

Exotic hadrons with heavy quarks

Part 1: Introduction

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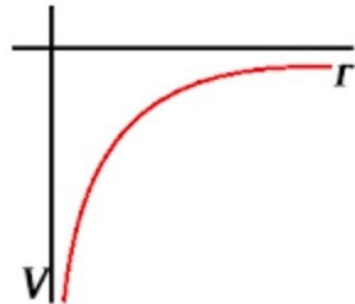
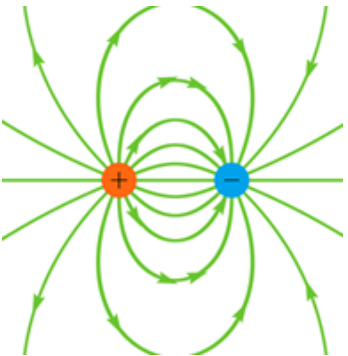
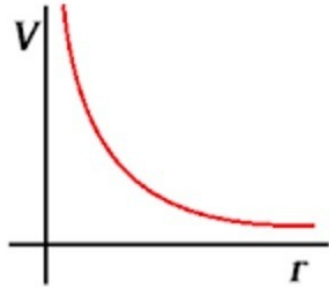
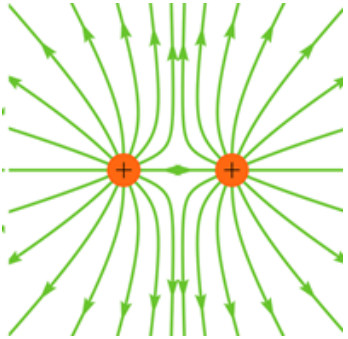


Standard Model of fundamental interactions

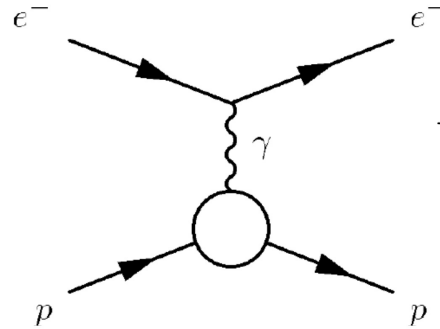
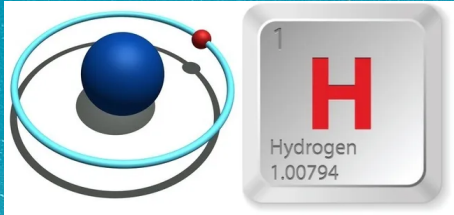
Standard Model of Elementary Particles								
three generations of matter (elementary fermions)			three generations of antimatter (elementary antifermions)			interactions / force carriers (elementary bosons)		
	I	II	III	I	II	III		
mass	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	$\approx 2.2 \text{ MeV}/c^2$	$\approx 1.28 \text{ GeV}/c^2$	$\approx 173.1 \text{ GeV}/c^2$	0	$\approx 124.97 \text{ GeV}/c^2$
charge	$\frac{2}{3}$	$\frac{2}{3}$	$\frac{2}{3}$	$-\frac{2}{3}$	$-\frac{2}{3}$	$-\frac{2}{3}$	0	0
spin	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	0
QUARKS	u up	c charm	t top	$\bar{u}$ antiup	$\bar{c}$ anticharm	$\bar{t}$ antitop	g gluon	H higgs
	d down	s strange	b bottom	$\bar{d}$ antidown	$\bar{s}$ antistrange	$\bar{b}$ antibottom	γ photon	
	e electron	μ muon	τ tau	e^+ positron	$\bar{\mu}$ antimuon	$\bar{\tau}$ antitau	Z Z ⁰ boson	
LEPTONS	ν_e electron neutrino	ν_μ muon neutrino	ν_τ tau neutrino	$\bar{\nu}_e$ electron antineutrino	$\bar{\nu}_\mu$ muon antineutrino	$\bar{\nu}_\tau$ tau antineutrino	W^+ W ⁺ boson	W^- W ⁻ boson
	$< 2.2 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 18.2 \text{ MeV}/c^2$	$< 2.2 \text{ eV}/c^2$	$< 0.17 \text{ MeV}/c^2$	$< 18.2 \text{ MeV}/c^2$	$\approx 80.39 \text{ GeV}/c^2$	$\approx 80.39 \text{ GeV}/c^2$
	0	0	0	0	0	0	1	-1
	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	$\frac{1}{2}$	1	1
							GAUGE BOSONS VECTOR BOSONS	
							SCALAR BOSONS	

Quantum Electrodynamics

$$L_{\text{QED}} = -\frac{1}{4}F_{\mu\nu}F^{\mu\nu} + \bar{\psi}(i\gamma_{\mu}\partial^{\mu} - m - e\gamma_{\mu}A^{\mu})\psi$$



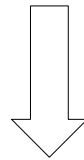
Bound states in QED



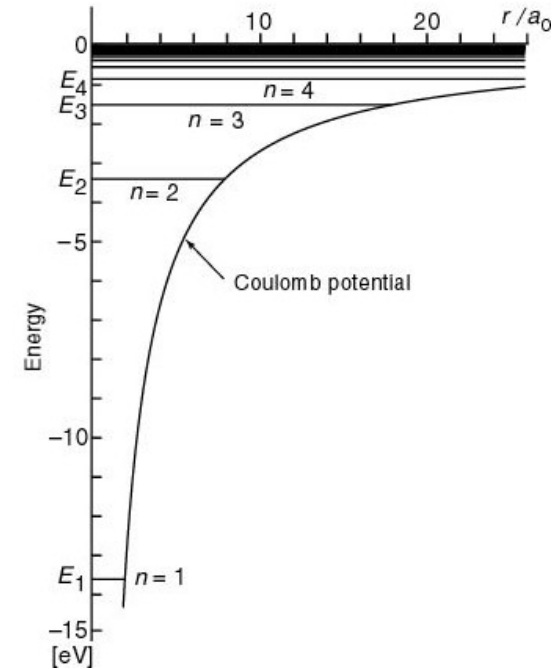
$$V(r) = -\frac{\alpha}{r}$$

$$\alpha = \frac{e^2}{4\pi\hbar c} \approx \frac{1}{137} \ll 1$$

$$E = \frac{m}{\sqrt{1 + \frac{\alpha^2}{\left(n_r + \sqrt{(j + \frac{1}{2})^2 - \alpha^2}\right)^2}}}$$

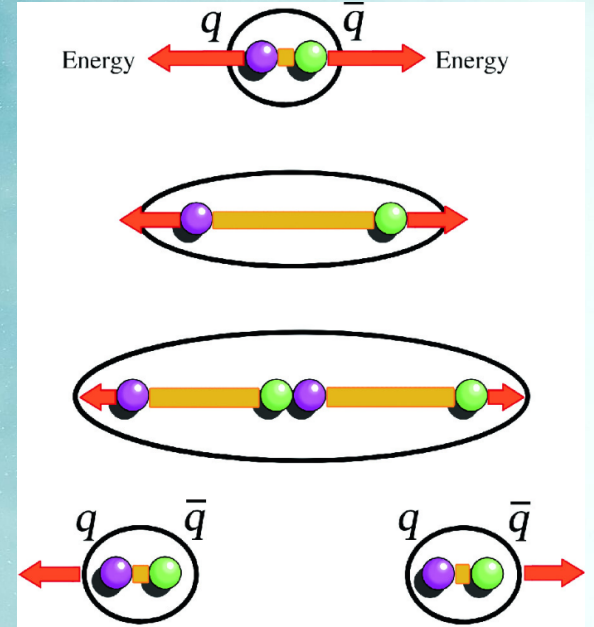
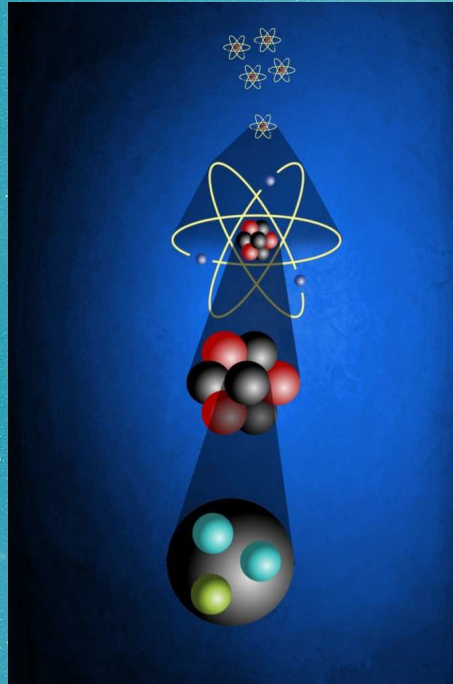
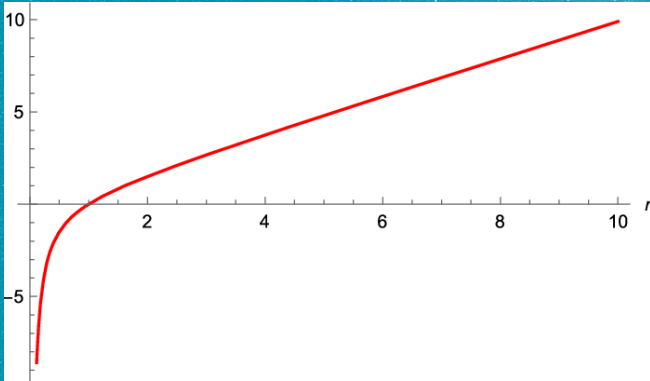
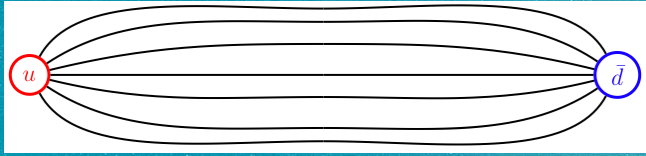


$$E = m - \frac{m\alpha^2}{2(n_r + j + \frac{1}{2})^2} = m - \frac{m\alpha^2}{2n^2}$$



