

Velocity Acoustic Oscillations on Cosmic Dawn 21cm power spectrum as a probe of Axion Dark Matter

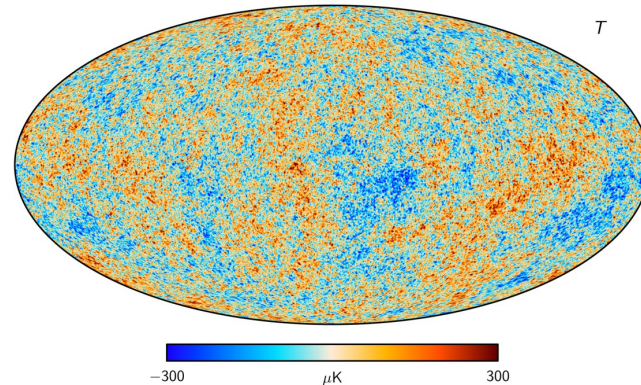
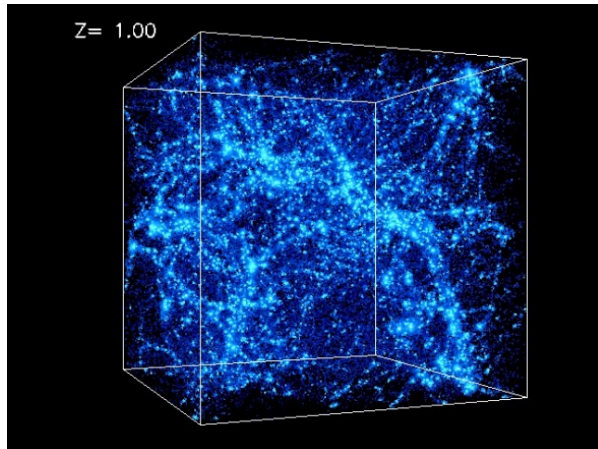
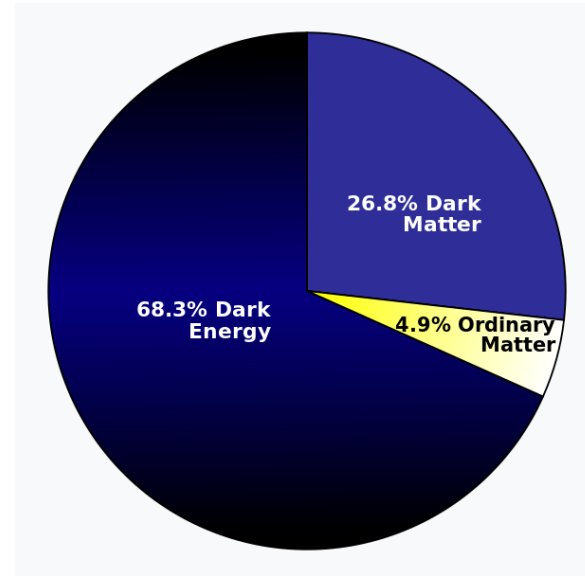
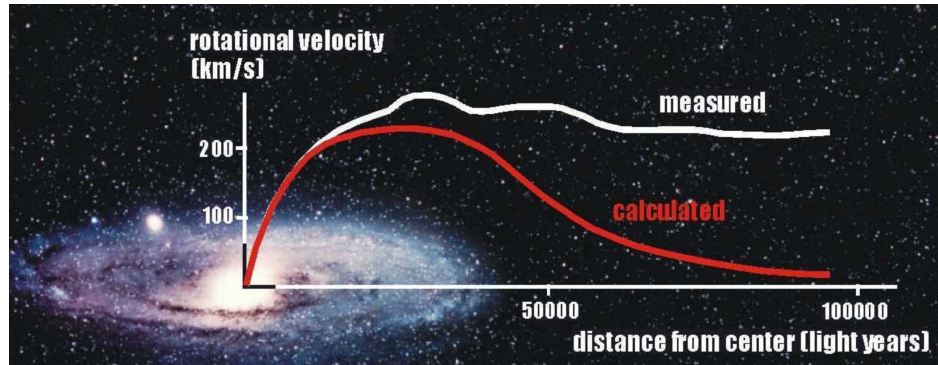


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Based on 2401.14234, The Astrophysical Journal, Volume 964, Number 1
In collaboration with Hengjie Lin, Meng Zhang, Bin Yue, Yan Gong, Yidong Xu,
and Xuelei Chen

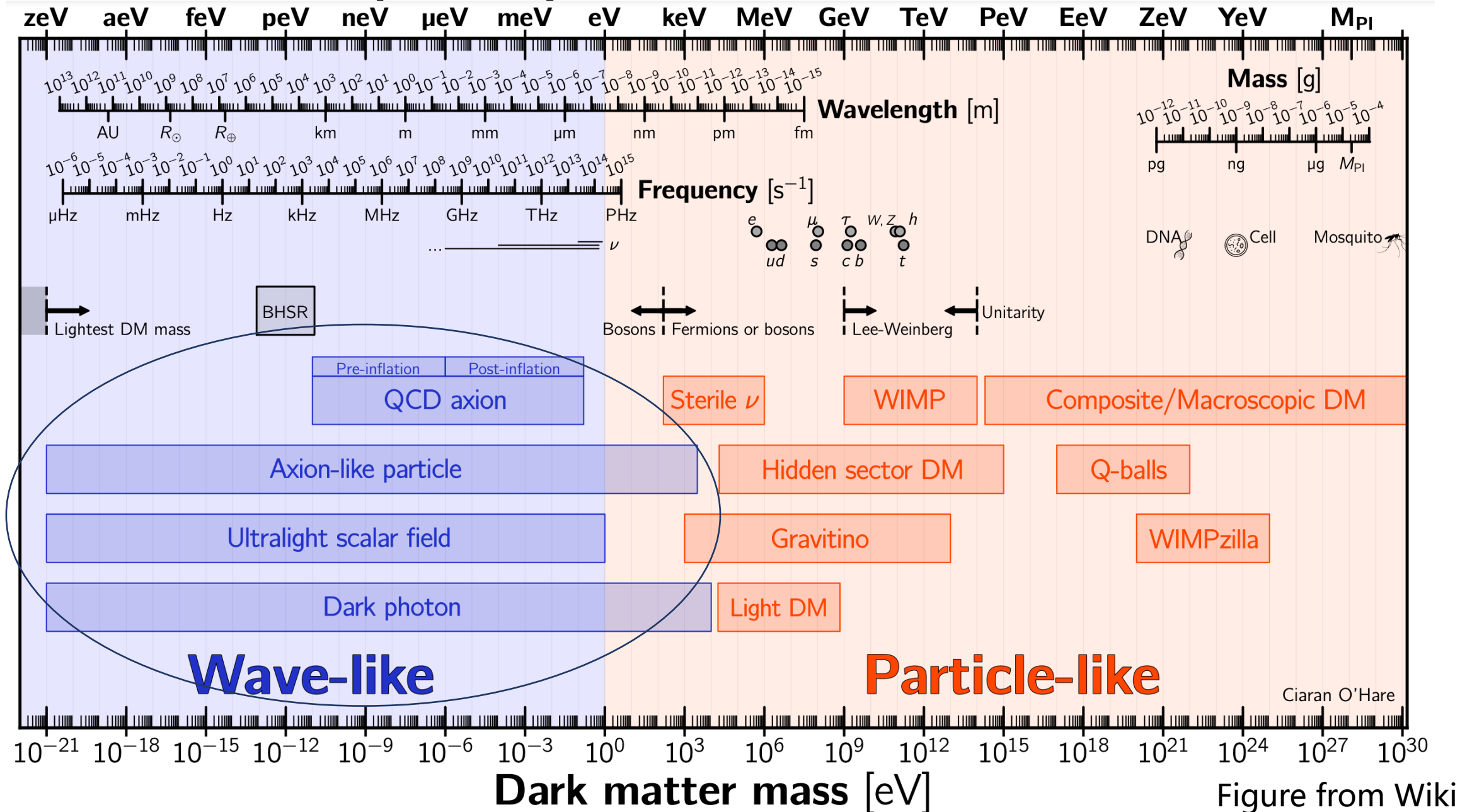
Axion2024 Zhangjiajie, Hunan Province, July22, 2024

We have plenty of indirect evidences for Dark Matter



All the indirect evidences of dark matter are from gravitational effects of astronomical observation

We also have plenty of Dark Matter candidate



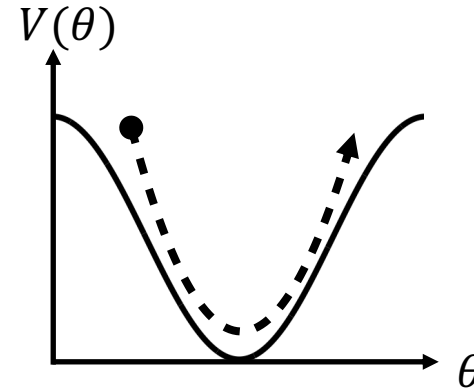
Motivation of Axion

New physics beyond the standard model (SM) in particle physics:

1. Strong CP problem $|\theta| \lesssim 10^{-10}$

Peccei-Quinn mechanism

$$\mathcal{L}_\theta = \theta \frac{A\alpha}{8\pi} G^{\mu\nu,i} \tilde{G}_{\mu\nu}^i \longrightarrow \theta \equiv \frac{a}{f_a} \longrightarrow$$

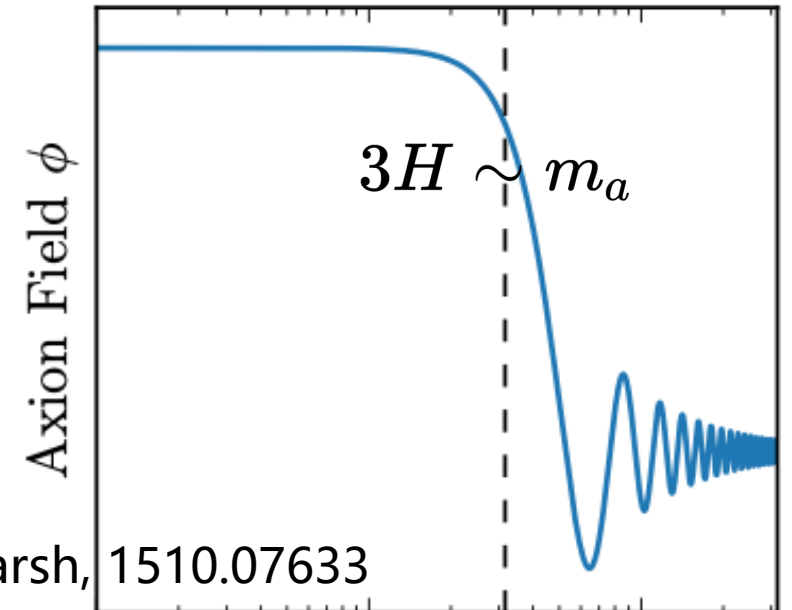


Misalignment mechanism (Axion DM non-thermal production)

$$\ddot{a} + 3H\dot{a} + \frac{dV_{\text{eff}}(a)}{da} = 0,$$

$$V_{\text{eff}}(a) = \frac{1}{2}m_a^2 a^2 + \dots \quad H \equiv \frac{\dot{R}}{R} = \text{Hubble rate}$$

2. Generic prediction of the new physics of string
and M-theory compactifications



D. J. E. Marsh, 1510.07633

Small scale structure anomaly in cosmology

1.cusp-core problem

CDM predicts that the center of dwarf galaxies are cusp-like, observation shows they are core-like

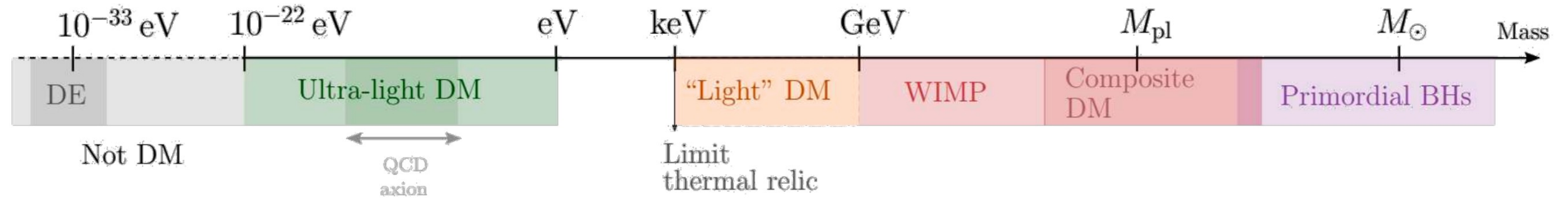
2.missing satellite problem

CDM predicts more satellite galaxies than we observed.

3. too-big-to-fail problem

Many of the satellites are so big that there must be enough stars in it so that we can see them.

Axion as Dark Matter: Fuzzy Dark Matter(FDM)



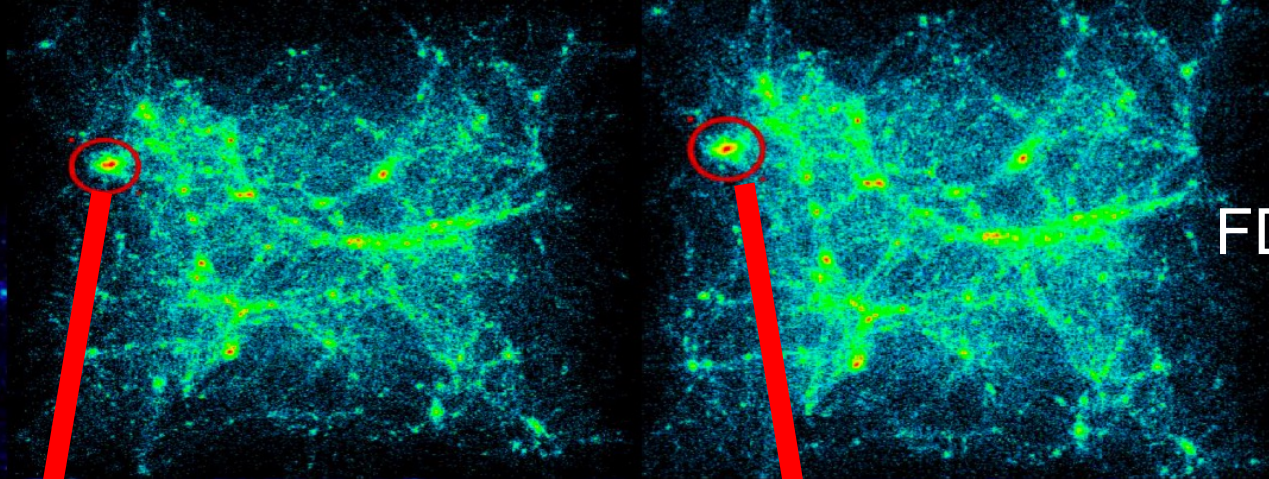
Many name in references: FCDM, BECDM, ULDM, ψ DM, quantum-wave DM, (ultra-light) axion(-like) DM (ULA, UALP)...

Fuzzy Dark Matter: DM as an extremely light(pseudo) scalar particle.

Very tiny mass \Rightarrow large de Broglie wavelength ($\lambda \sim 1/m$)

Macroscopic quantum (wave) effects on kpc scales(Galaxy scale):
Structures resist collapse below quantum wavelength

