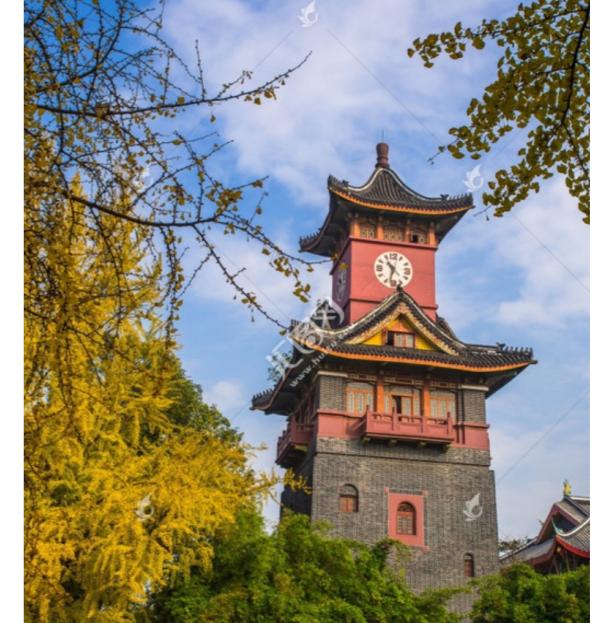
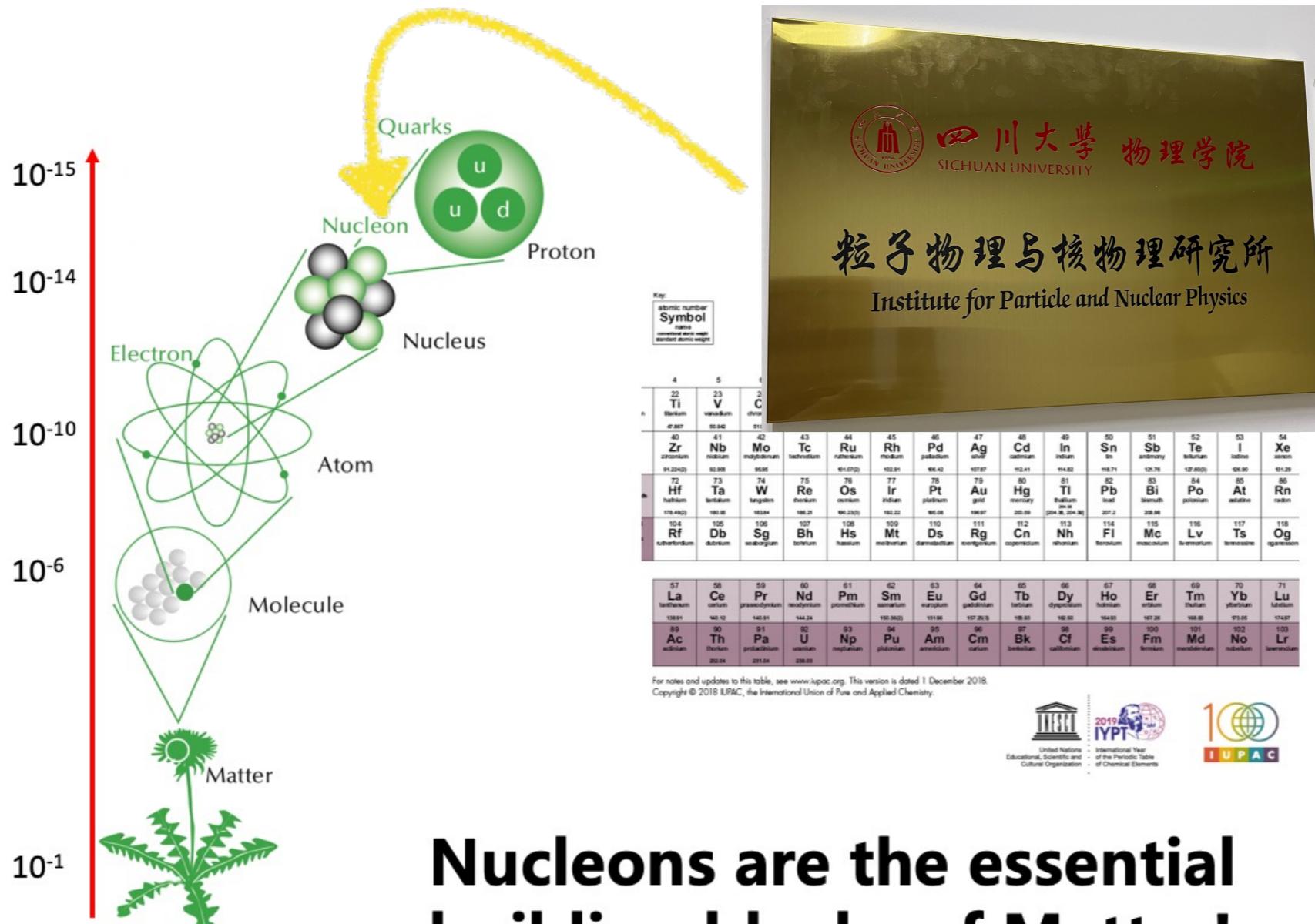


Nuclear weak currents from chiral effective field

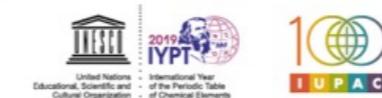
龙炳蔚
Bingwei Long



粒子物理与核物理结合，大有可为

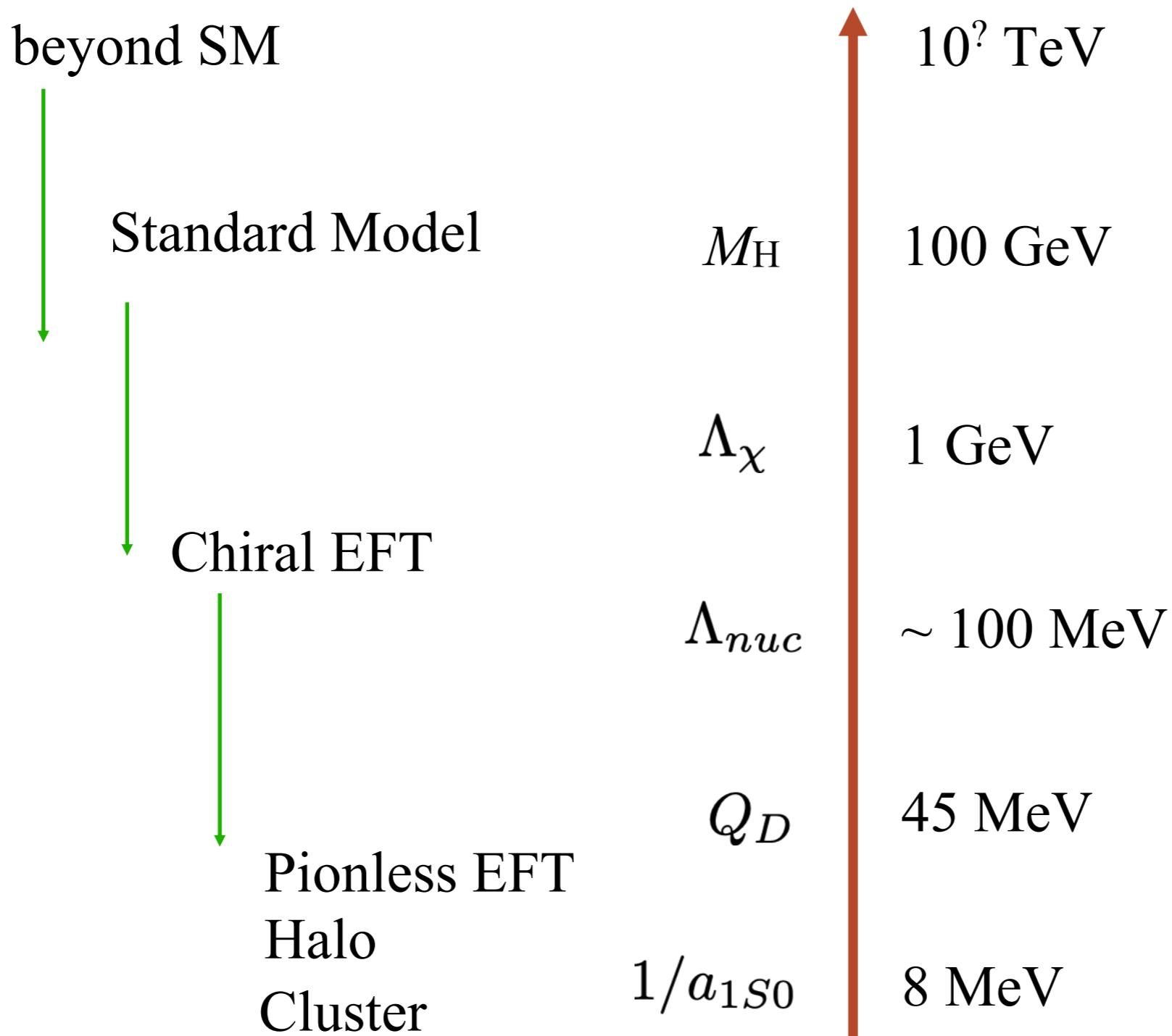


For notes and updates to this table, see www.iupac.org. This version is dated 1 December 2018.
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**Nucleons are the essential
building blocks of Matter!**

