

# Probing Polarized FFs in Unpolarized Collisions

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*H.C. Zhang, S.Y. Wei; PLB 839, 137821 (2023)*

*X.W. Li, Z.X. Chen, S. Cao, S.Y. Wei, PRD 109, 014035 (2024)*

*Z.X. Chen, H. Dong, S.Y. Wei, arXiv:2404.19202 (2024)*



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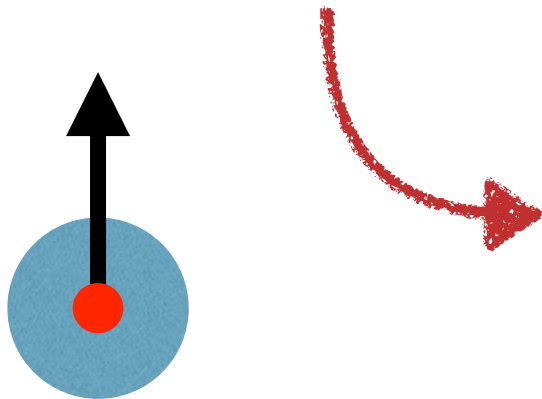
## QCD factorization

partonic interaction, perturbative

Cross Section = short distance  $\otimes$  long distance

non-perturbative, universal

TMD PDFs:  $\mathcal{FT} \langle p | \bar{\psi}(0) \psi(x^-, \vec{x}_\perp) | p \rangle$



$$\not{n}_+ \left[ f_1 - \frac{(\hat{e}_p \times \mathbf{k}_T) \cdot \mathbf{S}_\perp}{M} f_{1T}^\perp \right] + \gamma_5 \not{n}_+ \left[ \lambda g_{1L} + \frac{k_T \cdot \mathbf{S}_\perp}{m} g_{1T}^\perp \right] +$$

$$\frac{i[k_T, \not{n}_+]}{2m} h_1^\perp + \frac{1}{2} [\not{S}_\perp, \not{n}_+] \gamma_5 h_{1T} + \frac{[k_T, \not{n}_+] \gamma_5}{2m} \left[ \lambda h_{1L}^\perp + \frac{k_T \cdot \mathbf{S}_\perp}{m} h_{1T}^\perp \right]$$

TMD FFs:  $\mathcal{FT} \langle 0 | \psi(0) | hX \rangle \langle hX | \bar{\psi}(x^-, \vec{x}_\perp) | 0 \rangle$



$$\not{n}_- \left[ D_1 + \frac{(\hat{e}_j \times \mathbf{p}_T) \cdot \mathbf{S}_\perp}{zM} D_{1T}^\perp \right] + \gamma_5 \not{n}_- \left[ \lambda G_{1L} + \frac{p_T \cdot \mathbf{S}_\perp}{zM} G_{1T}^\perp \right] +$$

$$\frac{i[\not{p}_T, \not{n}_-]}{2M} H_1^\perp + \frac{1}{2} [\not{S}_\perp, \not{n}_-] \gamma_5 H_{1T} + \frac{[\not{p}_T, \not{n}_-] \gamma_5}{2M} \left[ \lambda H_{1L}^\perp + \frac{p_T \cdot \mathbf{S}_\perp}{M} H_{1T}^\perp \right]$$

## QCD factorization

		Baryons		
		Unpolarized	L	T
Quarks	Unpolarized	$D_1$		$D_{1T}^\perp$
	L		$G_{1L}$	$G_{1T}^\perp$
	T	$H_1^\perp$	$H_{1L}^\perp$	$H_{1T}, H_{1T}^\perp$

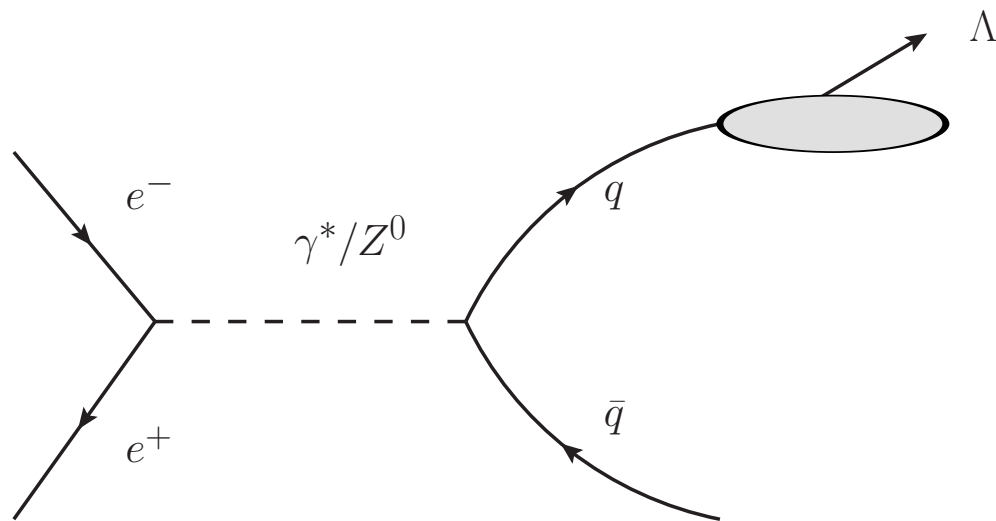
☒  $G_{1L}$ , aka, the longitudinal spin transfer

Number density of longitudinally polarized hadrons produced from longitudinally polarized quarks.

polarized beams  
or  
weak interaction



## Single Inclusive $\Lambda$ Production in $e^+e^-$ Annihilation Experiment



Final state quarks gain polarization through weak interaction

spin transfer

$$\mathcal{P}_L^\Lambda = \lambda_q \frac{G_{1Lq}^\Lambda}{D_{1q}^\Lambda}$$

quark polarization

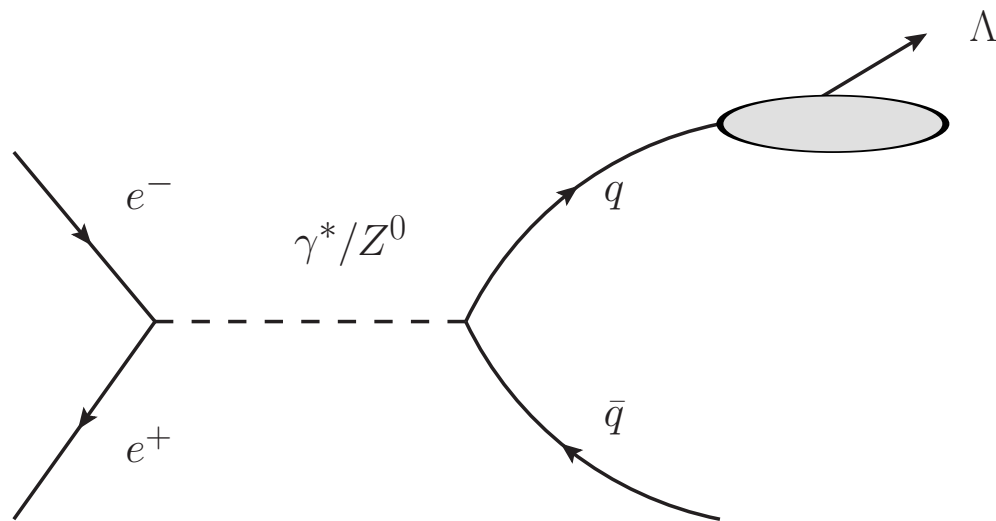
$$\frac{d\sigma}{dPS} = \sigma_0 \left[ D_{1q}^\Lambda(z) + \lambda_q \lambda_\Lambda G_{1Lq}^\Lambda(z) \right]$$