Large scale

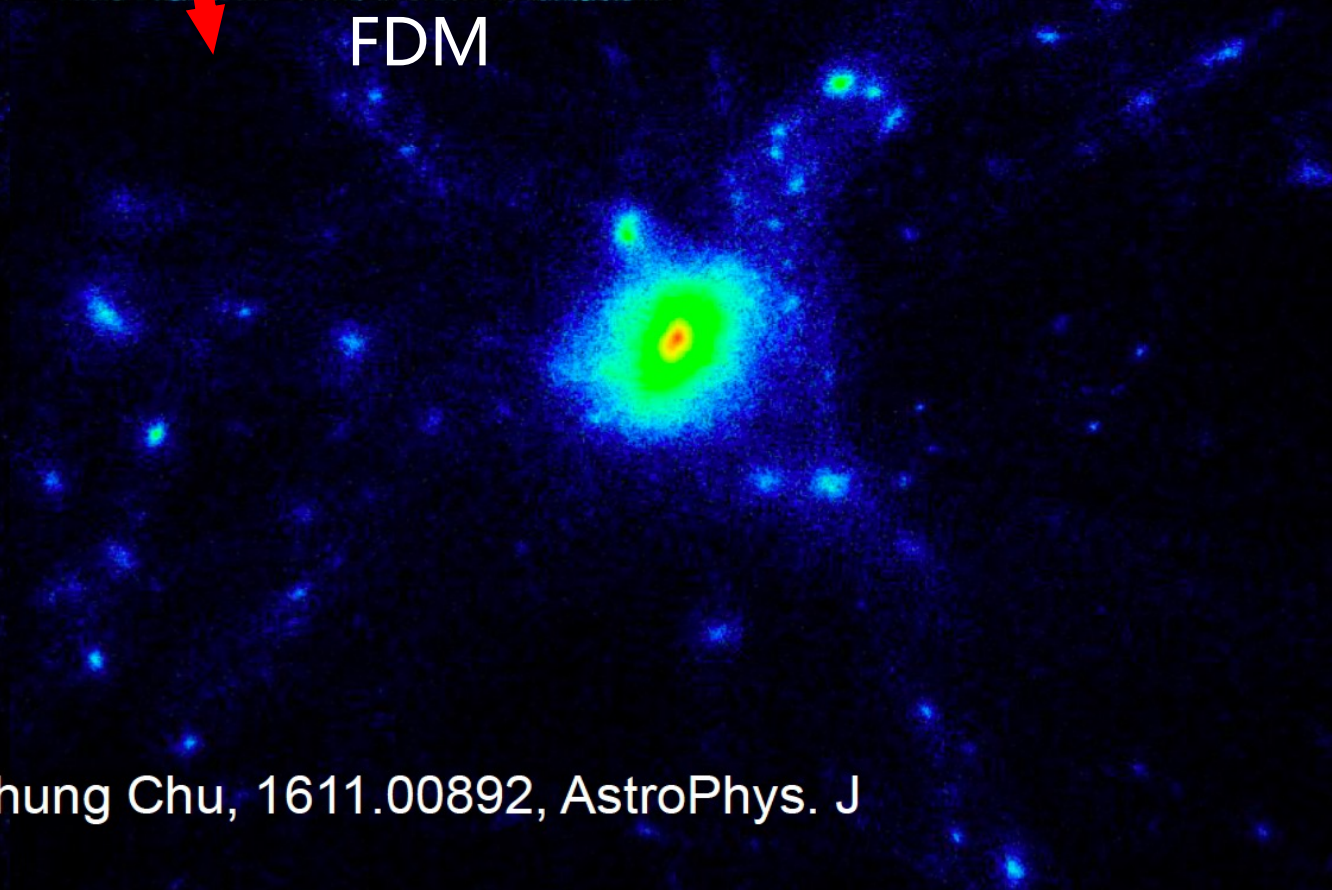
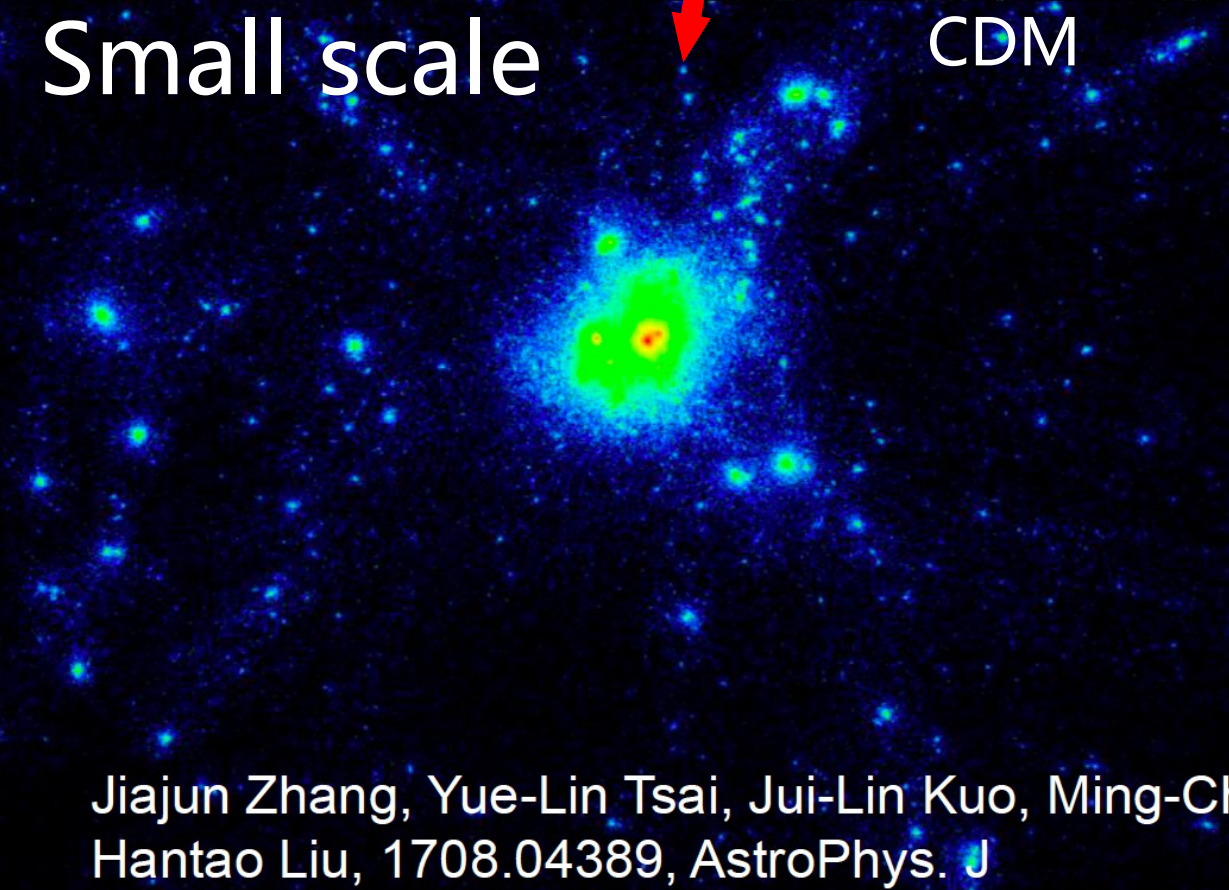


FDM equivalent to CDM

Small scale

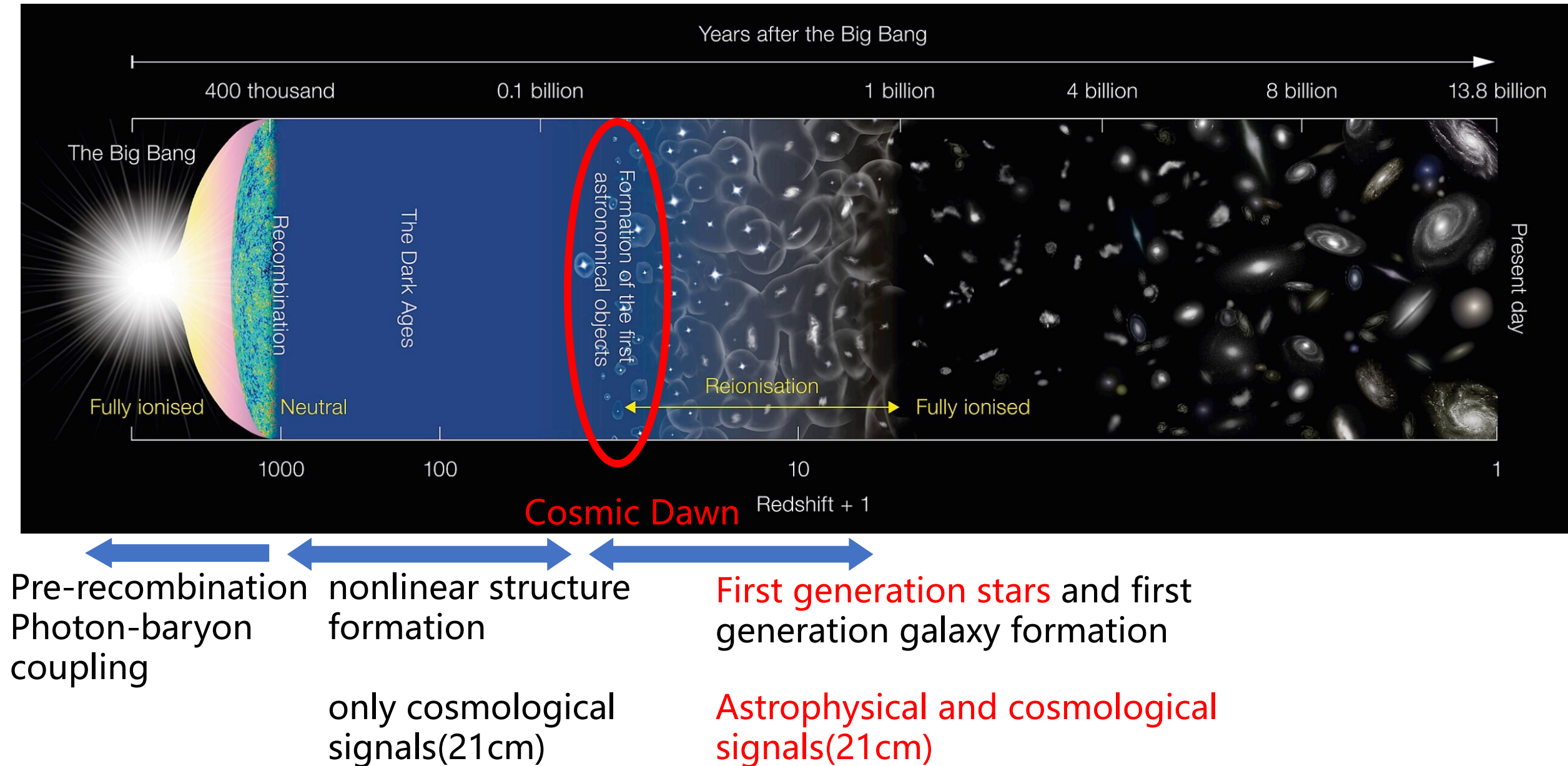
CDM

FDM



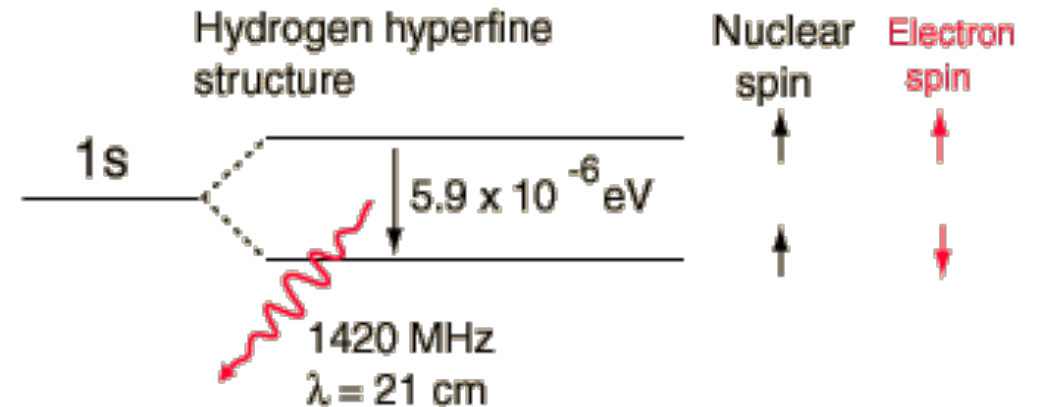
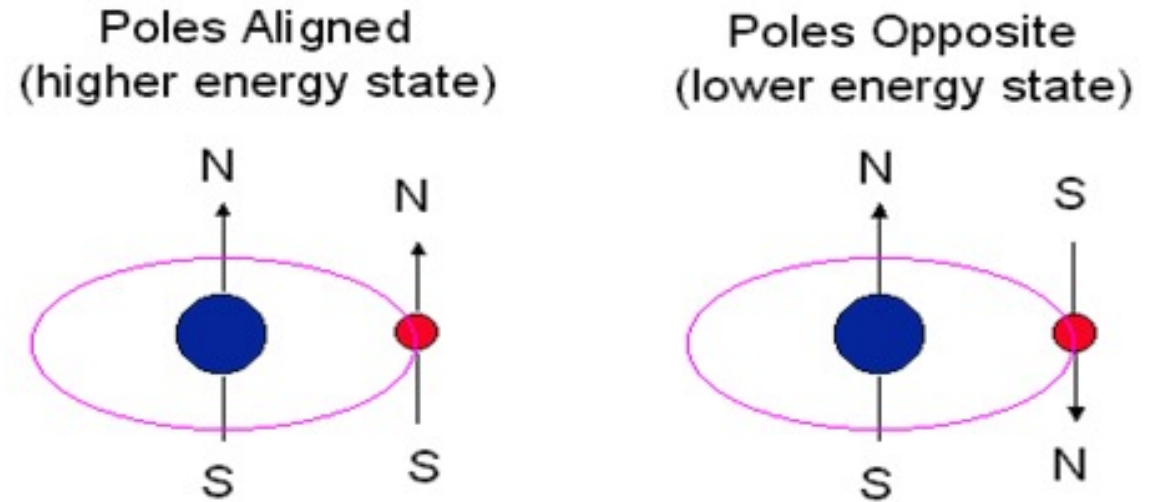
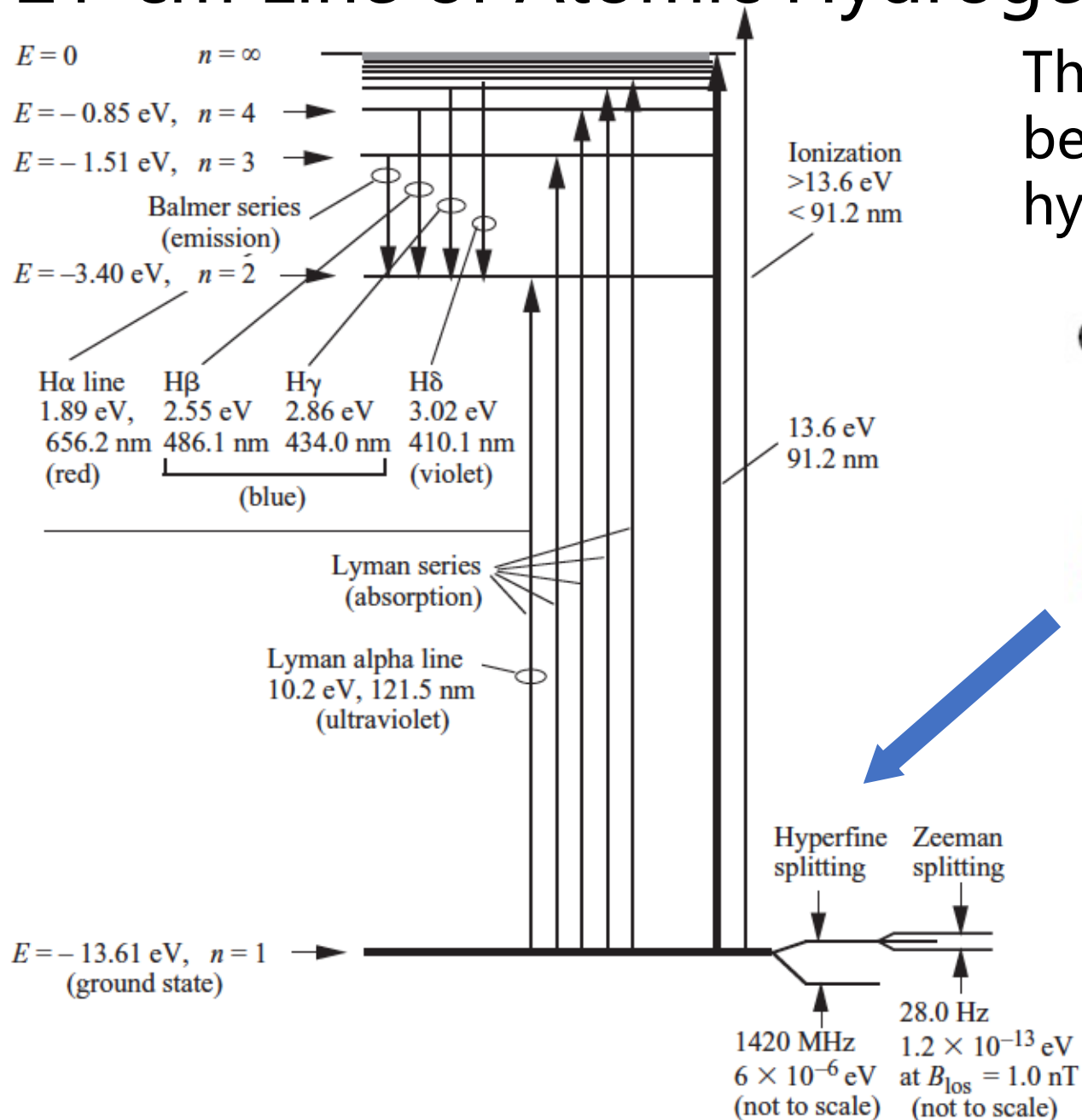
Jiajun Zhang, Yue-Lin Tsai, Jui-Lin Kuo, Ming-Chung Chu, 1611.00892, AstroPhys. J
Hantao Liu, 1708.04389, AstroPhys. J

Evolution history of the universe:



21-cm Line of Atomic Hydrogen

The 21cm line comes from the transition between two hyperfine states of the hydrogen atom 1s state.



21cm spin temperature and brightness temperature

Radiative transfer equation: $\frac{dI_\nu}{d\xi} = -\kappa_\nu I_\nu + j_\nu$

Rayleigh-Jeans limit: $I_\nu \equiv \frac{2k_B\nu^2}{c^2} T_b(\nu)$

The formal solution of radiative transfer equation:

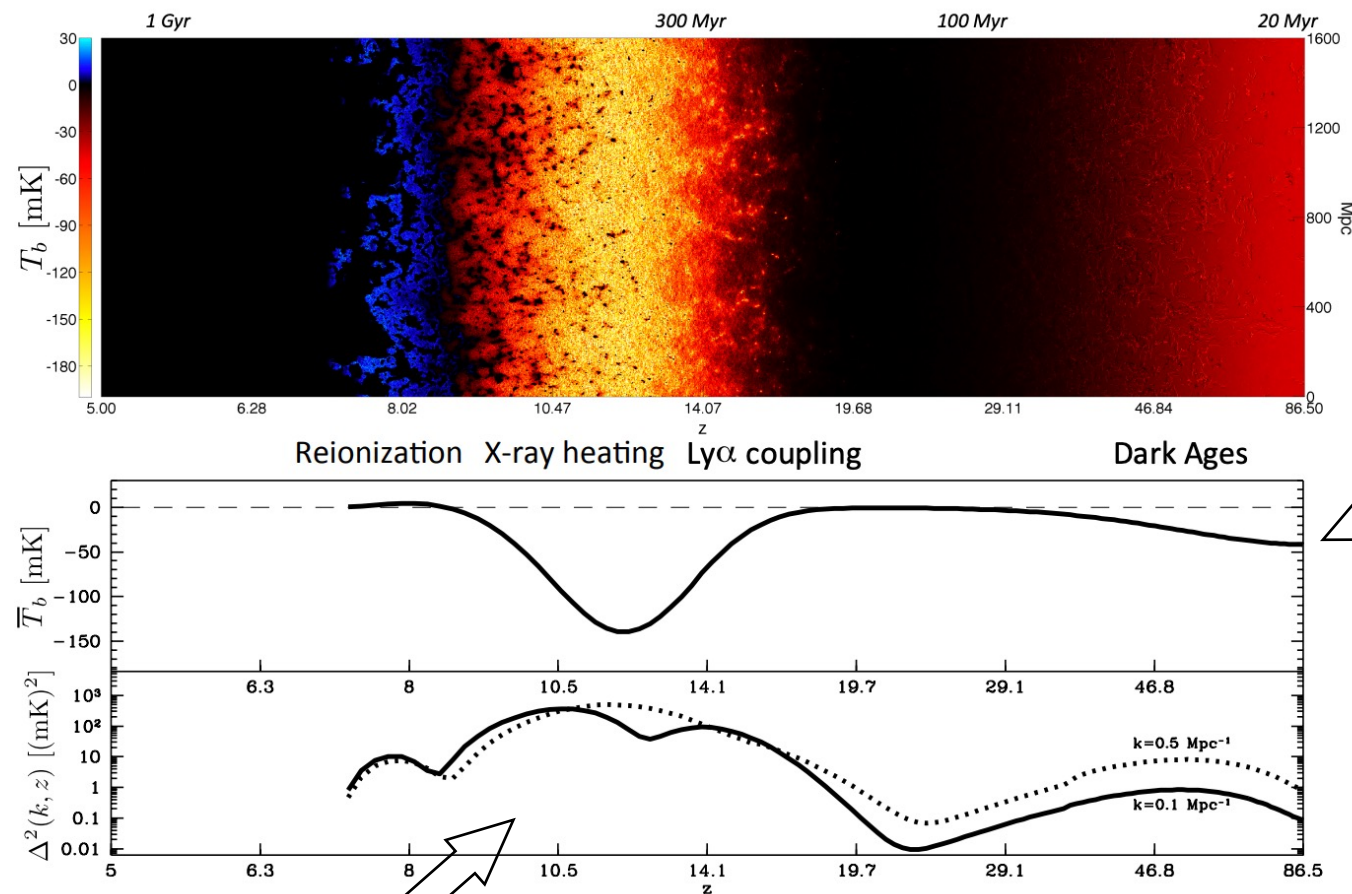
$$\delta T_b(\nu_{\text{obs}}) \equiv T_b(\nu_{\text{obs}}) - T_{\text{CMB},0} = a[T_s(\mathbf{x}, z) - T_{\text{CMB}}(a)](1 - e^{-\tau_{\nu_{\text{obs}}}})$$