Hierarchy of EFTs



Relay: from quarks & gluons to Uranium

Low-energy constants



Chiral EFTs

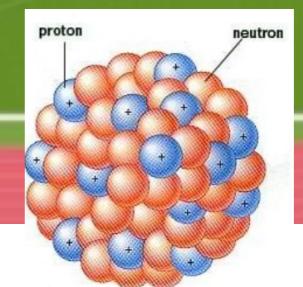
Lattice QCD

A: number of nucleons \uparrow

Light nuclei
 $: A = 4 \sim 12$

Many-body methods

$^2\text{H} \& ^3\text{H}$



EFT is a strategy



- Predictive power is w/ the SM (QCD)
- If EFTs have many LECs, so be it
- Job of EFTs:
 - where LECs lie;
 - *not* the values of LECs

QCD → Chiral Lagrangian

External sources: v, a, s, ps

$$\mathcal{L}_{QCD} = \bar{q}_R i\gamma_\mu D^\mu q_R + \bar{q}_L i\gamma_\mu D^\mu q_L - \frac{1}{4} G_{\mu\nu}^a G^{a\mu\nu} + \bar{q}\gamma^\mu(v_\mu + \gamma_5 a_\mu)q - \bar{q}(s - i\gamma_5 p_s)q$$

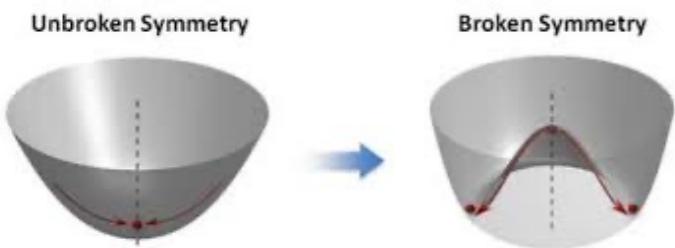
$$r_\mu = v_\mu + a_\mu = -eQA_\mu,$$

$$l_\mu = v_\mu - a_\mu = -eQA_\mu - \frac{g}{\sqrt{2}}(W_\mu^\dagger T_+ + h.c),$$

$$s = \mathcal{M} \equiv \text{diag}(m_q, m_q),$$

Incomplete!

$$p_s = 0.$$

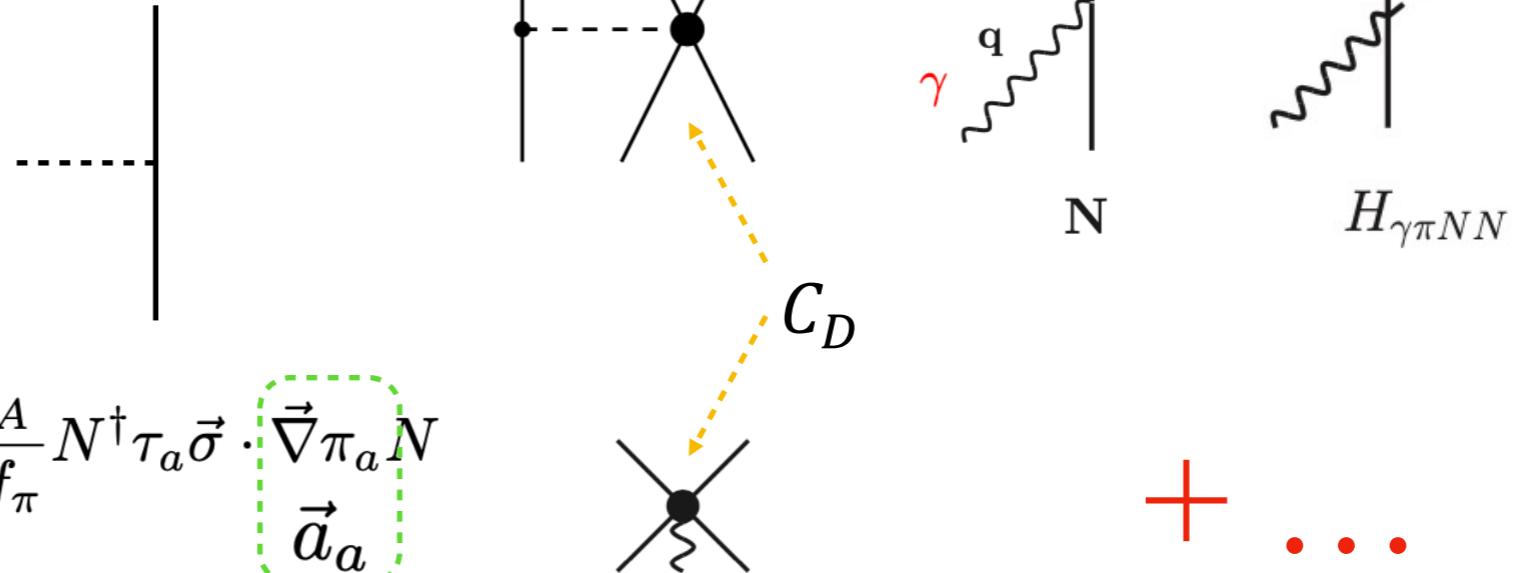
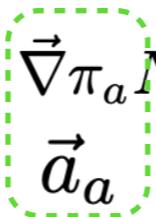


nonlinear realization
CCWZ; Weinberg; ...



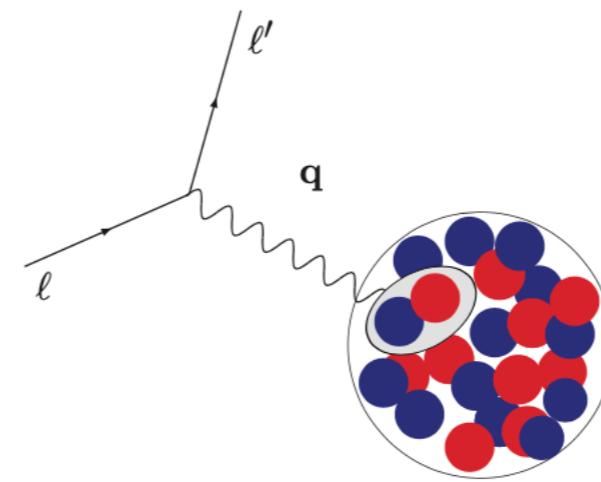
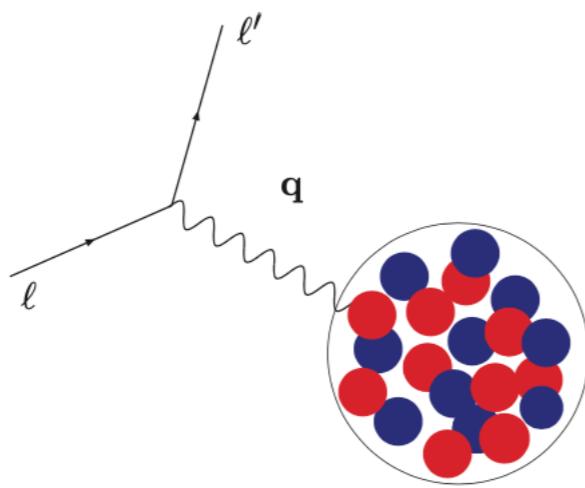
Nucleonic operators

$$-\frac{g_A}{2f_\pi} N^\dagger \tau_a \vec{\sigma} \cdot \left(\vec{\nabla} \pi_a N \right)$$

