

# Deep inelastic neutrino scatterings at high and ultra-high energies

The First Workshop on Neutrino Scattering: Theory, Experiment, Phenomenology

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### Outline

### ♦ 1. Introduction

### ♦ 2. NNLO calculations for Deep inelastic neutrino scatterings

### ♦ 3. Predictions for ultra-high energy neutrino scatterings



### Deep inelastic neutrino scattering

consisting of quarks and gluons (QPM), and lead to the establishment of QCD

#### **DIS** kinematics



$W^2 = M^2 + 2M\nu + q^2 =$	= Л
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W=M $\rightarrow$ elastic scattering
$1 < W < 2 \text{ GeV} \rightarrow \text{inelastic}$
$W > 2 \text{ GeV} \rightarrow \text{deep inelast}$

lab frame - proton at rest before collision:  $p_2 = (M, 0, 0, 0)$ 

lorentz invariant form

energy loss of  
ncoming particle 
$$\nu = E_1 - E_3$$
  $\nu = \frac{F^2 4}{M}$   
Bjorken x  $x = \frac{Q^2}{2M\nu}$  x in [0,1]  $x = \frac{Q^2}{2p_2}$   
fractional energy loss  $x = \frac{E_3}{2m}$  wip [0,1]  $\nu = \frac{p_2}{2m}$ 

of incoming particle

y in [0, 1]

4-momentum transver

$$q^2 = (p_1 - p_3)^2$$
  $Q^2 = -q^2$   $q^2 = (p_1 - p_3)^2$ 

$$F_1(x,Q^2) \qquad F_2(x,Q^2)$$

• Deep inelastic lepton/neutrino scatterings on fixed-target reveal the internal structure of nucleons,

Bjorken scaling and QPM





# QCD collinear factorization

DIS cross sections



• QCD collinear factorization ensures universal separation of long-distance and short-distance contributions in high energy scatterings involving initial/final state hadrons, and enables predictions on

$$\frac{V(\bar{\nu})}{Q^2} = \frac{G_F^2}{4\pi x (1 + Q^2/M_{W,Z}^2)^2} \left[ Y_+ F_2 - y^2 F_L \pm Y_- x F_3 \right], \quad Y_\pm = 1 \pm (1 - y)^2$$

$$F_{2}(x,Q^{2}) = \sum_{i=q,\bar{q},g} \int_{0}^{1} d\xi C_{2}^{i}(x/\xi,Q^{2}/\mu_{r}^{2},\mu_{f}^{2}/\mu_{r}^{2},\alpha_{s}(\mu_{r}^{2})) \times f_{i/h}(\xi,\mu_{f})$$
 [Collins, Soper, Sterman, 1989]

## Global analysis of PDFs

productions at fixed-target and collider experiments, with increasing weight from LHC



- with large momentum effective theory or pseudo-PDFs [2004.03543]

◆ PDFs are usually extracted from global analysis on variety of data, e.g., DIS, Drell-Yan, jets and top quark

[JG, Harland-Lang, Rojo, 1709.04922; 1905.06957 for recent review articles]

\* diversity of the analysed data are important to ensure flavor separation and to avoid theoretical/experimental bias; possible extensions to include EW parameters and possible new physics for a self-consistent determination

\* alternative approach from lattice QCD simulations, for various PDF moments or PDFs directly calculated in x-space

# Impact of neutrino DIS data

for separations of quark and anti-quarks and also quark flavors, especially strange quark PDF

Data set	Nucleus	$E_{\nu/\bar{\nu}}(\text{GeV})$	# pts	Corr.sys.	Ref.
CDHSW $\nu$	Fo	23 - 188	465	No	[48]
CDHSW $\bar{\nu}$	LE		464		
CCFR $\nu$	Fo	35 - 340	1109	No	[50]
CCFR $\bar{\nu}$	$\mathbf{L}_{\mathbf{C}}$		1098		
NuTeV $\nu$	Fe	35 - 340	1170	Yes	[23]
NuTeV $\bar{\nu}$			966		
Chorus $\nu$	Ph	25 - 170	412	Yes	[27]
Chorus $\bar{\nu}$	1.0		412		
CCFR dimuon $\nu$	Fo	110 - 333	40	No	[10]
CCFR dimuon $\bar{\nu}$	L,G	87 - 266	38	INU	
NuTeV dimuon $\nu$	Fo	90 - 245	38	No	[10]
NuTeV dimuon $\bar{\nu}$	$\mathbf{L}\mathbf{G}$	79 - 222	34	INU	

#### representative neutrino DIS data





• In probing partonic structure of nucleons, neutrino CC DIS are complementary to the neutral current DIS