Spin temperature: $\frac{n_1}{n_0} \equiv \frac{g_1}{g_0} e^{-\Delta E_{10}/k_B T_s} = 3e^{-T_\star/T_s} \quad T_s^{-1} = \frac{T_\gamma^{-1} + x_\alpha T_\alpha^{-1} + x_c T_k^{-1}}{1 + x_\alpha + x_c}$

$$\widehat{\delta T}_b(z) = 23.88 \left(\frac{\Omega_b h^2}{0.02} \right) \sqrt{\frac{0.15}{\Omega_M h^2} \frac{1+z}{10}} \bar{x}_{\text{HI,m}}(z) \text{ mK}$$

$$\delta T_b(\nu_{\text{obs}}) = \widehat{\delta T}_b(z) [1 + \delta_{\rho_{\text{HI}}}] \left[1 - \frac{T_{\text{CMB}}(a)}{T_s} \right],$$

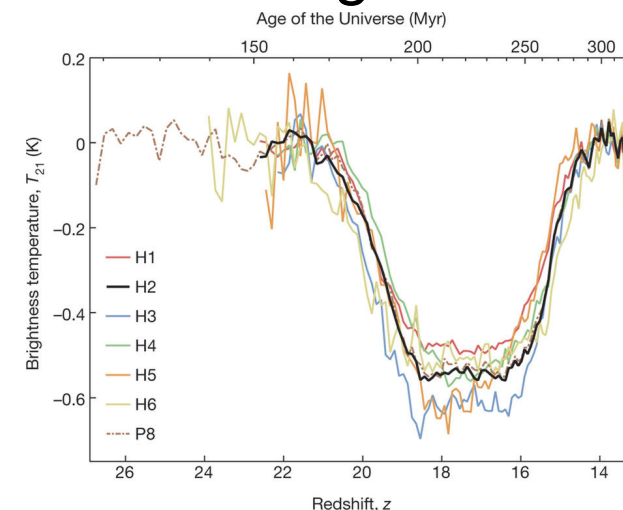
Measure the 21cm brightness temperature



Global spectrum:

$$\bar{T}_b(\nu) = \int d\Omega \delta T_b(\hat{\mathbf{r}}, \nu)$$

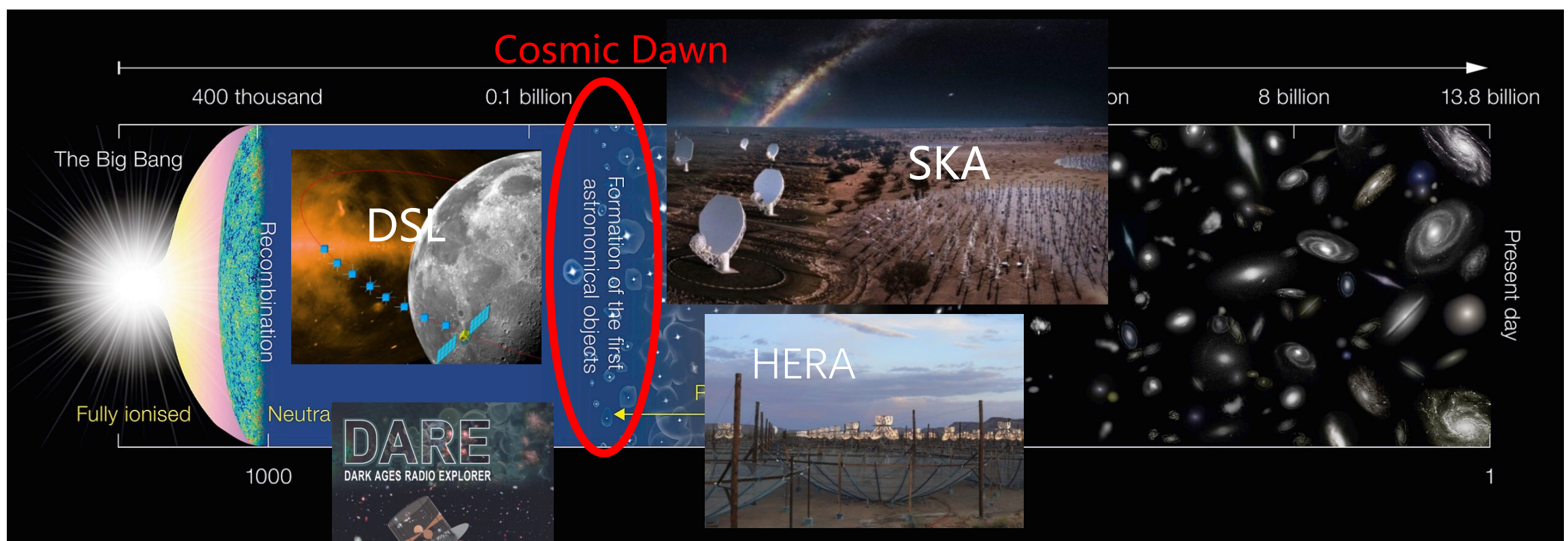
Measurement of the averaged T field over all angles on the sky



Power spectrum: $\delta T_b^2(\mathbf{r}) = \int_{-\infty}^{\infty} \frac{d^3 k}{(2\pi)^3} P(k) = \int_0^{\infty} d \ln k \Delta^2(k)$ Bowman et al.(2018)

Measurement of the variance of fluctuations in a T field as a function of k

Cosmic Dawn

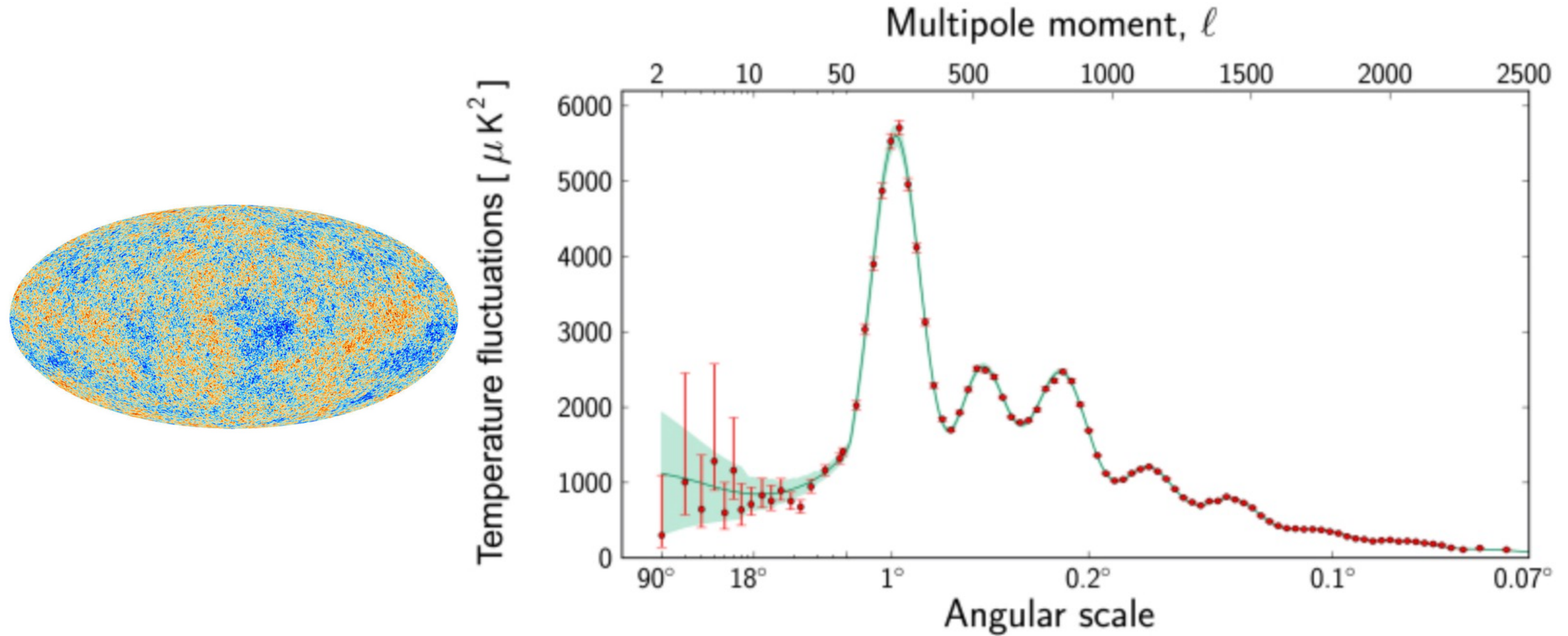


The cosmological signals in the dark age are more difficult to measure, which may require the next generation instruments:



The astrophysical and cosmological signals can be measured by current and near future instruments.

Baryon Acoustic Oscillations(BAO)




Planck 2018

ESA and the Planck Collaboration

Velocity Acoustic Oscillations(VAO)

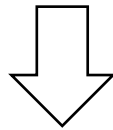
Under the interaction of photon pressure and gravity, both the baryon density (BAO) and velocity (VAO) different from dark matter.

$$v_{\text{db}} \sim 30 \left(\frac{1+z}{1081} \right) [\text{km/s}]$$

e.g. $z_{\text{rec}} = 1020$ $\mathcal{M} \equiv v_{\text{db}}/c_s \sim 5$  $k_{\text{vdb}} \equiv \frac{aH}{\langle v_{\text{db}}^2 \rangle^{1/2}} \bigg|_{\text{dec}} = \frac{k_J}{\mathcal{M}} \sim 40 \text{Mpc}^{-1}$

$$z_{\text{cosmic dawn}} = 20 \quad v_{\text{db}} \sim 1 [\text{km/s}]$$

Comparable with the circular velocities of dark matter halo has:

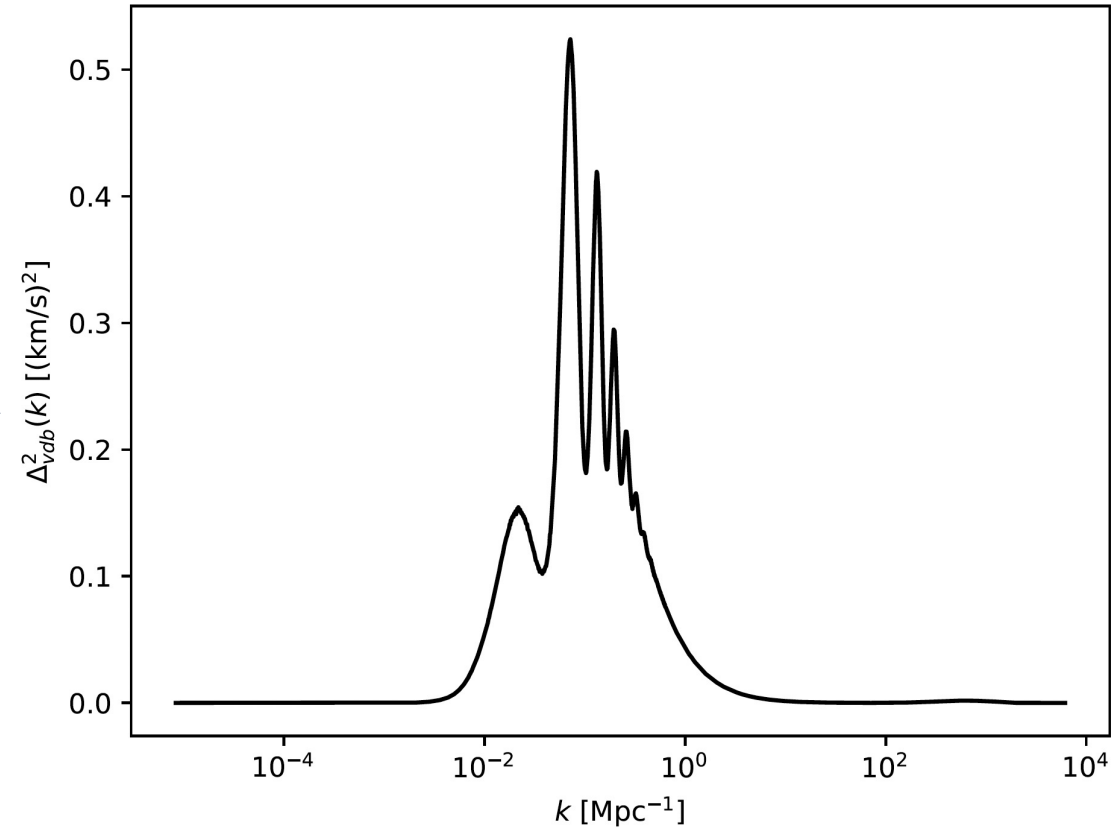
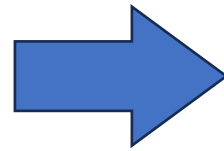
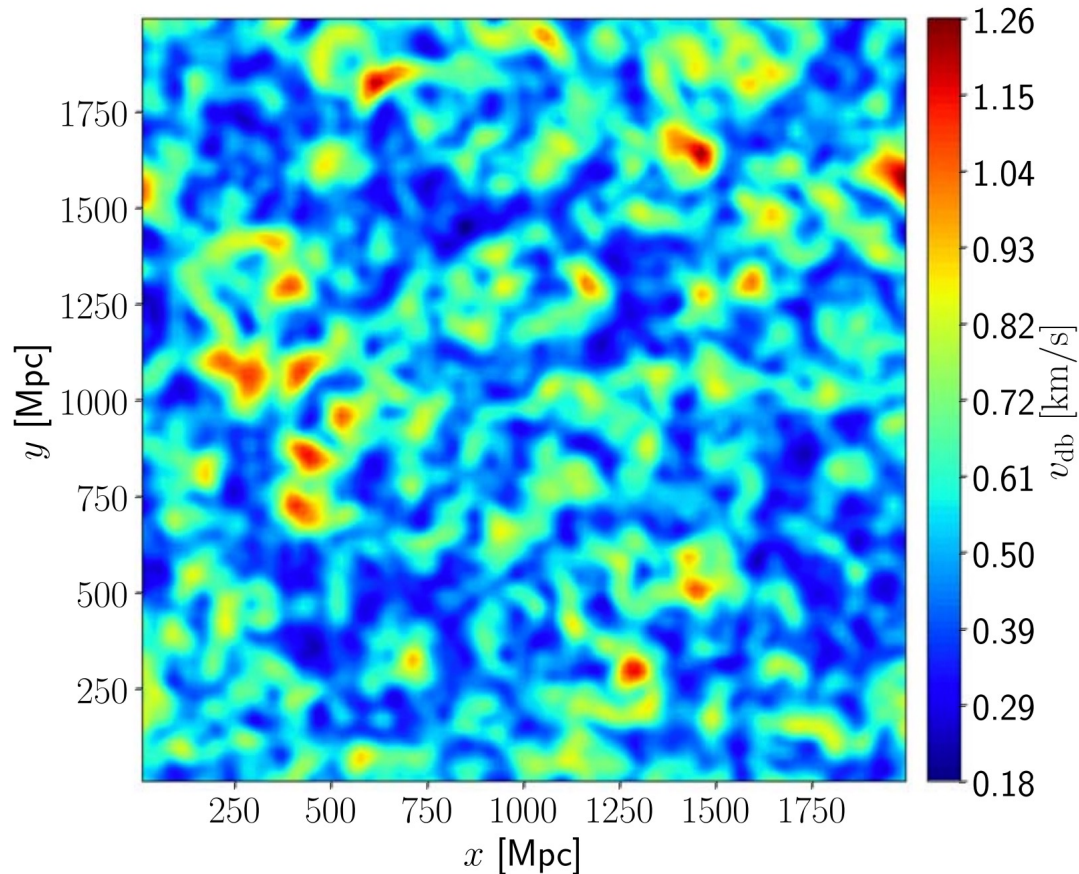


$$\sigma_{\text{rms}} \approx 0.6 \text{ km s}^{-1}$$

The relative velocity $\gtrsim 1\sigma_{\text{rm}}$ would be significantly suppressed the formation of the first generation stars, in turn impacts the 21 cm signal

Velocity Acoustic Oscillations(VAO)

Relative velocity field of baryon and dark matter: $\mathbf{v}_{\text{d/b}}(\mathbf{k}, a) = -i \frac{a\mathbf{k}}{k^2} \dot{\delta}_{\text{d/b}}(\mathbf{k}, a) = -i \frac{a\mathbf{k}}{k^2} \dot{T}_{\text{d/b}}(\mathbf{k}, a) \delta_{\text{pri}}(\mathbf{k})$



Relative velocity field of baryon and dark matter

$$v_{\text{db}}^2(\mathbf{x}) = \int \frac{dk}{k} \Delta_{v_{\text{db}}}^2(k) \quad \text{power spectrum}$$

The Big Picture

DM-Baryon relative velocity

(the large-scale coherently oscillating structures)

$$f_{\text{coll}}(z | \delta_{\text{cell}}, M_{\text{cell}}, v_{\text{db}}) = \frac{1}{\bar{\rho}_m} \int_{M_{\text{crit}}}^{\infty} dM M \frac{dn}{dM}(M, z | \delta_{\text{cell}}, M_{\text{cell}})$$

Minihalo distribution (critical mass) allowing the formation of first-generation stars

$$\rho_{\text{cell}}^* \sim f_* \frac{\Omega_b}{\Omega_m} \rho_m f_{\text{coll}}$$