Quantum Chromodynamics

$$L_{\text{QCD}} = -\frac{1}{4}F_{\mu\nu}^a F_a^{\mu\nu} + \bar{\psi}_\alpha (i\gamma_\mu \partial^\mu - m - g\gamma_\mu A_a^\mu t^a)_\beta^\alpha \psi^\beta$$



Bound states in QCD - Hadrons

1964 – Invention of constituent Quark Model



Murray Gell-Mann and George Zweig



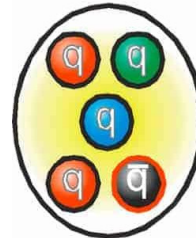
Quark-antiquark meson



Three-quark baryon

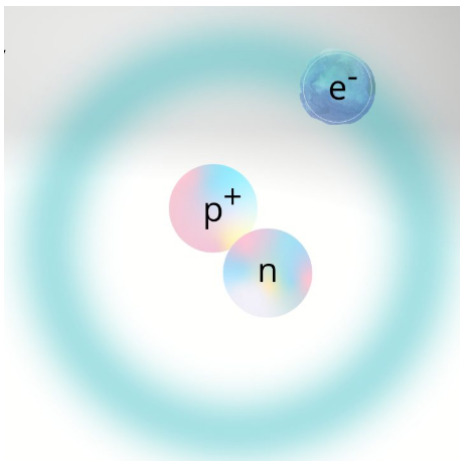
} Ordinary
hadrons

Exotic hadrons:



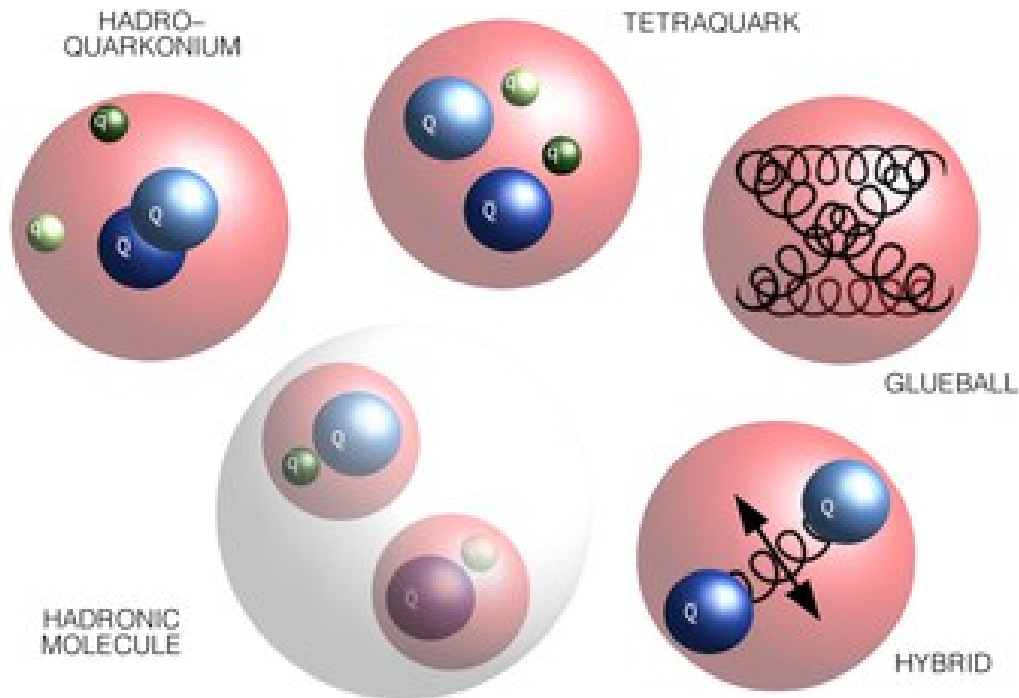
Hadronic Exotica

Deuterium

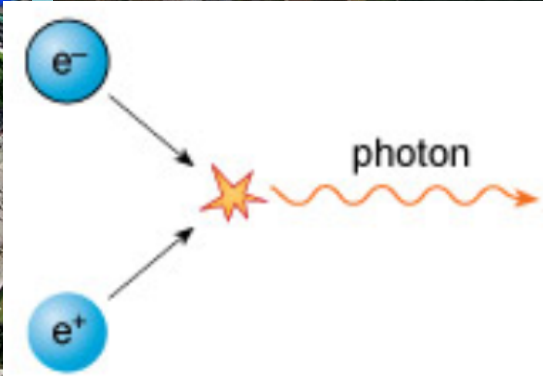
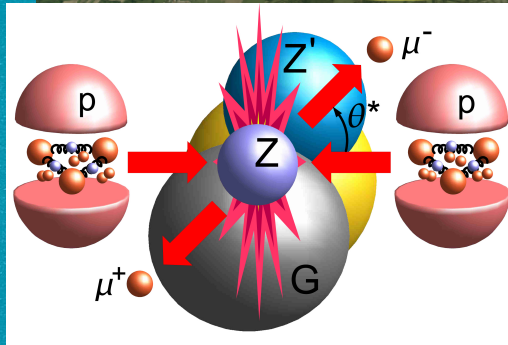
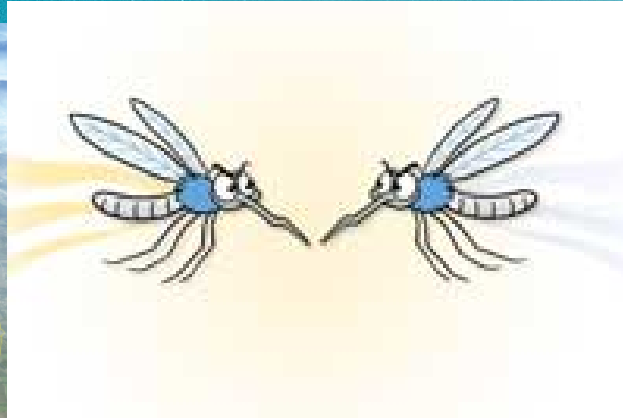
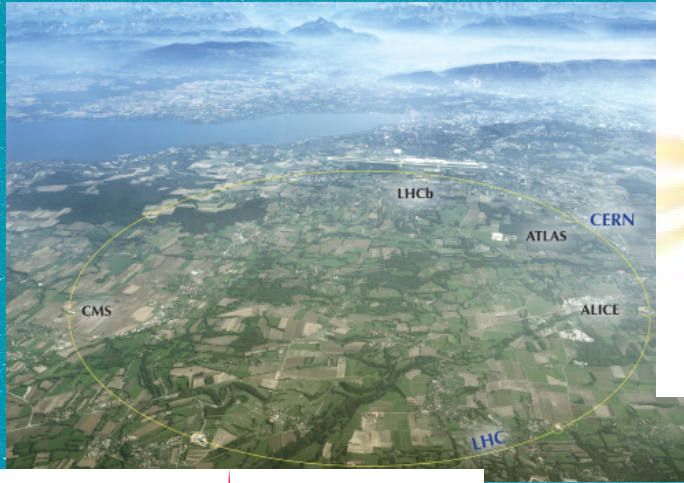


Nucleus is deuteron:
bound state of
proton & neutron

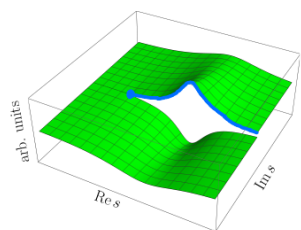
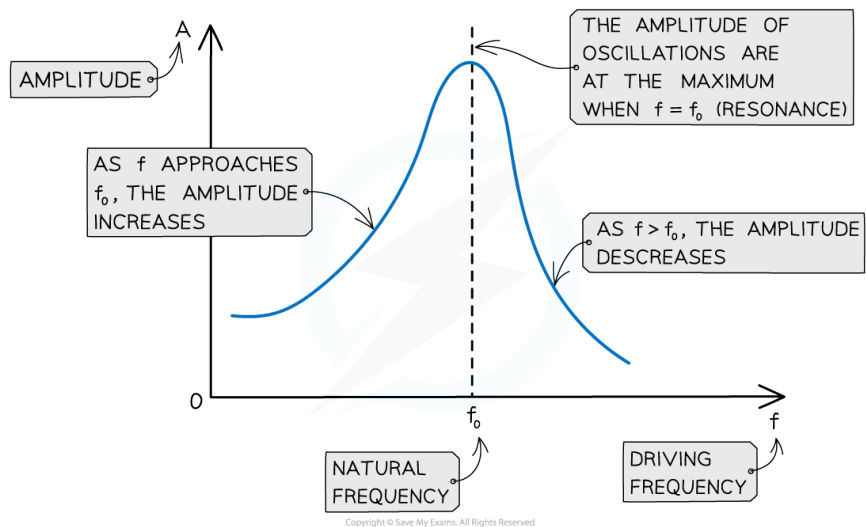
Known for nearly 100 years



