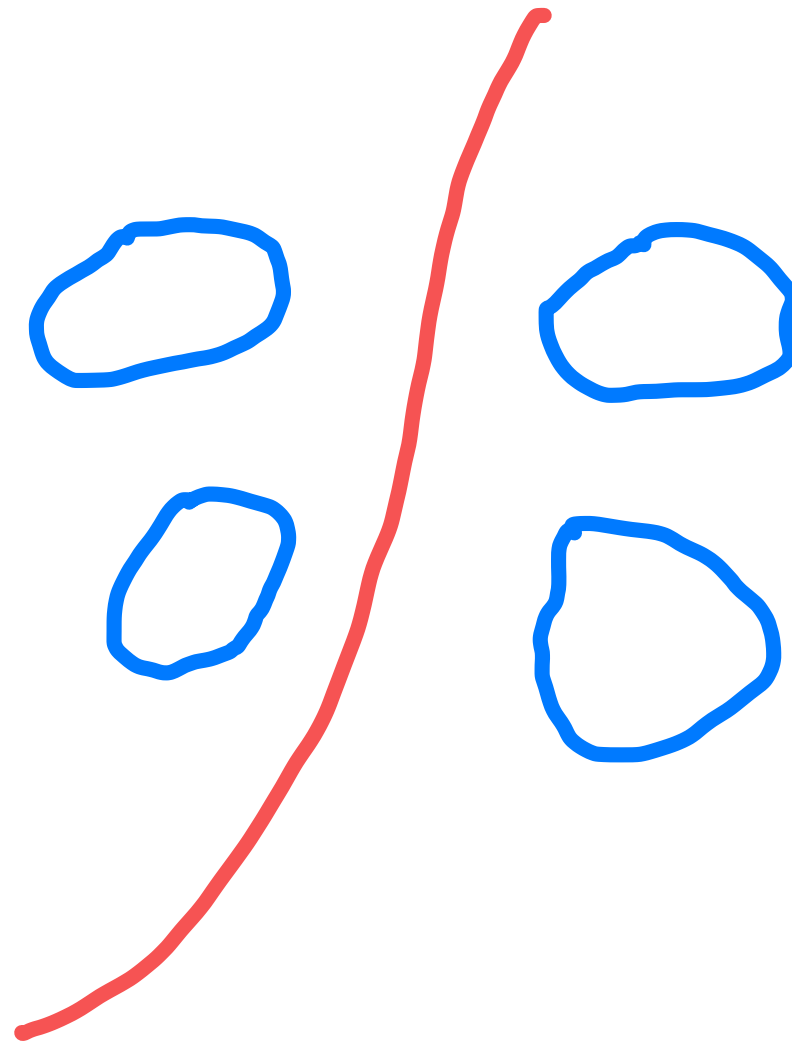
$(1,n)$ string in the second phase transition

- $(1, \textcolor{red}{n})$ string $\rightarrow ?$



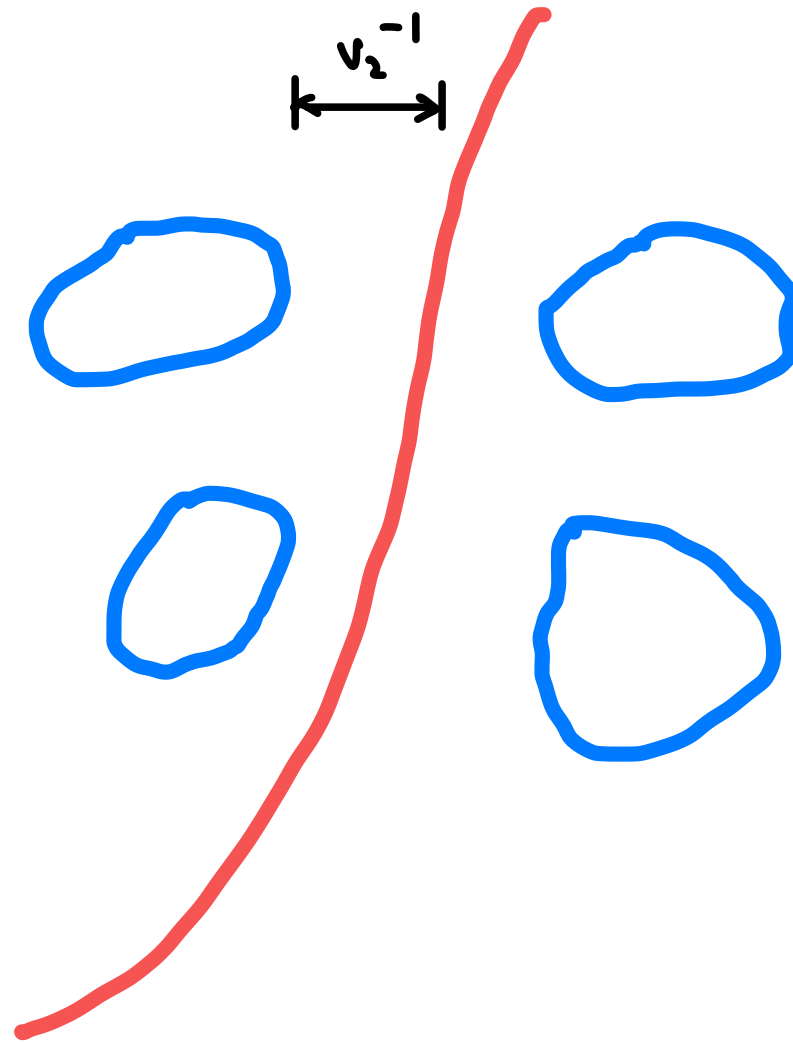
$(1,n)$ string in the second phase transition