#### CT18 proton PDF fit

W>3.5 GeV, Q>2 GeV

dimuon data on strange PDFs

[Hou, JG+, 1912.10053]

#### nCTEQ nuclear PDF fit

Inclusive and dimuon data on PDFs of Pb

[2204.13157]





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# Charm quark production in CC DIS

and NOMAD; especially the NOMAD dimuon data have a high quality

Dimuon production from SIDIS



♦ In semi-inclusive CC DIS, charm-quark production with subsequent leptonic decays (so-called dimuon) events) can serve as a direct probe of the strange-quark parton distribution, e.g., as from NuTeV, CCFR,

Differential cross sections

$$\begin{aligned} \frac{d\sigma_c^{\nu}}{dxdy} &= \frac{G_F^2 M E}{\pi (1 + Q^2 / M_W^2)^2} \left[ \left( 1 - y - \frac{M x y}{2E} \right) F_{2,c}^{\nu}(x, Q^2) + \right. \\ &\left. + \frac{y^2}{2} F_{T,c}^{\nu}(x, Q^2) + y \left( 1 - \frac{y}{2} \right) x F_{3,c}^{\nu}(x, Q^2) \right] \end{aligned}$$

- Finite charm-quark mass corrections are important for kinematic region of neutrino DIS (E $\nu \leq$ hundreds GeV), known at LO via the slow scaling variable  $\chi \sim x_B(1+m_h^2/Q^2)$
- Charm-quark production has been calculated to NLO in QCD by T. Gottschalk (1981) and M. Gluck+(1997) in a closed analytic form
- No one bother to calculate the NNLO corrections for more than 30 years since the first NLO result

## Production cross section at NNLO

subtractions of infrared/collinear divergences



\* unresolved region: factorization derived in soft-

The two-loop production cross sections with exact charm-quark mass dependence are calculated numerically with MC integrations and with a phase-space slicing (on beam thrust) method to render [Berger, JG, Li, Liu, Zhu, 1601.05430; JG, 1710.04258]

Corrections with CCFR data kinematics

# GM-VFNs at NNLO for inclusive CC DIS



- If actorization on heavy-quark contribution to DIS was proved by Collins and result in ACOT-like HQ schemes
- the coefficient functions for CC DIS in simplified ACOT-**x**, a GM-VFN, scheme have been calculated to NNLO with full mass dependence

#### • General-mass(GM) VFN schemes are frameworks to predict full heavy-mass dependence, interpolating between FFNs and ZM-VFNs, have been widely used in PDF fits, e.g., ACOT-like, FONLL, RT schemes

[JG, Hobbs, Nadolsky, Sun, Yuan, 2107.00460]

application to total cross sections of neutrino CC DIS on iso-scalar target



Inite mass corrections are about 2% at low Ev comparing GM and ZM both at N2LO; scale variations are 1~3% for GM N2LO, and <1% for N3LO' at large Ev



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# Predictions for ultra-high energy neutrinos

IceCUBE) and are dominated by regions of deeply inelastic scatterings (DIS)

total cross sections vs. energy



and various uncertainties are required

 Neutrino scattering cross sections on nucleus, both charged-current (CC) and neutral-current (CC) are key inputs for calculating absorption rate and detection efficiency of ultra-high energy neutrinos (e.g.,



\* an important observation is DIS of ultra-high energy (UHE) neutrinos involves kinematic regions with rather small-x (<10<sup>-5</sup>) that are unexplored by all other experiments; careful examination on small-x extrapolation



# Uncertainty: Perturbative QCD calculations

results available at NNLO and partial results known at N3LO for CC and NC DIS



◆ Perturbative calculations on the coefficient functions can be improved order by order in QCD, with full



charm mass corrections

\* perturbative calculations show good convergence with remaining unc. at most a few percents

![](_page_13_Figure_0.jpeg)

al major analysis groups (CT, MSHT, NNPDF, ABM, HERAPDF, ATLASpdf, ant heavy-quark schemes, selections of data, and methodologies

- CT18 represents the state of art PDFs of proton with a faithful/ conservative estimation of uncertainties
- PDF unc propagated into neutrino cross sections at the level of <10%in general for energy upto EeV

## Uncertainty: Nuclear modifications

for charged-current scatterings

general feature of nuclear modifications

![](_page_14_Figure_3.jpeg)

[1401.2345]

◆ DIS structure functions or cross sections per nucleon are different for free nucleons and for nucleus, and depends on the kinematic regions (x); those nuclear modifications are especially uncertain at small-x and

![](_page_14_Figure_6.jpeg)

![](_page_14_Figure_7.jpeg)

spread of predictions from different groups can be as large as 20% for tens PeV neutrinos