Belle  
Energy

LEP  
Energy

## Single Inclusive $\Lambda$ Production in $e^+e^-$ Annihilation Experiment

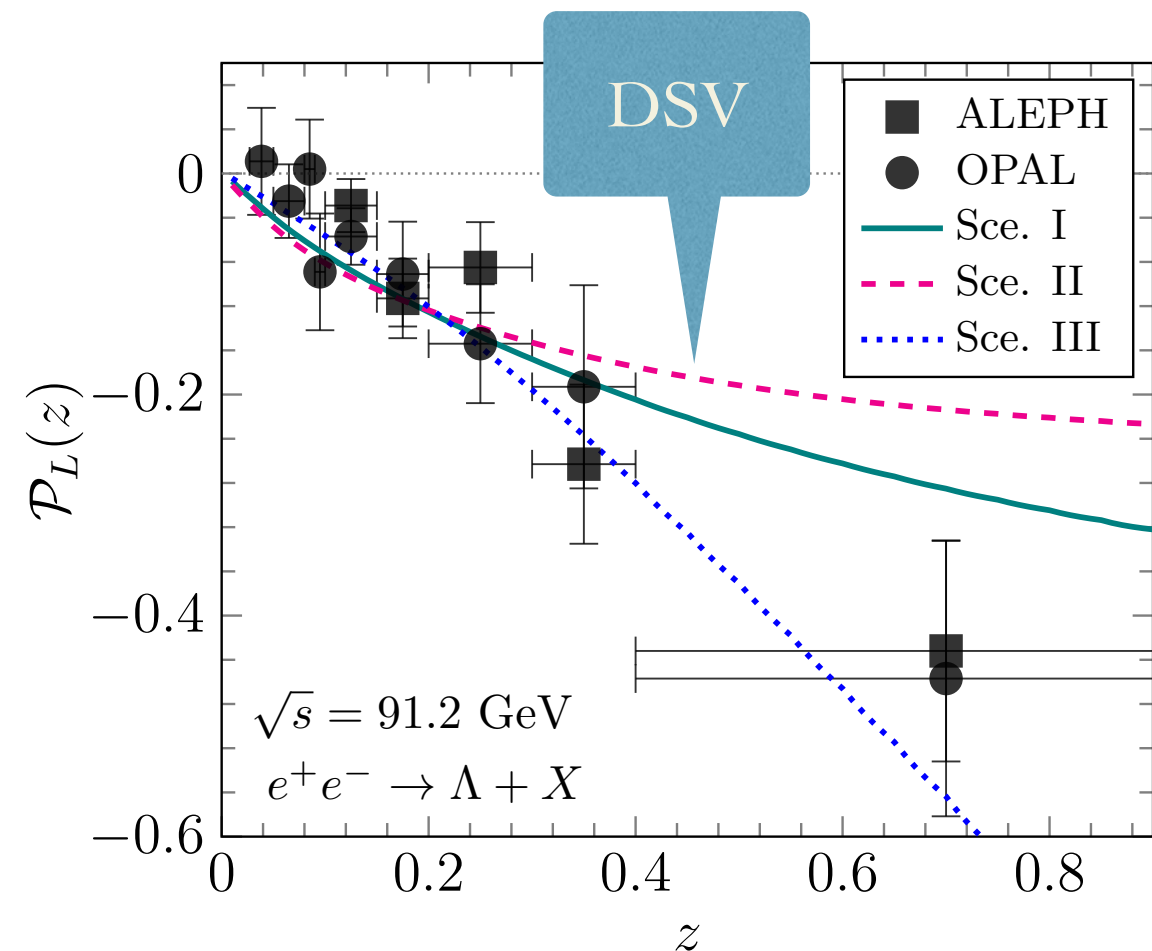
Final state quarks gain polarization through weak interaction



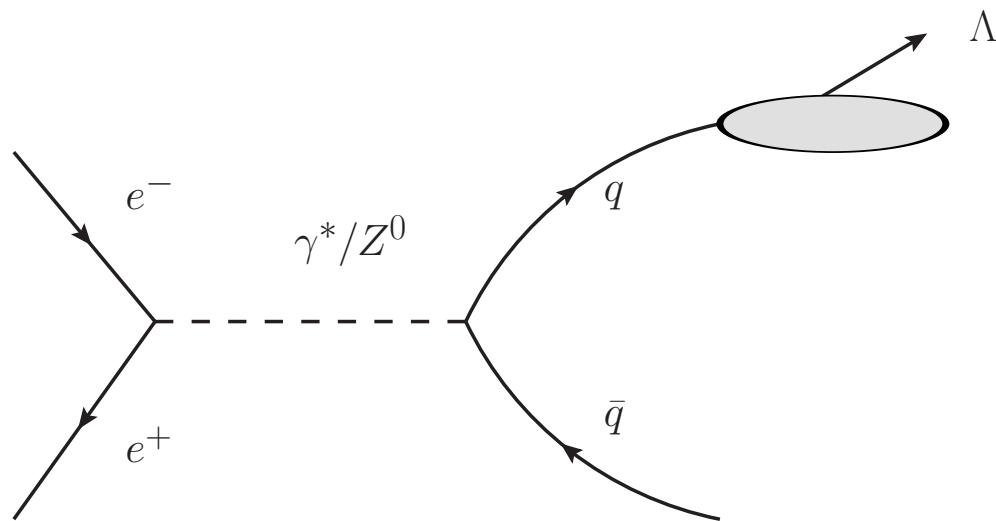
spin transfer

$$\mathcal{P}_L^\Lambda = \lambda_q \frac{G_{1Lq}^\Lambda}{D_{1q}^\Lambda}$$

quark polarization



## Single Inclusive $\Lambda$ Production in $e^+e^-$ Annihilation Experiment



Final state quarks gain polarization through weak interaction

spin transfer

$$\mathcal{P}_L^\Lambda = \lambda_q \frac{G_{1Lq}^\Lambda}{D_{1q}^\Lambda}$$

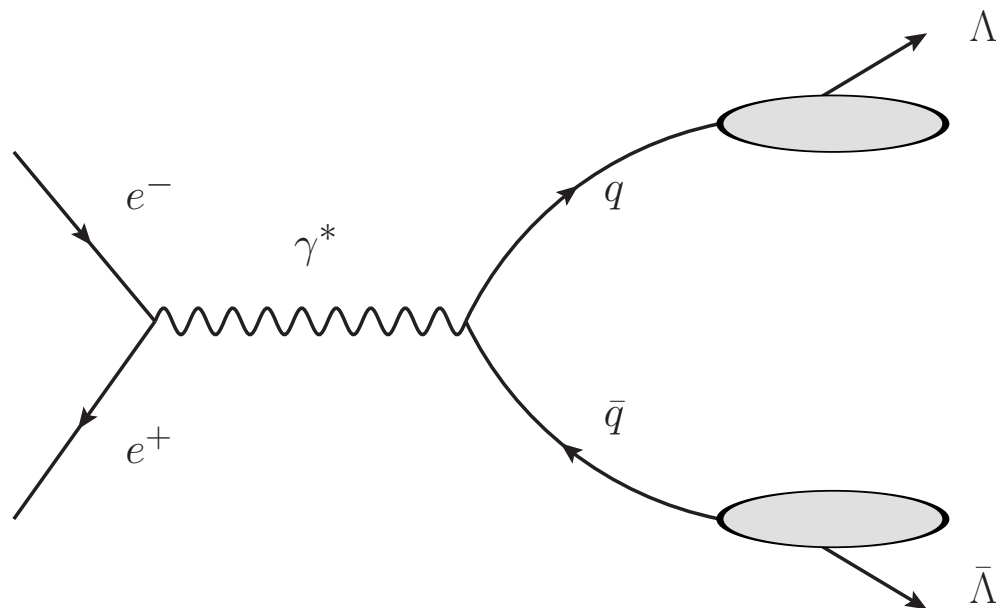
quark polarization

$$\frac{d\sigma}{dPS} = \sigma_0 \left[ D_{1q}^\Lambda(z) + \lambda_q \lambda_\Lambda G_{1Lq}^\Lambda(z) \right]$$

Belle  
Energy

LEP  
Energy

## $\Lambda\bar{\Lambda}$ -pair Production in $e^+e^-$ Annihilation Experiment



Belle  
Energy

$$\frac{d\sigma}{dPS} = \sigma_0 \left[ D_{1q}^{\Lambda}(z_1) D_{1\bar{q}}^{\bar{\Lambda}}(z_2) - \lambda_{\Lambda} \lambda_{\bar{\Lambda}} G_{1Lq}^{\Lambda}(z_1) G_{1L\bar{q}}^{\bar{\Lambda}}(z_2) \right]$$

### ☑ Helicity Conservation

$q$  and  $\bar{q}$  are on the same fermion line. They must have opposite helicities.

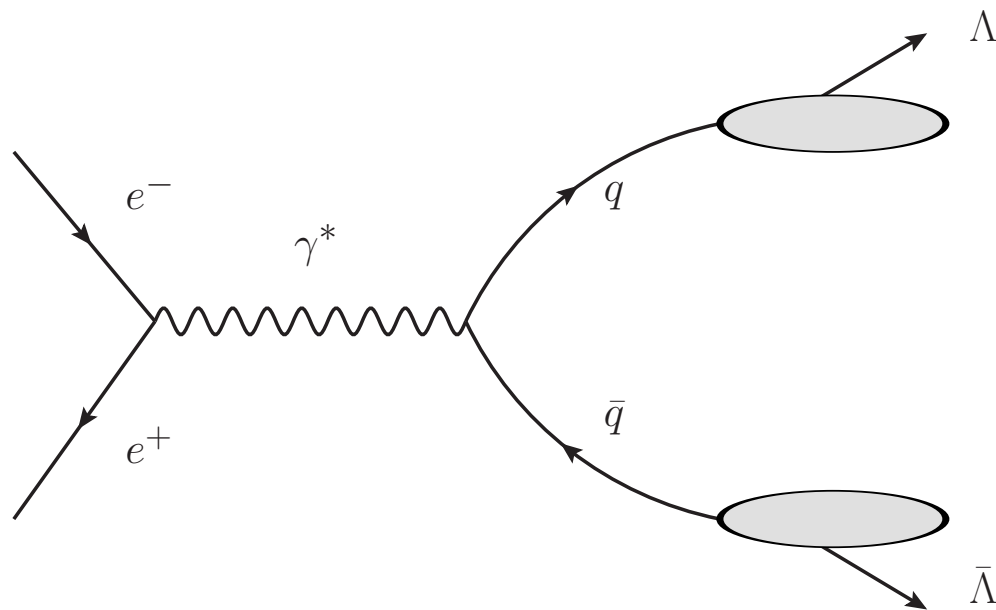
### ☑ Helicity Correlation

A novel probe to the spin-dependent fragmentation functions

*H.C. Zhang, SYW; PLB 839 (2023) 137821*  
see also Nucl. Phys. B 445 (1995) 380.