- $(1, \textcolor{red}{n})$ string $\rightarrow ?$



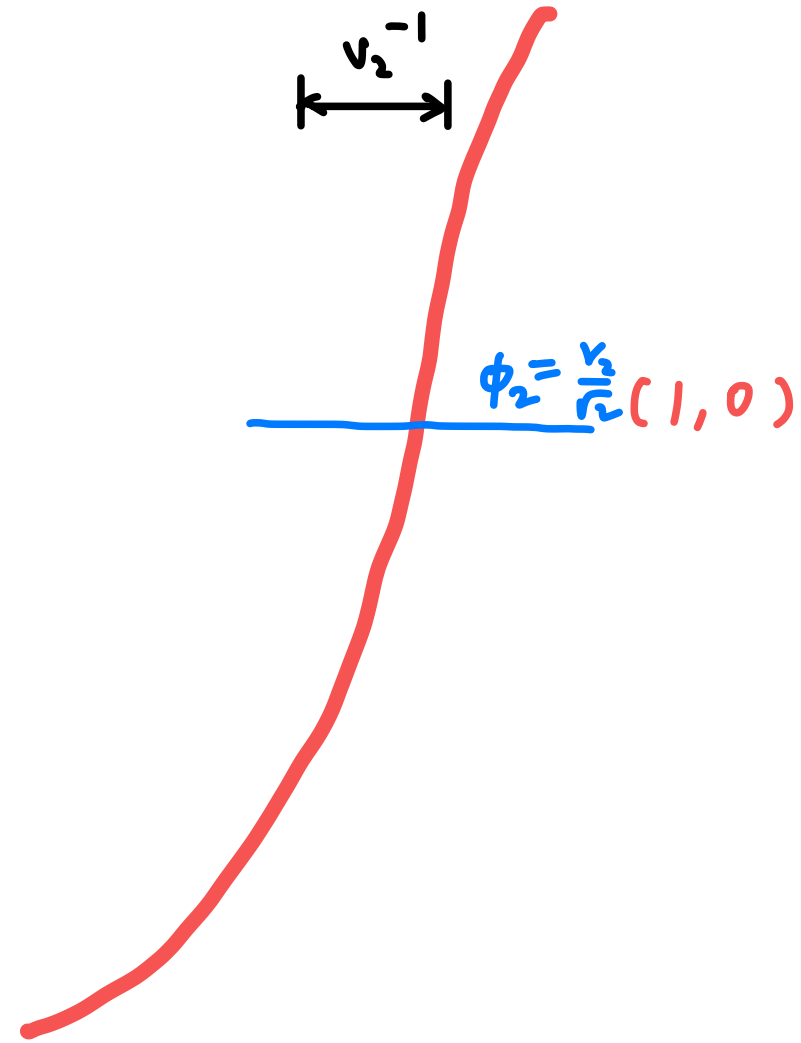
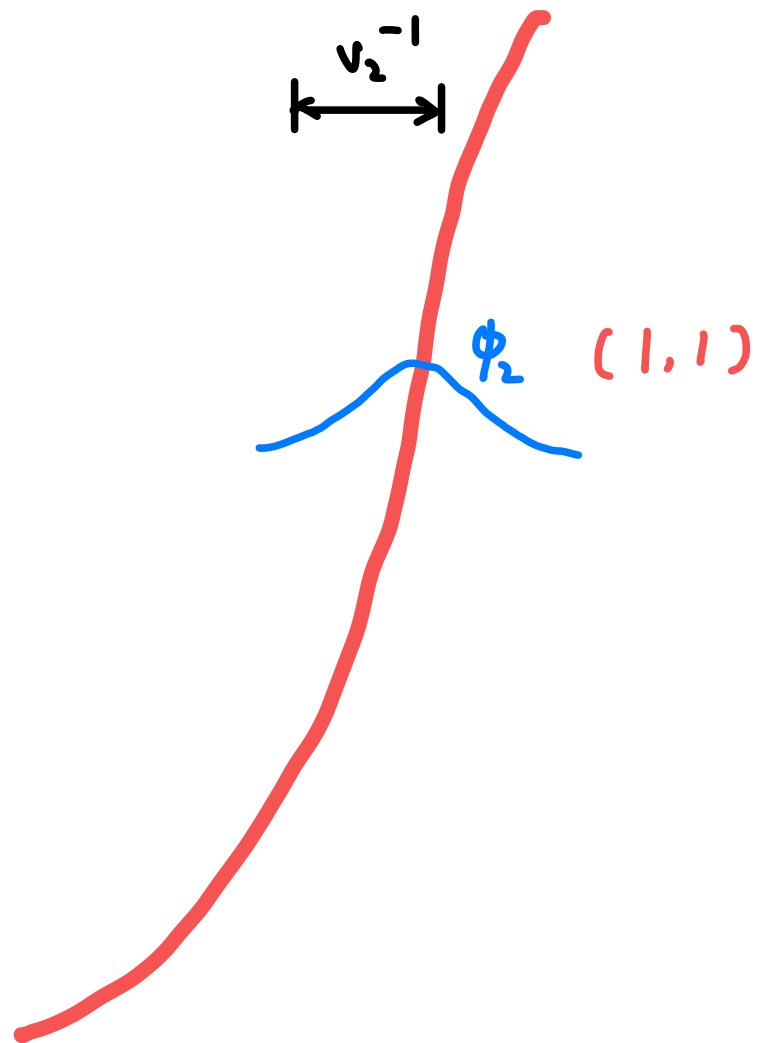
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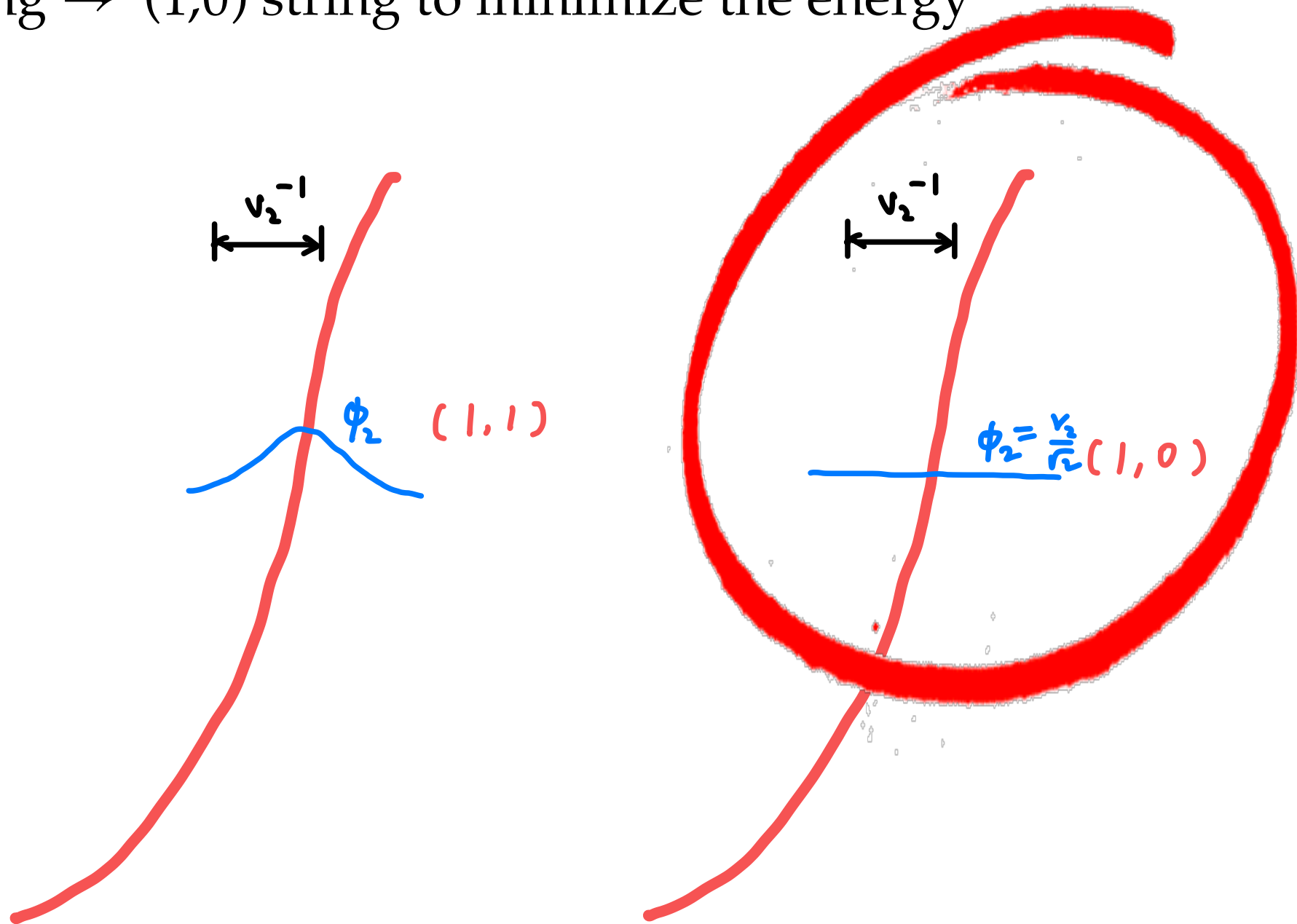
$(1,n)$ string in the second phase transition