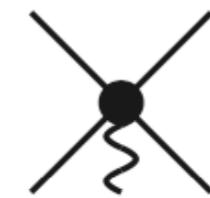


1N and 2N currents

- NN correlations may be enhanced by long-distance scale of NN correlations

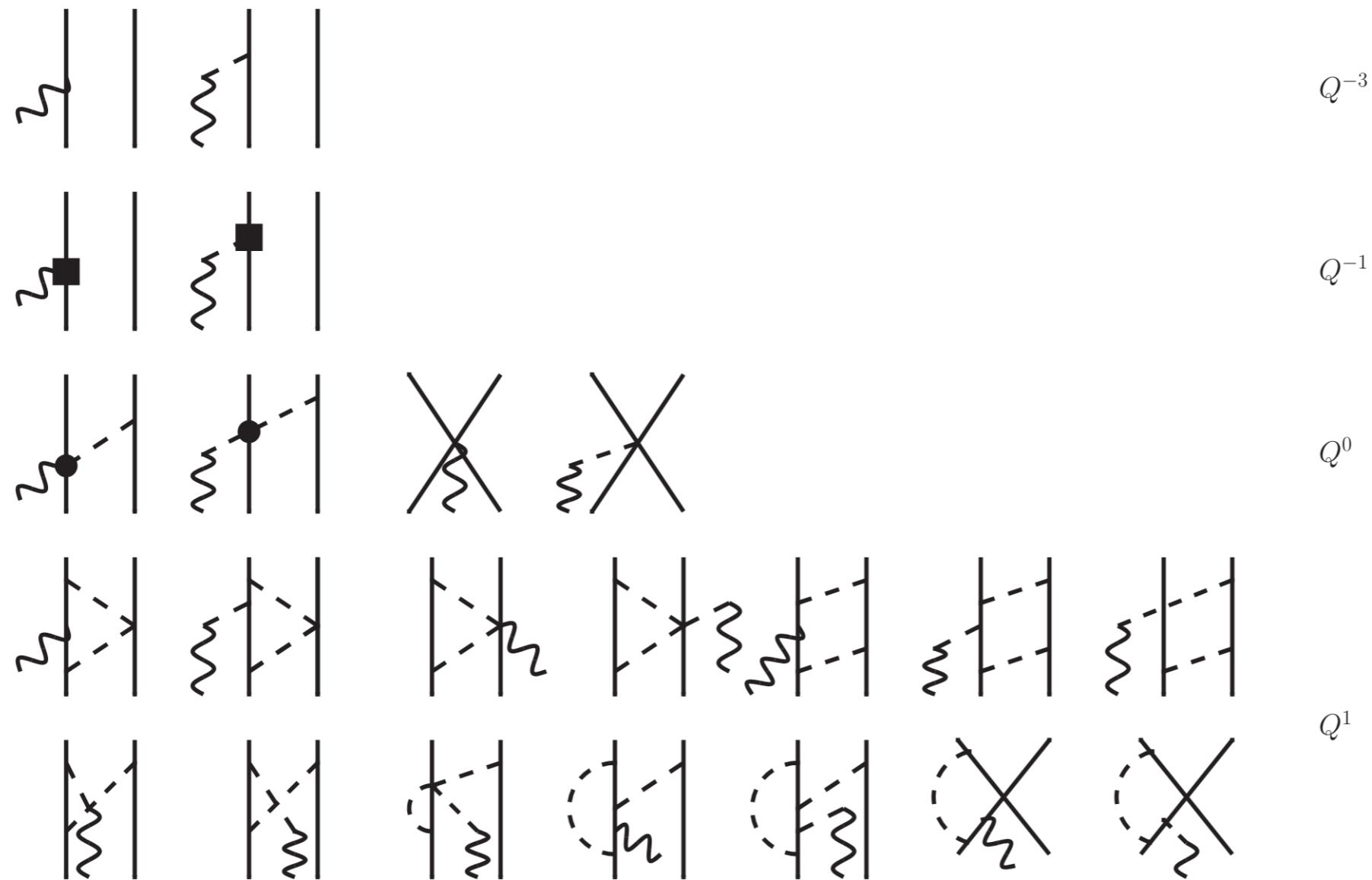


From Pastore's talk



Axial current diagrams

From Girlanda's talk



Riska & Schiavilla, arxiv:1603.01253

Krebs, arxiv:2008.00974

NN potentials from ChEFT

- Contact forces are doing heavy lifting

24 contact terms
in NDA up to Q^4