## Helicity Amplitude Approach



$\sigma_{\lambda_q \lambda_{\bar{q}}}$  denotes the differential X of  $q\bar{q}$ -pair production

$$\sigma_{+-} = \sigma_{-+} = \sigma_0/2$$

$$\sigma_{++} = \sigma_{--} = 0$$

$\mathcal{D}$  denotes the helicity dependent fragmentation function

$$\mathcal{D}(\lambda_q, \lambda_{\Lambda}, z) = D_{1q}(z) + \lambda_q \lambda_{\Lambda} G_{1Lq}(z)$$

Physical interpretation:

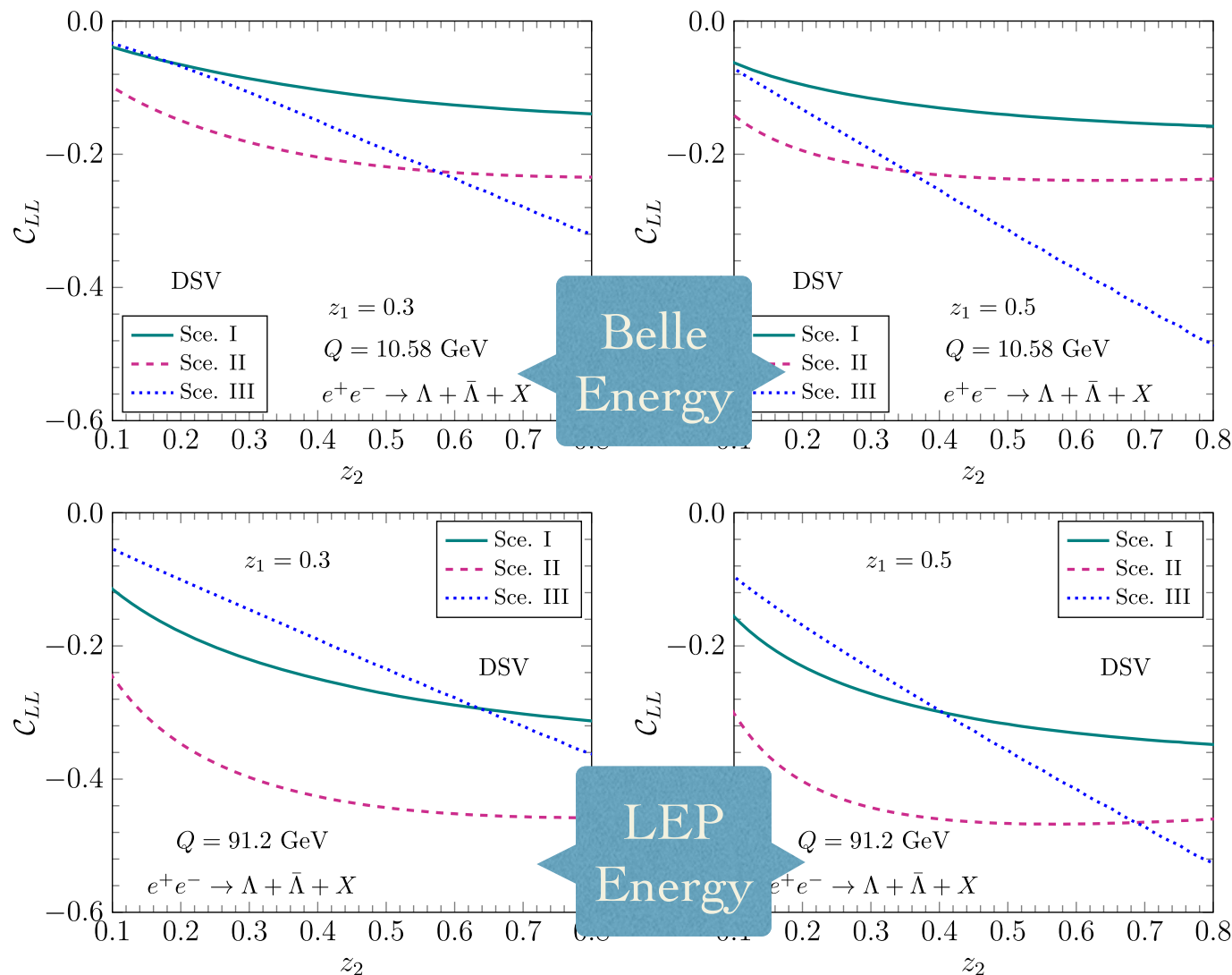
$$\begin{aligned} \frac{d\sigma}{dPS} &= \sigma_{+-} \otimes \mathcal{D}_q(+, \lambda_{\Lambda}, z_1) \otimes \mathcal{D}_{\bar{q}}(-, \lambda_{\bar{\Lambda}}, z_2) + \sigma_{-+} \otimes \mathcal{D}_q(-, \lambda_{\Lambda}, z_1) \otimes \mathcal{D}_{\bar{q}}(+, \lambda_{\bar{\Lambda}}, z_2) \\ &= \sigma_0 \left[ D_{1q}^{\Lambda}(z_1) D_{1\bar{q}}^{\bar{\Lambda}}(z_2) - \lambda_{\Lambda} \lambda_{\bar{\Lambda}} G_{1Lq}^{\Lambda}(z_1) G_{1L\bar{q}}^{\bar{\Lambda}}(z_2) \right] \end{aligned}$$

*H.C. Zhang, SYW; PLB 839 (2023) 137821*  
see also Nucl. Phys. B 445 (1995) 380.

# Dihadron helicity correlation

## Helicity Correlation of $\Lambda\bar{\Lambda}$ -pair

$$C_{LL} = \frac{\text{same signs} - \text{opposite signs}}{\text{total cross section}} = \frac{\sum_q \sigma_0 G_{1Lq}^\Lambda(z_1) G_{1L\bar{q}}^{\bar{\Lambda}}(z_2)}{\sum_q \sigma_0 D_{1q}^\Lambda(z_1) D_{1\bar{q}}^{\bar{\Lambda}}(z_2)} \propto \langle \cos \theta_1^* \cos \theta_2^* \rangle$$



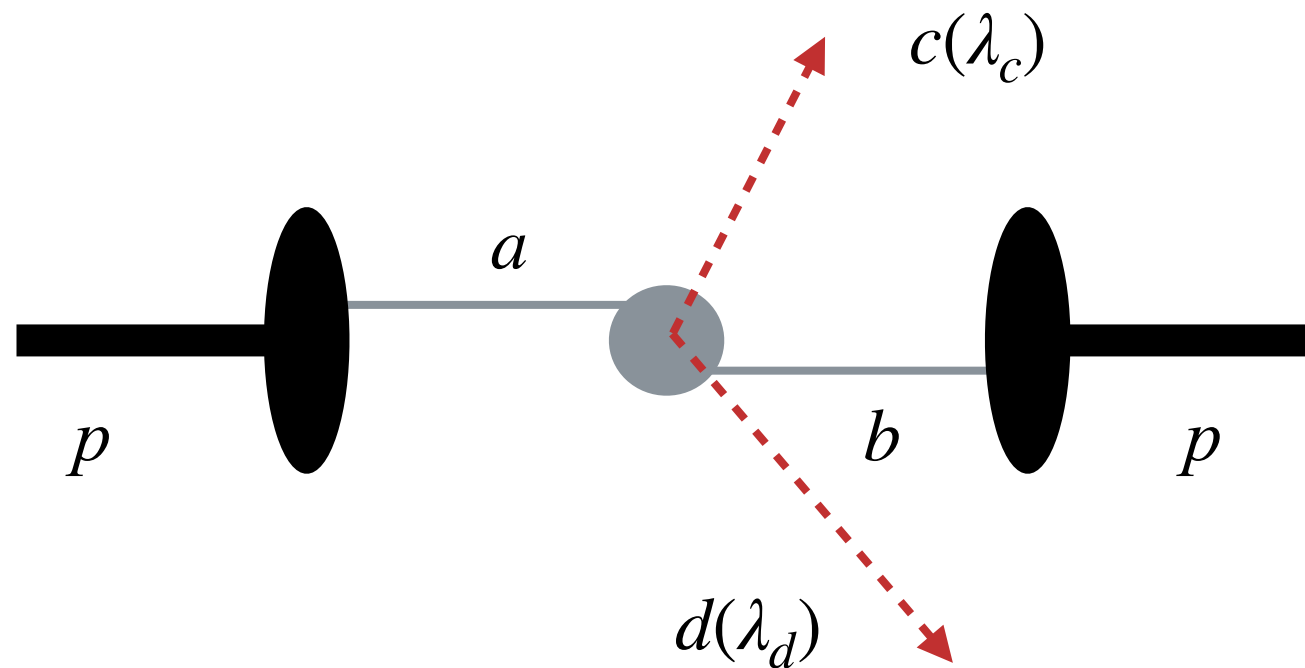
✓ The helicity correlation at the Belle energy has a similar magnitude with that at the LEP energy.

✓ It is now possible to extract the longitudinal spin transfer at Belle experiment.