- $(1, n)$ string $\rightarrow ?$



(1,n) string in the second phase transition

- (1, **n**) string \rightarrow (1,0) string to minimize the energy

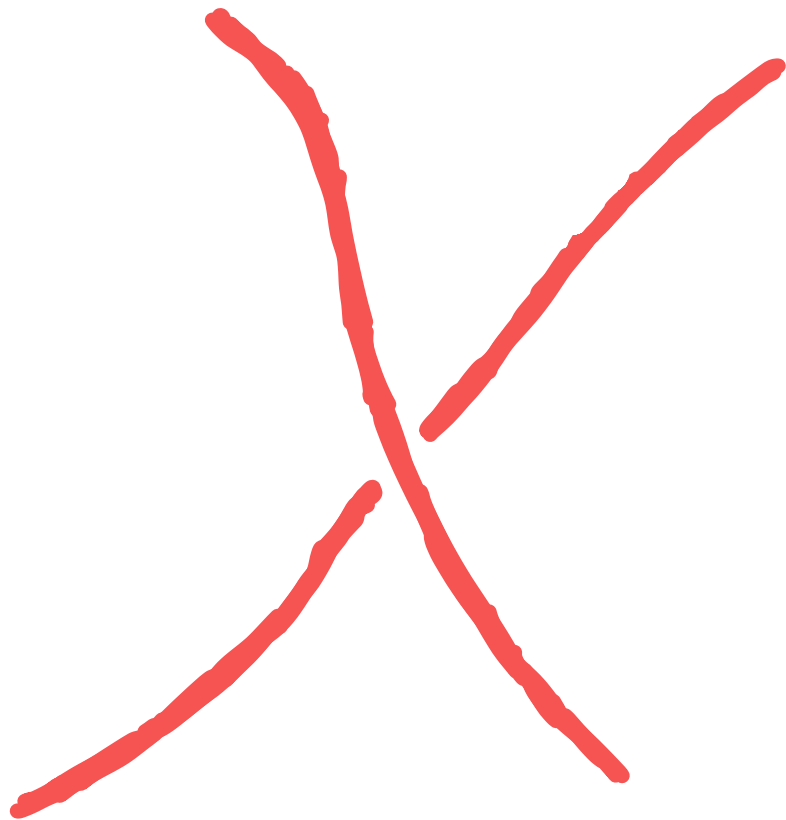


String network evolution

- $(1,0)$ string encounters $(1,0)$ string

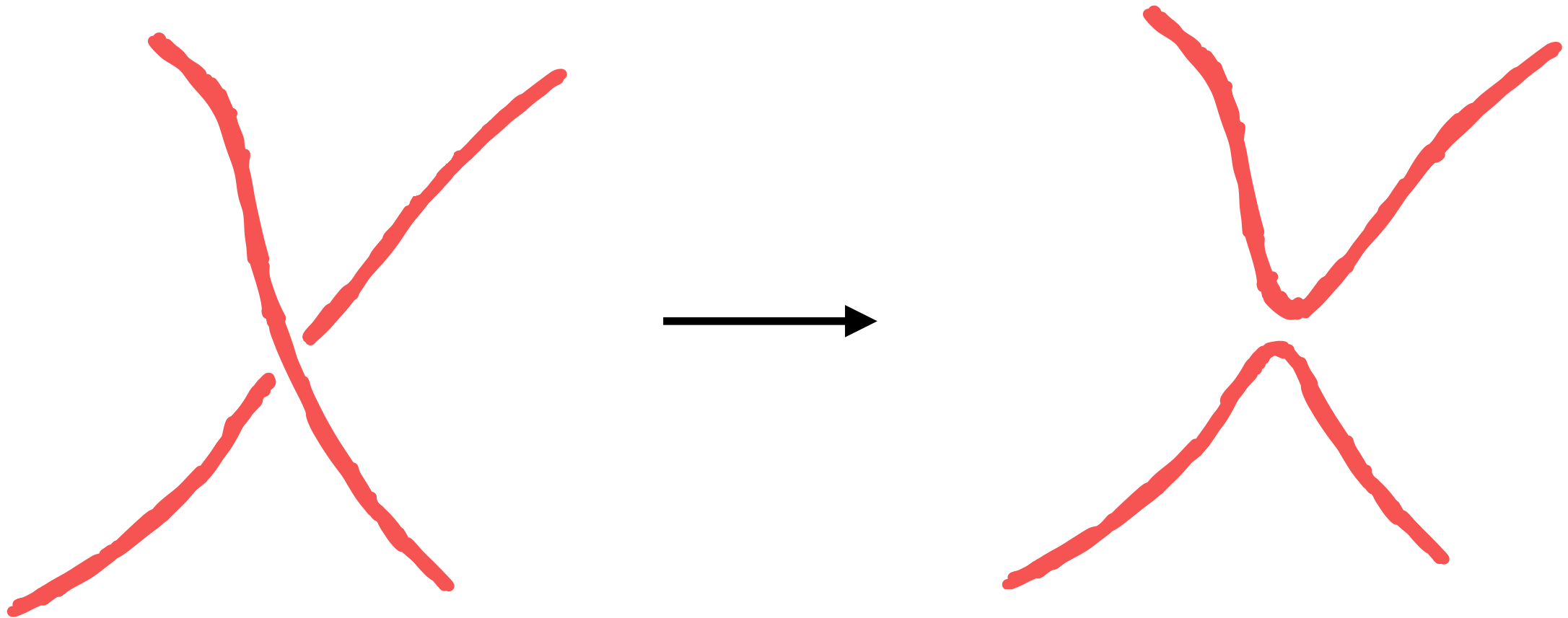
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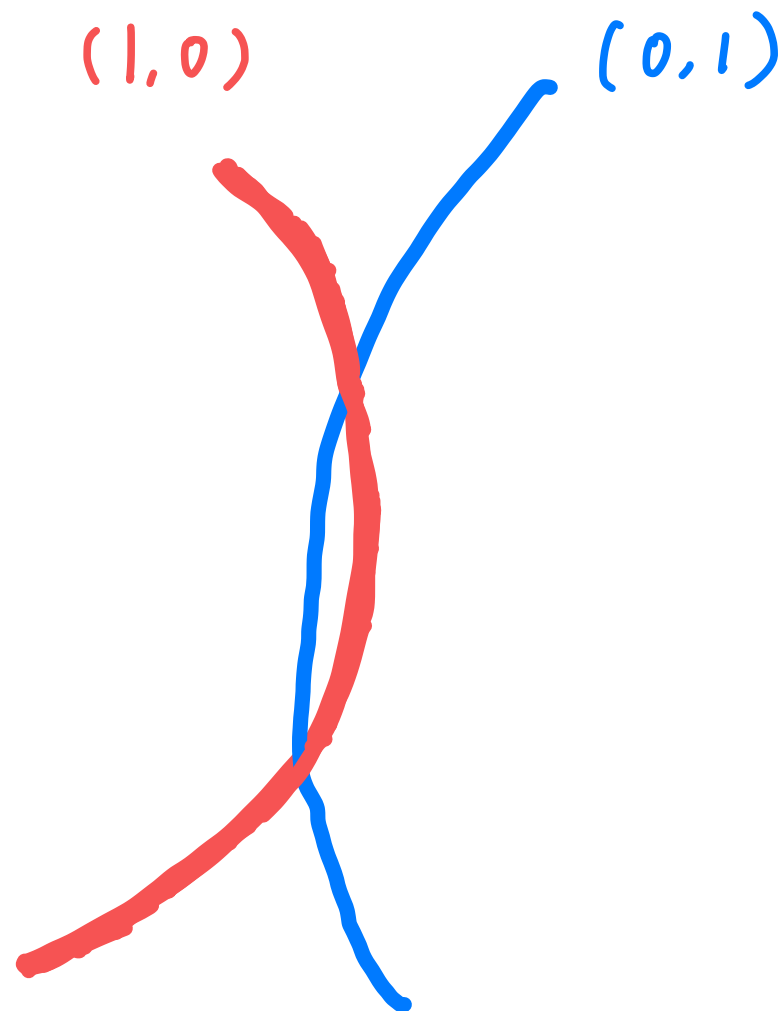
String network evolution

- $(1,0)$ string encounters $(0,1)$ string \rightarrow $(1,1)$ bound state
Y-junctions

String network evolution

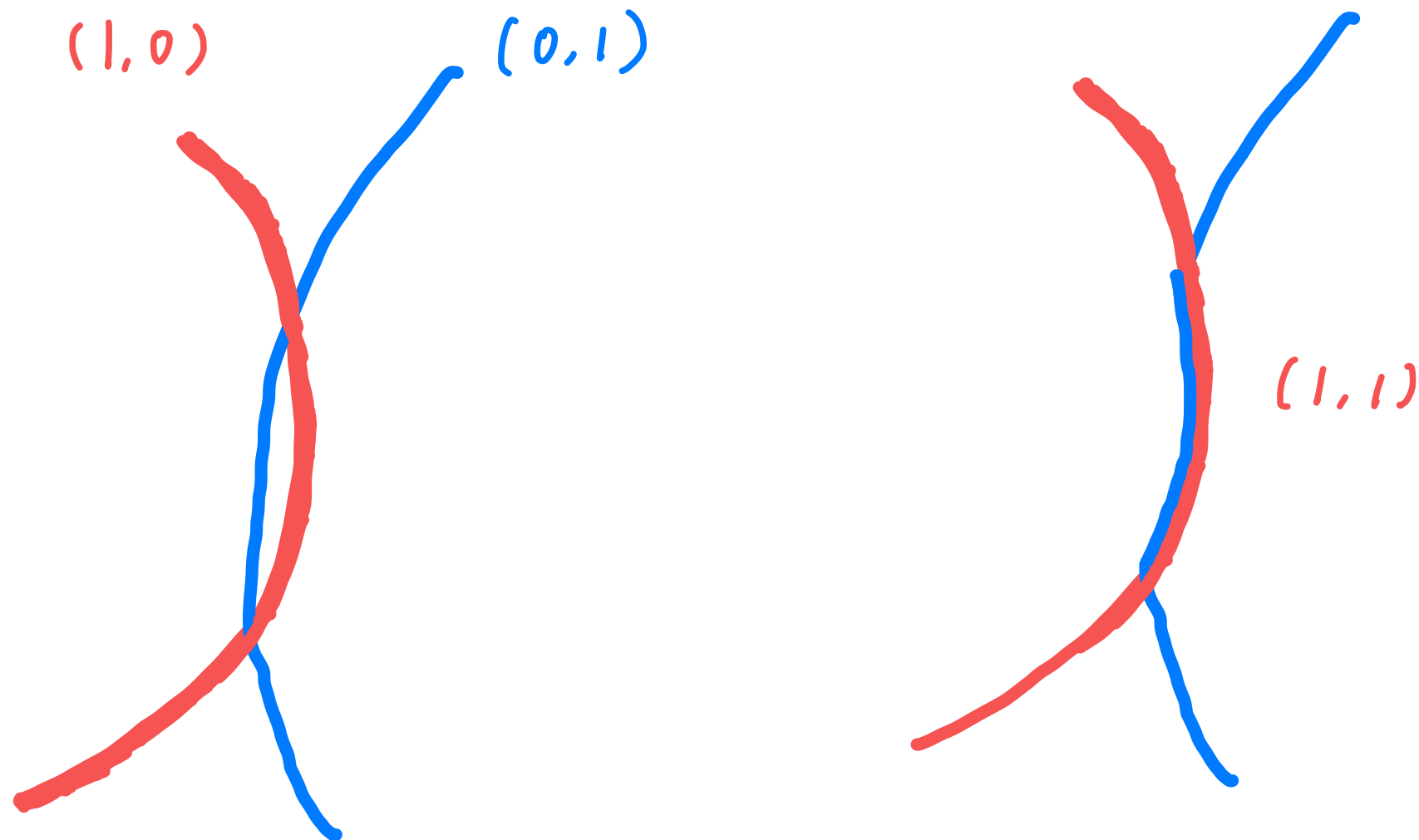
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Y-junctions