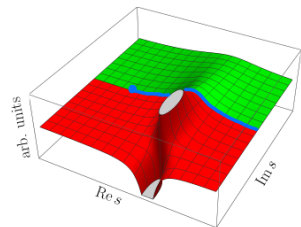
Contemporary collider experiments



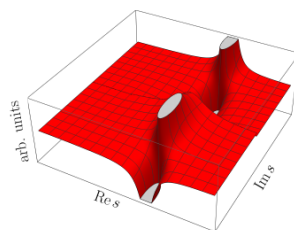
Resonance



first Riemann sheet



transition from first to second Riemann sheet



second Riemann sheet

$J/\psi(1S)$

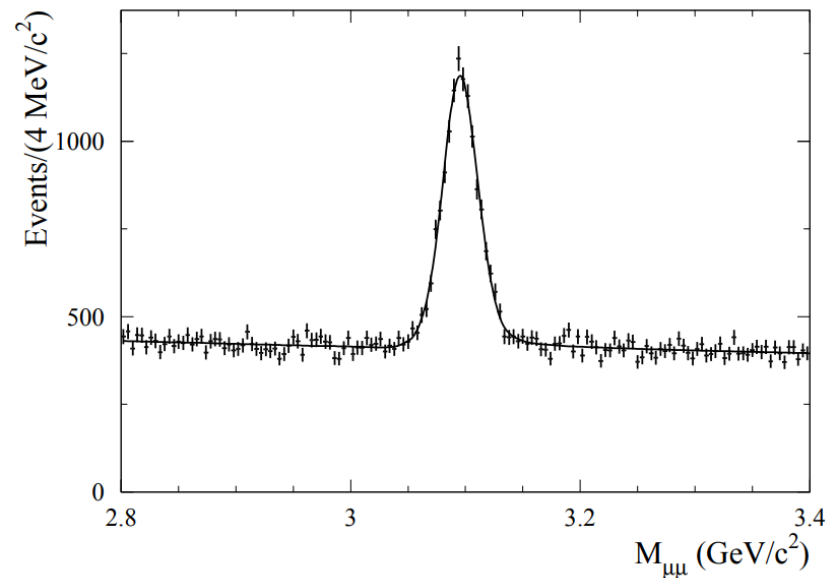
$$I^G(J^{PC}) = 0^-(1^{--})$$

$J/\psi(1S)$ MASS

$$3096.900 \pm 0.006 \text{ MeV}$$

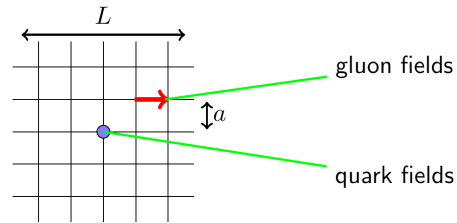
$J/\psi(1S)$ WIDTH

$$92.6 \pm 1.7 \text{ keV } (S = 1.1)$$

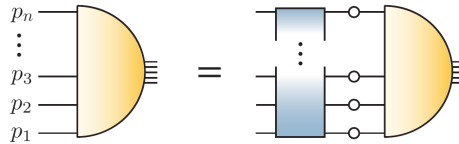


Theoretical Approaches

QCD on lattice



Functional Methods



Quark Models



$$H_0 = \sqrt{p^2 + m_1^2} + \sqrt{p^2 + m_2^2} + V_0(r) + V_{SD}(r)$$

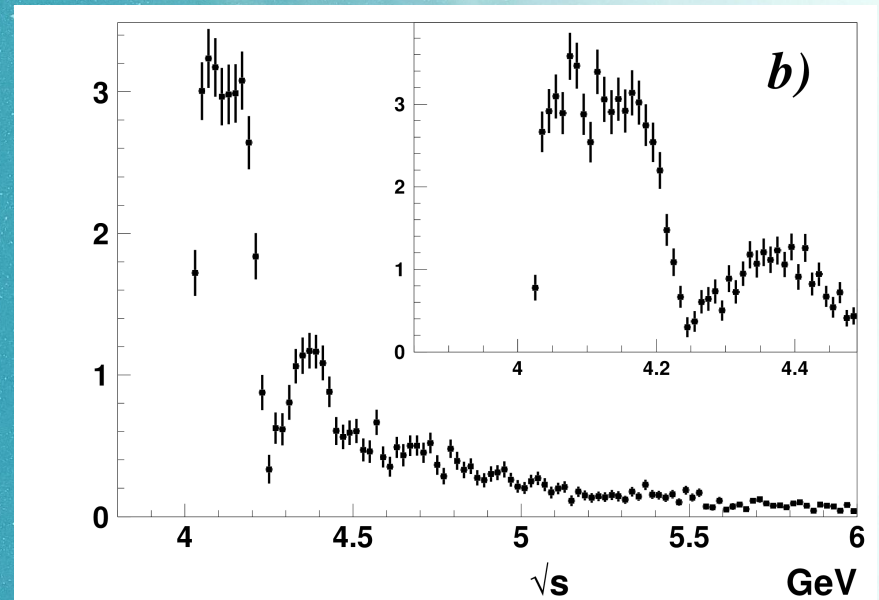
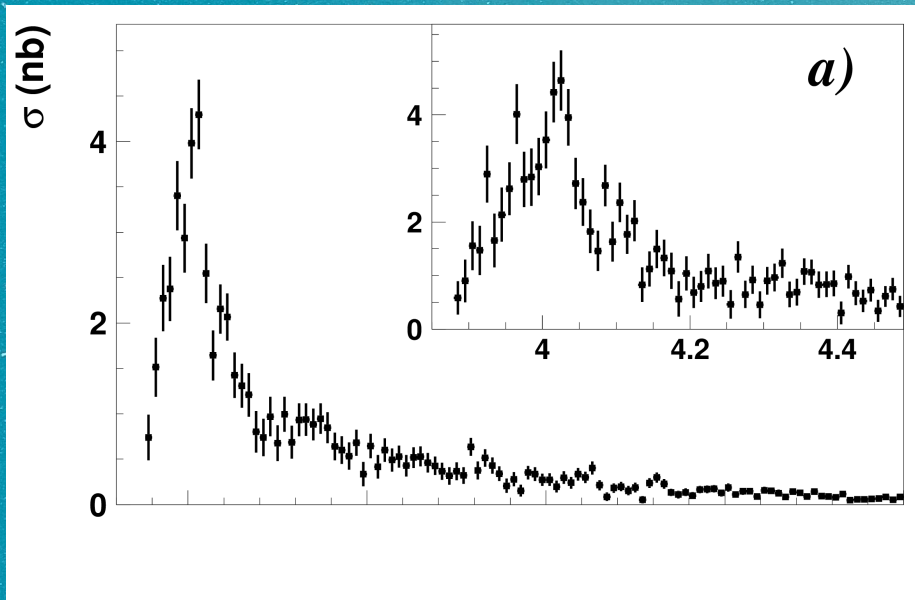
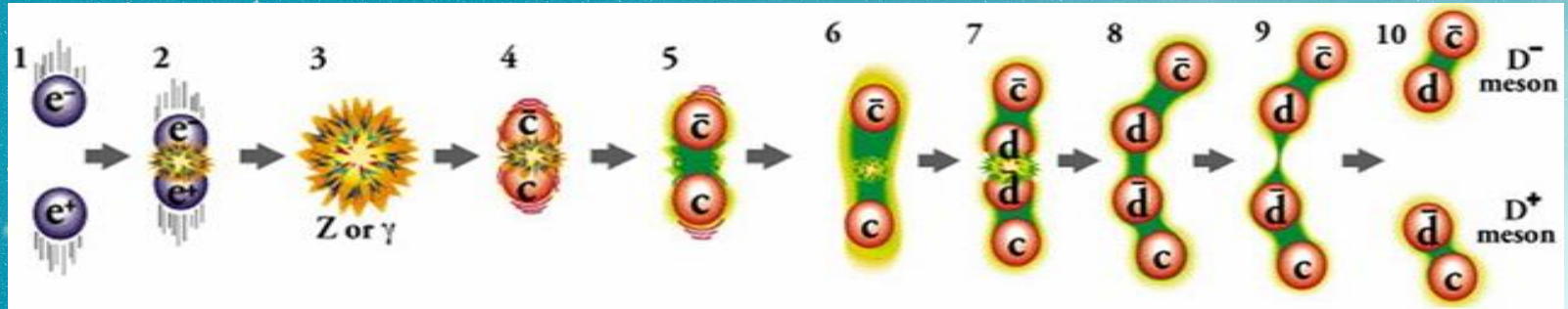
$$V_0(r) = \sigma r - \frac{\frac{4}{3}\alpha_s}{r} + C_0$$

XYZ states

In particle physics, XYZ particles (also referred to as XYZ states) are recently-discovered heavy mesons whose properties do not appear to fit the standard picture of charmonium and bottomonium states. They are therefore types of exotic mesons. The term arises from the names given to some of the first such particles discovered: X(3872), Y(4260) and Zc(3900), although the symbols X and Y have since been deprecated by the Particle Data Group.