# Dihadron helicity correlation

Applying to the unpolarized pp collisions

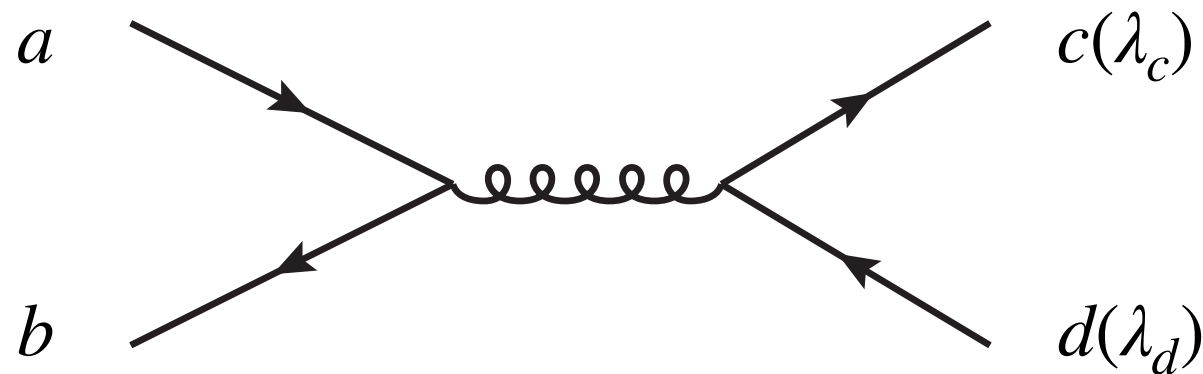


$$a + b \rightarrow c(\lambda_c) + d(\lambda_d)$$

☒ Are  $\lambda_c$  and  $\lambda_d$  correlated?

Yes!

“s-channel diagrams”: just like  $e^+e^-$  annihilation, maximum correlation



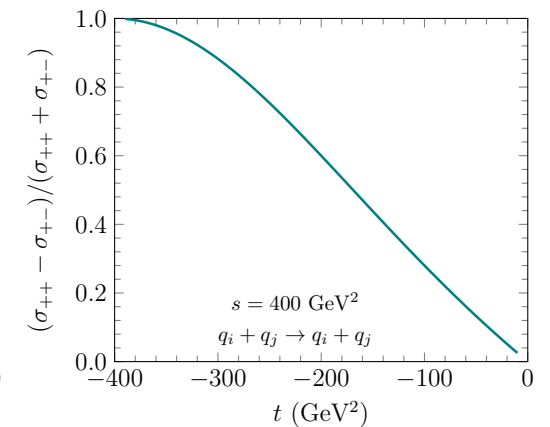
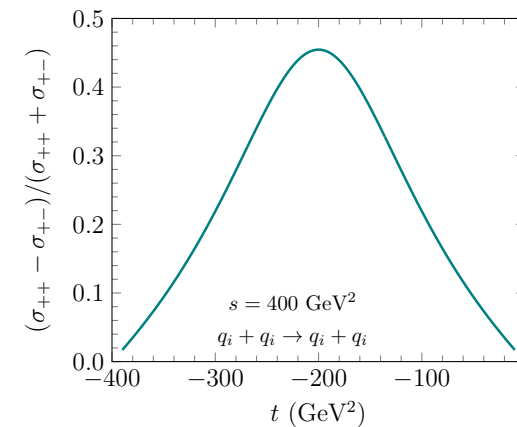
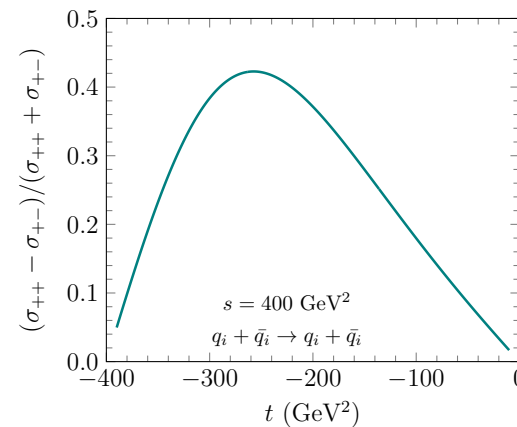
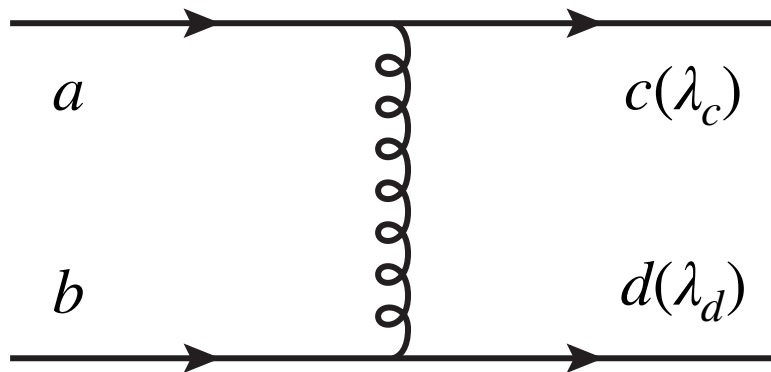
$$g + g \rightarrow q + \bar{q}$$