String network evolution

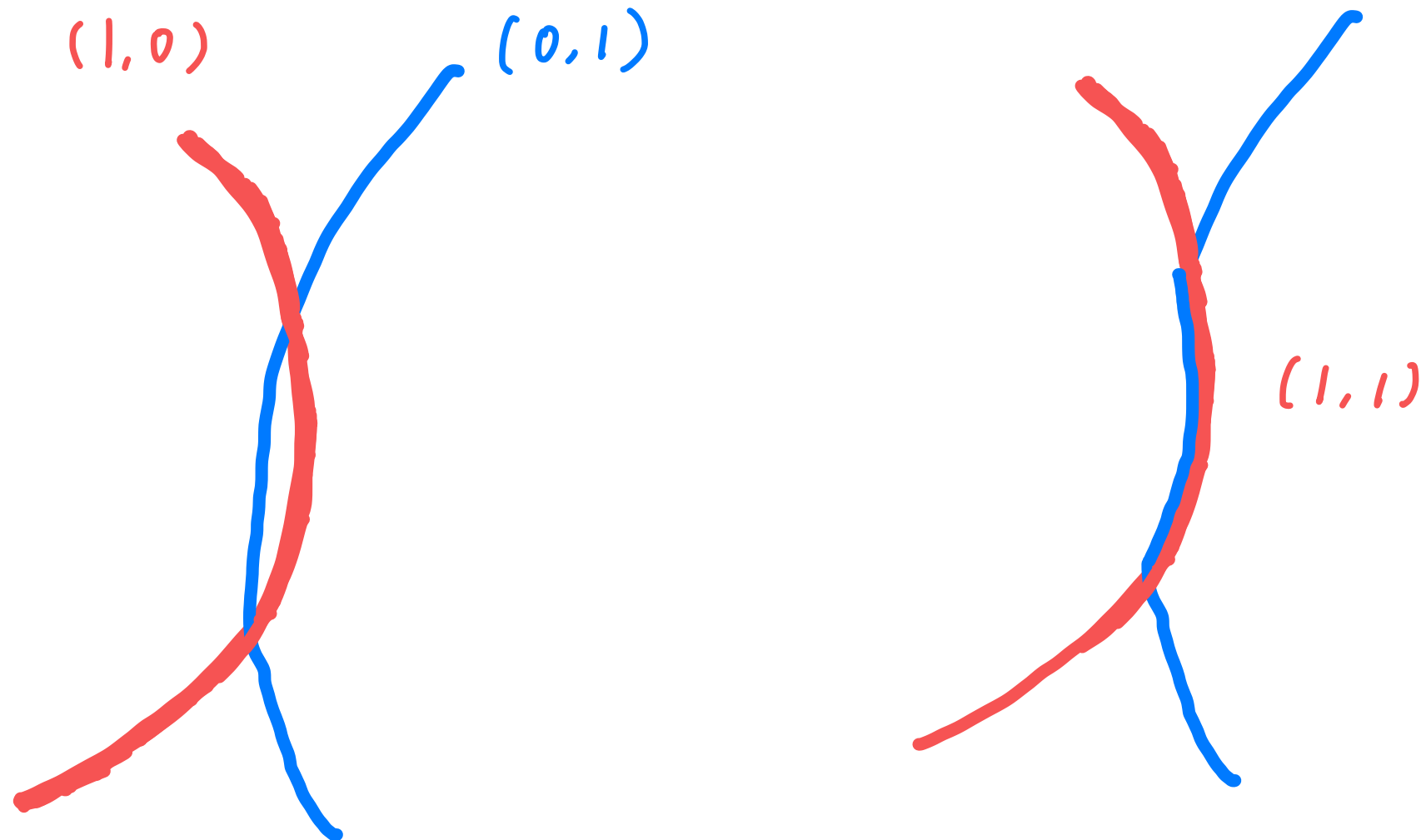
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String network evolution

- $(1,0)$ string encounters $(0,1)$ string \rightarrow $(1,1)$ bound state

Y-junctions



- Other works on simulations of Y-junctions
found 1) some fraction of Y-junctions remain
2) scaling solution

Urrestilla, Vilenkin JHEP(2008)

Rajantie, Skellariodou, Stoica, JCAP (2007)

Copeland, Saffin JHEP (2005)

...

QCD axion

QCD axion

- dark matter abundance

misalignment + string radiation + domain wall collapse

$$\rho_{a,0} = \rho_a^{\text{vac}}(t_0) + m_a n_a^{\text{str}}(t_0) + m_a n_a^{\text{DW}}(t_0)$$

QCD axion

- dark matter abundance

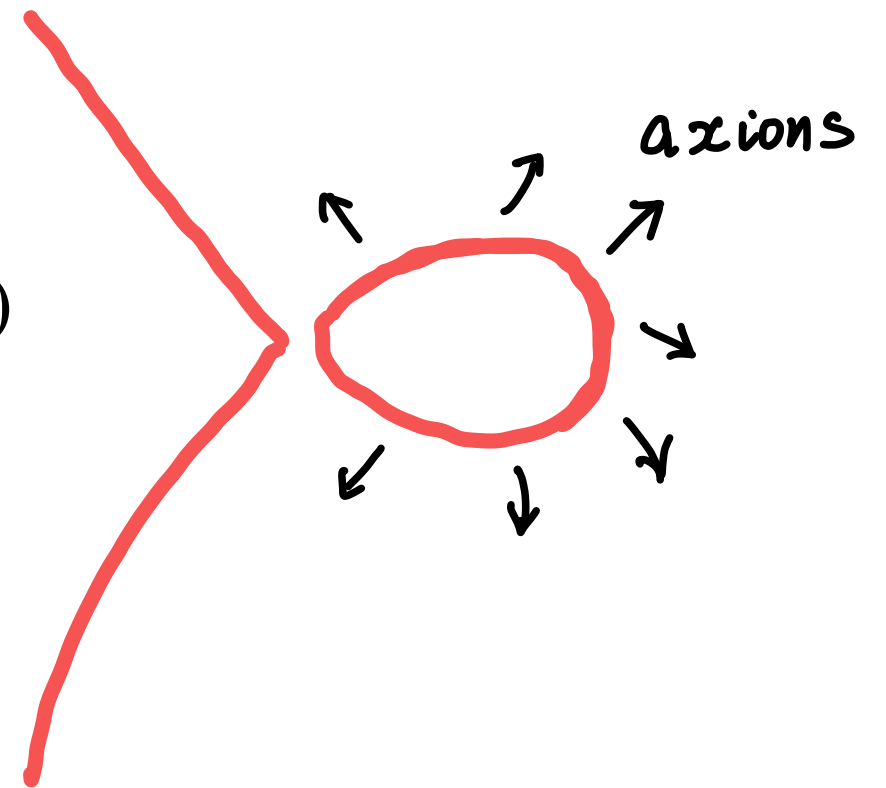
misalignment + **string radiation** + domain wall collapse

$$\rho_{a,0} = \rho_a^{\text{vac}}(t_0) + \textcolor{red}{m_a n_a^{\text{str}}(t_0)} + m_a n_a^{\text{DW}}(t_0)$$

- uncertainty from string radiation

Scenario A: IR spectrum $\frac{dE}{d\omega} \propto \delta(\omega - 2\pi t^{-1})$

Scenario B: flat spectrum $\frac{dE}{d\omega} \propto \frac{1}{\omega}$



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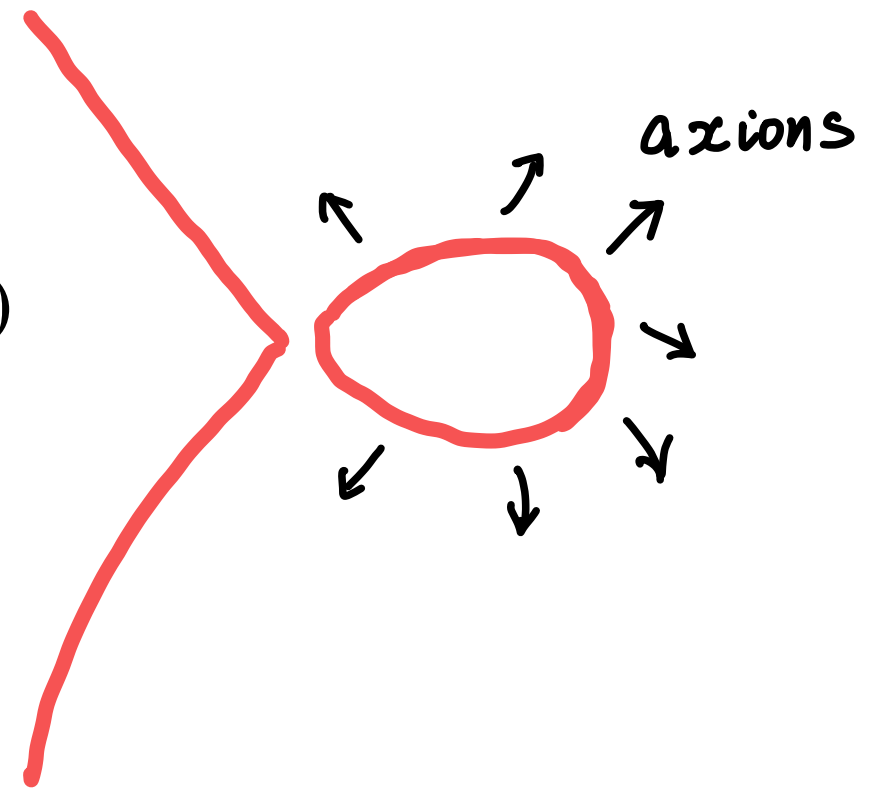
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- $\textcolor{red}{n_a^{\text{str}}} \propto \frac{1}{\langle \omega \rangle}$



Gauged global string

Gauged global string

- (0,1) string tension is the same as a standard QCD axion string
- **heavy core** of (1,0) string