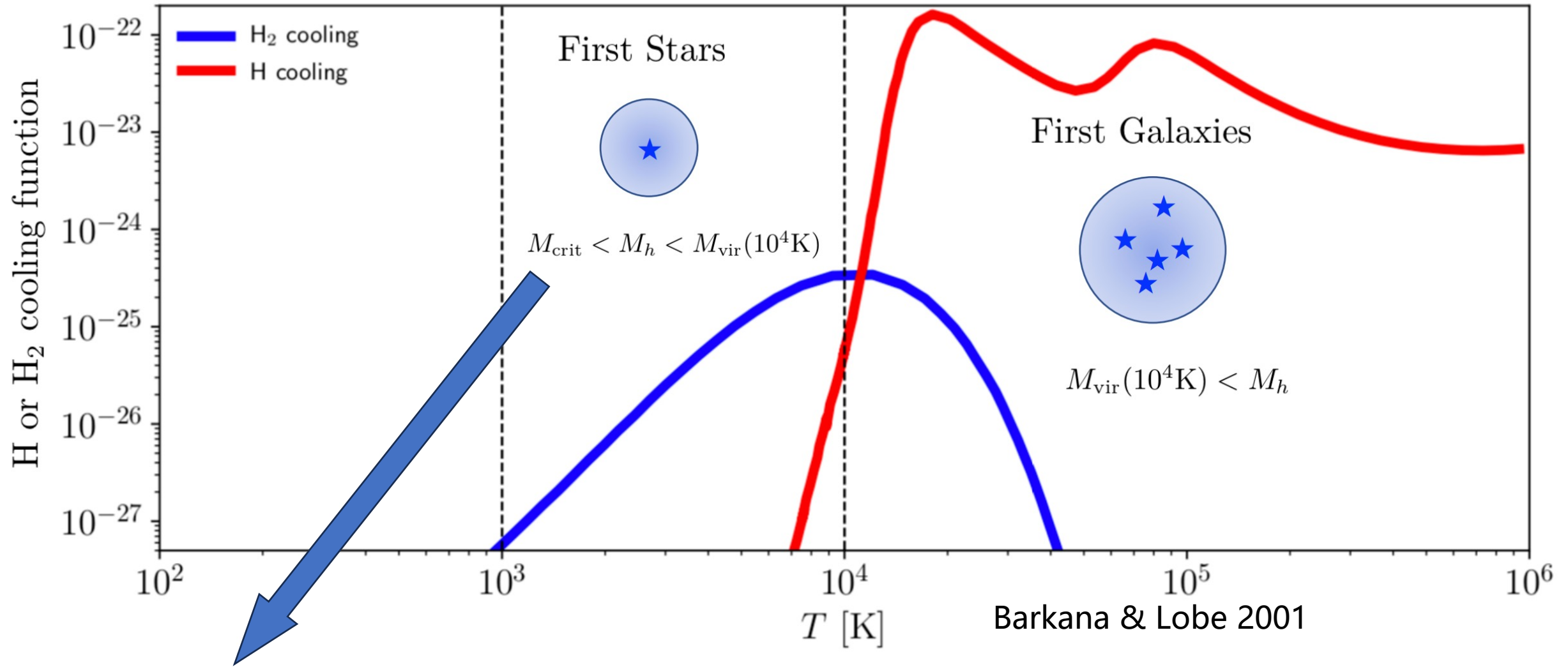
Radiation field produced by first-generation star formation

$$\begin{aligned} \epsilon_{\text{cell}}^{\alpha} &\propto \rho_{\text{cell}}^* & \epsilon_{\text{cell}}^X &\propto \rho_{\text{cell}}^* \\ T_s^{-1} &= \frac{T_{\gamma}^{-1} + x_{\alpha} T_{\alpha}^{-1} + x_c T_K^{-1}}{1 + x_{\alpha} + x_c} \end{aligned}$$

VAO feature on the 21 cm power spectrum

$$\delta T_b = 27 x_{\text{HI}} (1 + \delta_b) \left(\frac{\Omega_b h^2}{0.023} \right) \left(\frac{0.15}{\Omega_m h^2} \frac{1+z}{10} \right)^{1/2} \times \left(1 - \frac{T_{\gamma}}{T_s} \right) \left[\frac{\partial v / \partial r}{(1+z) H(z)} \right] \text{mK}$$

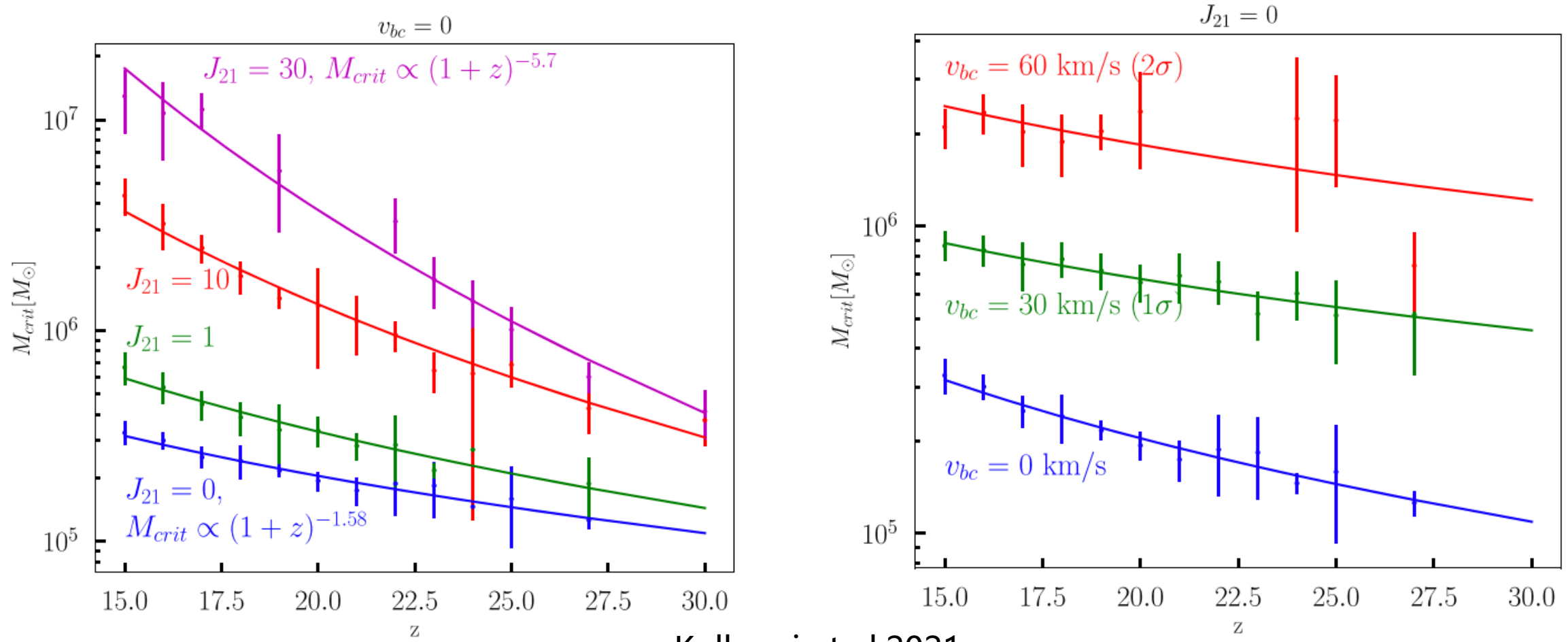
Critical Mass



The minimum mass of a dark matter halo allowing the formation of first-generation stars $M_{\text{crit}}(J_{\text{LW},21}, v_{\text{db}}^{\text{rec}}, z)$

Critical Mass

The critical mass depends on the relative velocity of baryons and dark matter, the cooling mechanism of gas, LW radiation, X-rays



Kulkarni et al.2021

$$M_{crit} (J_{LW,21}, v_{db}^{rec}, z) = M_{z=20}(J_{LW,21}, v_{db}^{rec}) \times \left(\frac{1+z}{21} \right)^{-\alpha(J_{LW,21}, v_{bc}^{rec})}$$

VAO signal

The VAOs statistically independent from density fluctuations sourced by overdensities so they can be linearly added to the usual (no-vdb) 21-cm power spectrum to obtain the total signal.

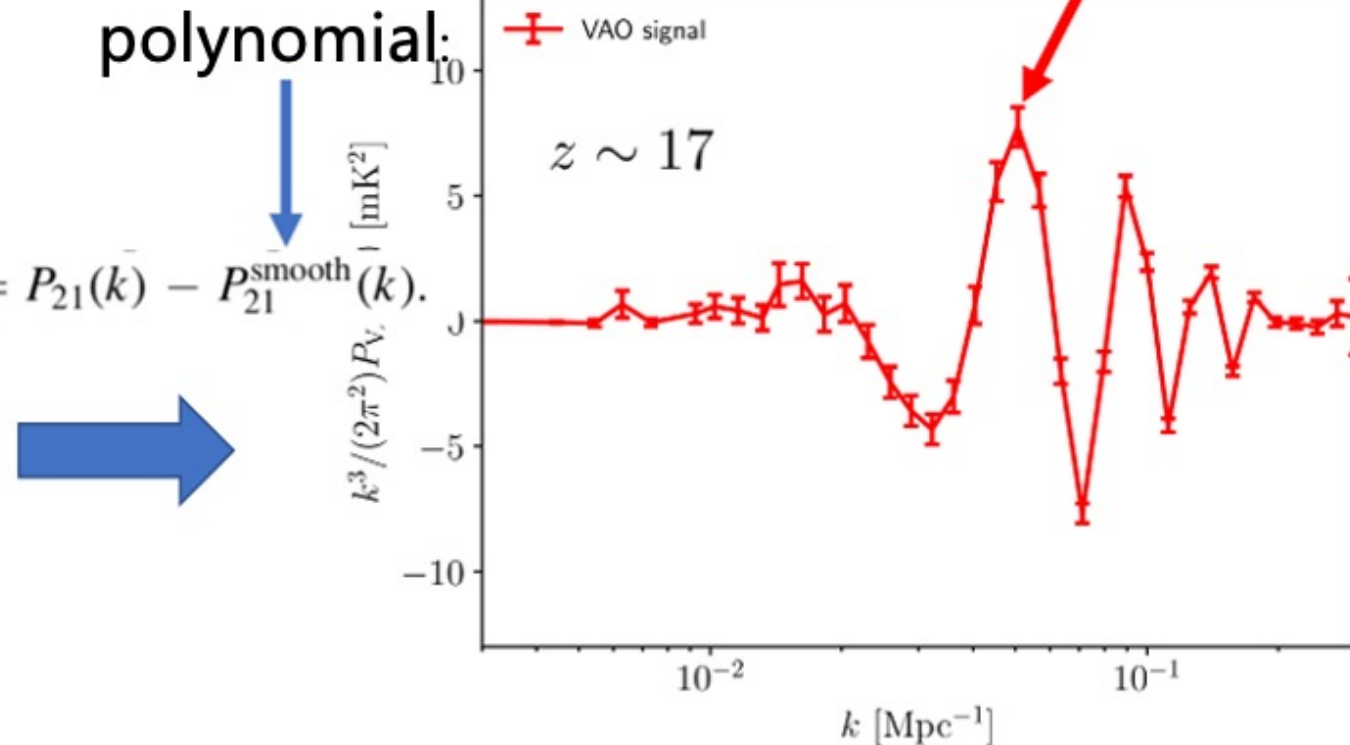
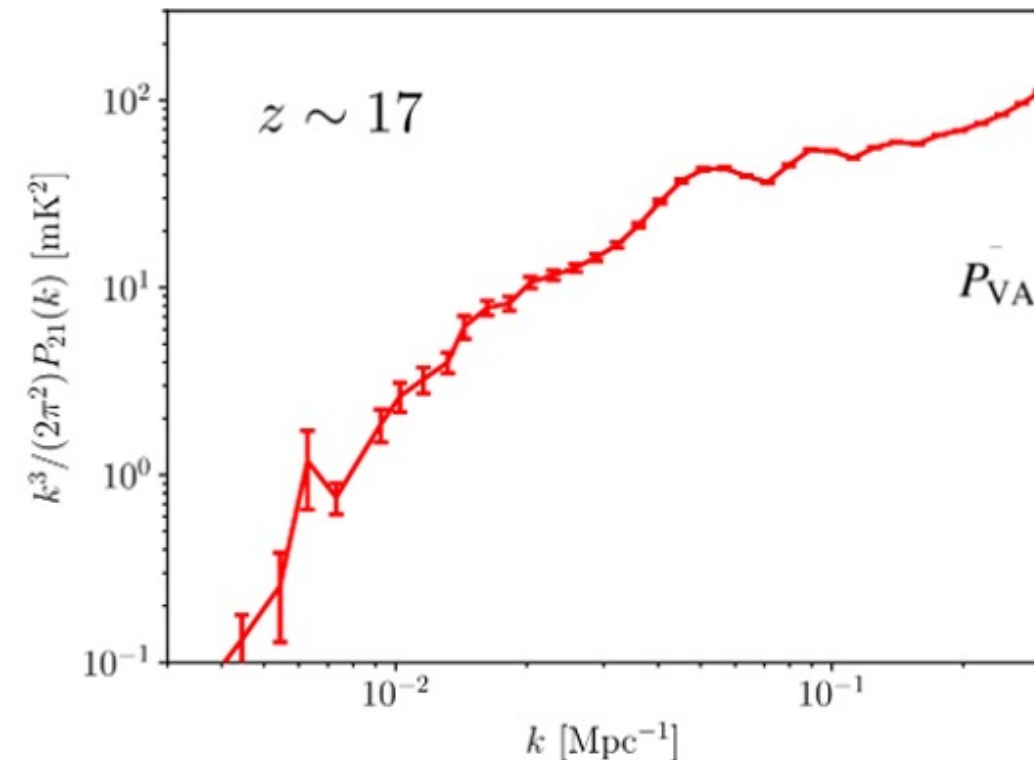
Muñoz, J. B. 2019a, PhRvD, 100, 063538

Muñoz, J. B. 2019b, PhRvL, 123, 131301

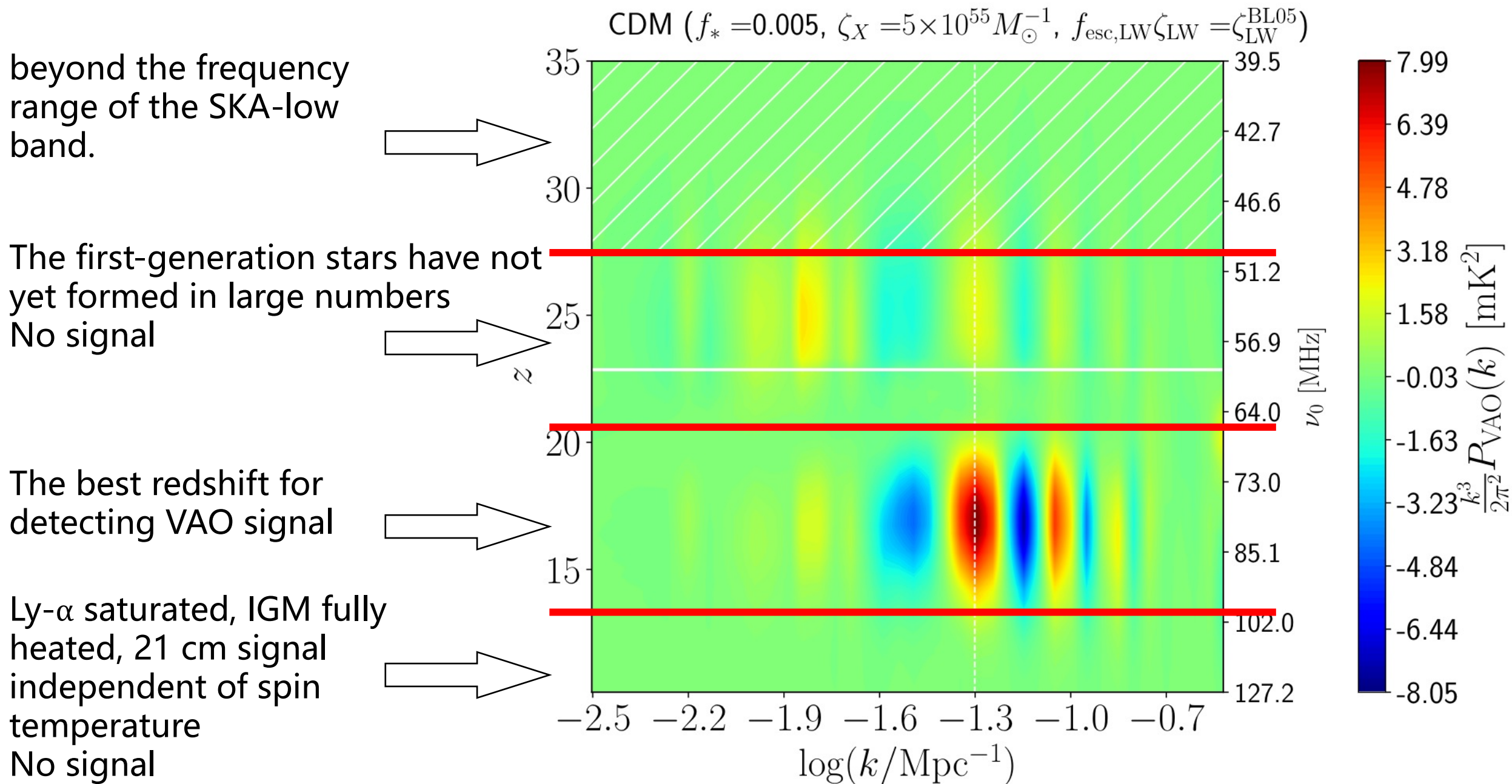
$$\Delta_{21}^2(k, z) = \Delta_{21, \text{vel}}^2(k, z) + \mathcal{P}_n(k, z)$$

$$\Delta_{21, \text{vel}}^2(k, z) = A_{\text{vel}}(z) \Delta_{v^2}^2(k) |W_i(k, z)|^2$$

first peak: $k \approx 0.05 \text{ Mpc}^{-1}$



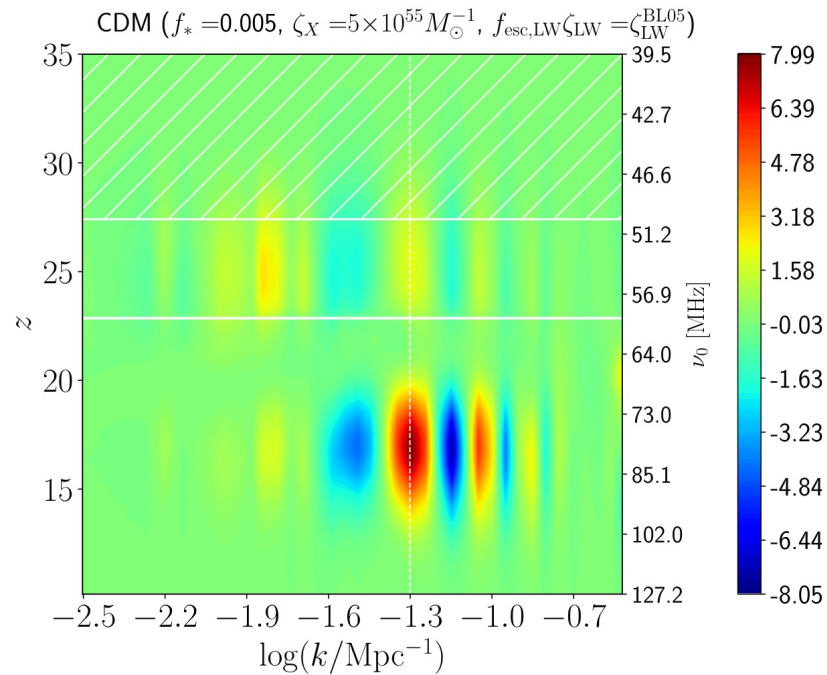
The evolution of VAO signal



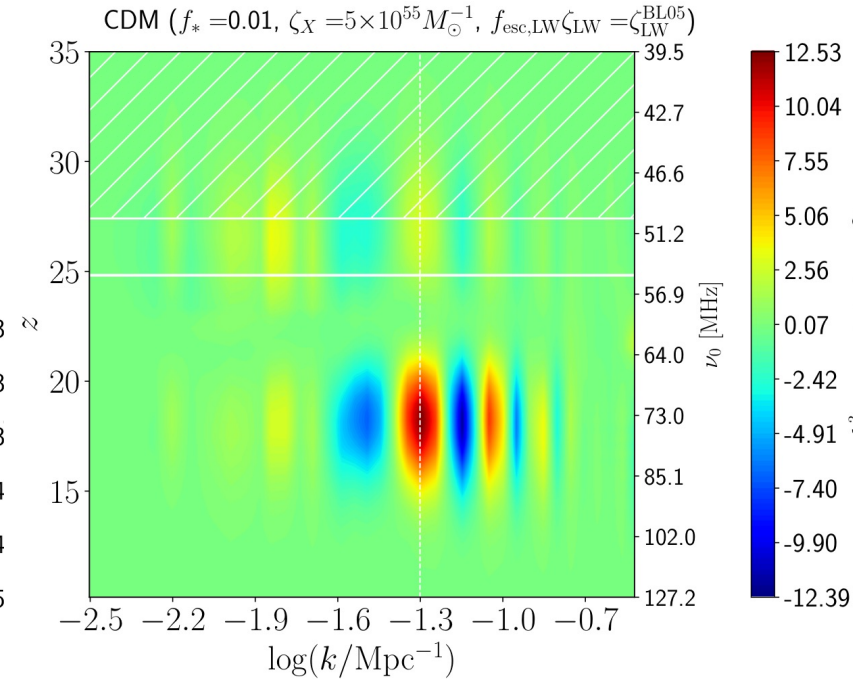
The evolution of VAO signal

Dependence on the efficiency of first-generation star formation :