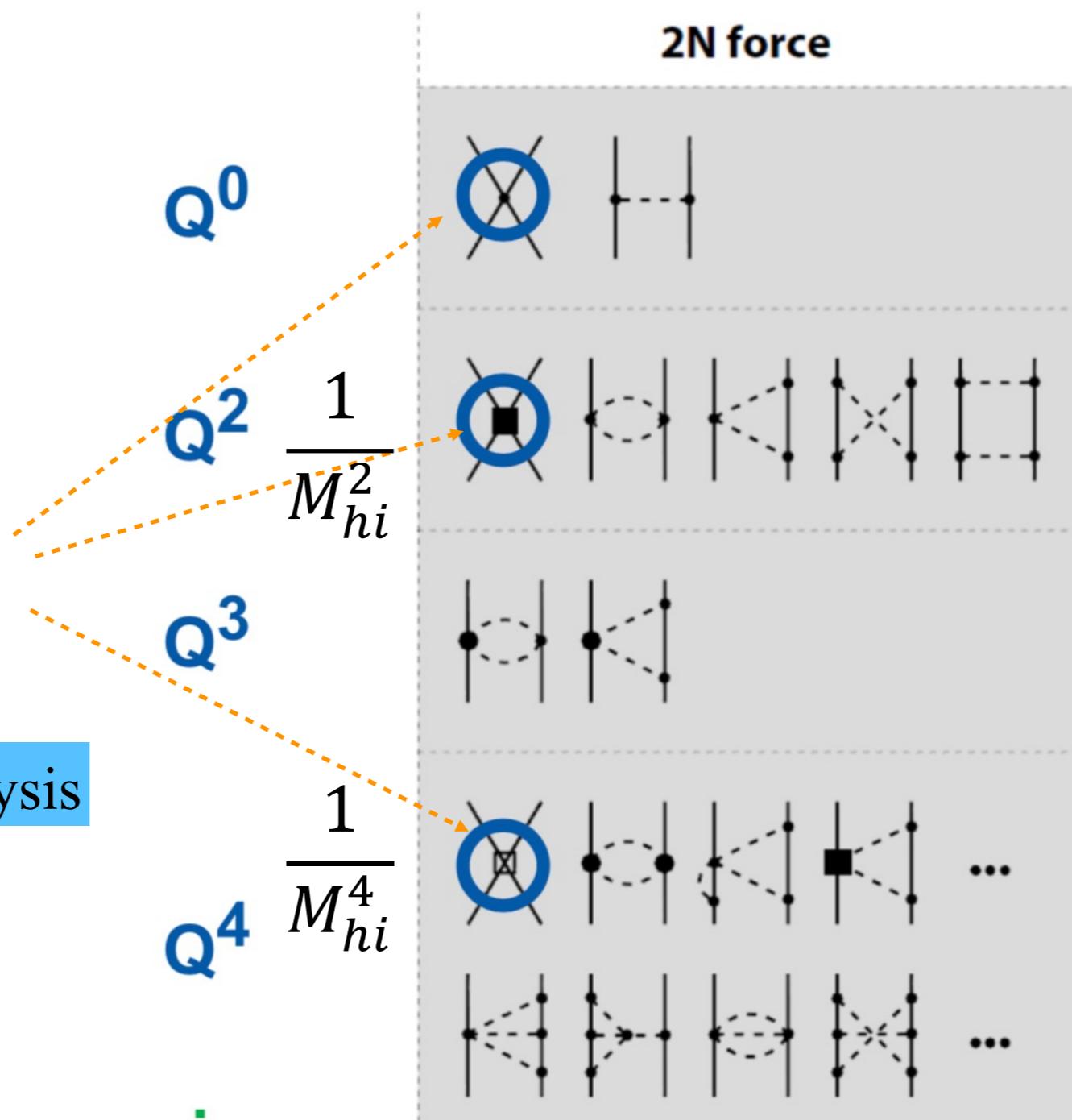
NDA = naive dim. analysis

van Kolck et al. since '92

Entem & Machleidt since '03

Epelbaum et al. since '99

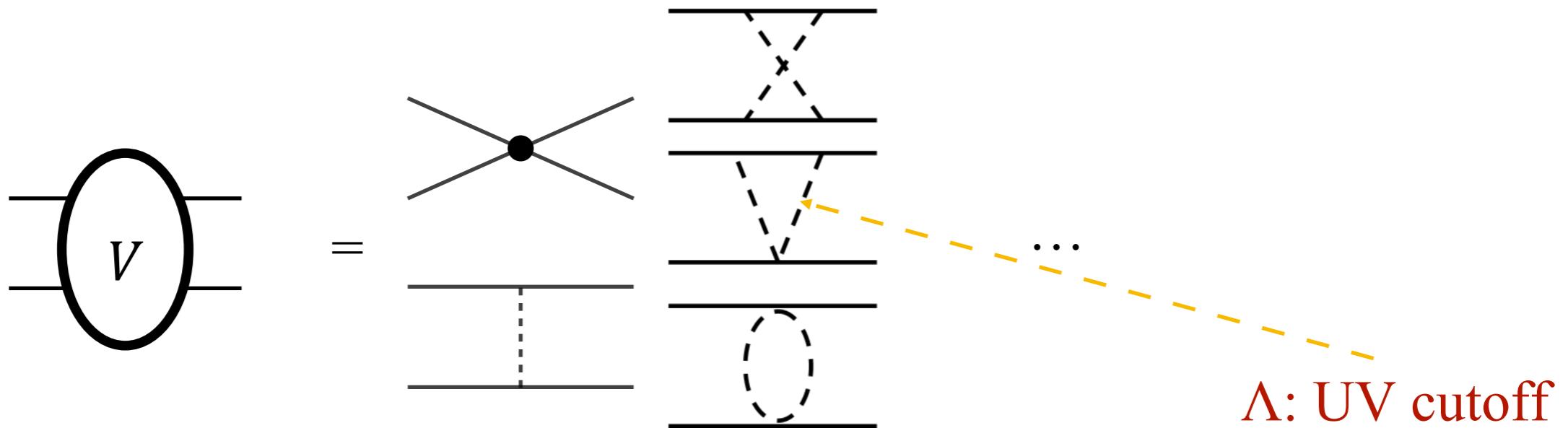
Pastore, Piarulli, et al. since '09



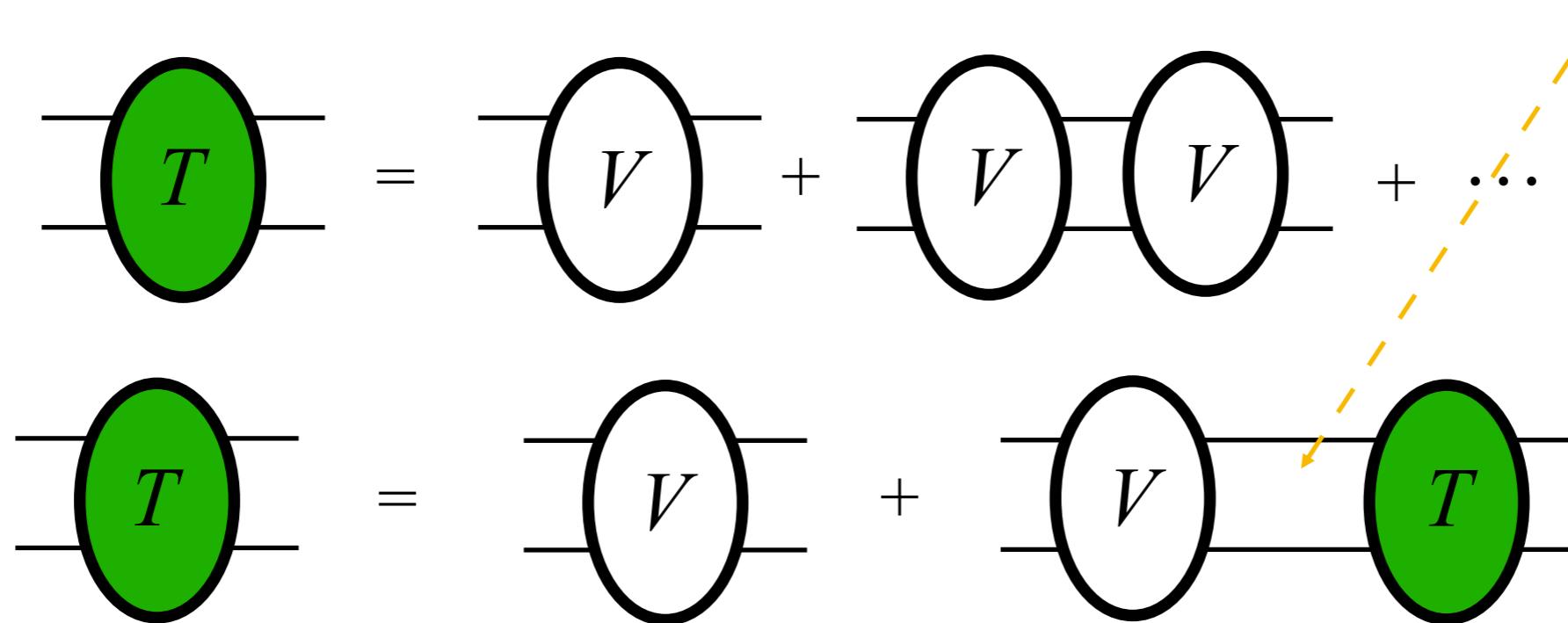
...

From potential to amplitude

- “ NN Potential”: two-nucleon irreducible diagrams



- Lippmann-Schwinger eqn (equivalent to Schrodinger eqn)