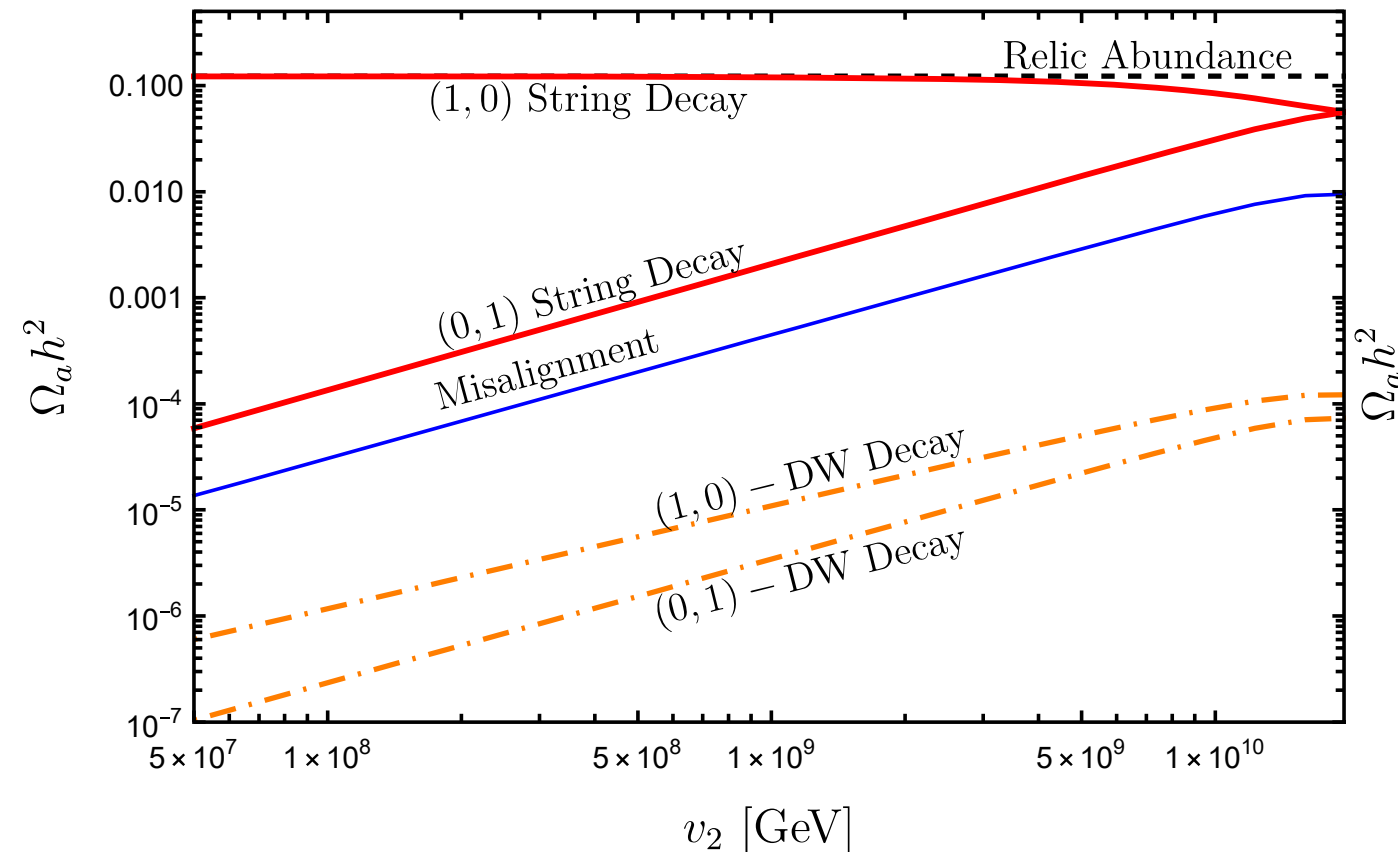
$$\mu_{(1,0)}(t) \simeq \pi v_1^2 \ln \left(\frac{m_1}{m_Z} \right) + \pi f_a^2 \ln \left(\frac{m_Z t}{2} \right)$$

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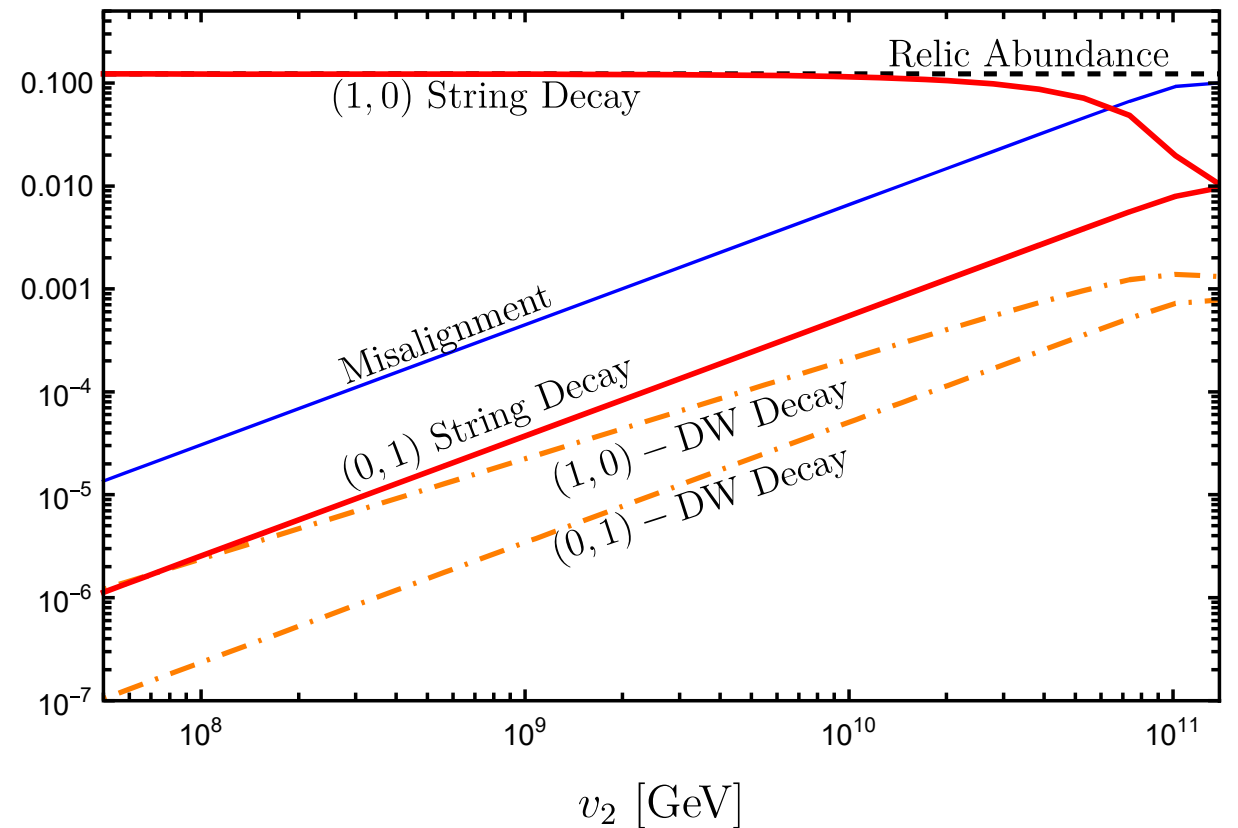
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Scenario A