Wikipedia

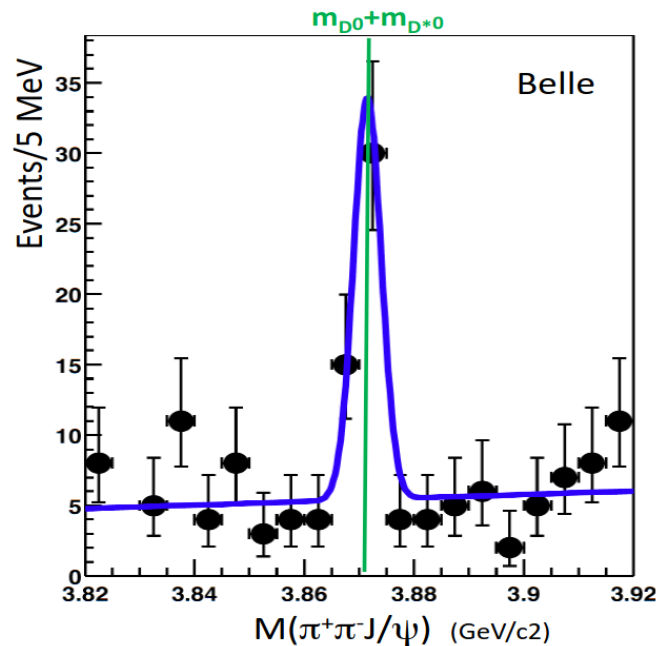
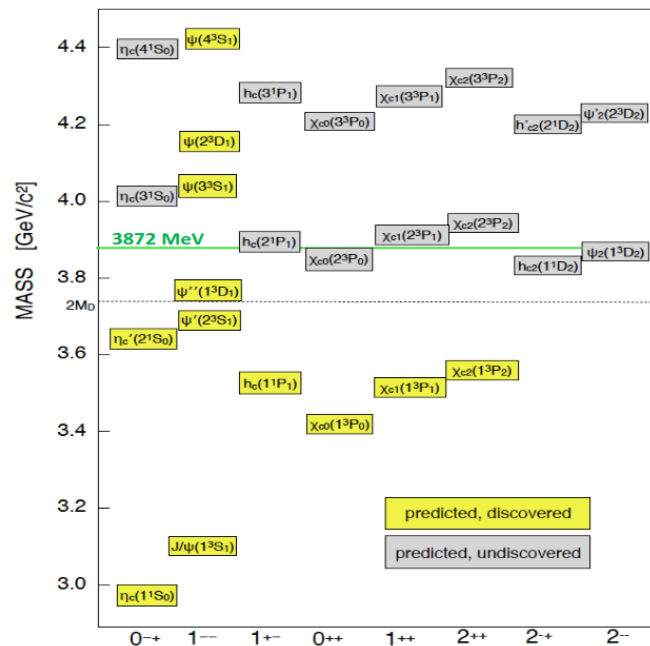
Open-flavour thresholds



X(3872)

$$M_{X(3872)} = 3871.65 \pm 0.06 \text{ MeV}$$

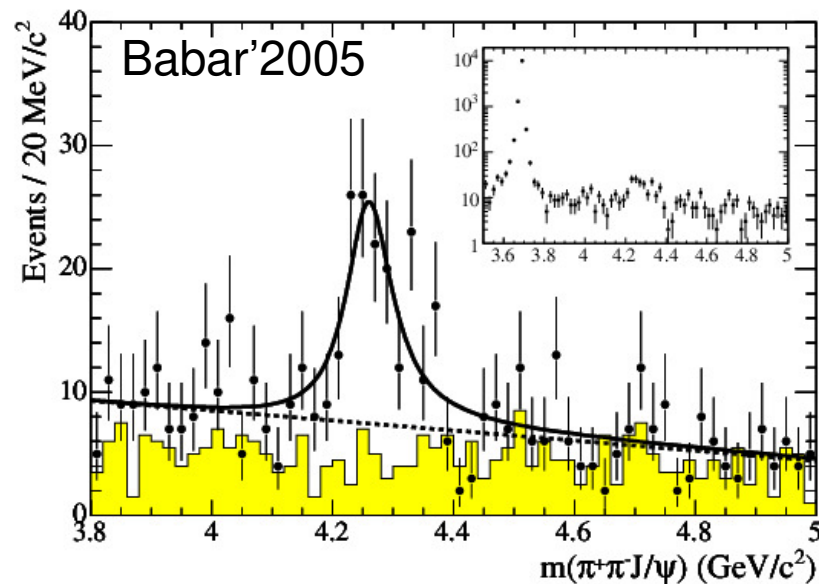
$$\Gamma_{X(3872)} = 1.19 \pm 0.21 \text{ MeV}$$



$$|X\rangle = \lambda|\bar{c}c\rangle + \chi(p)|D\bar{D}^*\rangle + \dots$$

$$\psi_{D\bar{D}^*} \underset{r \rightarrow \infty}{\sim} e^{-\kappa_B r} \quad \kappa_B = \sqrt{2\mu E_B}$$

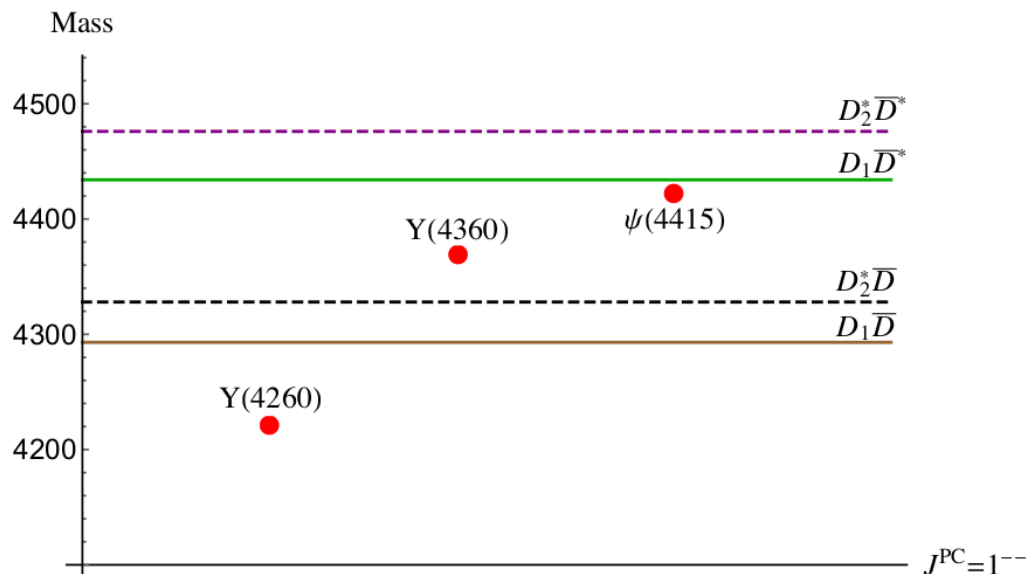
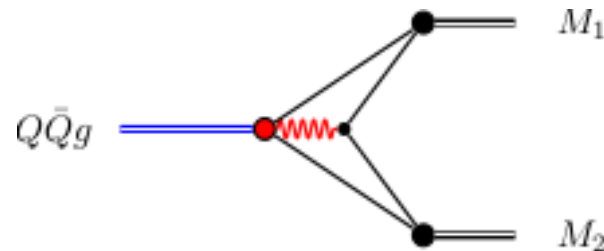
Y(4260)/Y(4230)



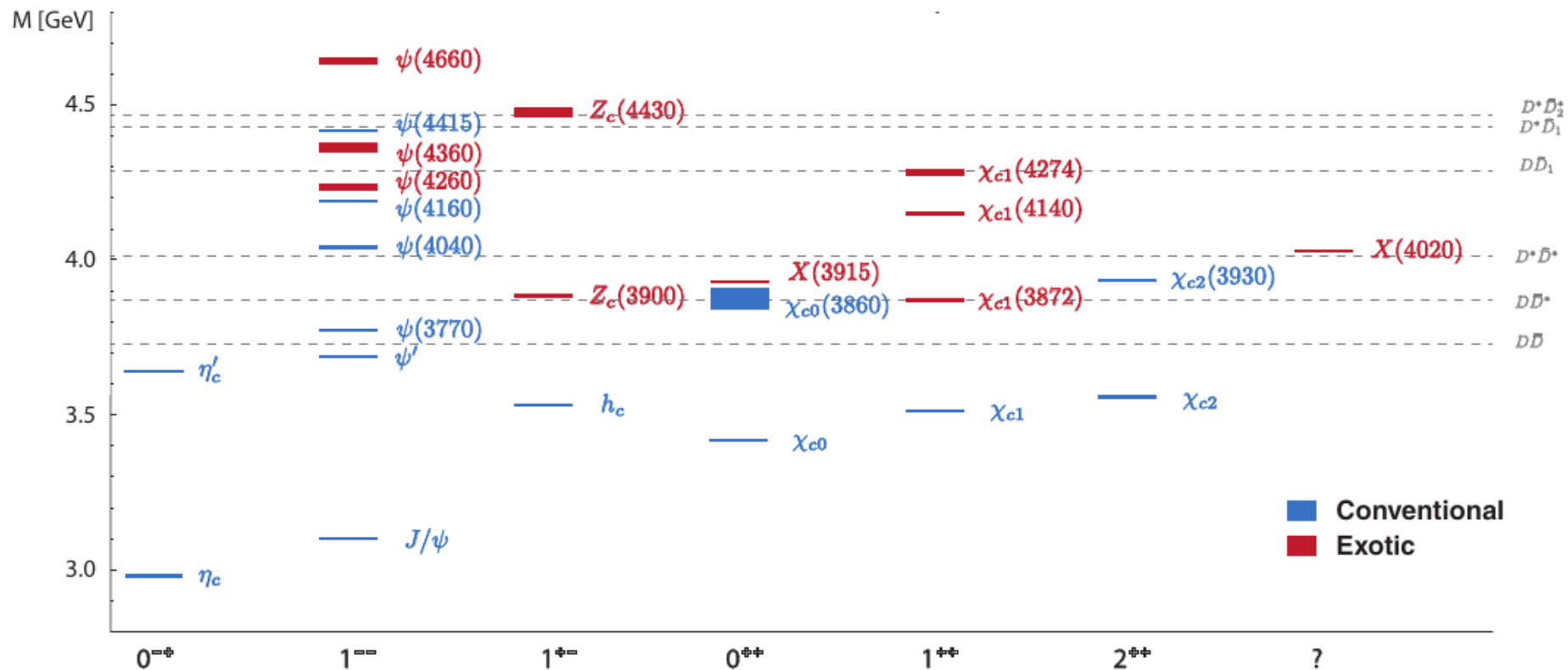
Molecule?