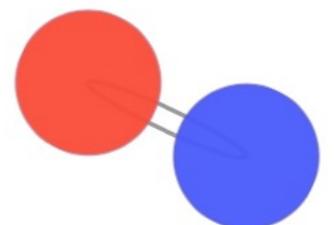


Few-body methods

- Equivalent to Schrodinger eqn
- Precise numerical solutions

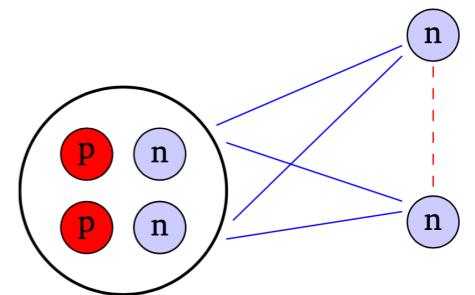
Easy

- Two-body: NN, N-cluster, cluster-cluster
Lippmann-Schwinger equation



Medium

- Three-body: NNN, 2n-cluster, 3alpha, ...
Faddeev equation



Hard

- Four-body: NNNN, 3n-alpha ...
Faddeev/Yakubowski equation



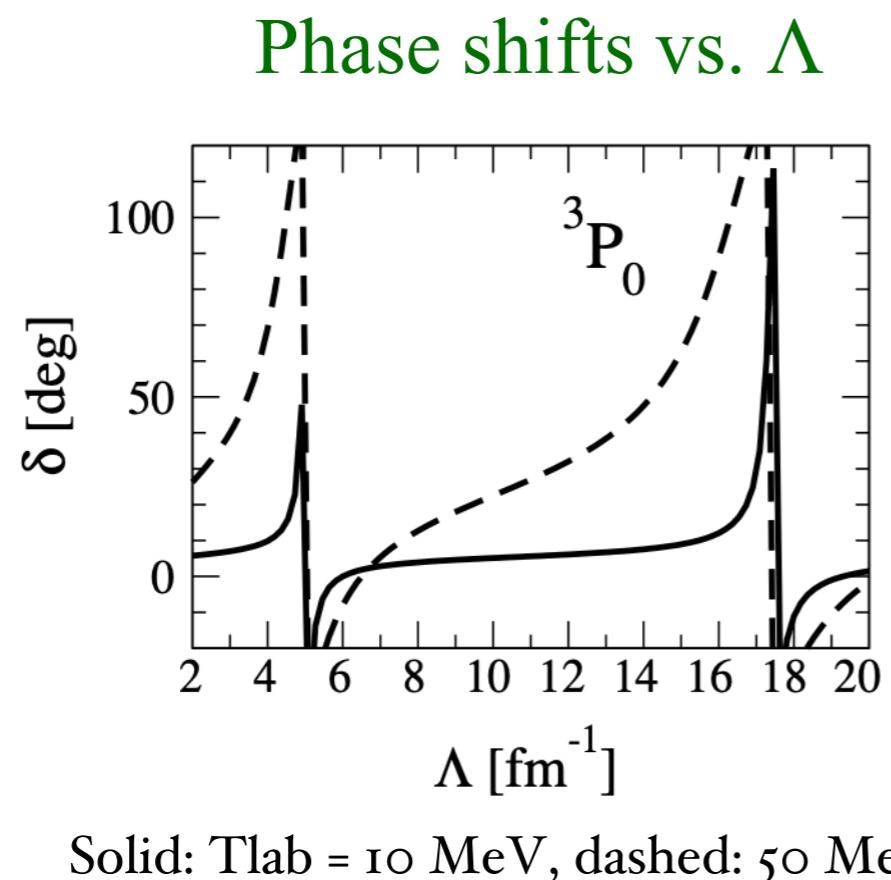
Strength of one-pion-exchange

$$\begin{array}{ccc} \left| \right. & Q \sim m_\pi & \left| \right. \\ \left| \right. \cdots \left| \right. & \sim \frac{1}{f_\pi^2} \frac{Q^2}{m_\pi^2 + Q^2} \sim \frac{1}{f_\pi^2} & \left| \right. \cdots \left| \right. \\ \left| \right. & & \left| \right. \\ & & \sim \frac{1}{f_\pi^2} \frac{m_N}{4\pi f_\pi} \frac{Q}{a_l f_\pi} \end{array}$$

- Strength of OPE $\sim a_l f_\pi$
- $a_l \sim 1$ for small l , $a_l \gg 1$ for large l
- $a_l f_\pi$ may have impact on contacts through renormalization
- coexistence of $a_l f_\pi$ and M_{hi} makes NDA no longer reliable

Renormalizing singular attraction

Nogga, Timmerman & van Kolck (2005)



Naive dimensional
analysis (NDA)

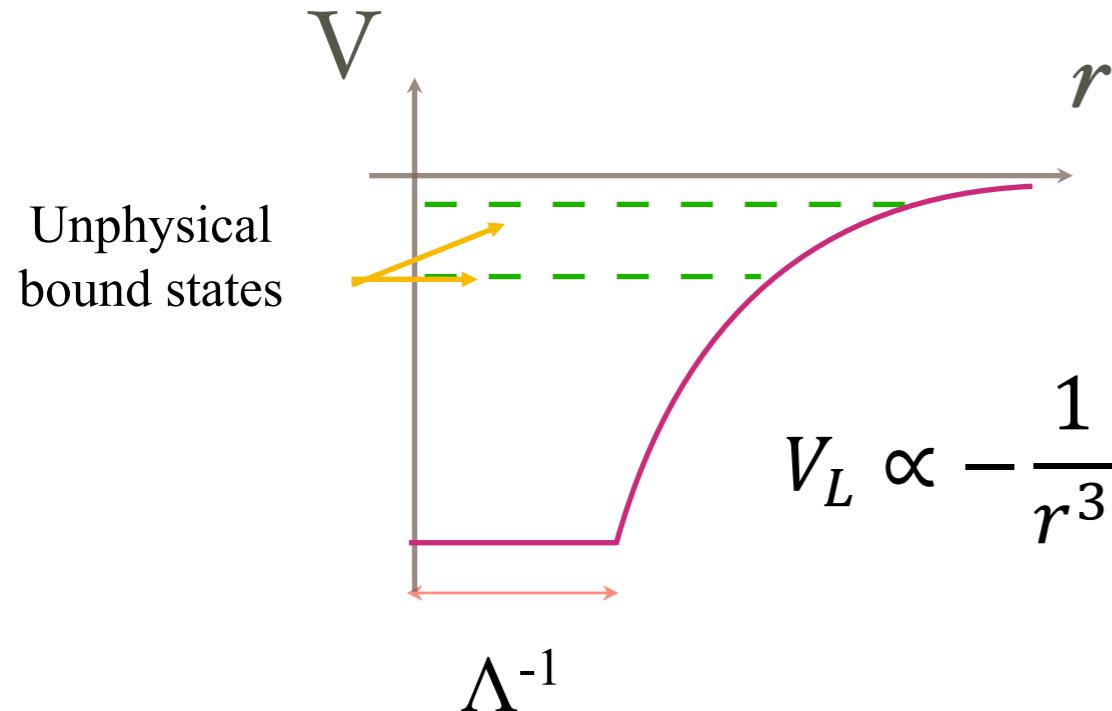
$$C_{3P0} \vec{p} \cdot \vec{p}' \sim \frac{Q^2}{m_{hi}^2} \quad \text{N2LO}$$

Renormalizing singular attraction

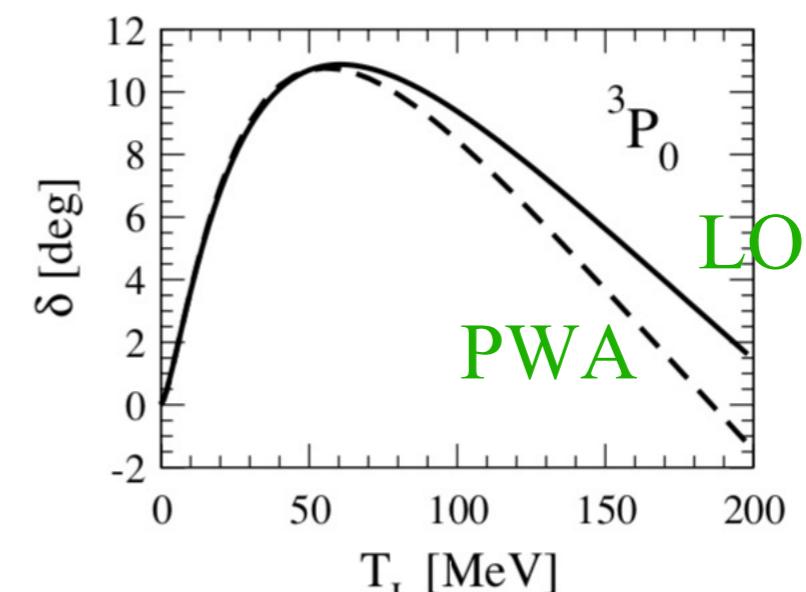
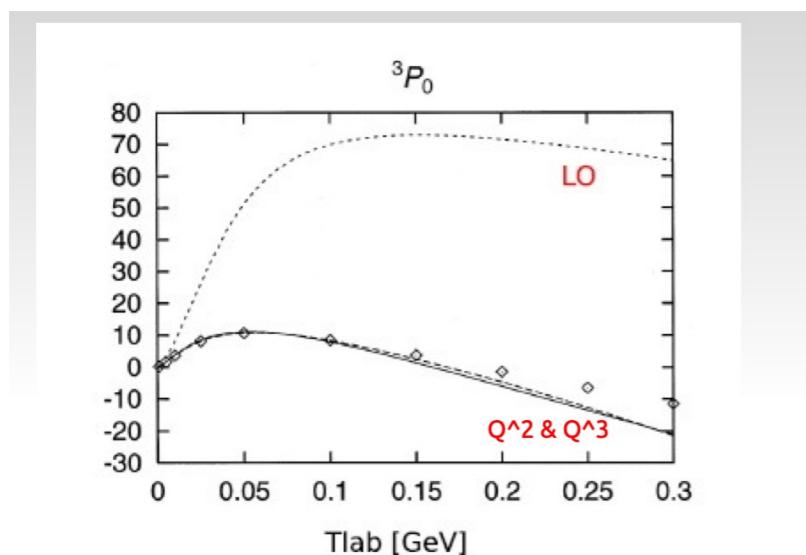
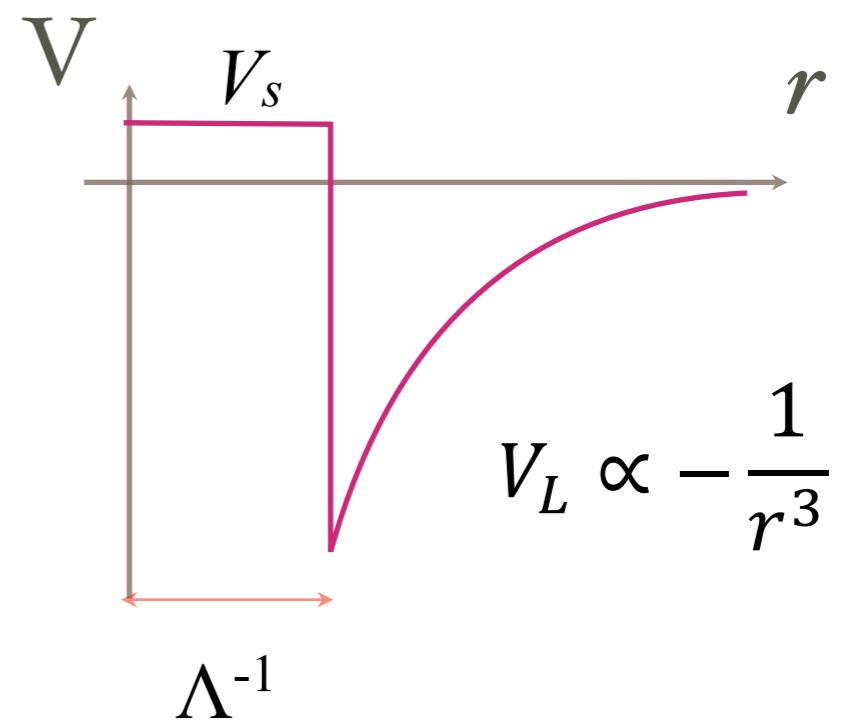
Beane et al ('01)