$$q_i + \bar{q}_i \rightarrow q_j + \bar{q}_j$$

$$q + \bar{q} \rightarrow g + g$$

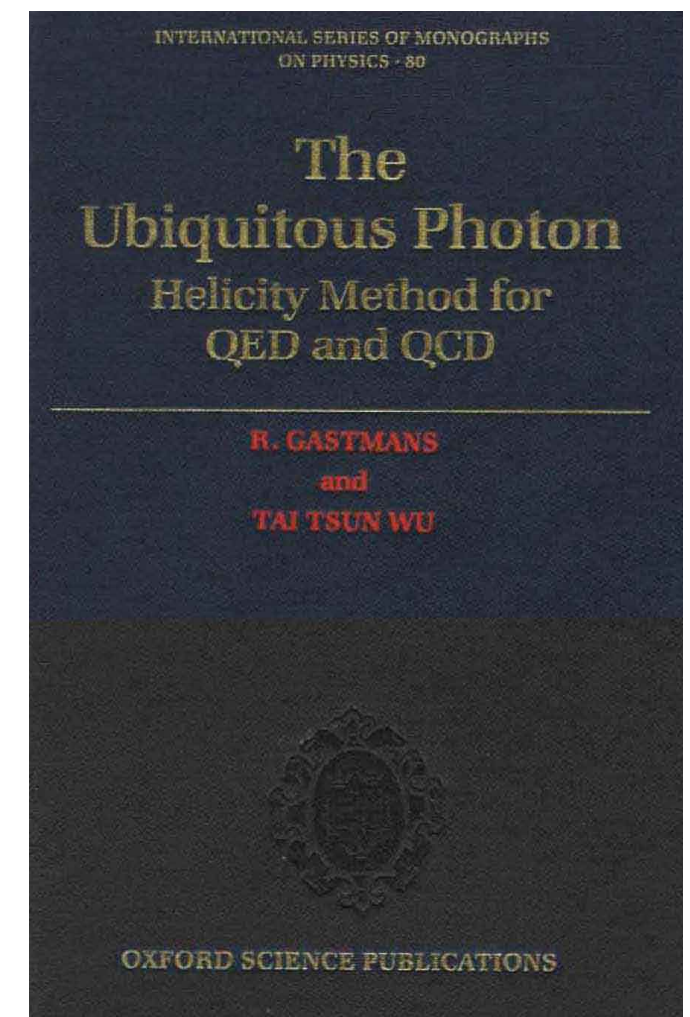
# Dihadron helicity correlation

“t-channel diagrams”: prefer same-sign correlation



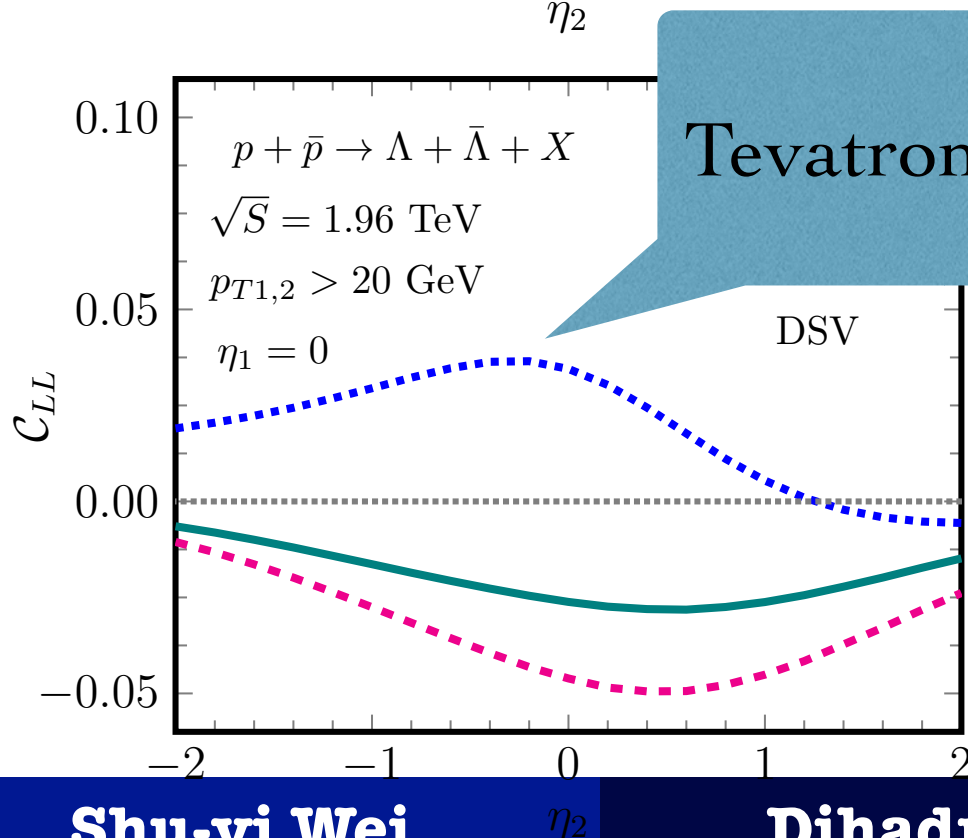
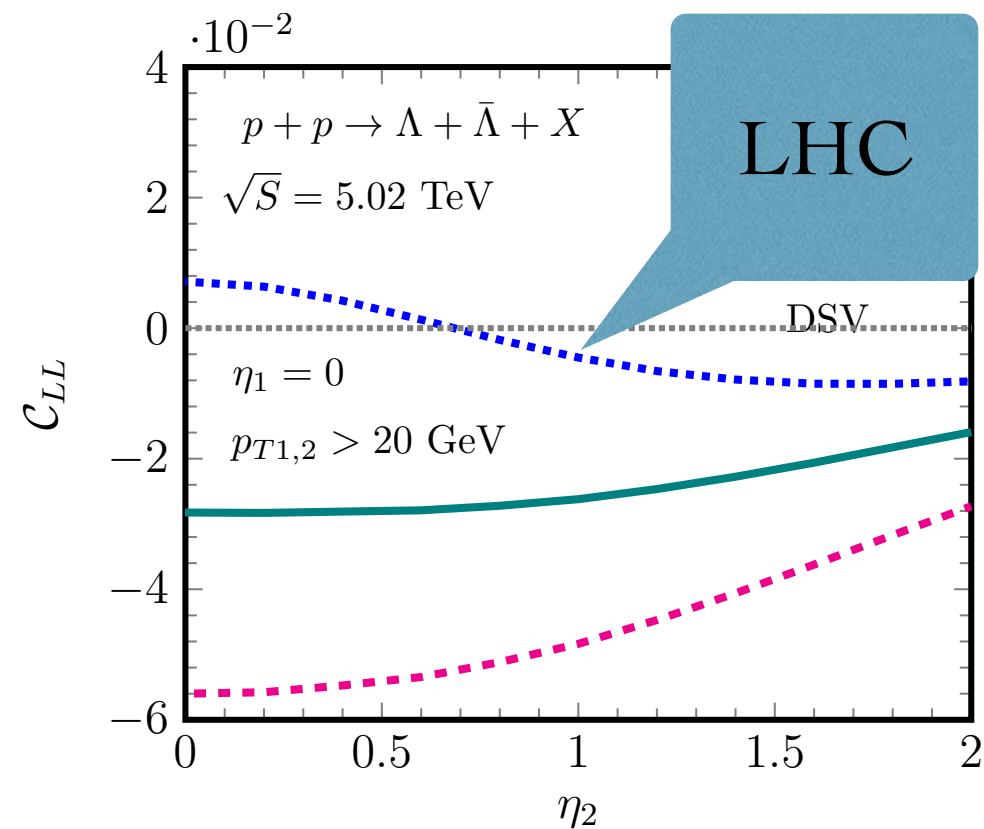
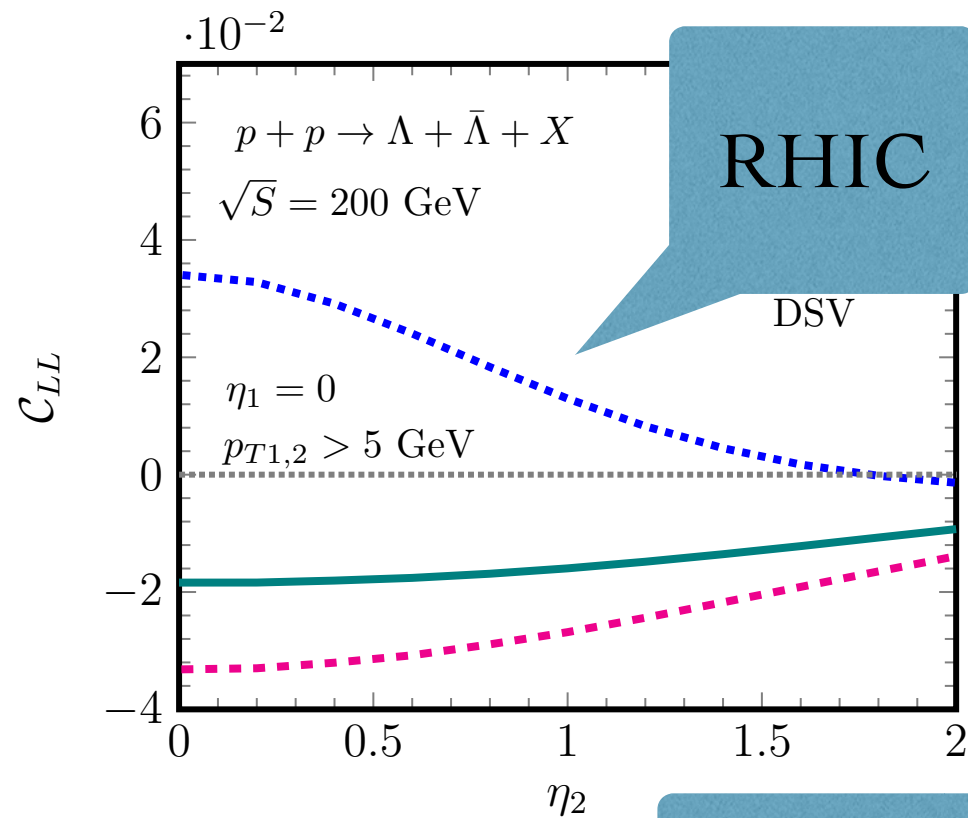
To summarize

- ☑ “s-channel”:  $\sigma_{+-} = \sigma_{-+} > \sigma_{++} = \sigma_{--} = 0$
- ☑ “t-channel”:  $\sigma_{++} = \sigma_{--} > \sigma_{+-} = \sigma_{-+} > 0$
- ☑ Probe polarized FF in unpolarized pp collisions
- ☑ Explore the circularly polarized gluon FF





## Polarization Correlation in unpolarized pp collisions



- ✓ Smaller, but none-zero
- ✓ Distinguish different scenarios
- ✓ Avoid contamination of polarized PDF
- ✓ Probe gluon spin transfer