Hybrid?

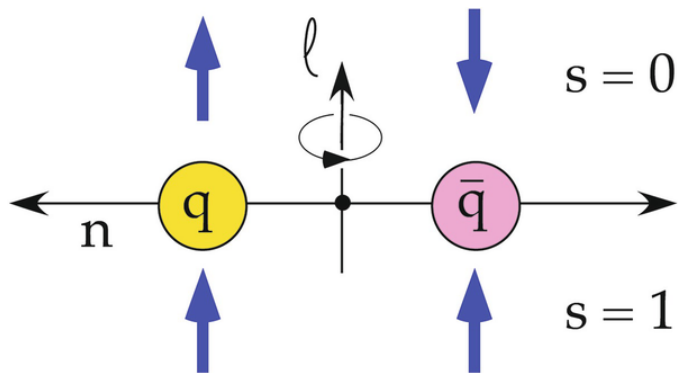


Charmonium

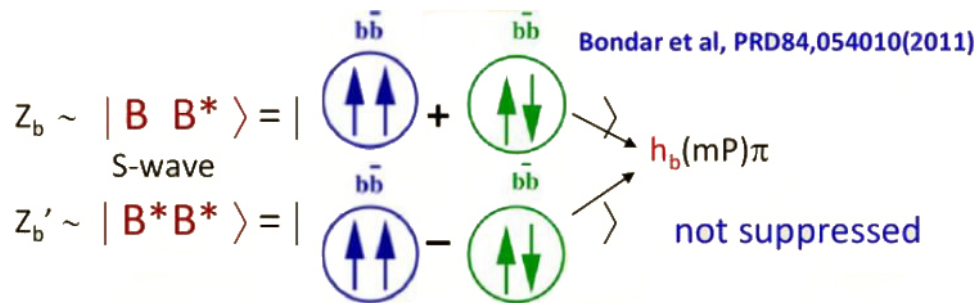
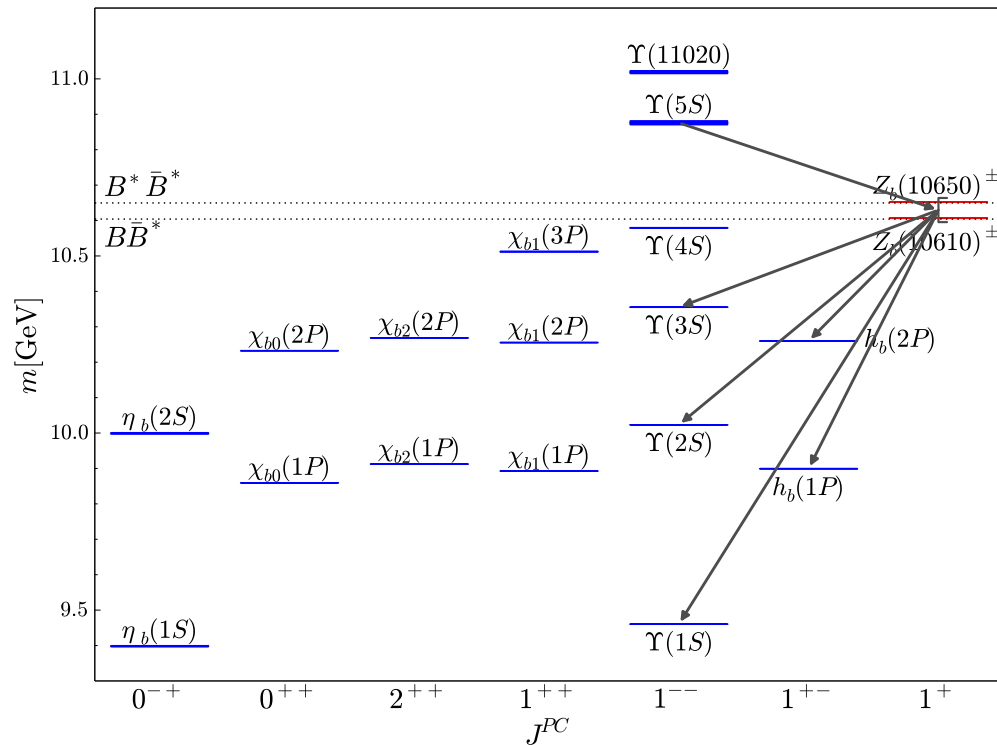


Charged Z states

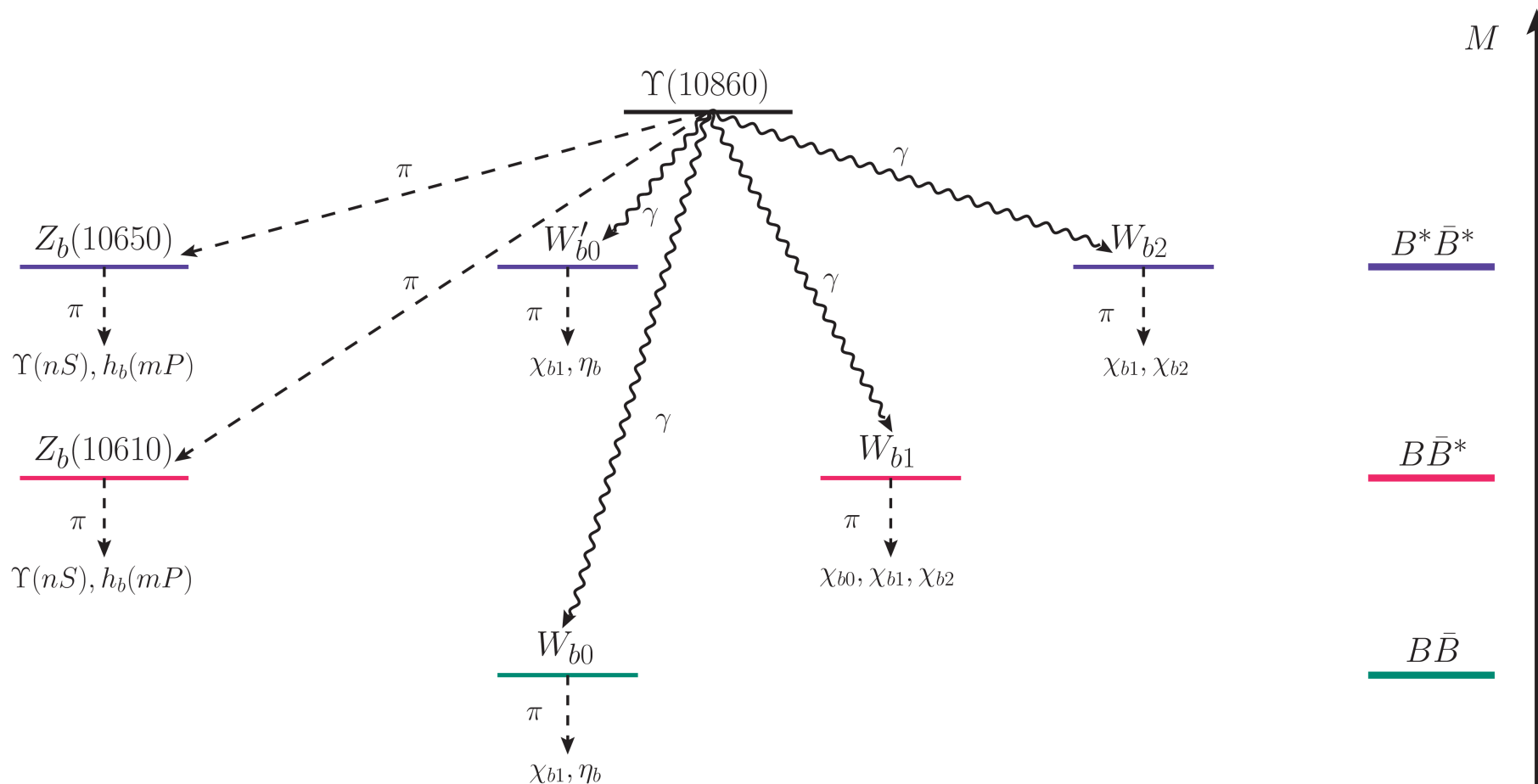
To flip or not to flip?