Pavon Valderrama & Ruiz Arriola ('05 ~ '07)

Nogga et al ('05)



Modify PC
⇒



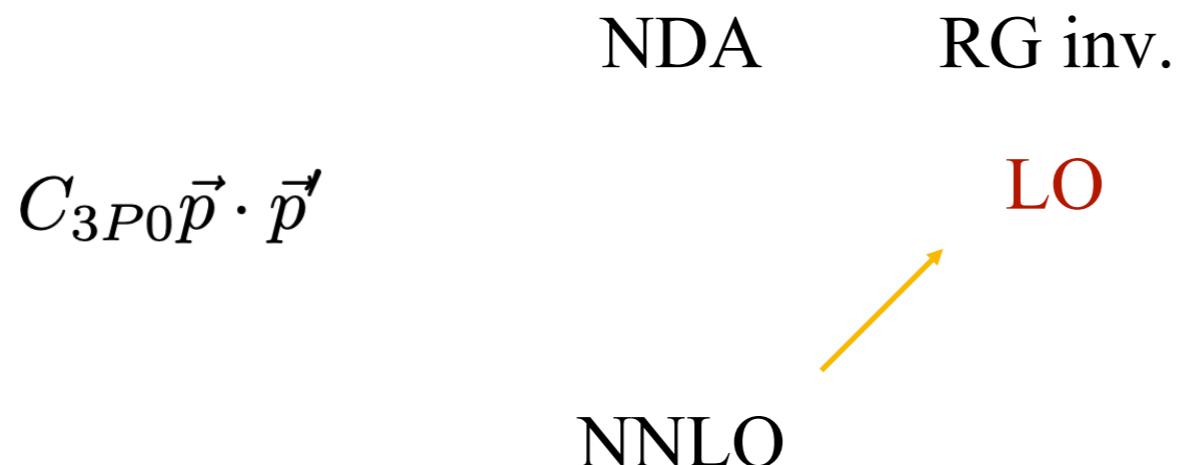
RG as diagnostic tool for PC

- The cutoff dependence of v th order must be of $v+1$ order or smaller

$$\frac{\Lambda}{T^{(\nu)}(Q, \Lambda)} \frac{dT^{(\nu)}(Q, \Lambda)}{d\Lambda} = \mathcal{O}\left(\frac{Q^{\nu+1}}{M_{hi}^\nu \Lambda}\right)$$

- A consistent power counting (PC) must be a solution to RGE

Modify PC if necessary!



Modified power counting for chiral nuclear forces

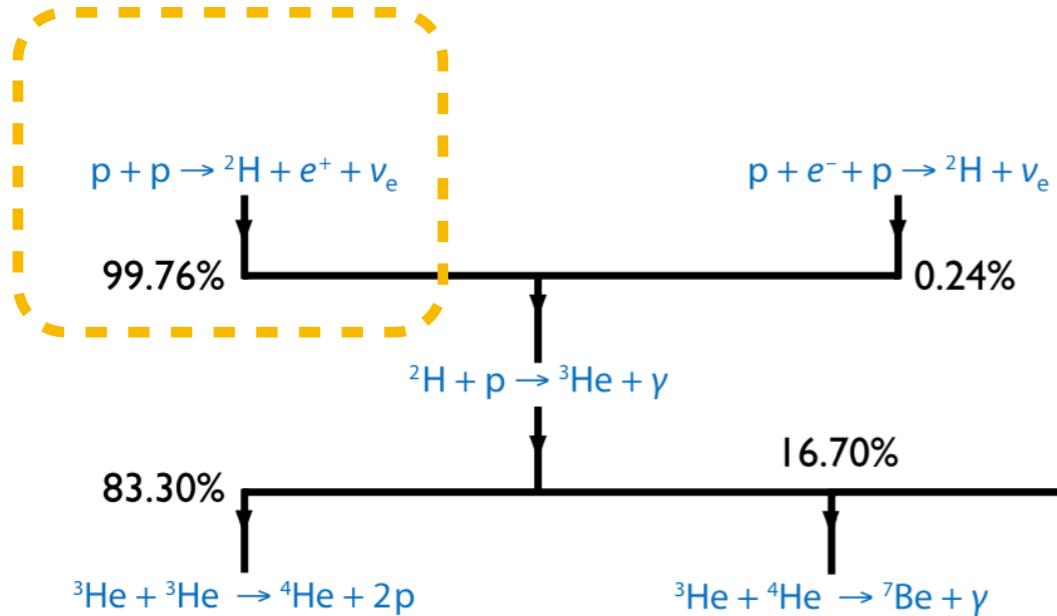
Nogga et al. '05

BwL & Yang '11, '12

Wu & BwL '19

- PC varies for different partial waves
- LO : (C + OPE) for 1S0, 3S1, **3P0** (perturbative OPE for most waves)
- NLO : **Q² C.T.** for 1S0; OPE for 1P1, 3P1, 3P2...
- N2LO: (**Q⁴ C.T.** + TPE) for 1S0; (Q² C.T. + TPE) for 3S1-3D1 and 3P0
- N3LO:

PP fusion



Park, Kubodera, Min & Rho '98
Marcucci, Schiavilla, Viviani '13
Acharya et al, '16
Valderrama & Phillips, '15

$$\langle \psi_d^M | A_-^i | \psi_{pp} \rangle = \delta_{Mi} \sqrt{\frac{32\pi}{\gamma^3}} g_A C_0 \Lambda_R(p)$$