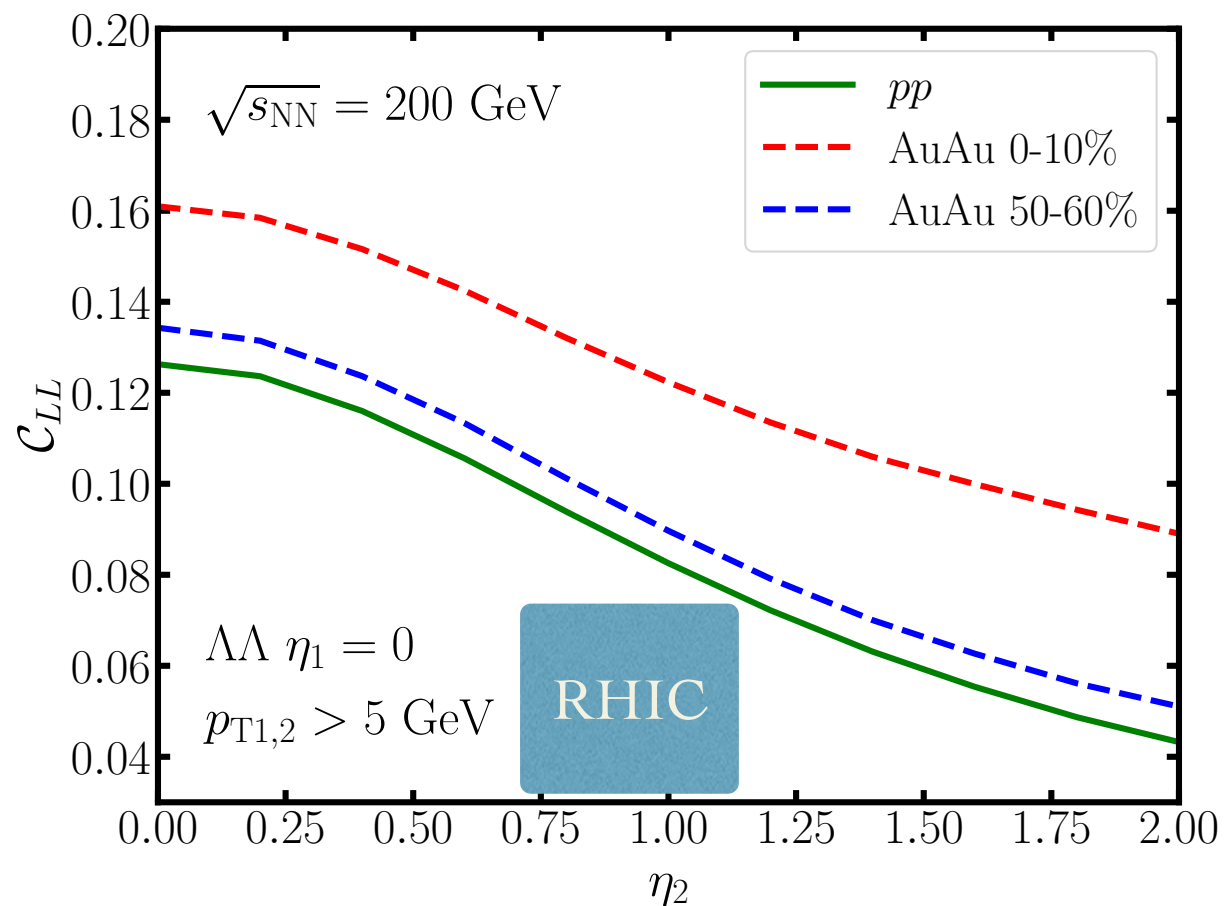
# Helicity Amplitude Approach

## Helicity Correlation in central and peripheral AA collisions

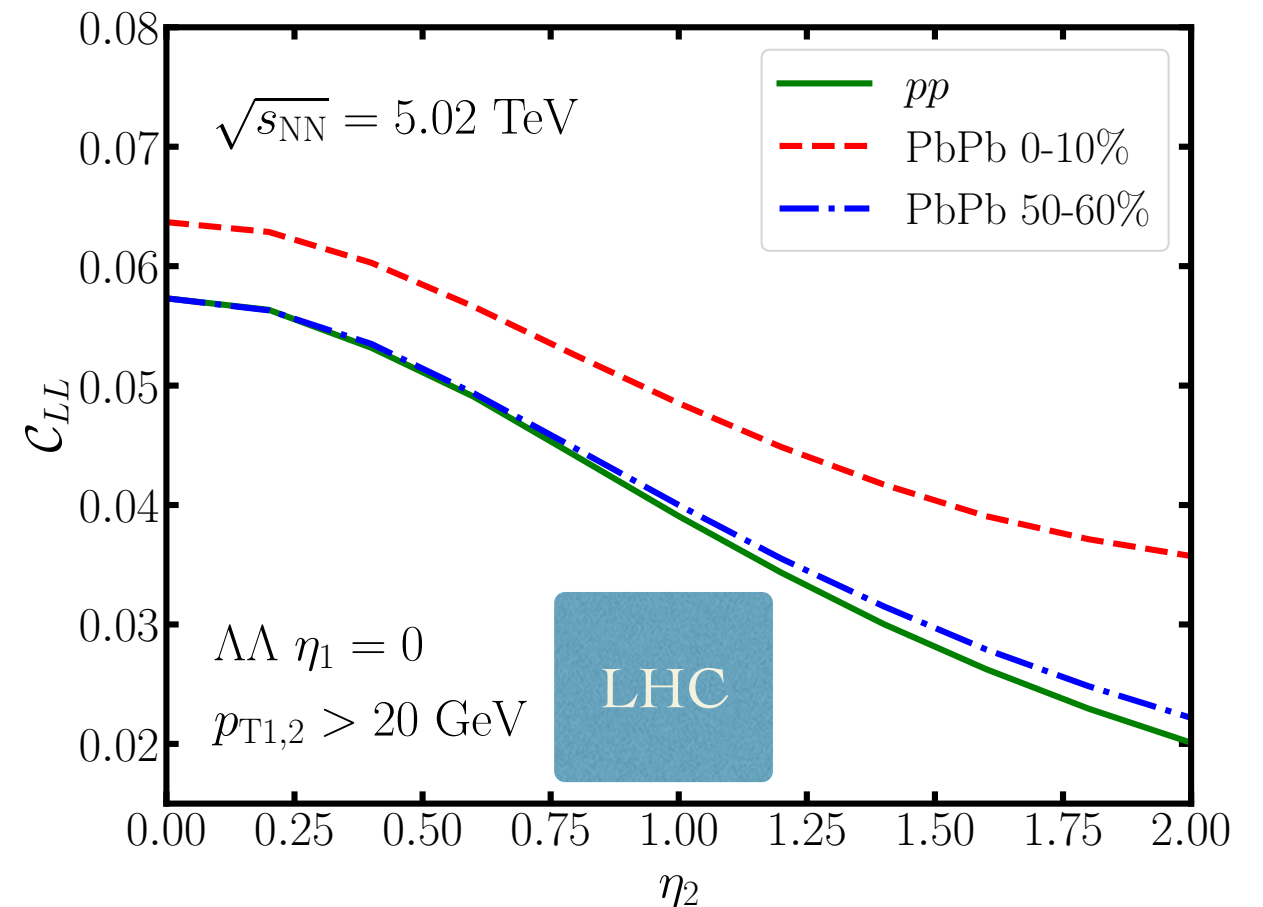
A toy model:  $\left. \frac{d\sigma}{dPS} \right|_{AA} = \text{Energy Loss} \otimes \left. \frac{d\sigma}{dPS} \right|_{pp}$