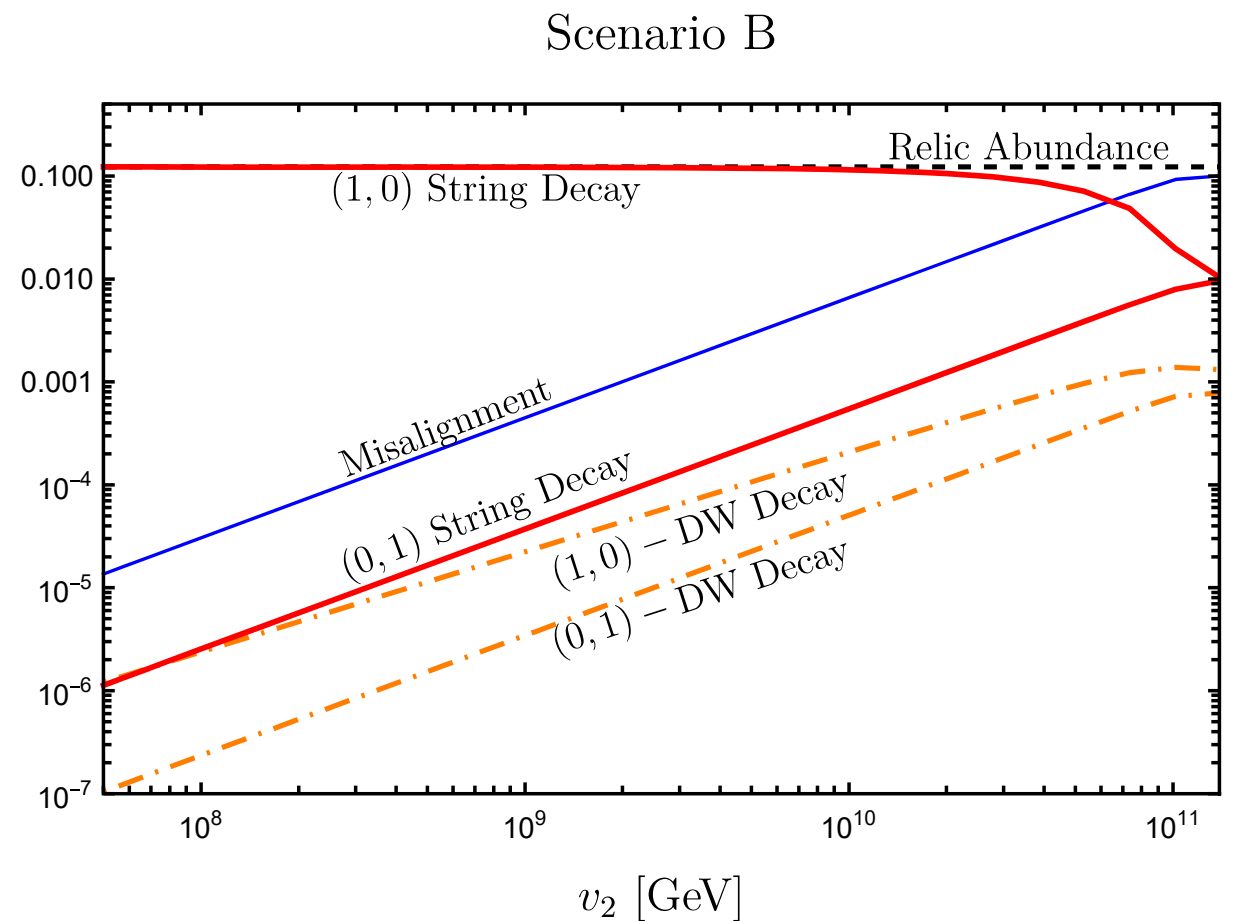
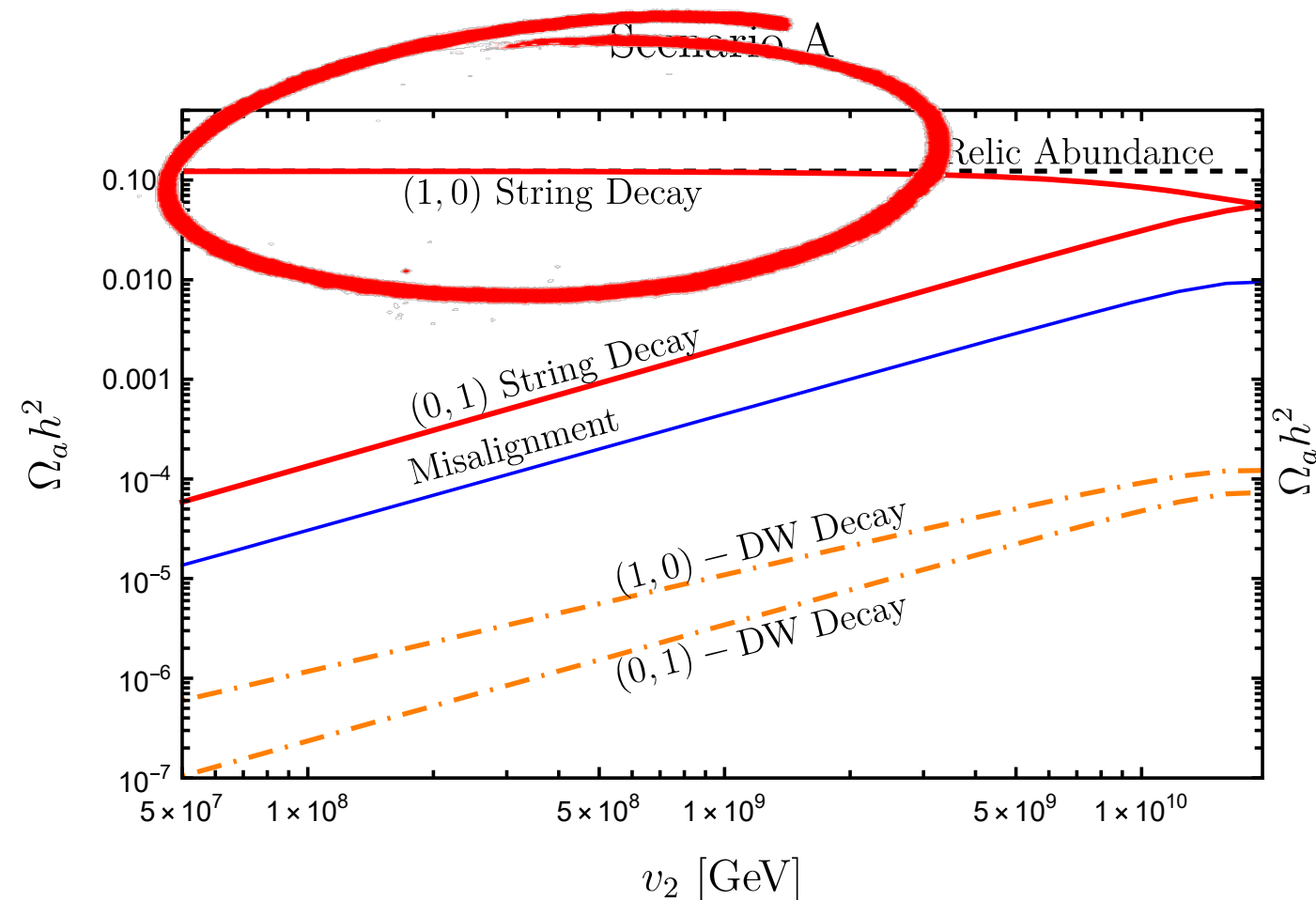
Scenario B