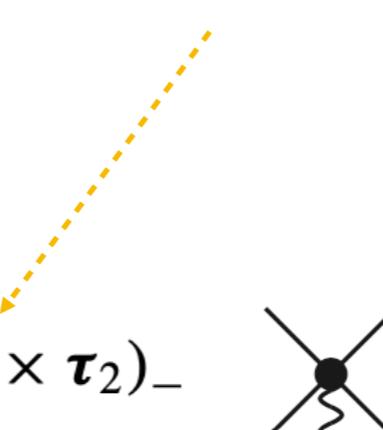
$$\vec{A}_-^{(0)}(\vec{p}', \vec{p}) = -g_A \sum_i \vec{\sigma}_i \tau_{i,-}$$

$$\vec{A}_-^{(2)}(\vec{p}', \vec{p}) = \frac{g_A}{2m_N^2} \sum_i [\vec{K}^2 \vec{\sigma}_i - (\vec{\sigma}_i \cdot \vec{K}) \vec{K}] \tau_{i,-}$$

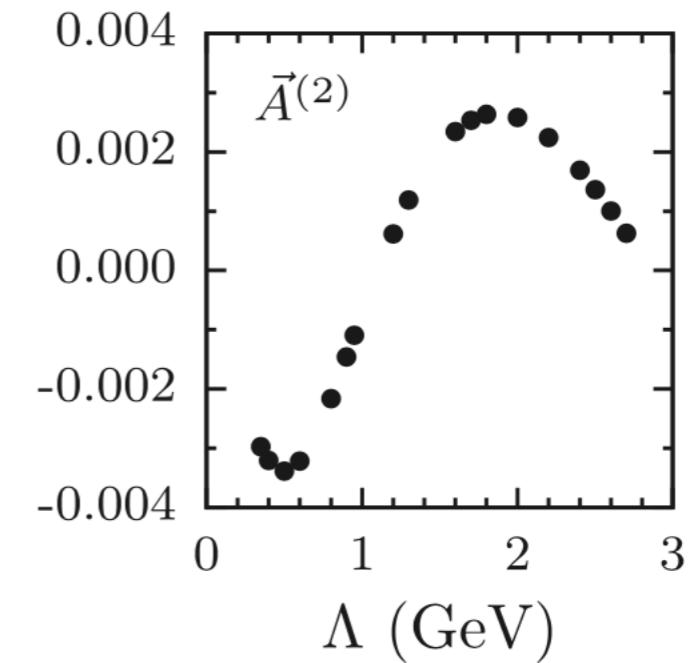
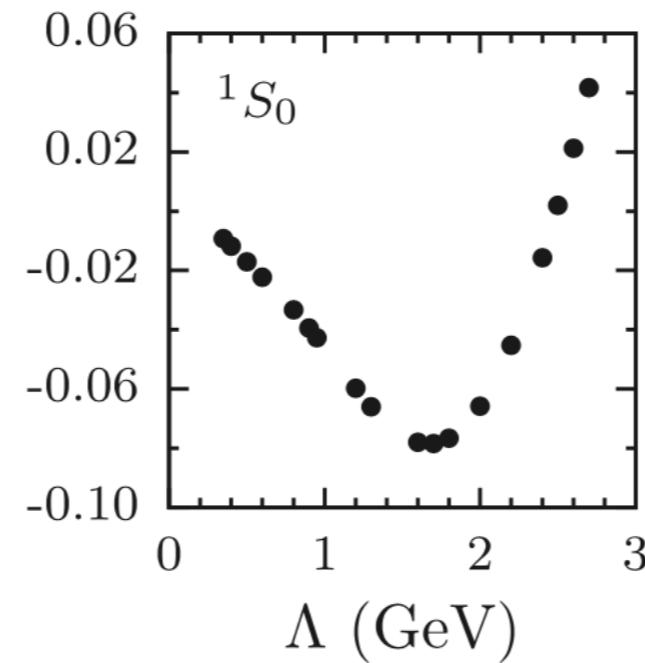
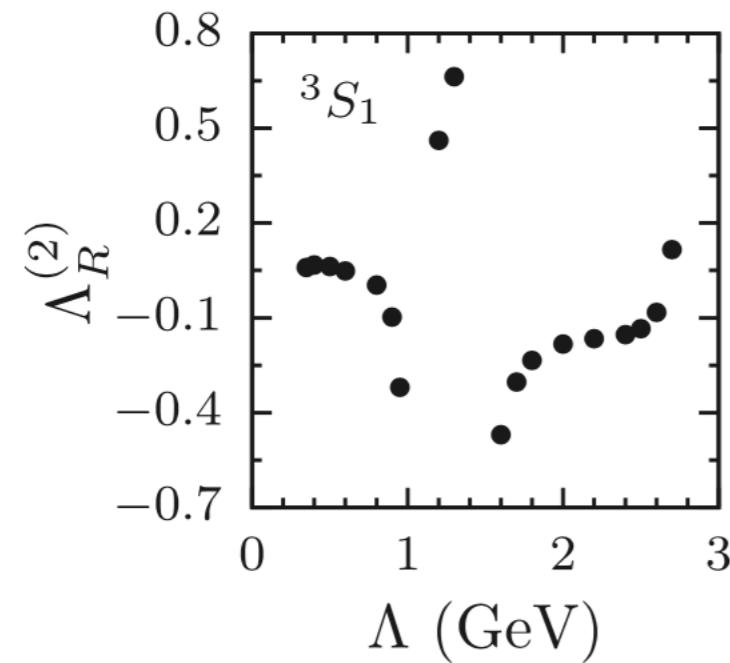
$$\vec{A}_{ct}(\vec{p}', \vec{p}) = \hat{d}_R \vec{\sigma}_1 \times \vec{\sigma}_2 (\boldsymbol{\tau}_1 \times \boldsymbol{\tau}_2)_-$$

- Two-body currents accounting for a few percents, an
- Promoted from N3LO to N2LO