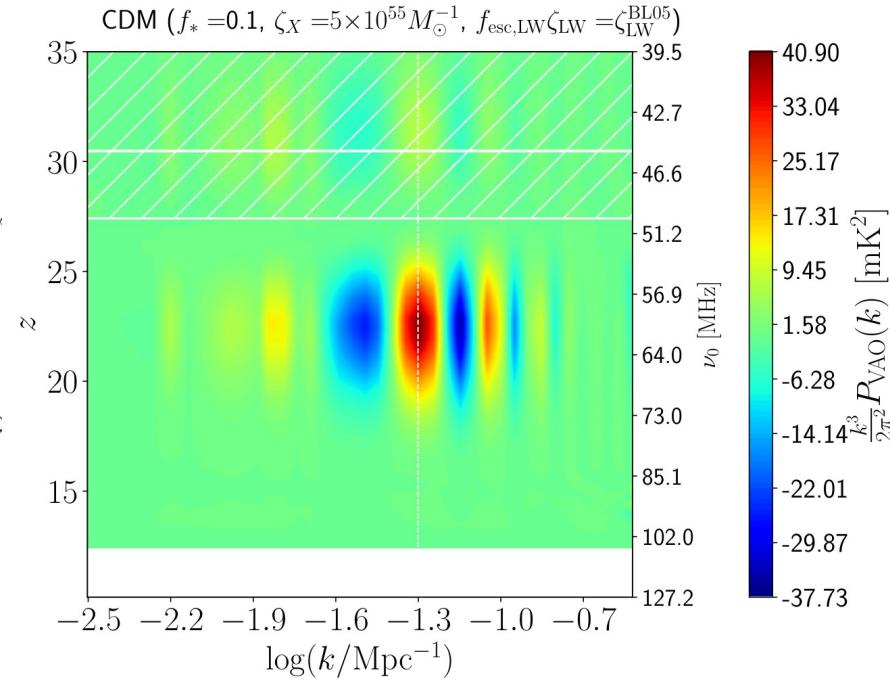
$$f_* = 0.005$$



$$f_* = 0.01$$



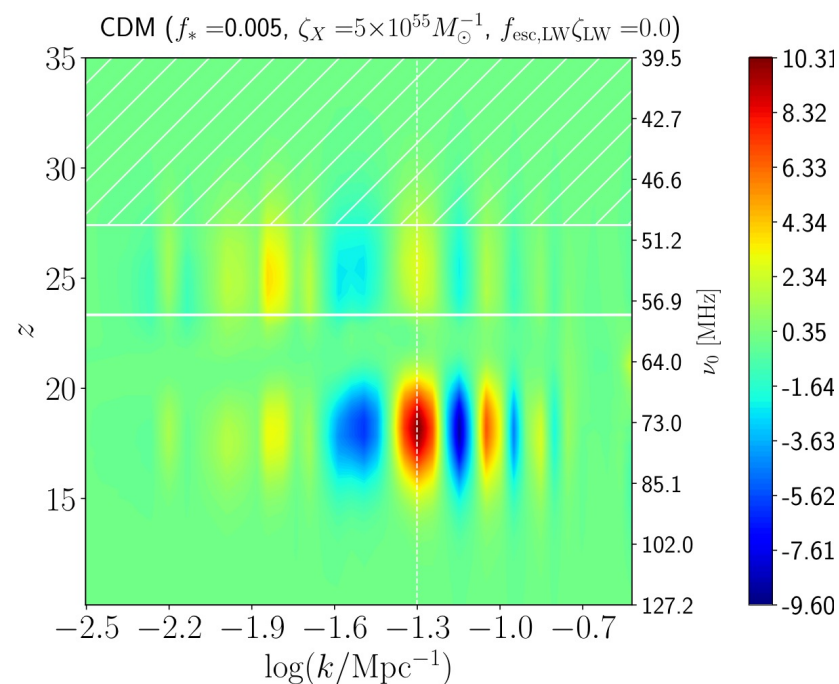
$$f_* = 0.1$$



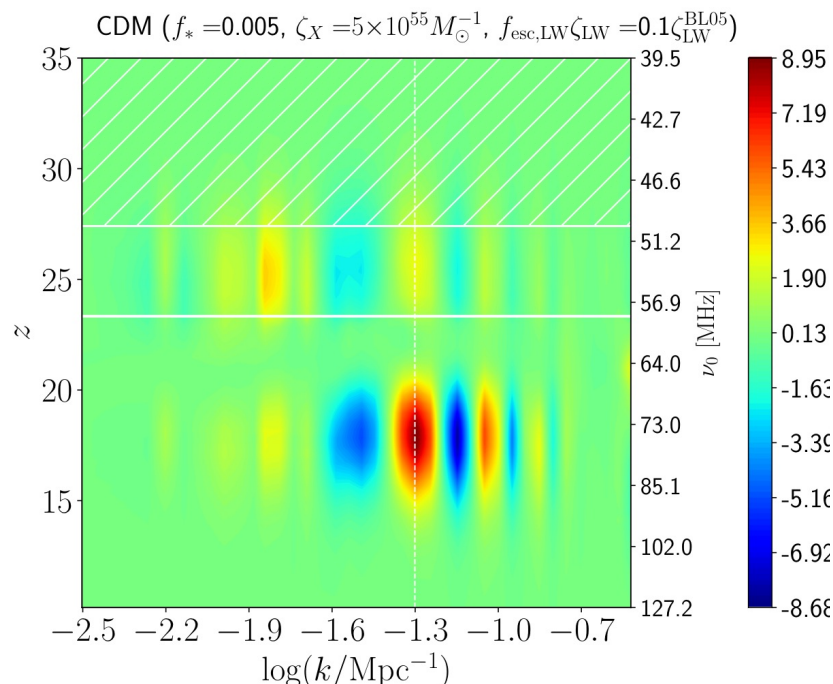
The higher the star formation efficiency, the stronger the signal, and the earlier the signal appears

Dependence on the LW feedback:

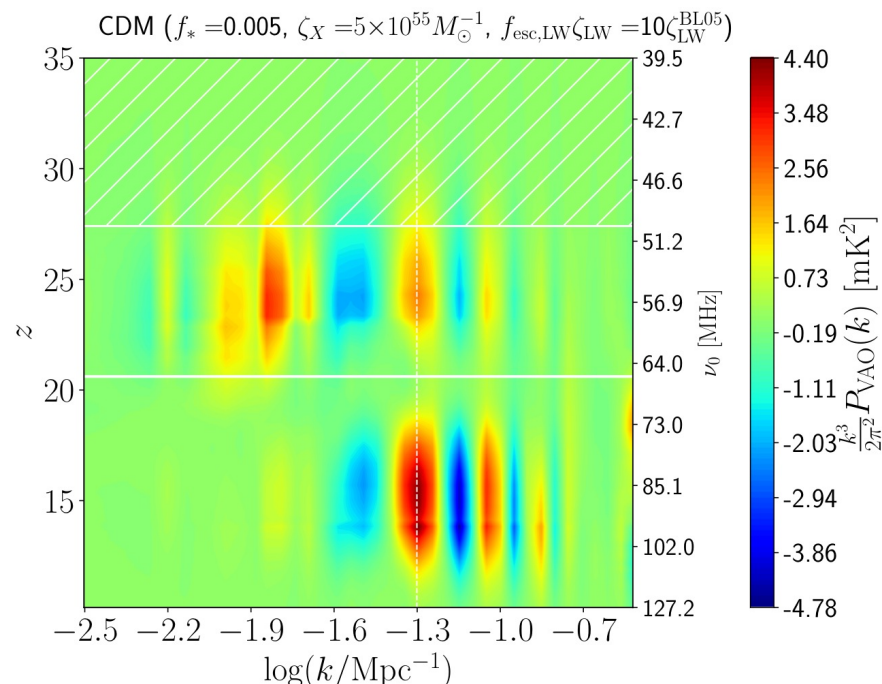
no LW feedback



weak LW feedback



Strong LW feedback

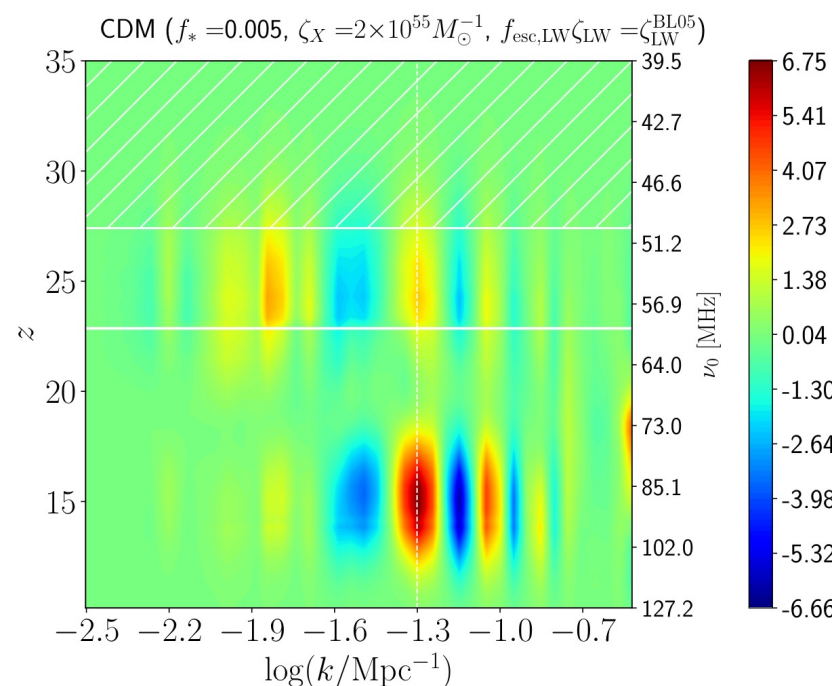


LW feedback effect will weaken VAO signal

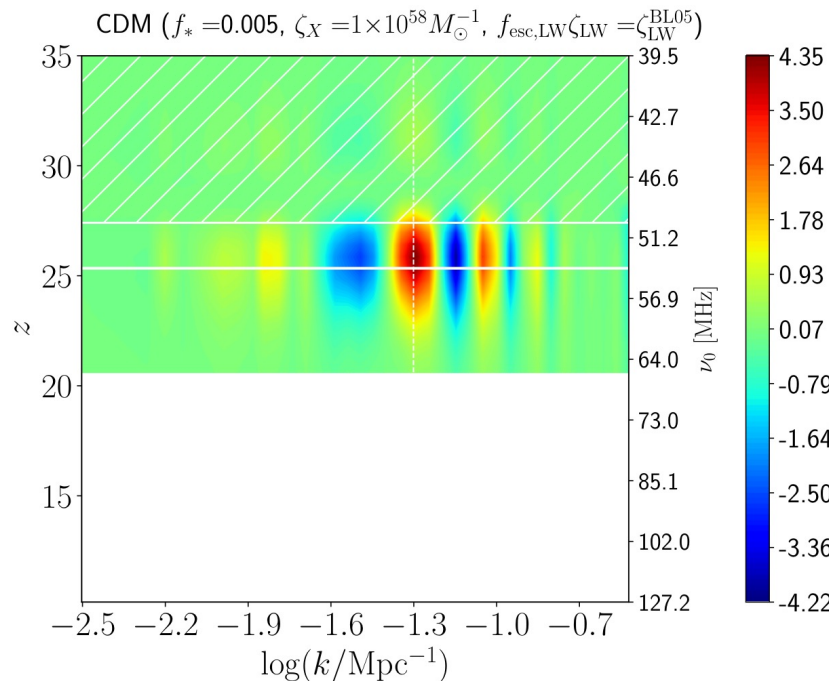
But it is not enough to make it disappear completely

Dependence on the X-rays:

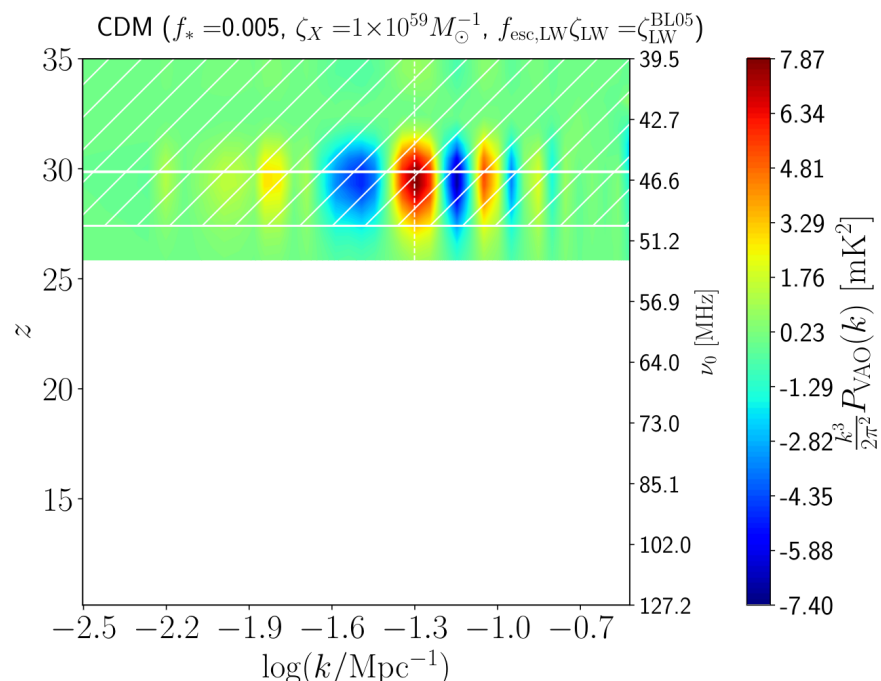
weak X-rays



strong X-rays



very strong X-rays



X-ray will weaken the VAO signal, but not monotonically decrease

X-ray is also the main Ly- α source, which can make the VAO signal appear in advance and enhance its amplitude

CDM vs FDM(Axion)

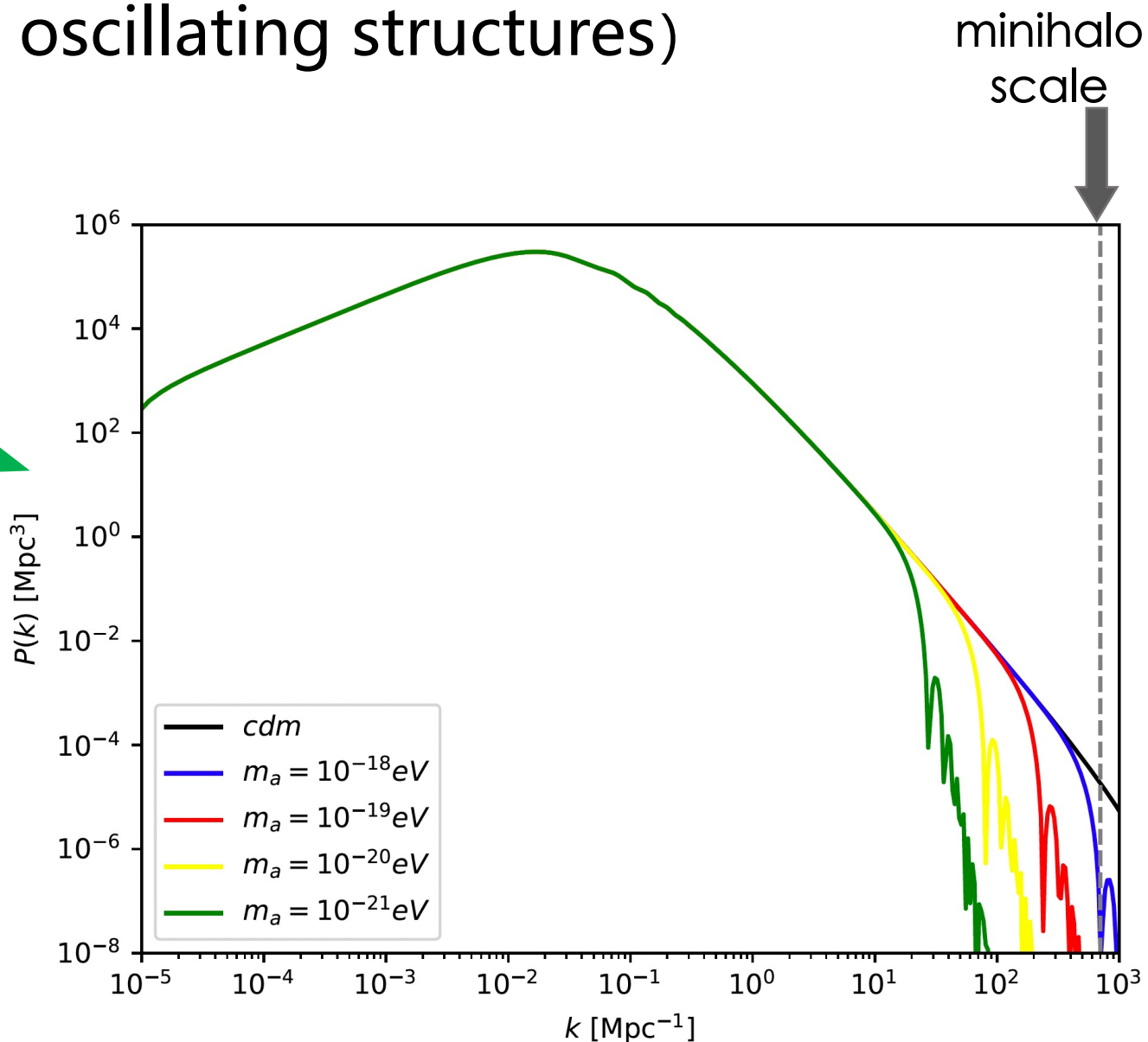
DM-Baryon relative velocity

(the large-scale coherently oscillating structures)