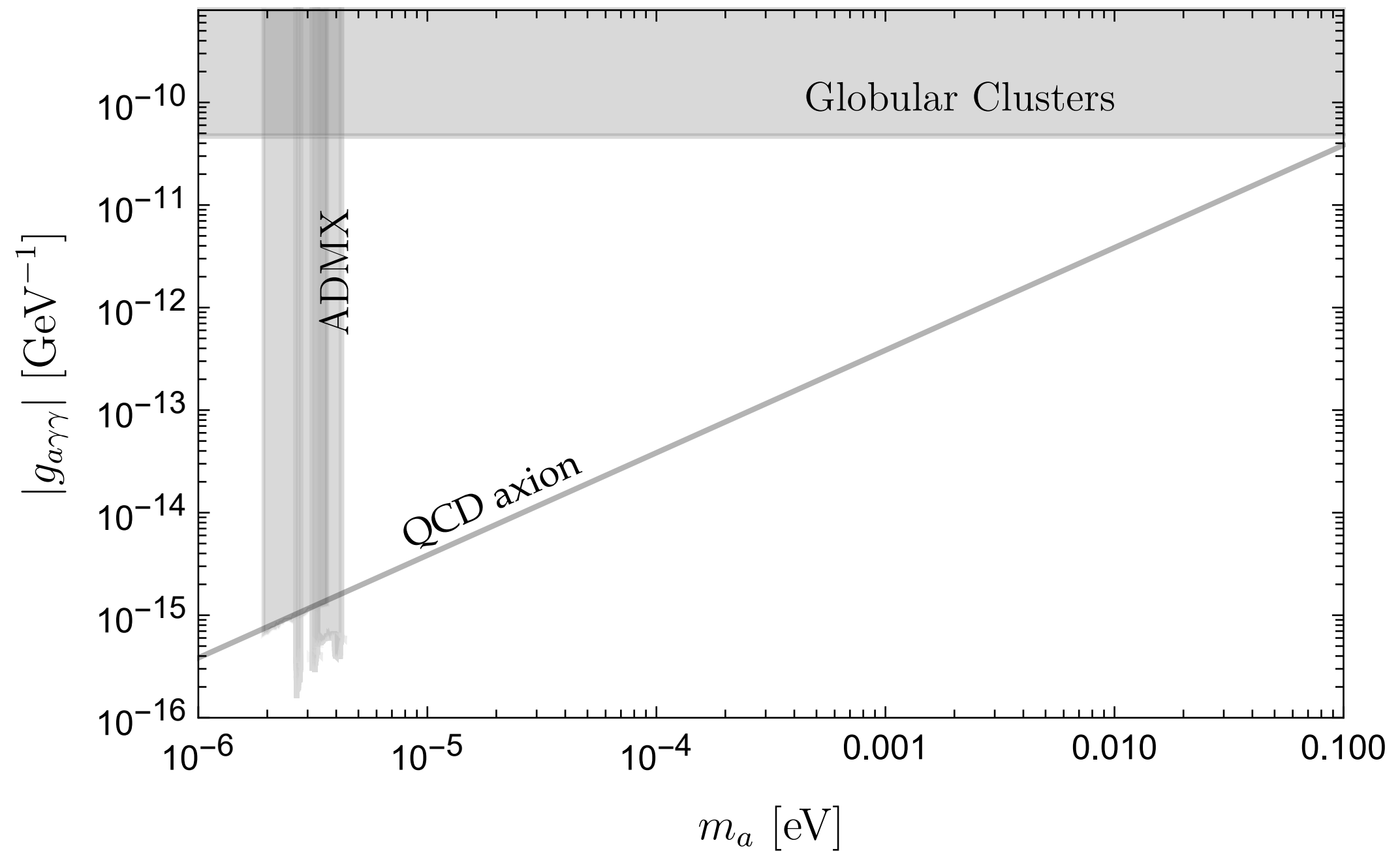
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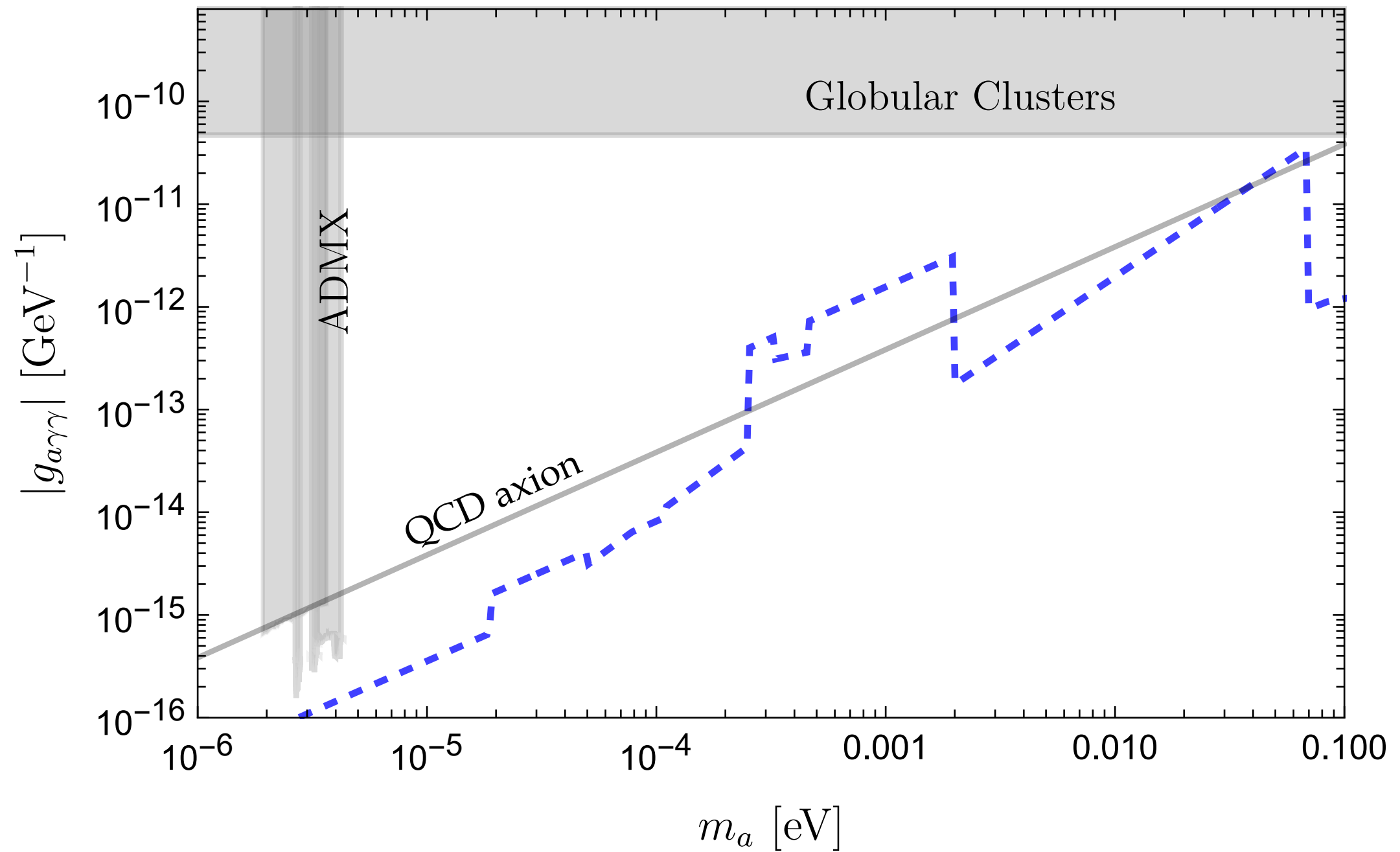
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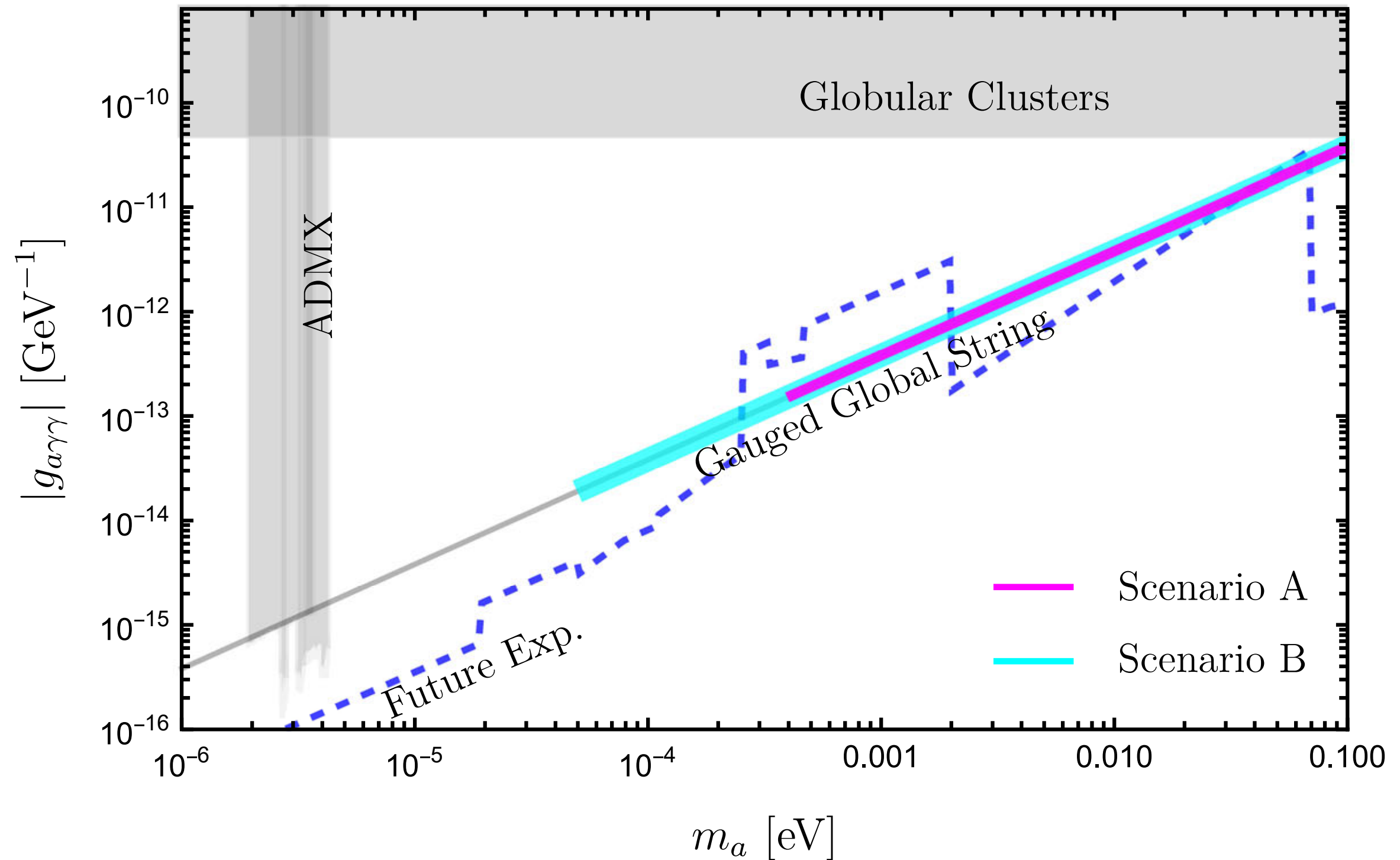
QCD axion window