Liu, Peng, Lyu & BwL '22



N2LO vs cutoff

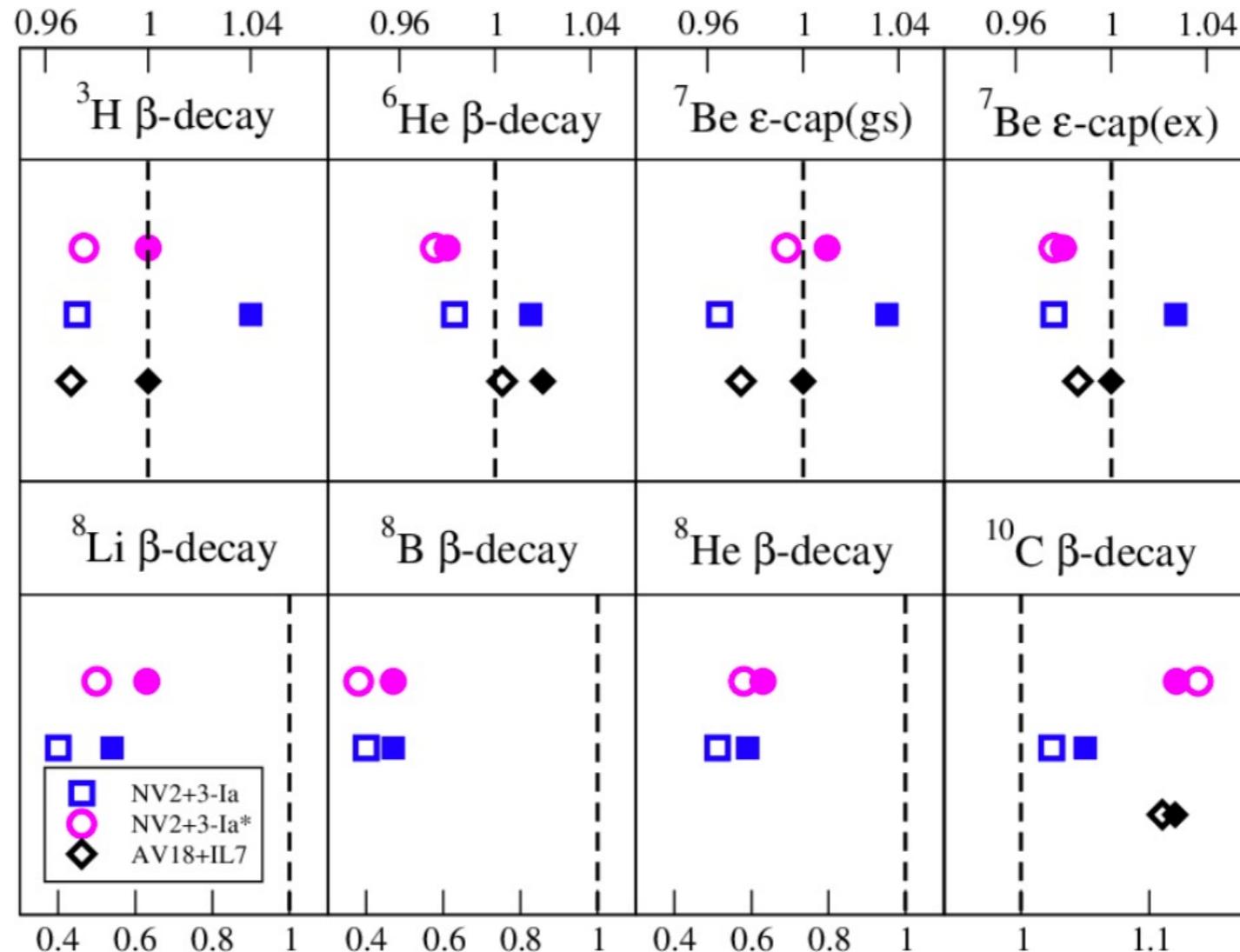


Beta decays by QMC

Baroni et al. '18

LO — empty

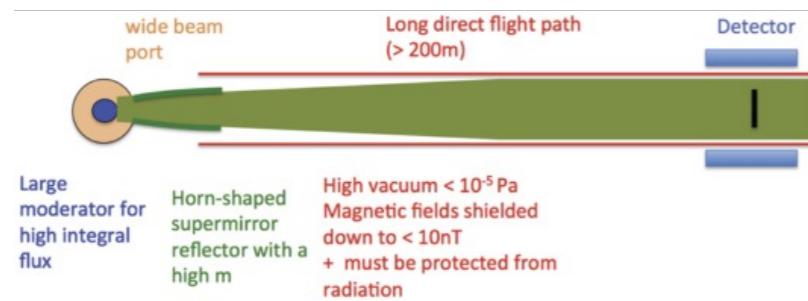
N3LO — filled



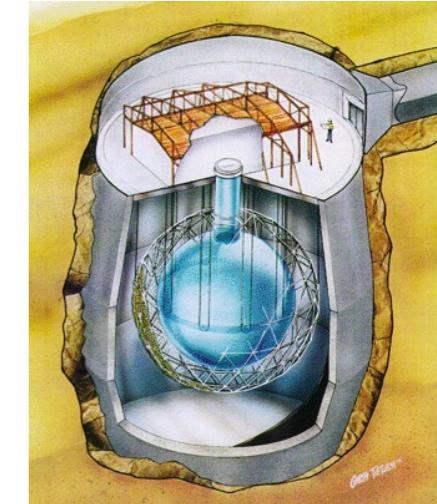
- One-body dominates for most nuclei
- gA quenching no longer an issue once many-body solvers improve
- For some nuclei, high-order corrections do not go right direction

Neutron-antineutron oscillation

- Some BSM models favor baryon-number violation $|\Delta B| = 2$
- Can explain baryon asymmetry of universe
- Stable nuclei become “unstable”
- Can we relate $\tau_{n-n\bar{b}ar}$ to Γ_d ?
- EFT helps disentangle different B-violating physics

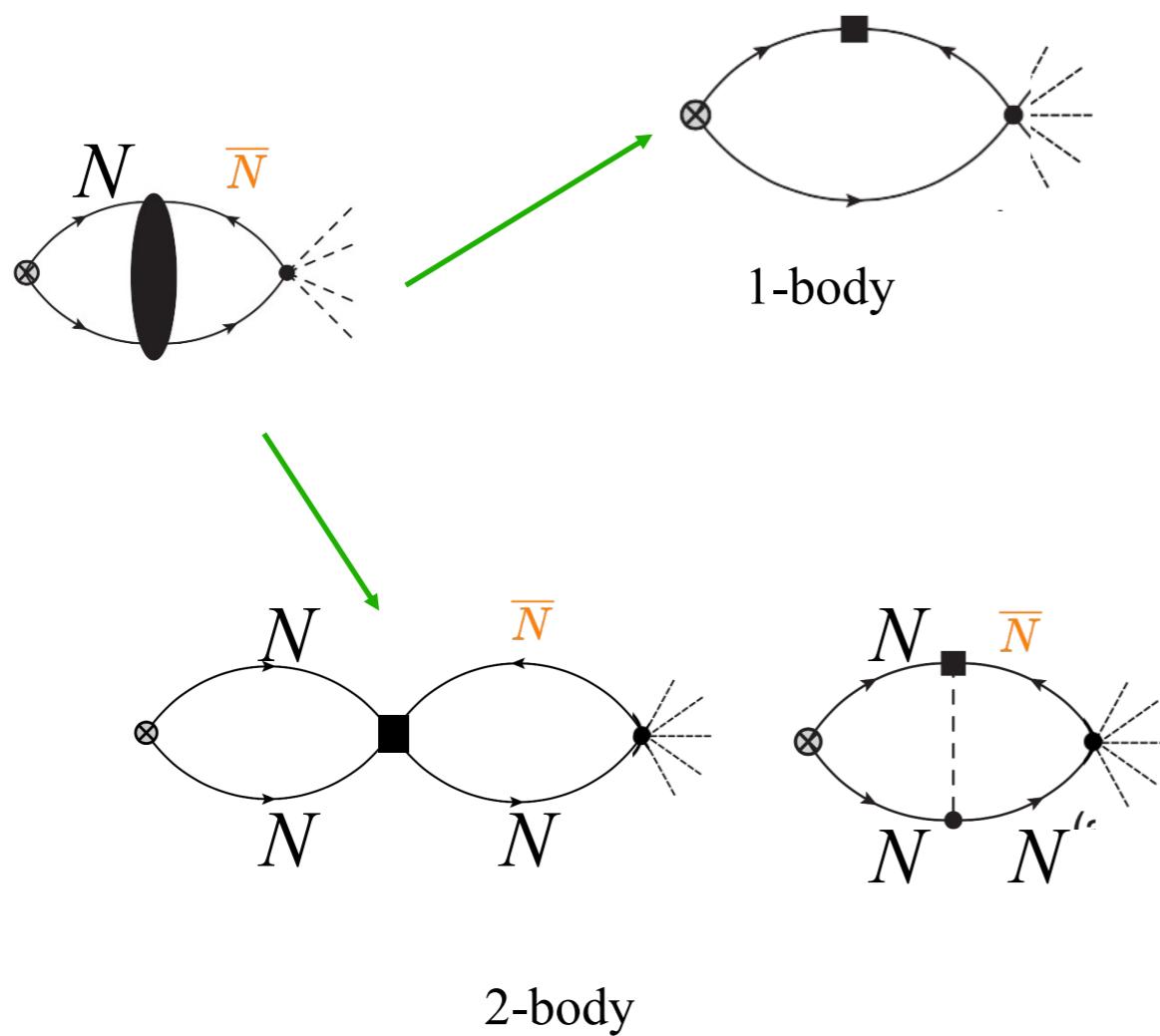


In vacuum (European spallation source)



Limit on deuteron lifetime (SNO)

$NN \rightarrow NNbar$



$$R_d \equiv \Gamma_d^{-1} / \tau_{n\bar{n}}^2$$

$$R_d = - \left[\frac{m_N}{\kappa} \text{Im} a_{\bar{n}p} (1 + 0.40 + 0.20 - 0.13 \pm 0.4) \right]^{-1}$$

$$= (1.1 \pm 0.3) \times 10^{22} \text{ s}^{-1}.$$

Uncertainty

Oosterhof, BwL, et al. PRL 122, 17

Summary

- QCD Lag + chiral symmetry + external
—> chiral Lagrangian w/ vector and axial vector sources
- Pion-exchange potentials and currents have been calculated to N4LO or higher
- Two-body currents may be enhanced due to nucleon-nucleon correlations at long distances
- Few-nucleon systems are testbeds for understanding of chiral forces and currents