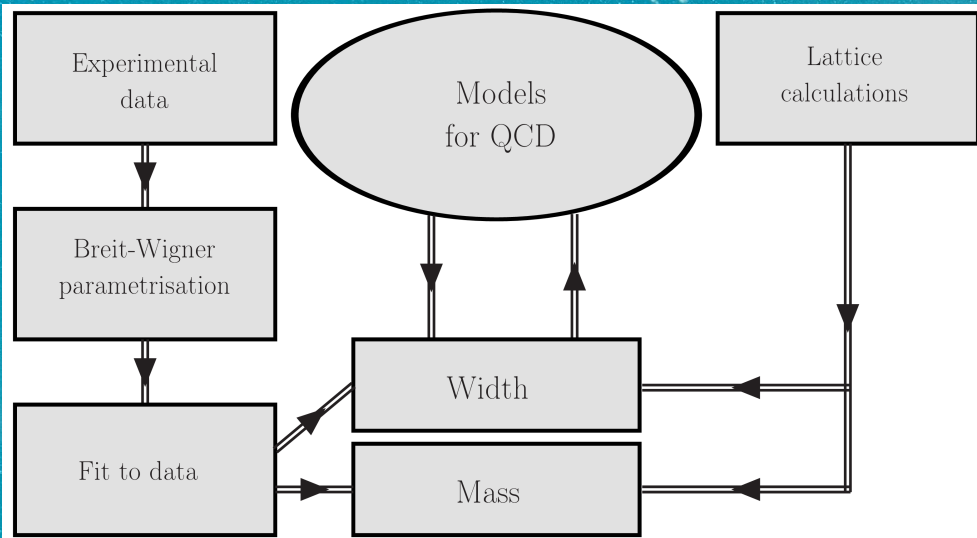


$$\delta H_{\text{flip}} \propto \frac{\Lambda_{\text{QCD}}}{m_b} \simeq \frac{0.25 \text{ GeV}}{4.5 \text{ GeV}} \ll 1$$

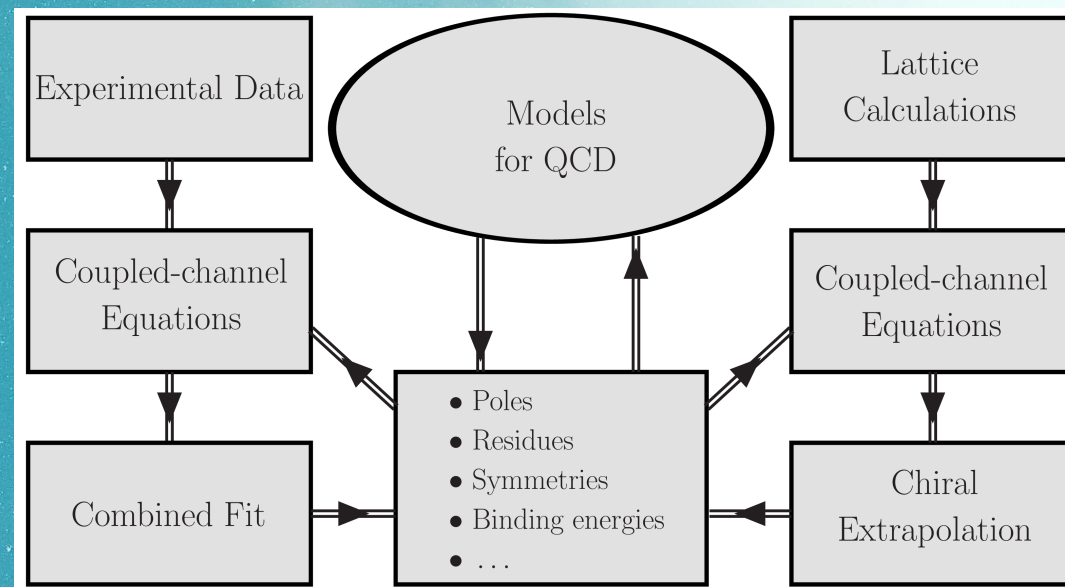


Spin partners

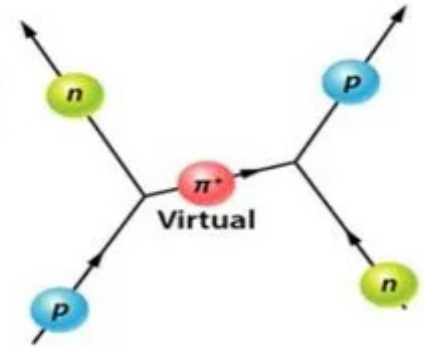
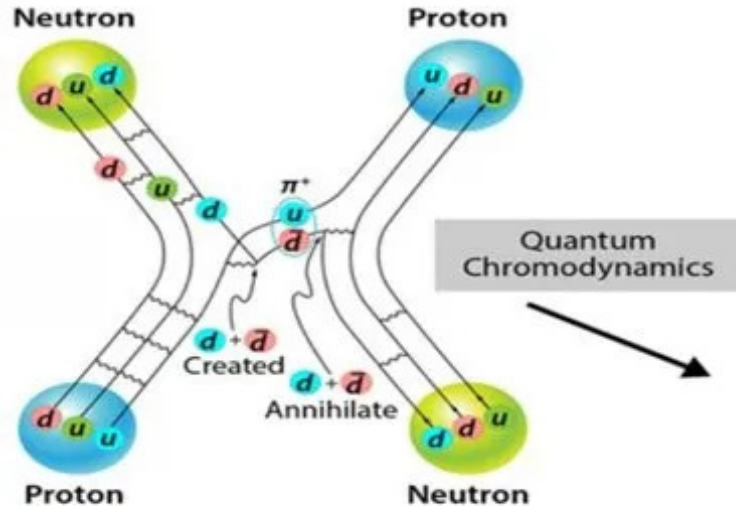