Future explorations



Gauged global string



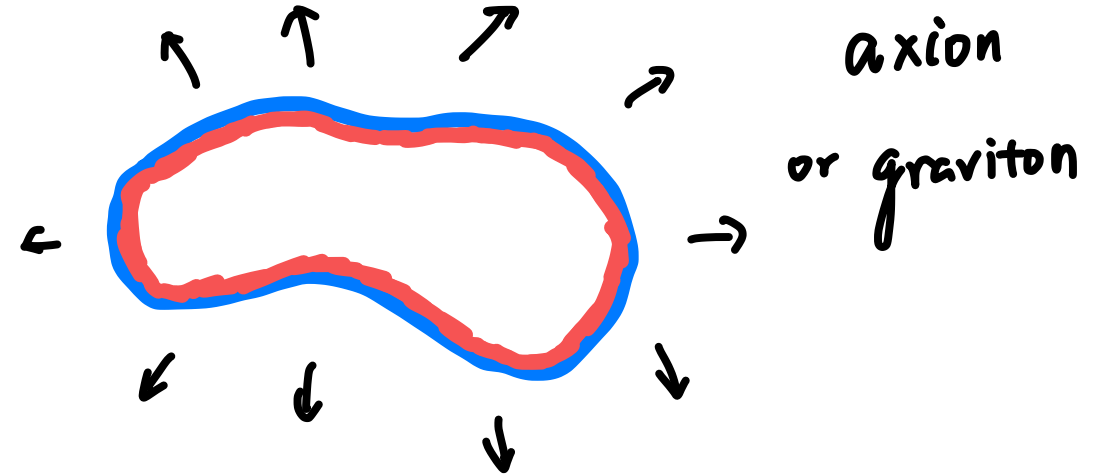
Gauge string (1,1) radiation

Gauge string (1,1) radiation

- gauge strings radiate gravitons

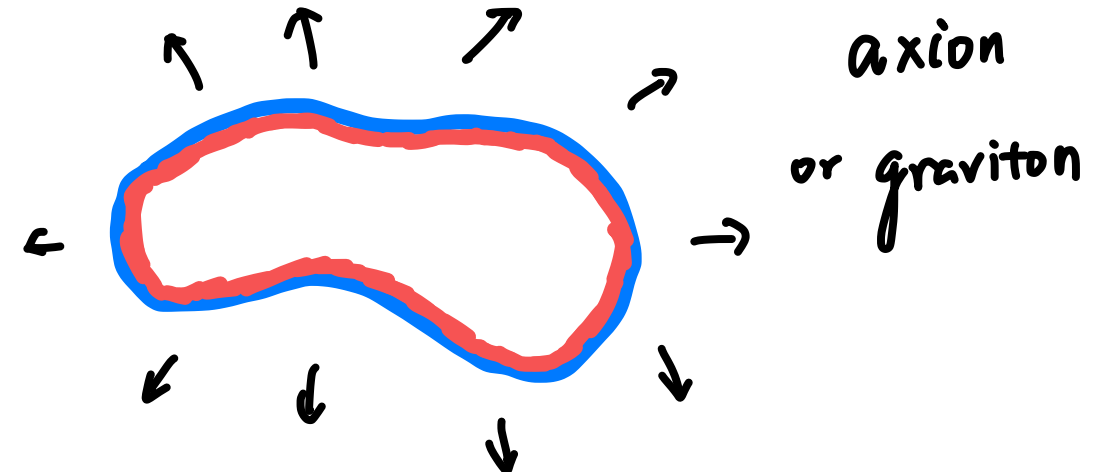
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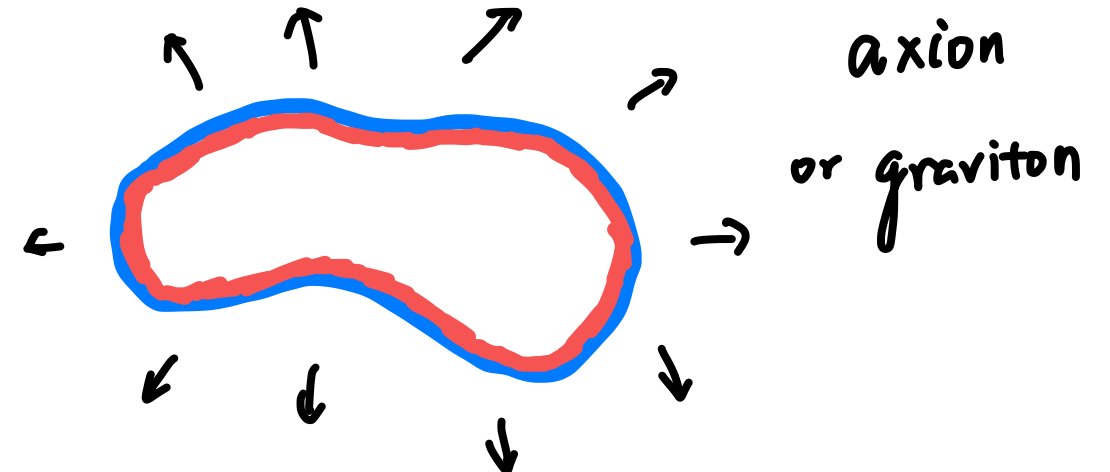


- (1,1) string is gauge string, but it also has axion as light d.o.f

$$\mathcal{L} = -\frac{1}{4}Z_{\mu\nu}Z^{\mu\nu} + \frac{1}{2}e^2(\phi_1^2 + \phi_2^2)Z_\mu^2 - g(\phi_1, \phi_2) e Z^\mu \partial_\mu a + \frac{1}{2} f(\phi_1, \phi_2) (\partial_\mu a)^2$$

$$g(\phi_1, \phi_2) = f_a \frac{\phi_1^2}{v_1^2} - f_a \frac{\phi_2^2}{v_2^2}$$

Gauge string (1,1) radiation



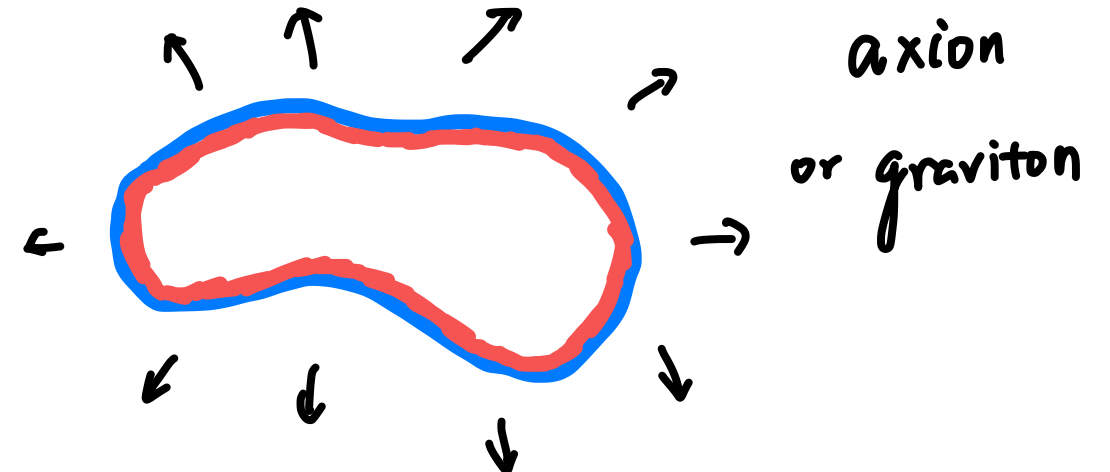
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- Kalb-Ramond field $B^{\mu\nu}$, $\partial_\mu a = \frac{1}{2} \epsilon_{\mu\nu\alpha\beta} \partial^\nu B^{\alpha\beta}$

Gauge string (1,1) radiation



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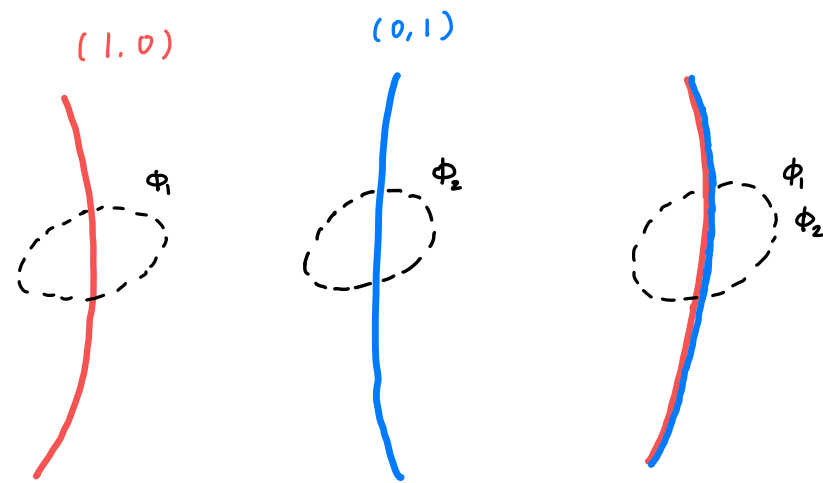
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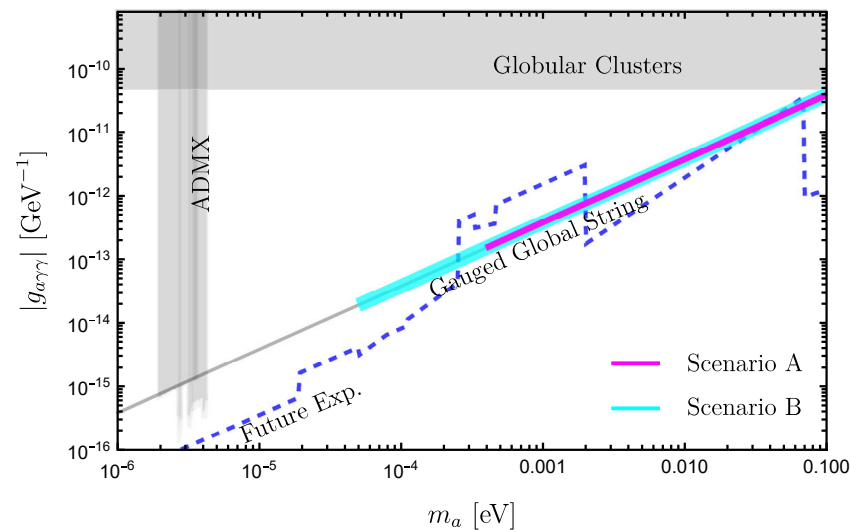
- radiation power

$$\frac{dP_a}{d\Omega} \sim e^2 f_a^2$$

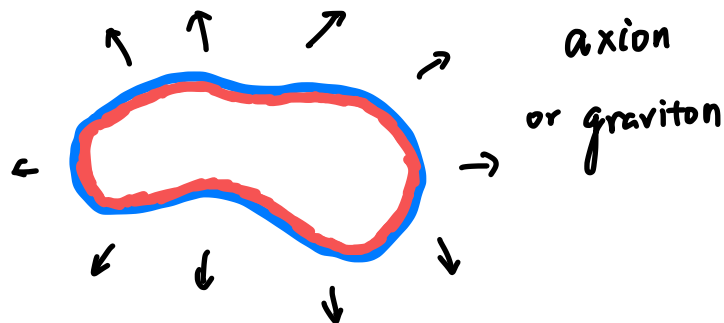
Conclusion



- $U(1)_Z \times U(1)_{PQ}$
(1,0), (0,1) and (1,1) strings



- Cosmology
Y-Junctions
opening QCD axion mass windows



- (1,1) gauge string
radiating axions and gravitons