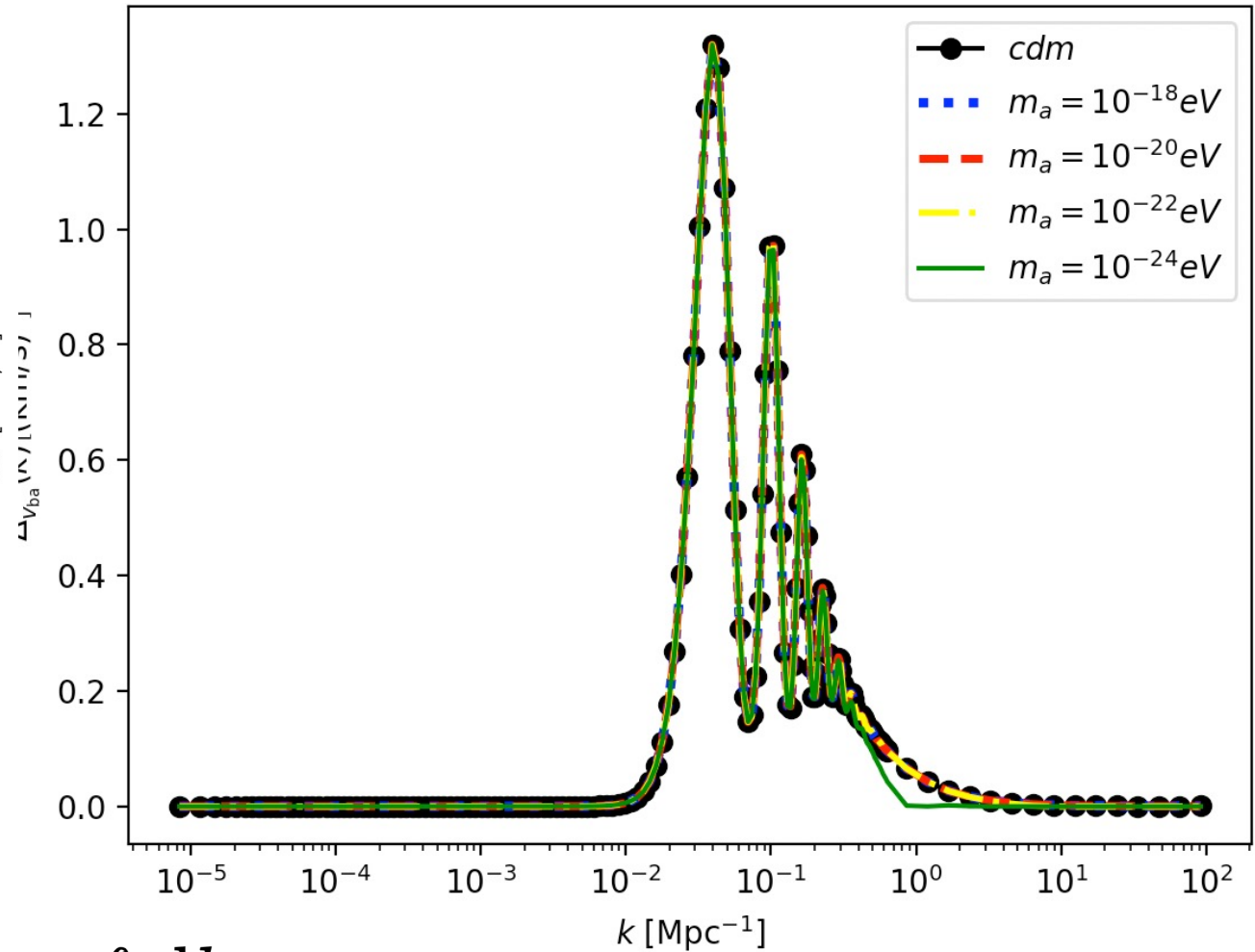
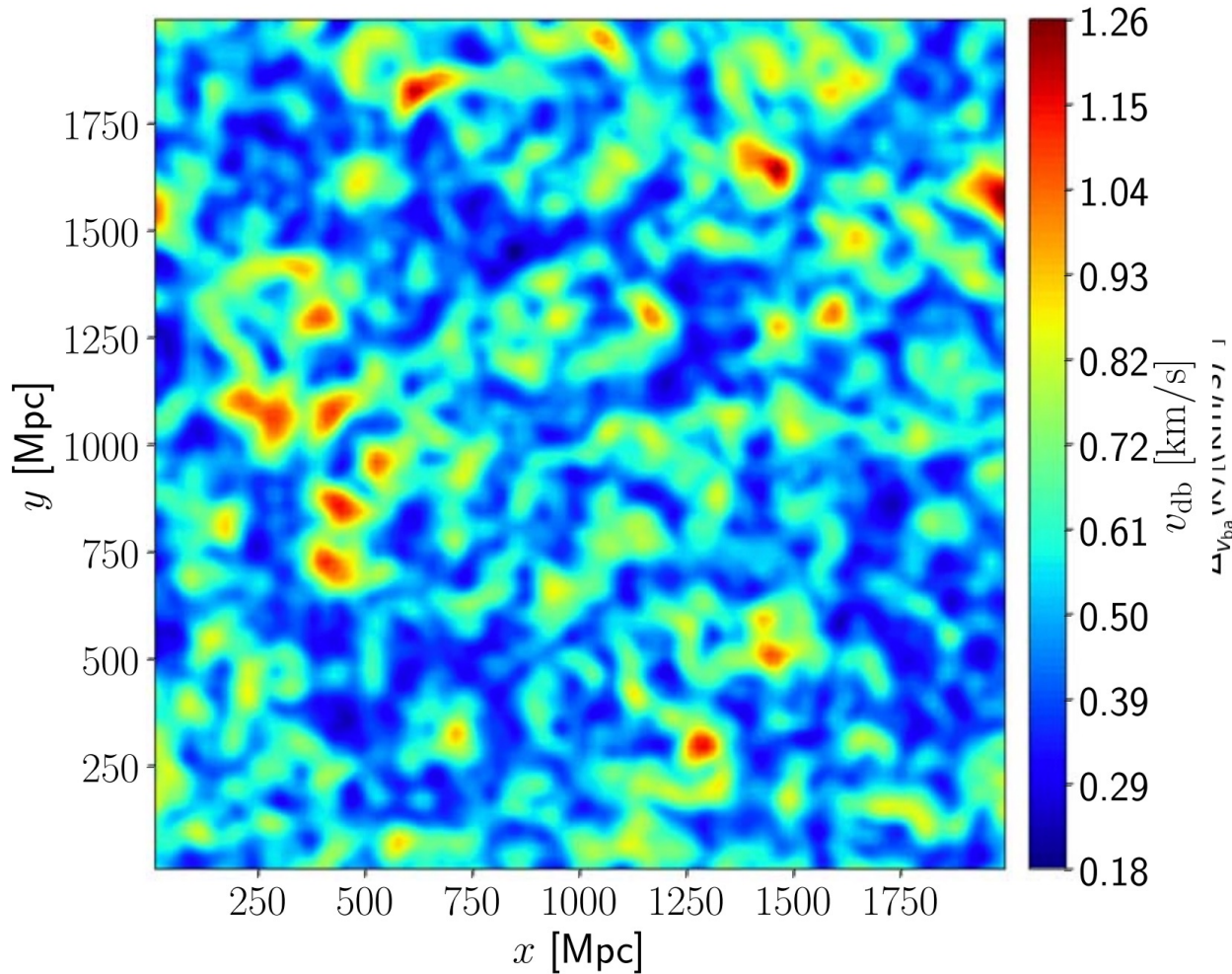
Minihalo distribution (critical mass) allowing the formation of first-generation stars

Radiation field produced by first-generation star formation

VAO feature on the 21 cm power spectrum



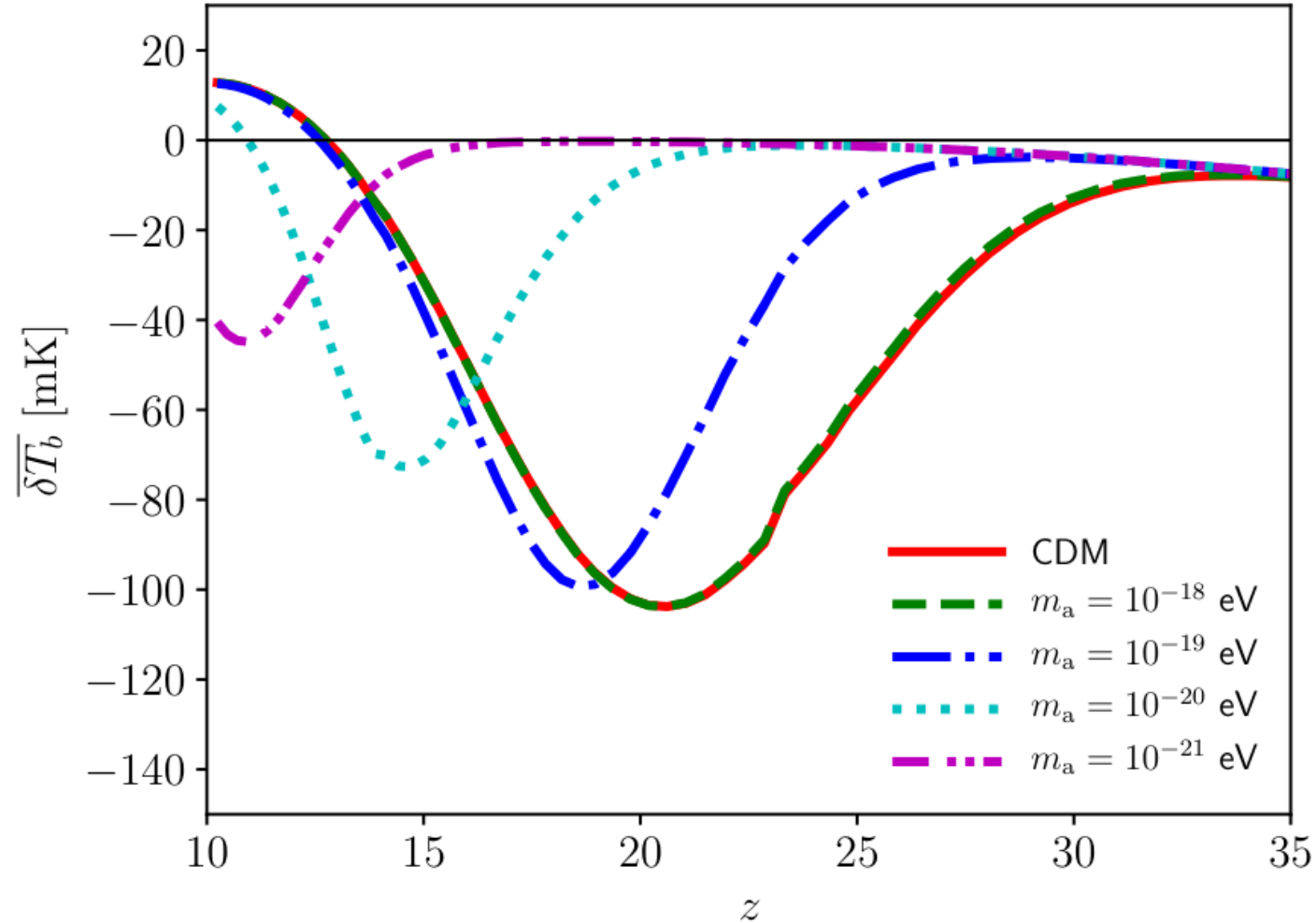
Velocity Acoustic Oscillations(VAO)



Relative velocity field of
baryon and dark matter

$$v_{\text{db}}^2(\mathbf{x}) = \int \frac{dk}{k} \Delta_{\text{vdb}}^2(k) \quad \text{power spectrum}$$

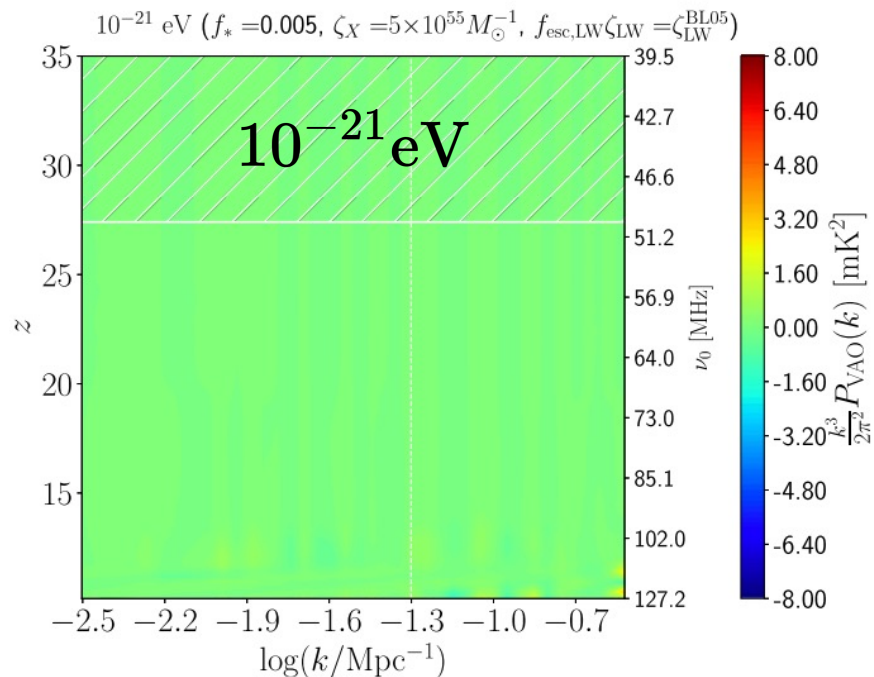
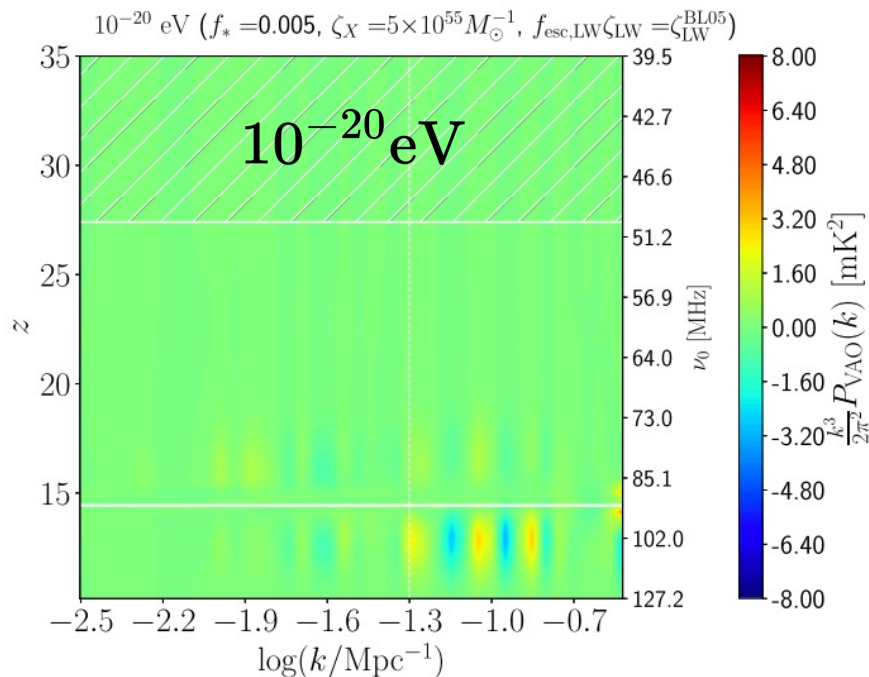
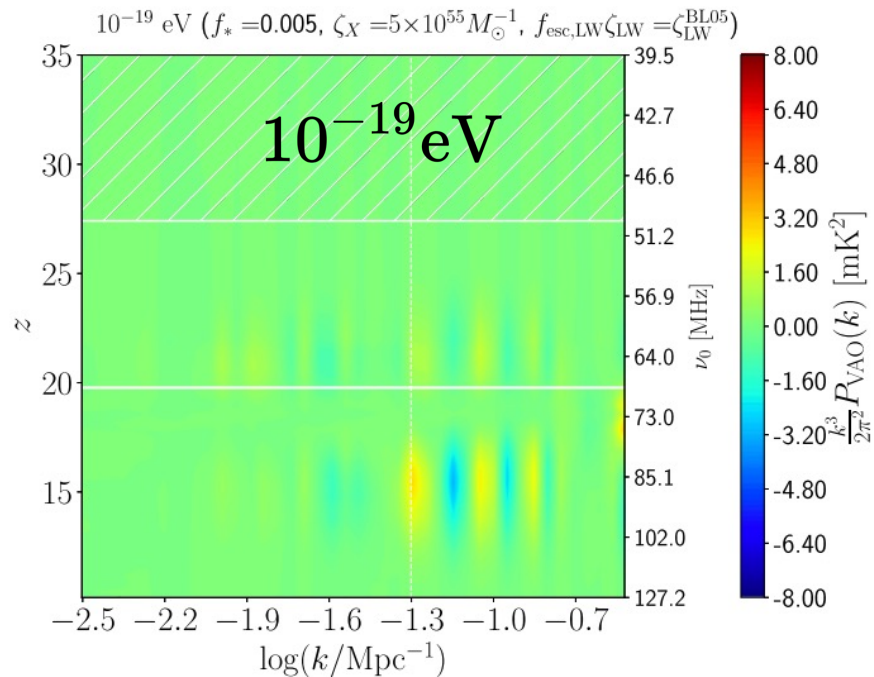
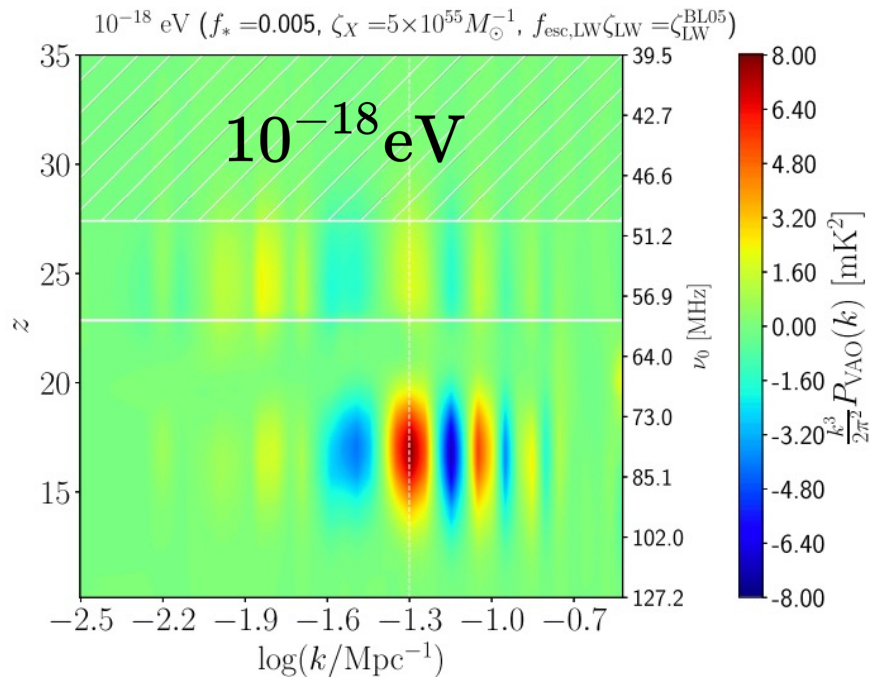
21cm Global spectrum in different dark matter model