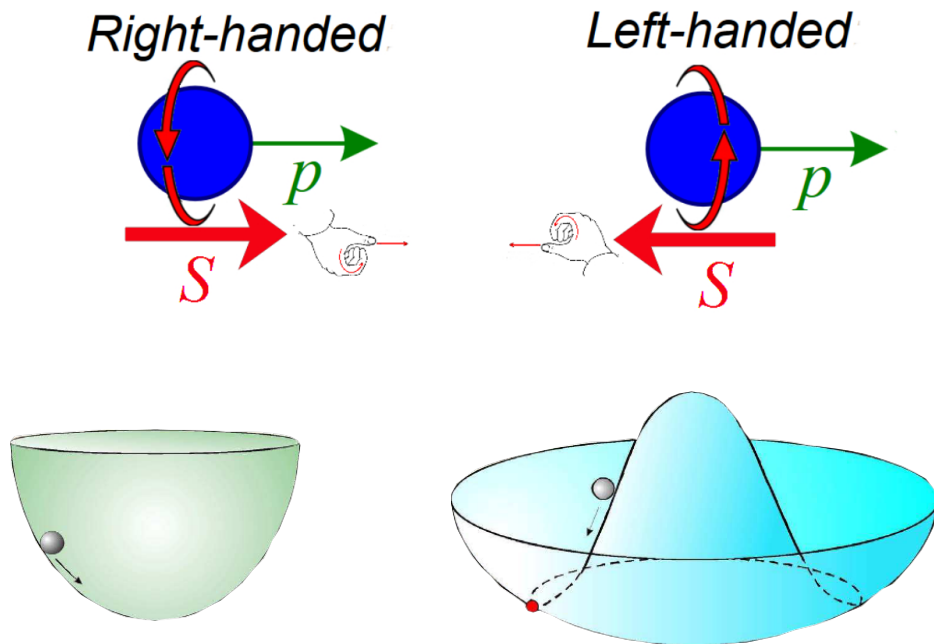
Above open-flavour threshold



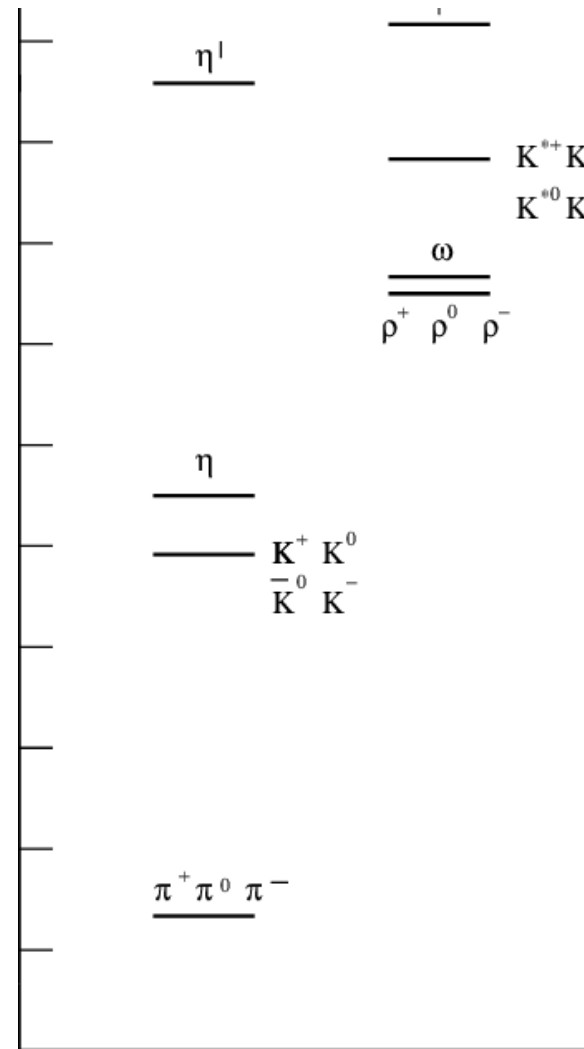
Effective Field Theory



Chiral Perturbation Theory

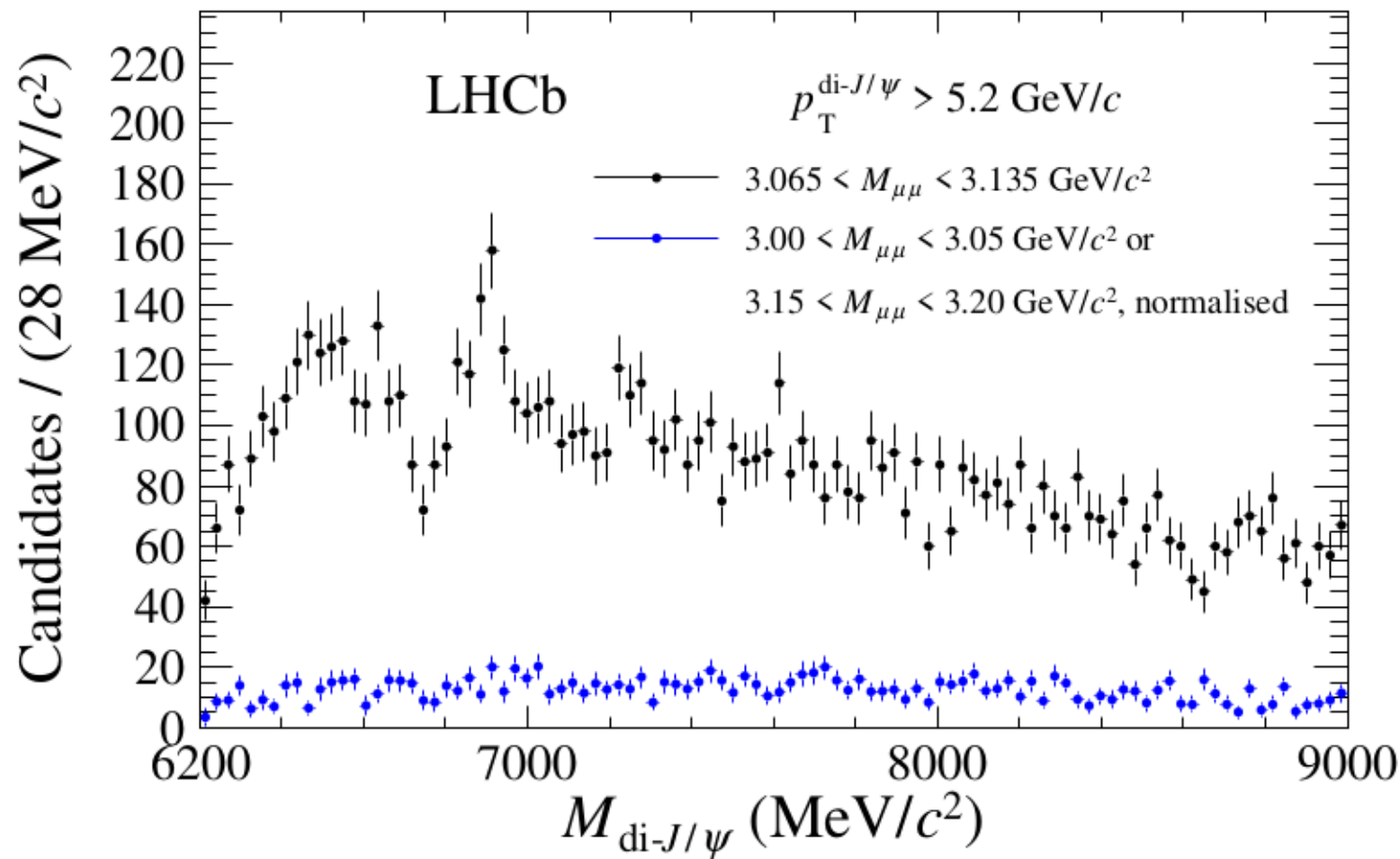


$$\mathcal{L}_{QCD} \xrightarrow{\text{low energy}} \mathcal{L}_{\pi\pi} + \dots$$

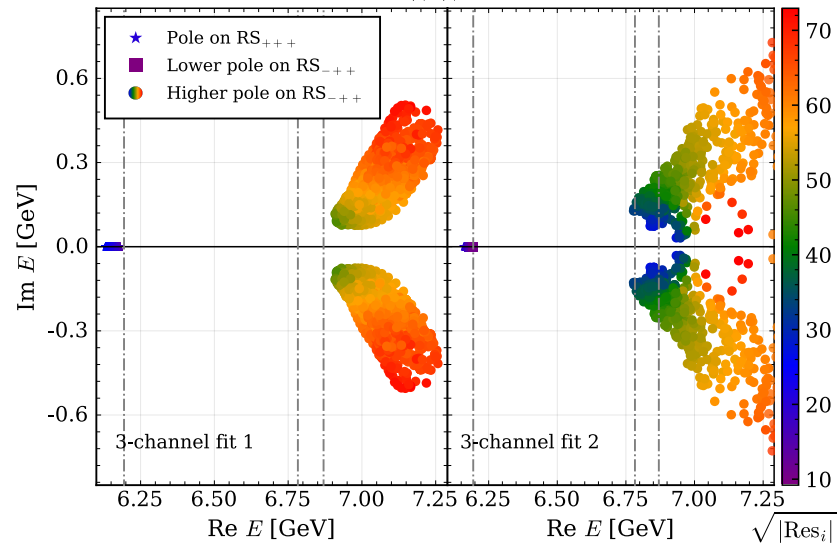
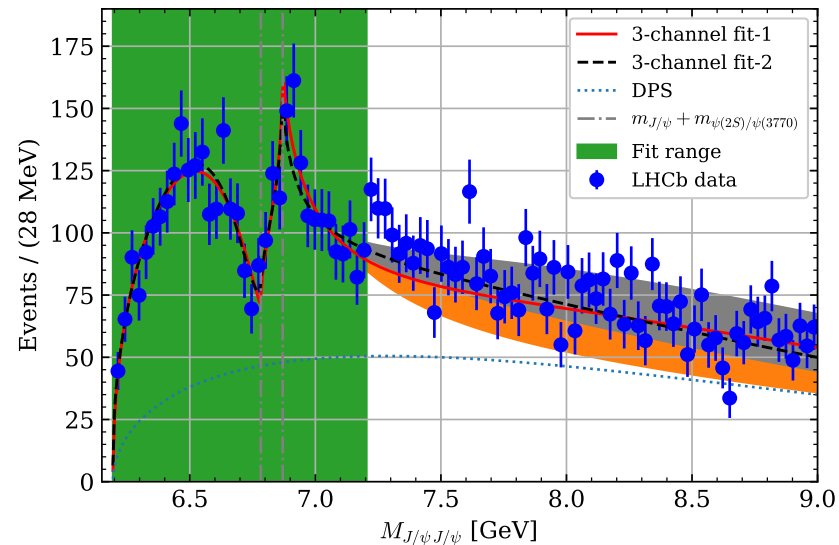
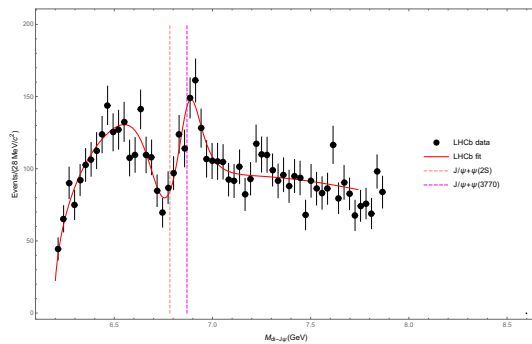
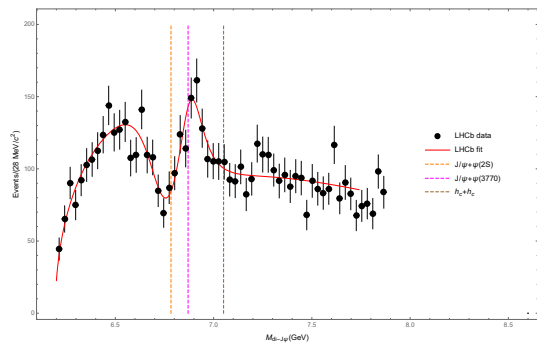
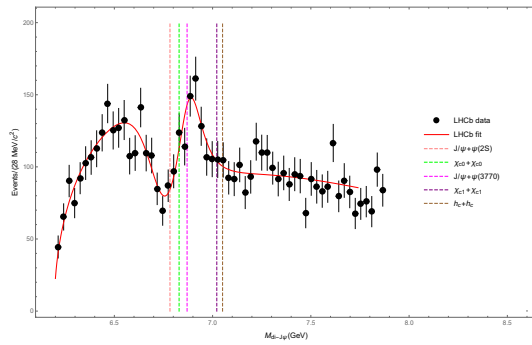
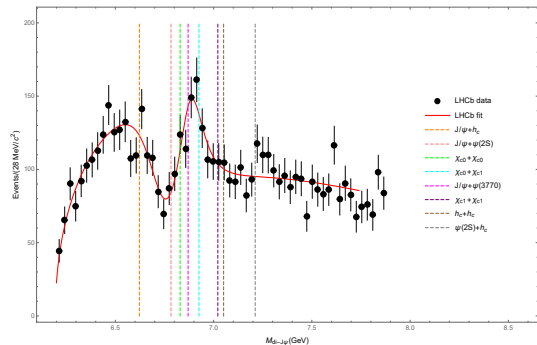