![](_page_14_Figure_10.jpeg)

### Comparison to CSMS results

collaboration ratio of DIS total cross sections

1.21.2Rato to CT8.0  $\operatorname{CI}$ Rato to 8.0  $\nu(W)$  $\bar{\nu}(W)$ -CT18 —CT18  $-CT18 \otimes H_2O(EPPS21)$  $-CT18 \otimes H_2O(EPPS21)$ 0.6 0.6-CSMS -CSMS  $10^5$   $10^6$   $10^7$  $10^9 \quad 10^{10} \quad 10^{11}$  $10^2 \quad 10^3 \quad 10^4 \quad 10^5 \quad 10^6$  $10^{8}$  $10^{3}$  $10^{4}$  $10^{2}$  $E_{\nu}$  [GeV] 1.2 1.2Rato to CT 8.0 R  $\operatorname{CT}$ Rato to 8.0  $\nu(Z)$  $\bar{\nu}(Z)$ -CT18 —CT18  $-CT18 \otimes H_2O(EPPS21)$  $-CT18 \otimes H_2O(EPPS21)$ 0.60.6 -CSMS —CSMS  $10^9 \quad 10^{10} \quad 10^{11}$  $10^3$   $10^4$  $10^{5}$  $10^6 \quad 10^7 \quad 10^8$  $10^2 \quad 10^3 \quad 10^4$  $10^{5}$  $10^{6}$  $10^{2}$ 

 $E_{\nu}$  [GeV]

◆ Our final predictions (from CT18+nuclear corrections) are compared to the CSMS predictions [Cooper-Sarkar, Mertsch, Sarkar 2011] which are the benchmarks of DIS cross sections used by the IceCube

![](_page_15_Figure_4.jpeg)

- three results are compared: CT18 with (w/o) nuclear corrections, CSMS
- central predictions differ especially for CC DIS by upto ~20%; can affect neutrino transmission probability and also the measured flux
- Iarger unc. from PDFs and nuclear corrections especially for UHE neutrinos

 $E_{\nu}$  [GeV]

### Revisit of QED corrections

![](_page_16_Figure_1.jpeg)

$$\left(\frac{\log\left(\frac{1}{m_{\ell}^2}\right)\left[\frac{1-z}{1-z}\right]_+ + \frac{1}{2}\delta(1-z)\right]$$

QED corrections lead to a systematic shift/smeaning of kinematic variable if using bare kinematics; similar effects also studied recently in the scenario of EIC using new factorization approach of QED correction in DIS [Qiu, Liu+, 2022]

#### 1.05

# DIS from collider neutrinos a

 New opportunities from forward detectors a probing region of ~ 1 TeV neutrino energies;

Detector	before cuts	after DIS and acceptance cuts	acceptance
Detector	$N_{\nu_e} + N_{\bar{\nu}_e}, \ N_{\nu_{\mu}} + N_{\bar{\nu}_{\mu}}$	$N_{\nu_e} + N_{\bar{\nu}_e}, \ N_{\nu_{\mu}} + N_{\bar{\nu}_{\mu}}$	$N_{\nu_e} + N_{\bar{\nu}_e}, \; .$
$\mathrm{FASER} u$	1.2k, 4.1k	610, 1.8k	51%, 4
SND@LHC	280,860	260, 700	92%, 8
$FASER\nu 2$	270k, 980k	170k, 510k	63%, 3
AdvSND-far	19k, 66k	18k, 56k	95%,8
FLArE10	65k, 202k	64k, 110k	98%, 3
FLArE100	427k, 1.3M	420k, 670k	98%, 3

#### Expected CC DIS event rates

![](_page_17_Figure_6.jpeg)

#### **DIS** kinematics

#### [2309.09581]

![](_page_17_Figure_9.jpeg)

next-generation proposals show potential strong constraint on PDFs<sup>1</sup> especially for strange quark, also possibility on pining down nuclear corrections

10<sup>-3</sup>

![](_page_17_Figure_14.jpeg)

![](_page_17_Picture_15.jpeg)

### Summary

- analyses of proton and nuclear parton distribution functions
- nuclear modifications, even more to expect with future EIC
- energies of a few tens GeV and below, and much more efforts are needed

• Neutrino DIS data play important role in development of QCD and still provide key inputs to global

Neutrino DIS cross sections can be well described by perturbative QCD approach within collinear factorization, with high precisions due to great efforts from loop calculations, fits of proton PDFs and

However, combining DIS predictions with other low-Q/non-DIS contributions are crucial for neutrino

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# Thank you for your attention!

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+ However, combining DIS predictions with other low-Q/non-DIS contributions are crucial for neutrino