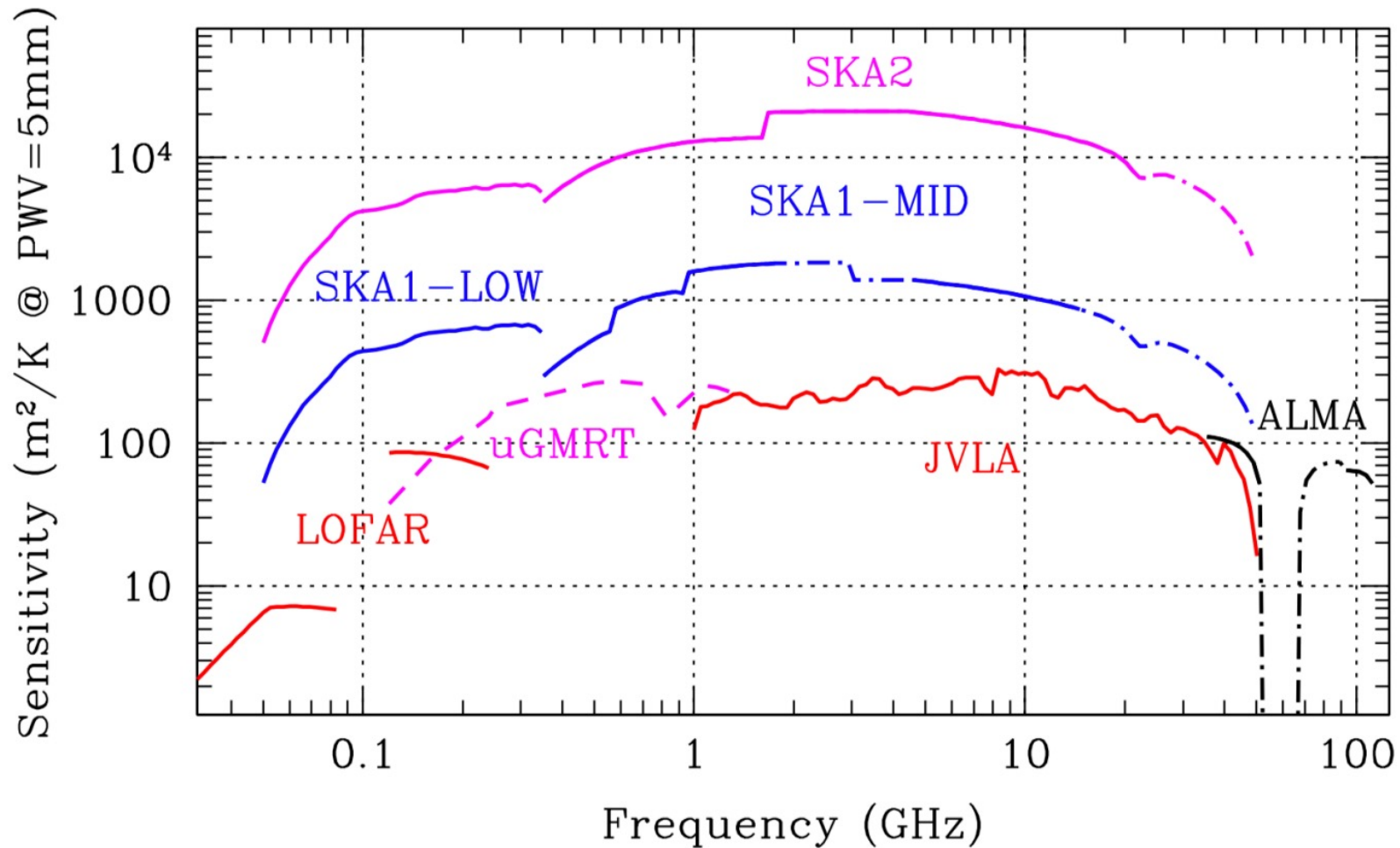
$$\delta T_b = 27 x_{\text{HI}} (1 + \delta) \left(\frac{T_s - T_{\text{CMB}}}{T_s} \right) \times \left(\frac{1+z}{10} \right)^{1/2} \left(\frac{0.15}{\Omega_m h^2} \right)^{1/2} \left(\frac{\Omega_b h^2}{0.023} \right) \text{mK}$$

$$T_s^{-1} = \frac{T_{\text{CMB}}^{-1} + x_\alpha T_k^{-1} + x_c T_k^{-1}}{1 + x_\alpha + x_c}$$

The evolution of VAO signal (For the Axion dark matter model)

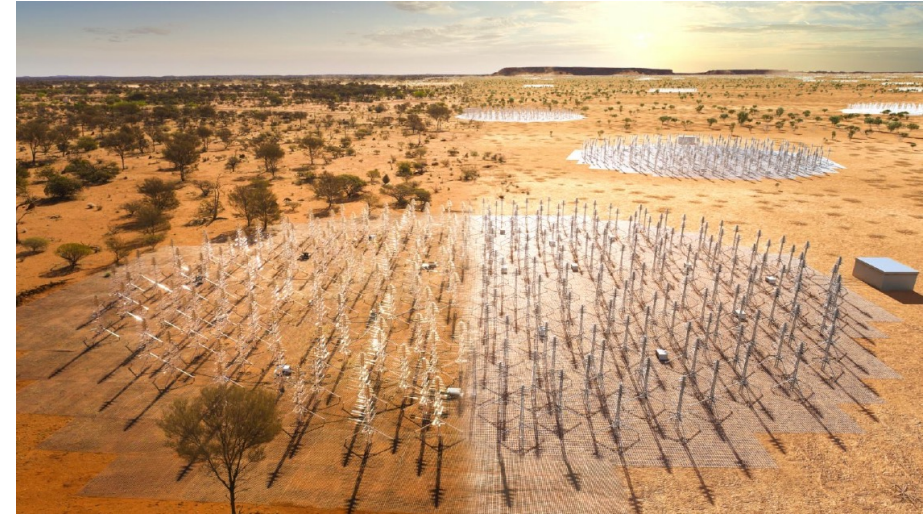


SKA-low: The detectability of VAO signal



Sensitivity parameters of SKA1 and SKA2

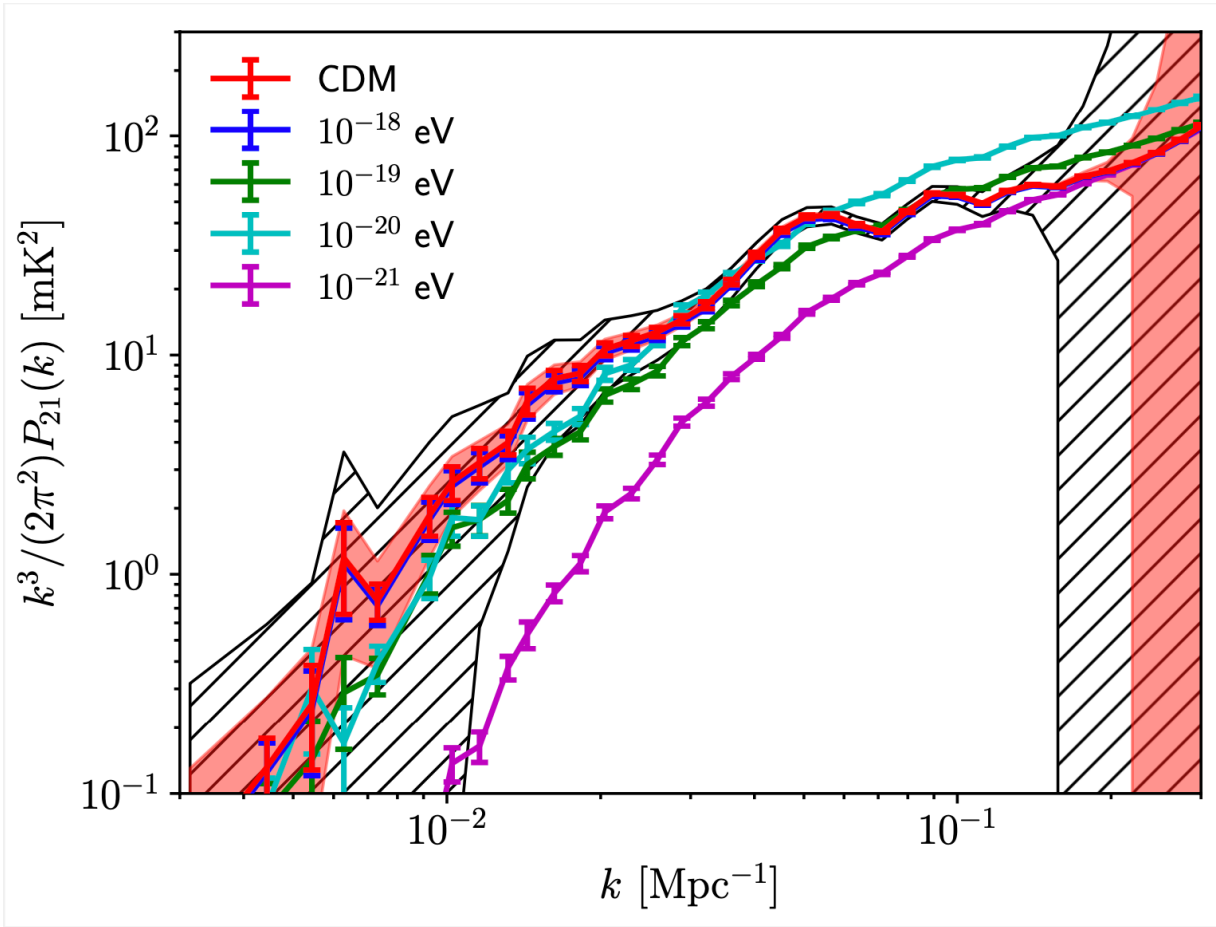
Braun et al. (2019)



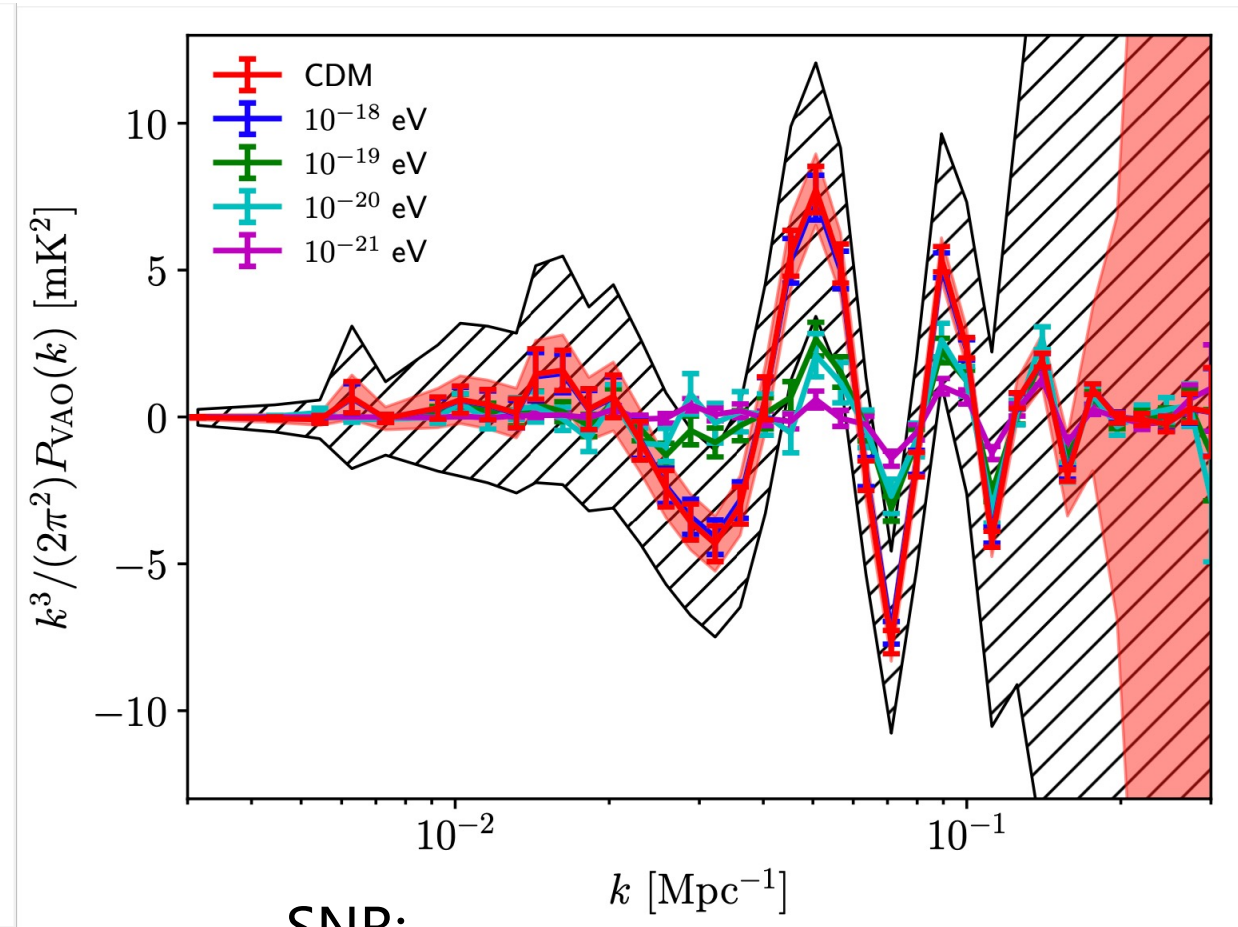
VAO signal and the expected measurement error of SKA

Fit the smooth component of the 21 cm power spectrum by a polynomial:

$$\ln \left(\frac{k^3}{2\pi^2} P_{21}^{\text{smooth}}(k) \right) = \sum_{i=0}^4 c_i(z) (\ln k)^i$$