☑ Much larger luminosity

☑ Jet Quenching + Polarization

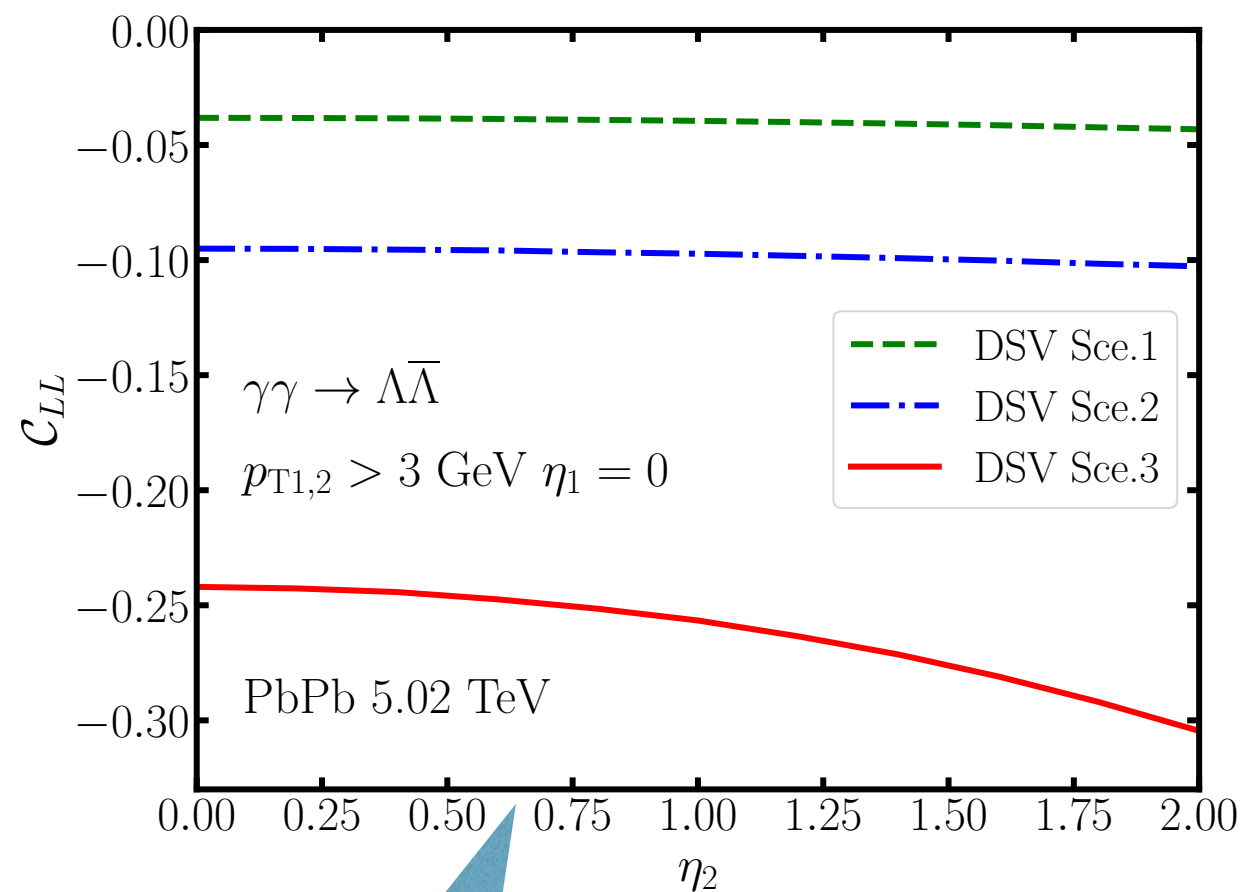


Clear Enhancement in central AA collisions





## Helicity Correlation in ultra-peripheral AA collisions: Type I

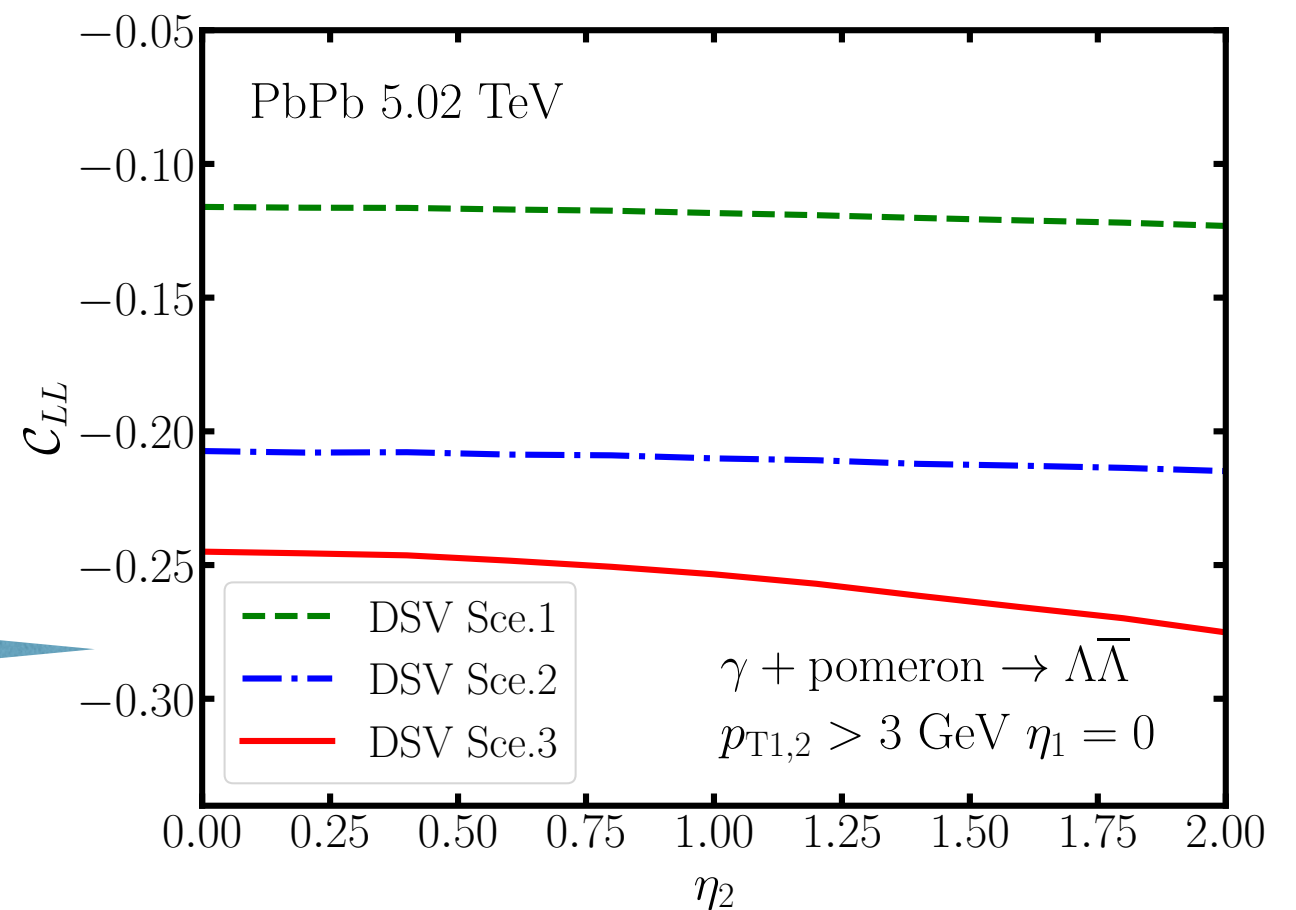


$\gamma + \gamma$

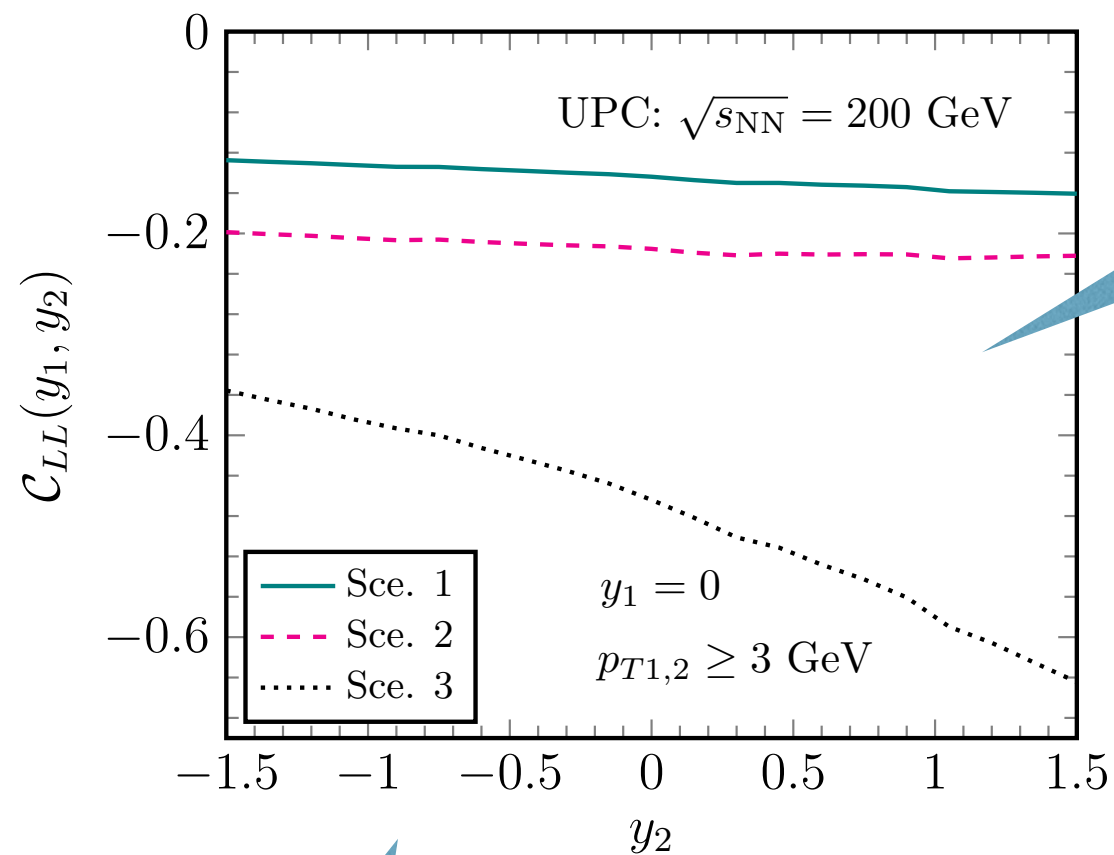
$\gamma + \mathbb{P}$

✓ Much larger luminosity

✓ Pomeron + Polarization



## Helicity Correlation in ultra-peripheral AA collisions: Type II

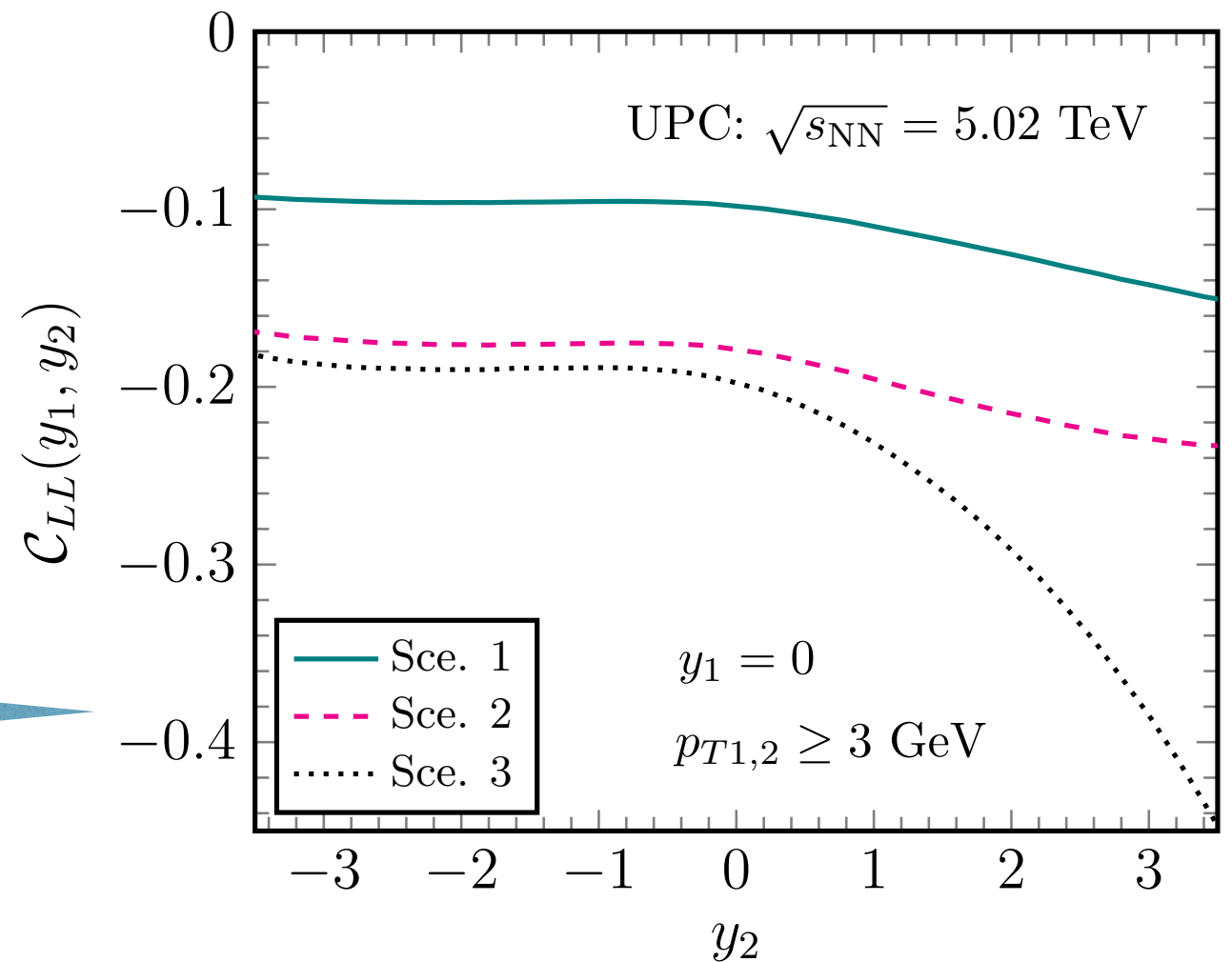


RHIC

LHC

$\gamma + q$  or  
 $\gamma + g$

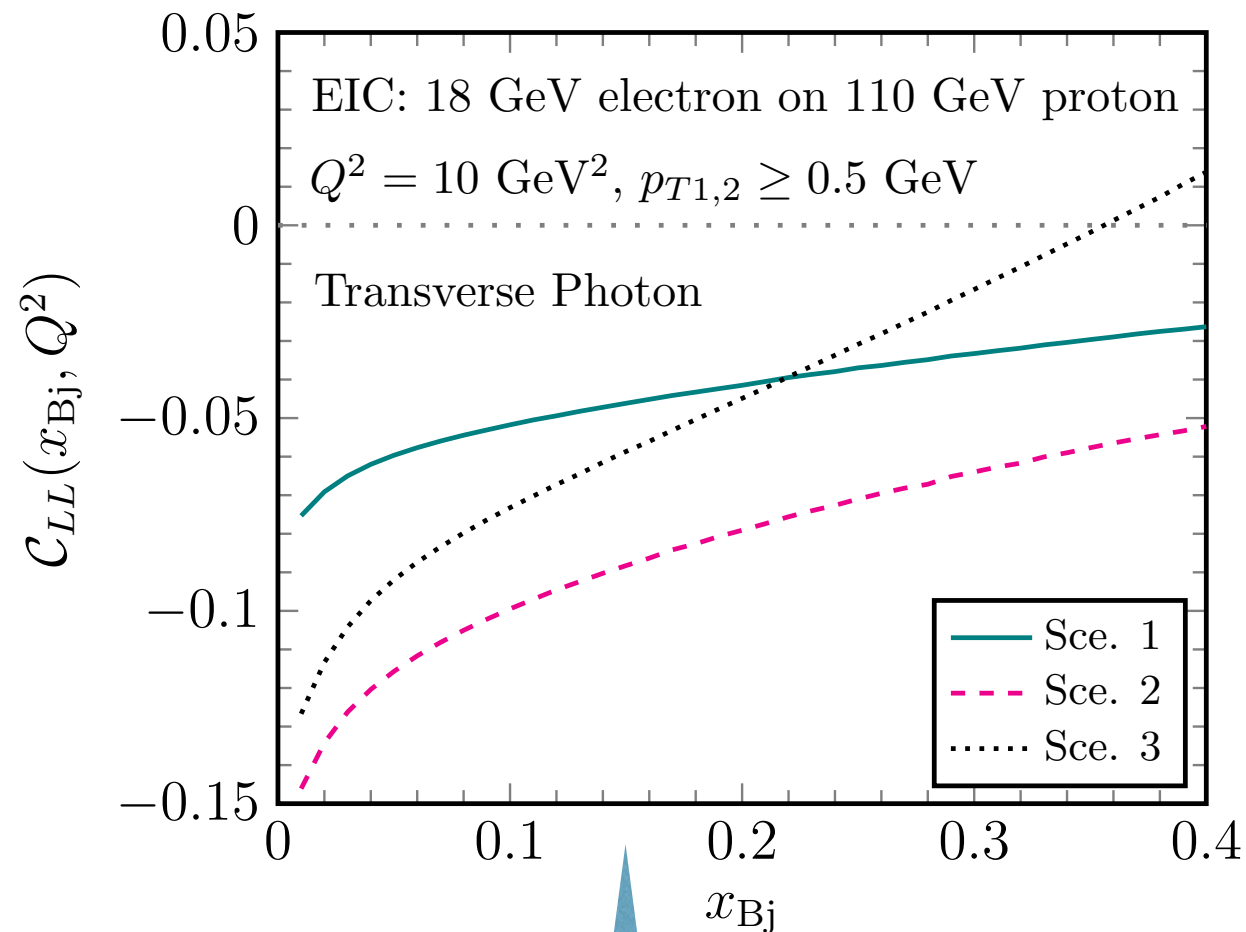
$\gamma + q$  or  
 $\gamma + g$