X(6200)/X(6900)

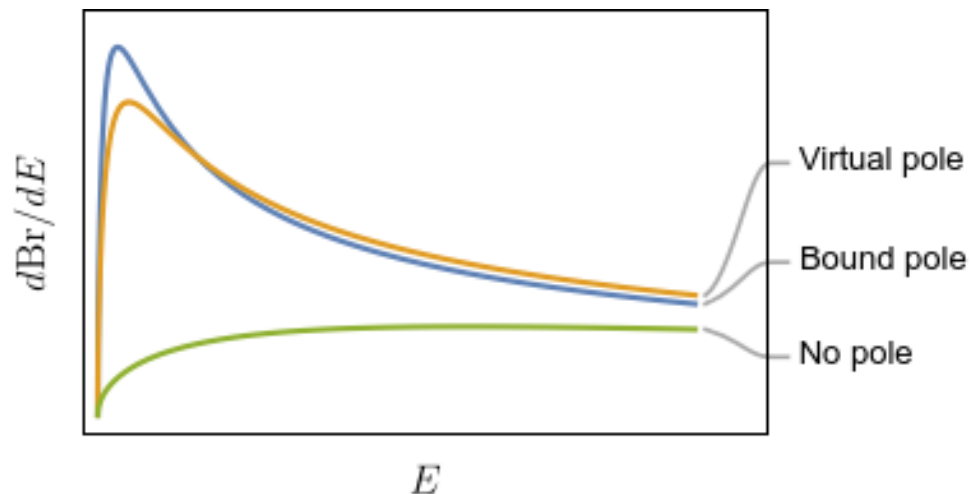
$$p + p \rightarrow J/\psi + J/\psi + \dots$$



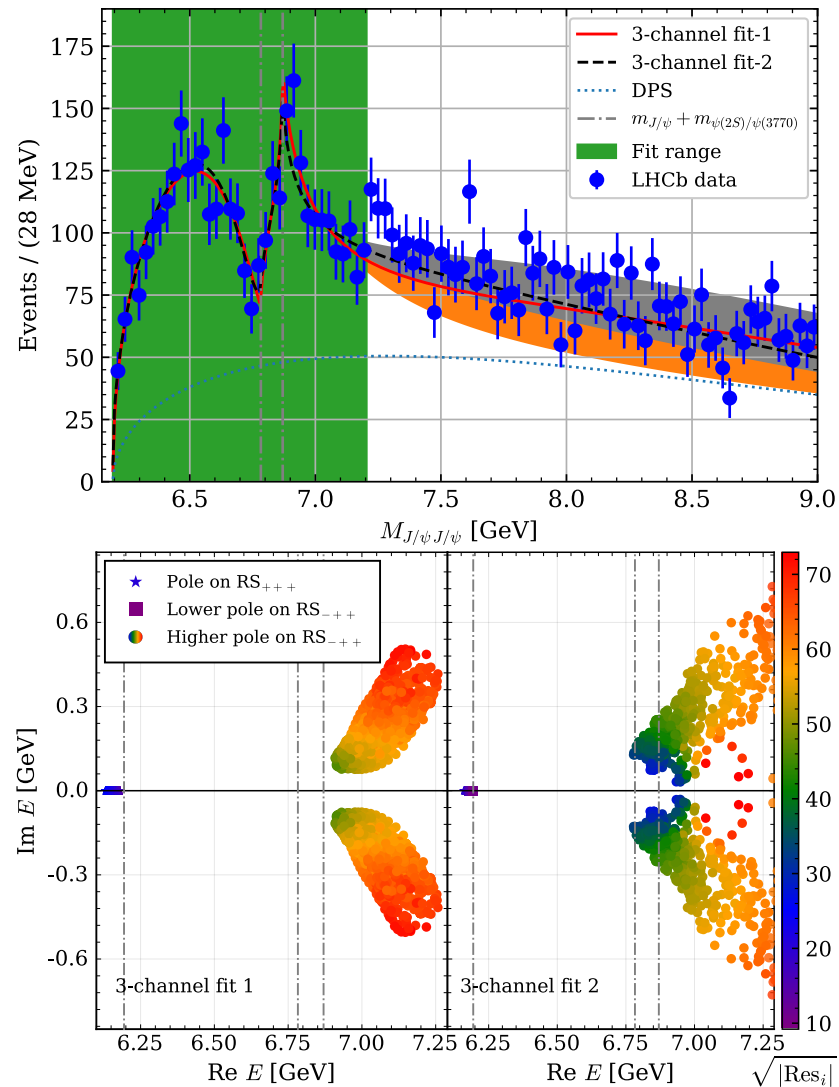
X(6200)/X(6900)



X(6200)/X(6900)

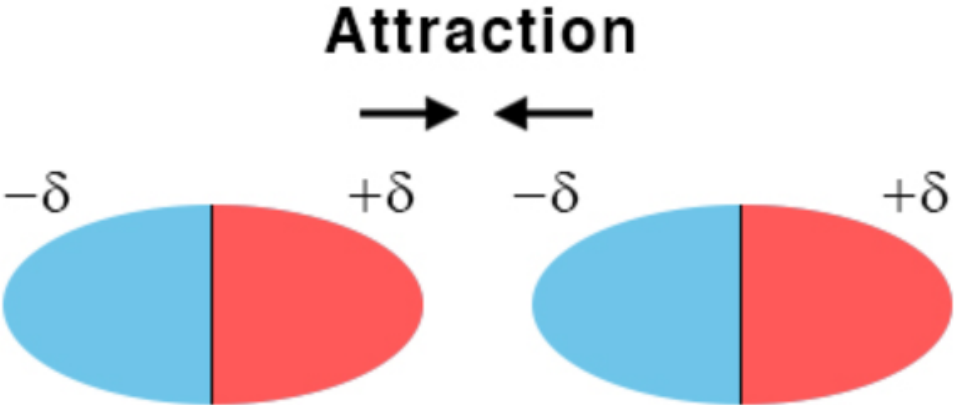
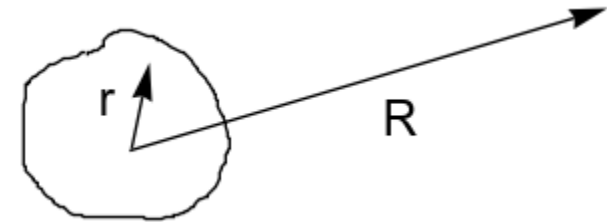


CMS & ATLAS Collaborations:
Additional BW structure near threshold
improves data description



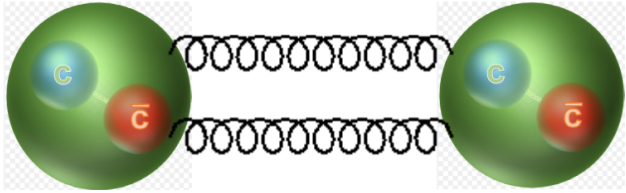
Polarisability in QED

$$\Phi(R) = \frac{Q}{R} + \text{Dipole} + \text{Quadrupole} + \dots$$



Strength of interaction of neutral objects at large distance depends on their *polarisability* that measures how easily they can be distorted by electric field

X(6200) as molecule



$$H_{\text{int}} \approx -\frac{1}{2} \zeta_a \vec{r} \cdot \vec{E}^a$$

$$\beta_{J/\psi} = \frac{1}{9} \langle J/\psi | \vec{r} \frac{1}{\hat{H}_O - M} \vec{r} | J/\psi \rangle = \frac{0.93}{\alpha_s^4 m_c^3}$$

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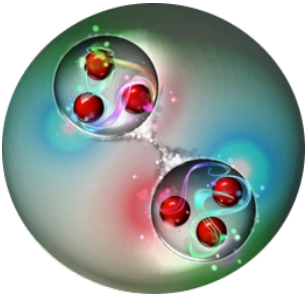
Dibaryon with Highest Charm Number near Unitarity from Lattice QCD

Yan Lyu,^{1,2,†} Hui Tong^{1,3,*}, Takuya Sugiura,³ Sinya Aoki,^{4,2} Takumi Doi,^{2,3}
Tetsuo Hatsuda,³ Jie Meng^{1,5} and Takaya Miyamoto²

PHYSICAL REVIEW LETTERS **130**, 111901 (2023)

Strongly Bound Dibaryon with Maximal Beauty Flavor from Lattice QCD

Nilmani Mathur^{1,*}, M. Padmanath^{2,†,§} and Debsubhra Chakraborty^{1,‡}



$$\beta_{\Omega_{ccc}} \approx 2.6 \times \beta_{J/\psi}$$

- Considerable progress in hadronic physics in last years
- Many new states with heavy quarks qualify as exotic hadrons
- Complementary theoretical approaches to exotic hadrons
- New insights into physics of strong interactions

In the forthcoming lectures:

Lecture 2. Exotic hadrons with heavy quarks: Methods

Lecture 3. Exotic hadrons with heavy quarks: Applications