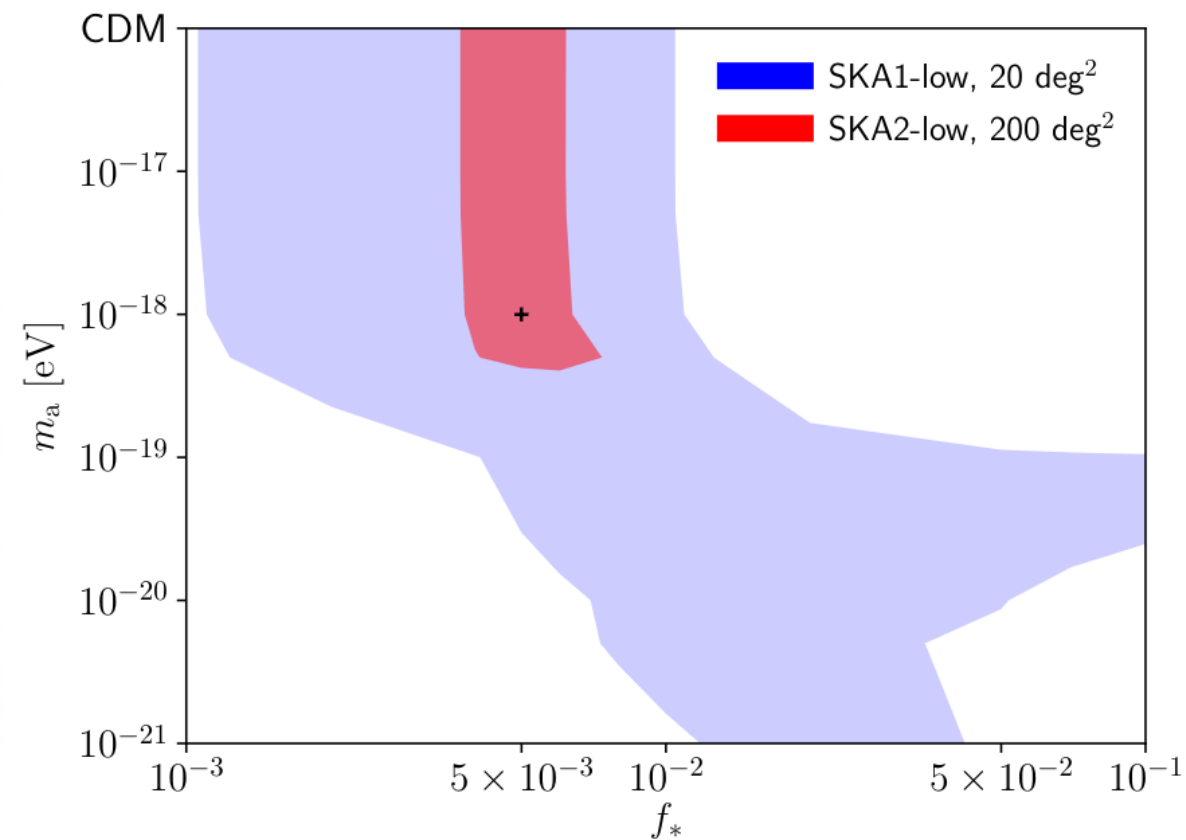
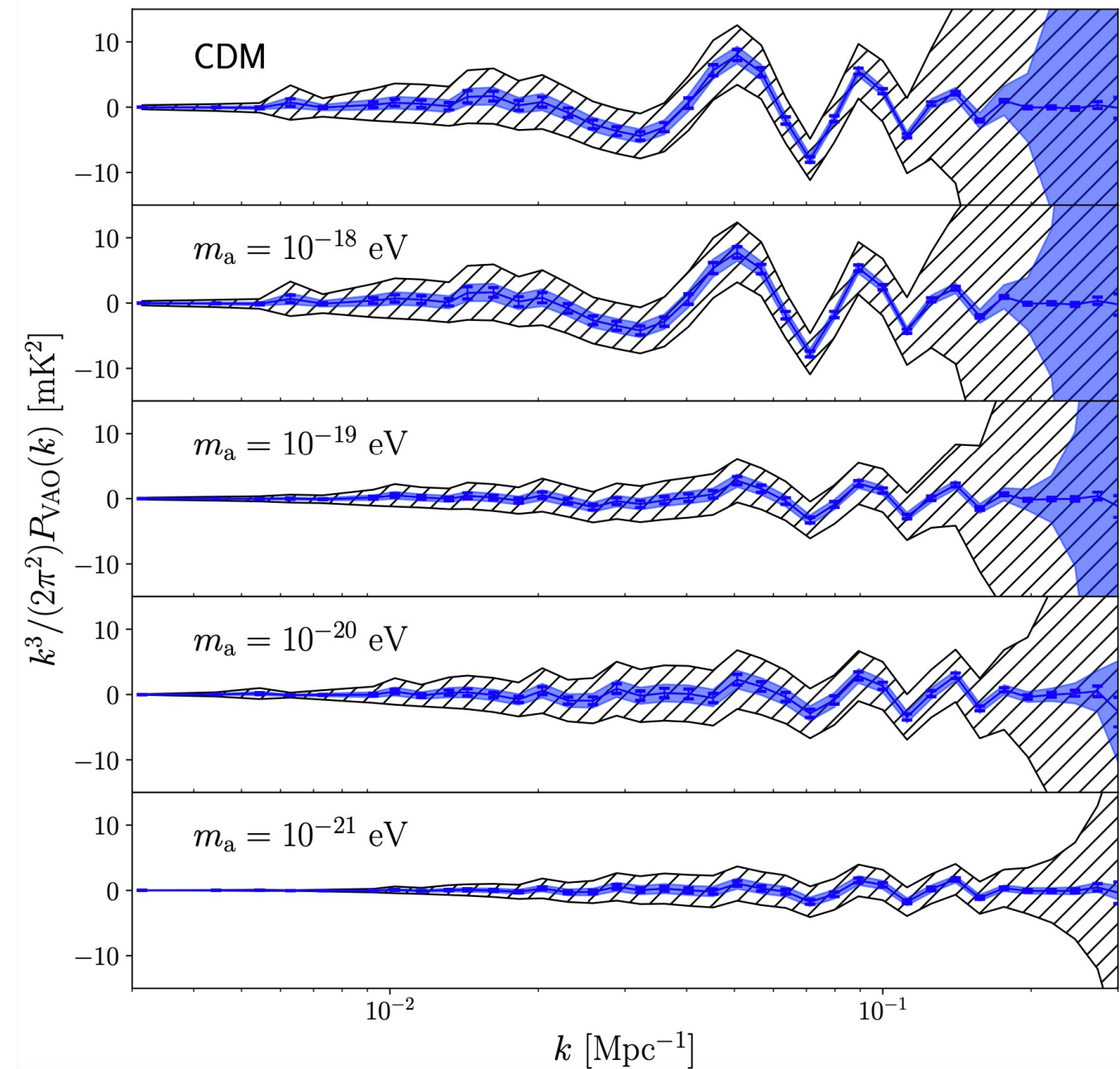


integration time: 2000 hour
survey area: SKA1-low 20deg^2 SKA2-low 200deg^2



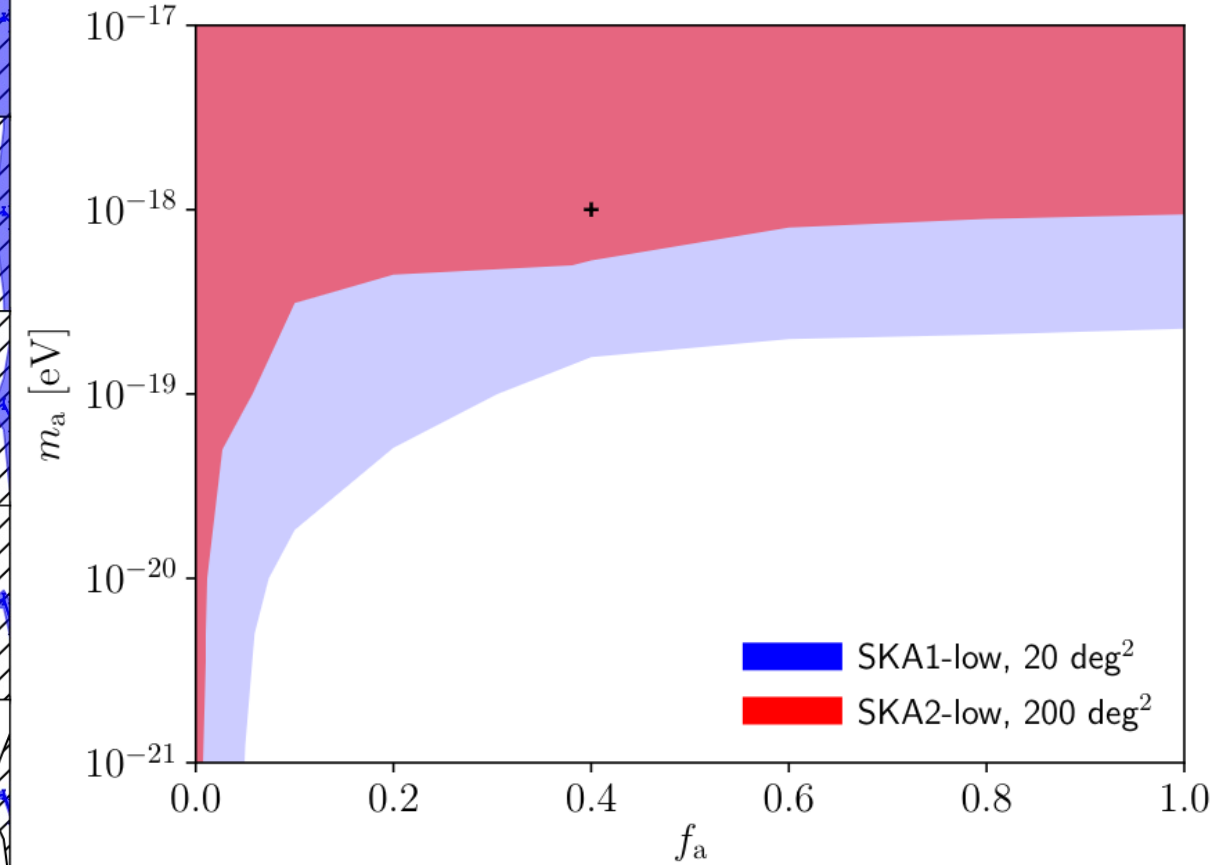
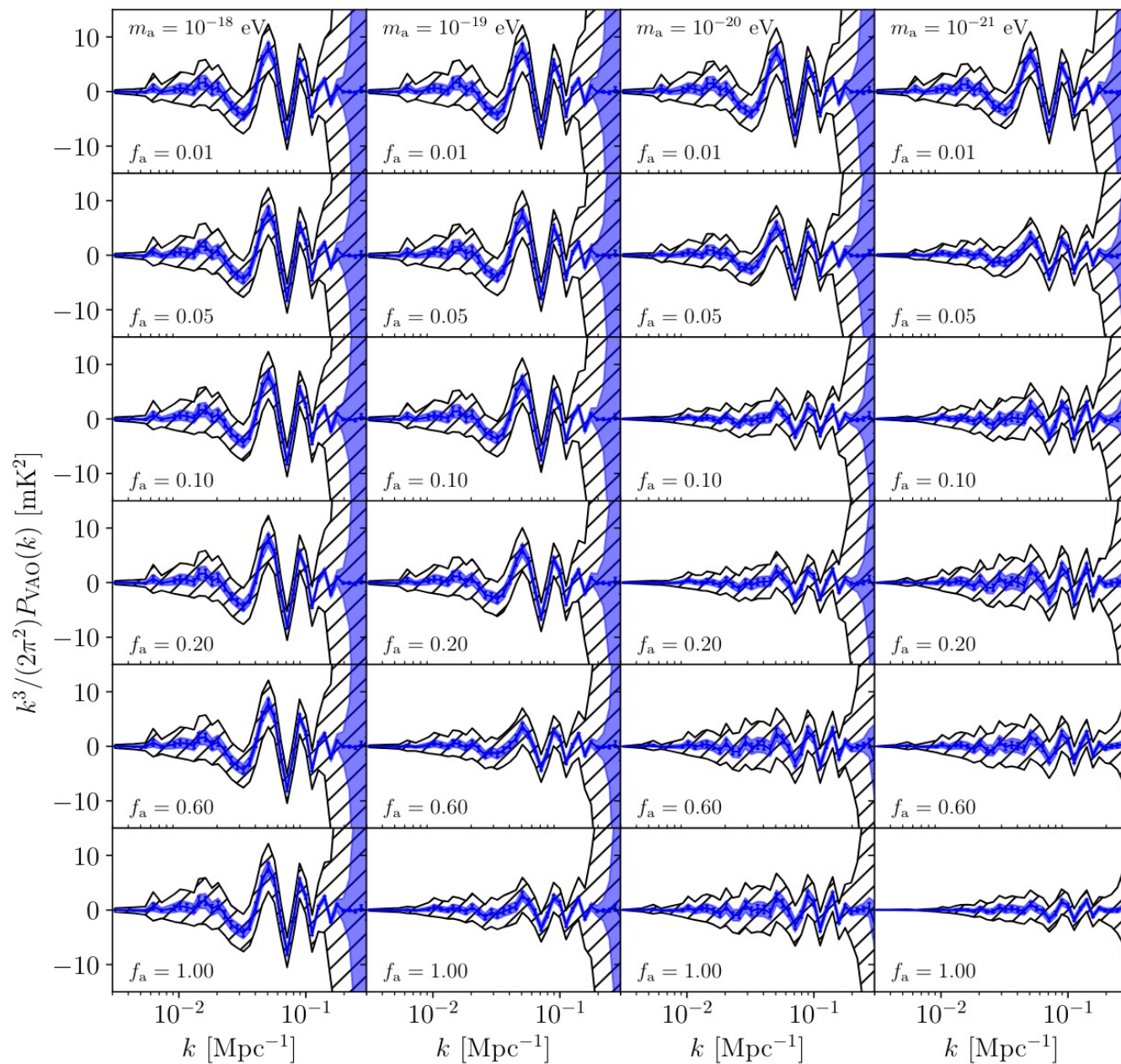
SNR:
SKA1-low: ~ 6
SKA2-low: ~ 22

The VAO signal in various dark matter models



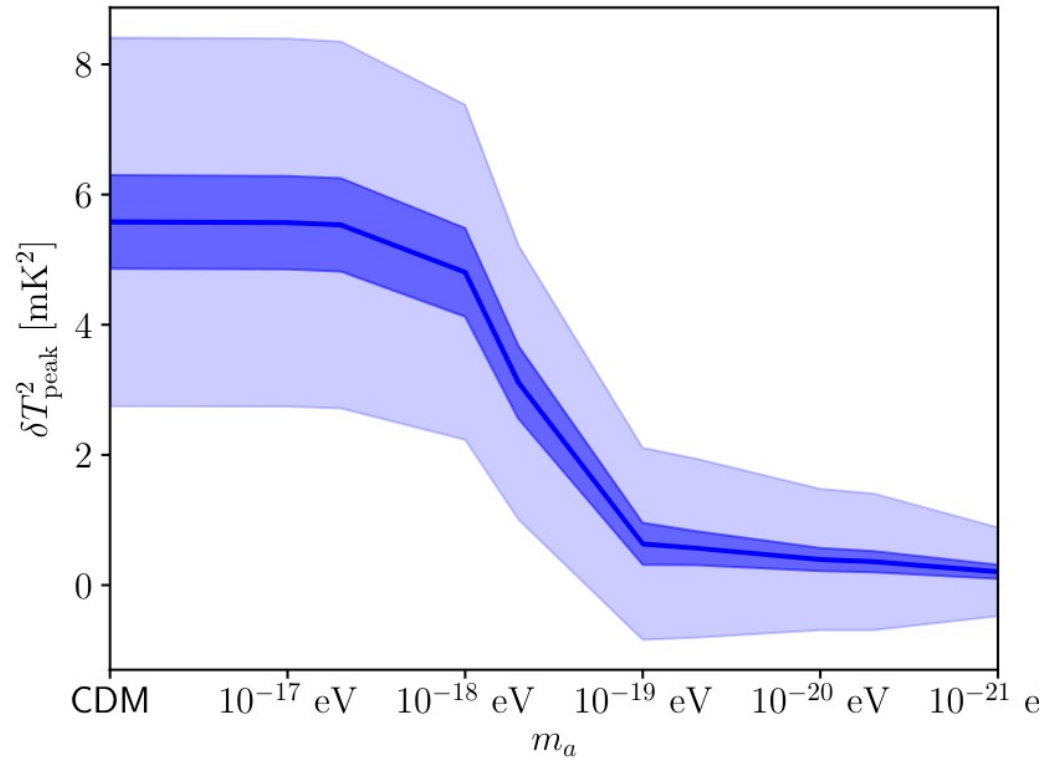
forecast the parameter constraints

The VAO signal in mixed dark matter model



forecast the parameter constraints

The amplitude of the strongest peak vs f_a , for different axion masses



We find this formula can be used to quickly estimate the VAO signal:

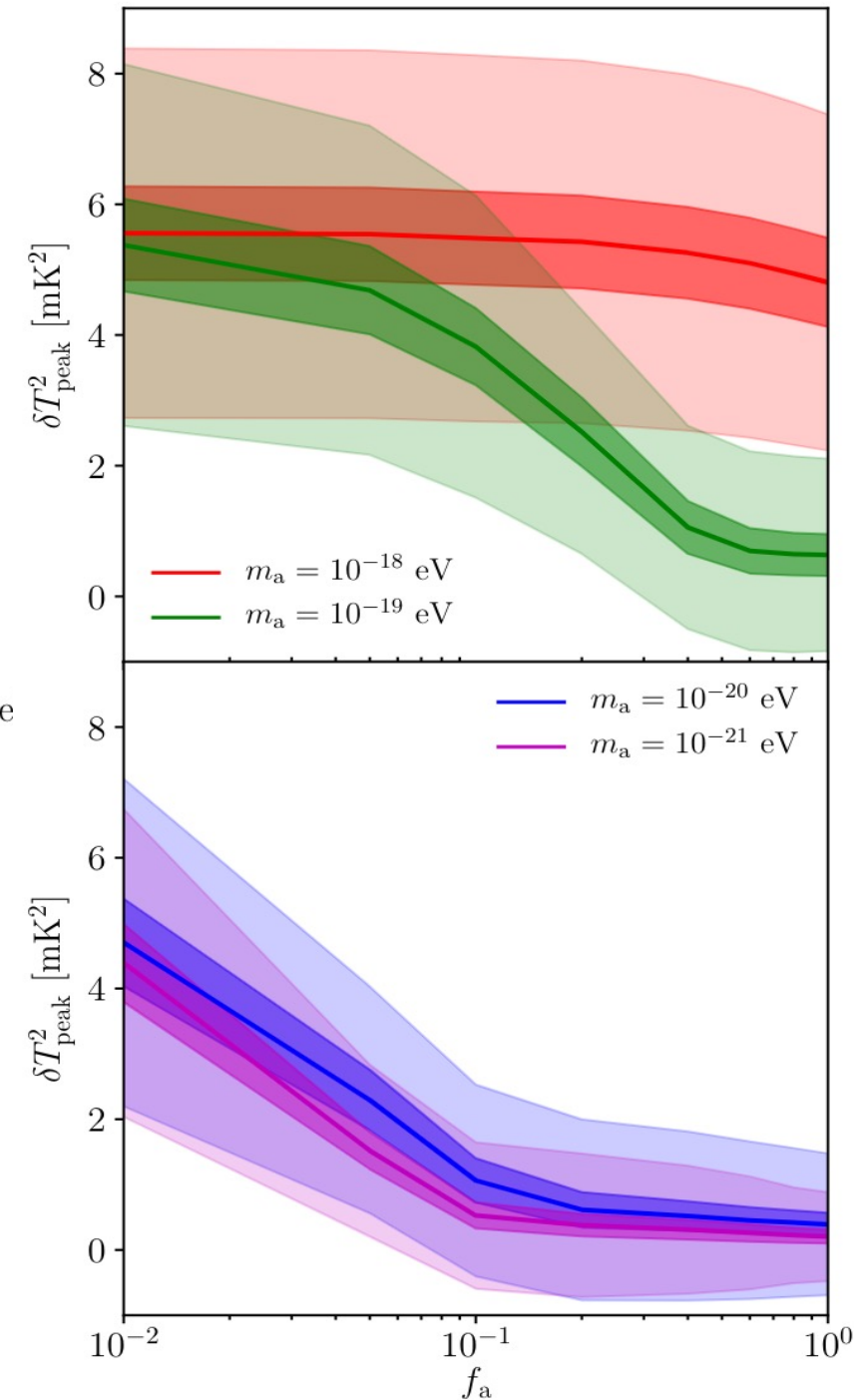
$$\frac{\delta T_{\text{peak}}^2}{(\text{mK})^2} \approx [5.6 - f(m_a)g(f_*)]\tilde{f}_{\text{voigt}}(f_a, \sigma_a, \gamma_a) + f(m_a)g(f_*),$$

$$f(m_a) = 2.8\{1 + \text{erf}[1.35(\log m_a + 18.5)]\},$$

$$g(f_*) = 7.34 + 4.12 \log f_* + 0.59(\log f_*)^2,$$

$$\sigma_a = \left[10^{(-18.3 - \log m_a)}\right]^{0.2} \exp\left[-10^{(-18.3 - \log m_a)}\right],$$

$$\gamma_a = 0.1 \exp\left[-2(-19.0 - \log m_a)^2\right] + 0.03.$$



Summary:

1. The 21cm signal at the cosmic dawn contains important information of cosmology and astrophysics.
2. With the formation of the first-generation stars at the cosmic dawn and the radiation field is build up, the VAO characteristics appear on the 21cm power spectrum in a specific window period.
3. VAO signals can be used as good tools to distinguish whether the dark matter has small scale structure(different dark matter model)
4. SKA1-low and SKA2-low may measure such special 21cm signal characteristic.

Thank You!

The CMB, cosmic gas temperature, and 21 cm spin temperature in the dark ages, cosmic dawn and the epoch of reionization