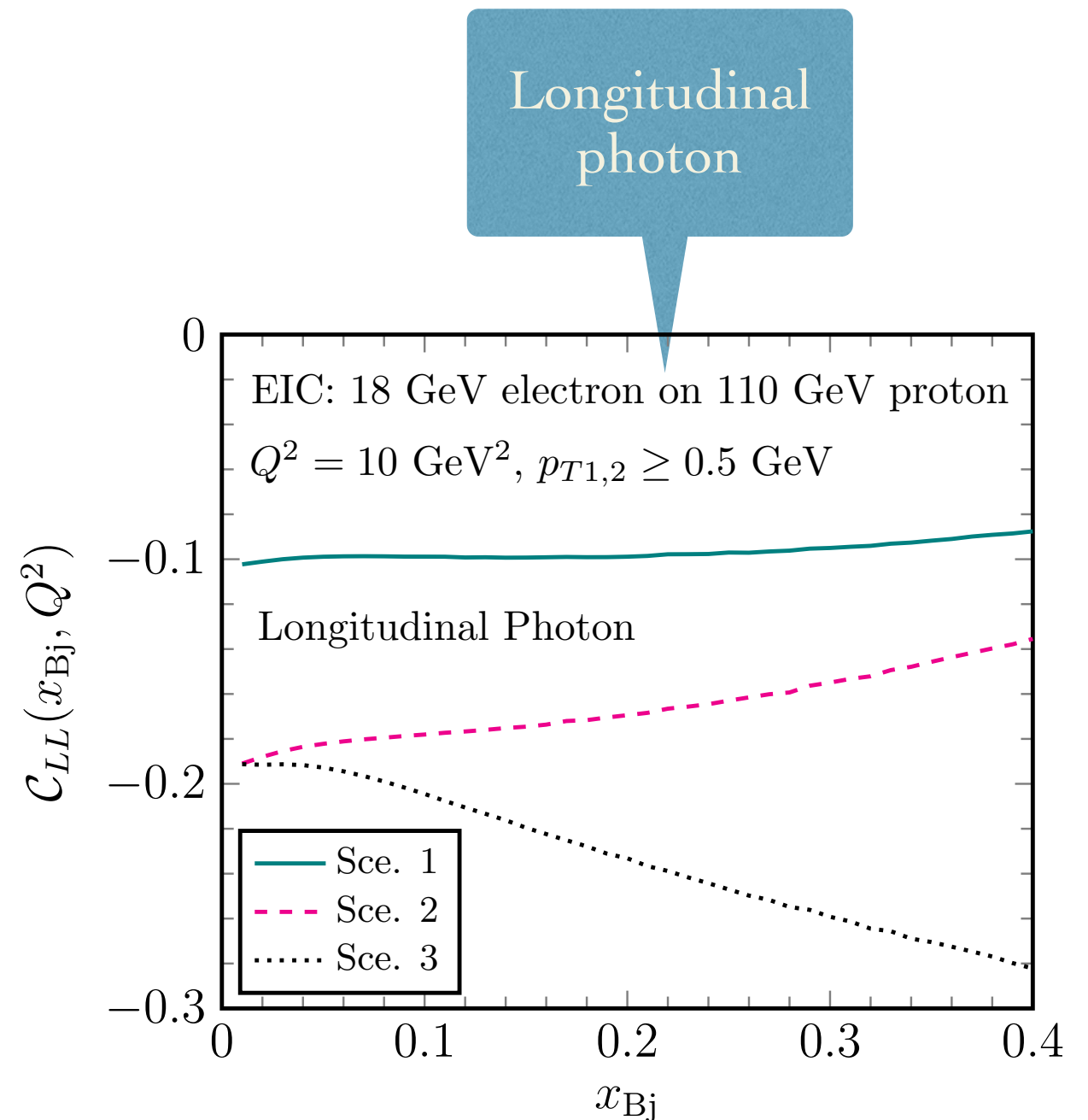


## Helicity Correlation at future EIC

$\gamma^* + q$  or  $\gamma^* + g$



Transverse  
photon



Longitudinal  
photon

- ☑ Spin effects can also be studied in unpolarized collisions.
- ☑ Dihadron helicity correlation in various high energy collisions offers a novel platform to investigate the flavor dependence of the  $G_{1L}$  fragmentation function.

Besides this talk, we also studied other spin effects in unpolarized collisions.

Phys.Lett.B 816, 136217 (2021).

Phys.Rev.D105, 034027 (2022).

Phys.Lett.B 850, 138509 (2024).

arXiv:2403.06133

Thanks for your attention!

The End



