

Interaction of fully heavy hadrons

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Long-term fruitful collaboration with my colleagues and co-authors is gratefully acknowledged!

- X. K. Dong, V. Baru, F. K. Guo, C. Hanhart, A. Nefediev,
“Coupled-Channel Interpretation of the LHCb Double- J/ψ Spectrum
and Hints of a New State Near the $J/\psi J/\psi$ Threshold,”
Phys. Rev. Lett. **126**, 132001 (2021)
- X. K. Dong, V. Baru, F. K. Guo, C. Hanhart, A. Nefediev, B. S. Zou,
“Is the existence of a $J/\psi J/\psi$ bound state plausible?,”
Sci. Bull. **66**, 2462 (2021)
- X. K. Dong, F. K. Guo, A. Nefediev, J. T. Castellà,
“Chromopolarizabilities of fully-heavy baryons,”
Phys. Rev. D **107**, 034020 (2023)

Motivation: Near-threshold fully heavy exotic states

Why to study?

- Encodes information on **internal structure** of heavy hadrons
- Relevant for **molecular** model
- Allows making **predictions** for exotic states

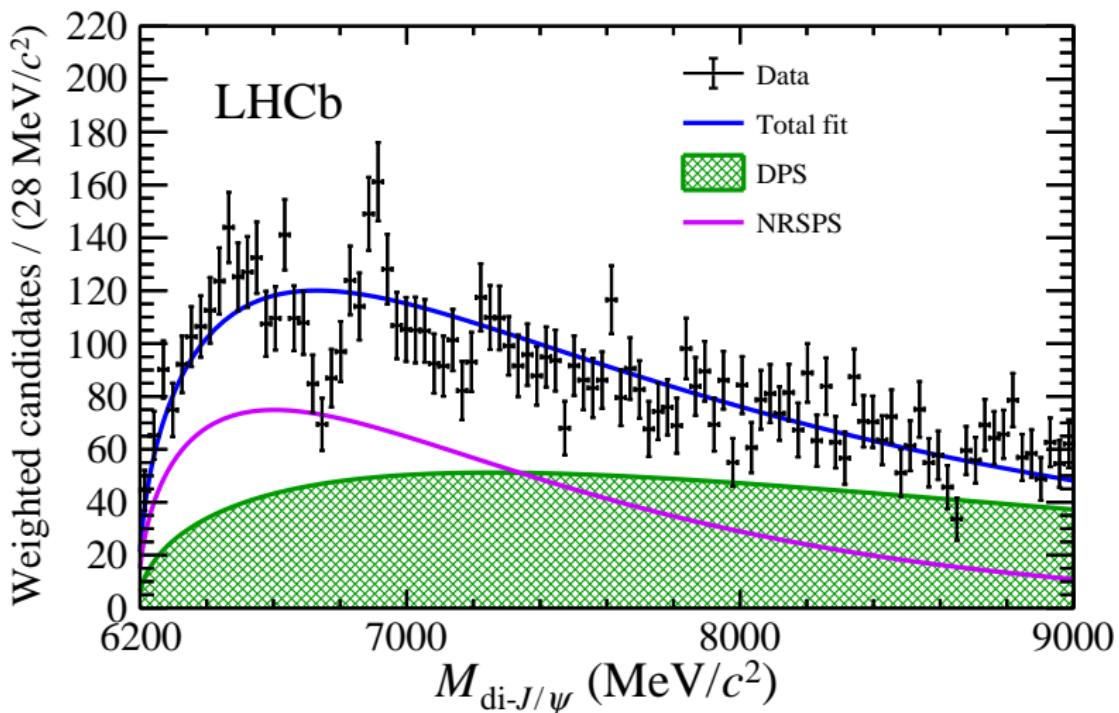
Do we have adequate theoretical approaches?

- Effective field theory
- Coupled channels
- Multipole expansion in QCD

Do we have promising candidates?

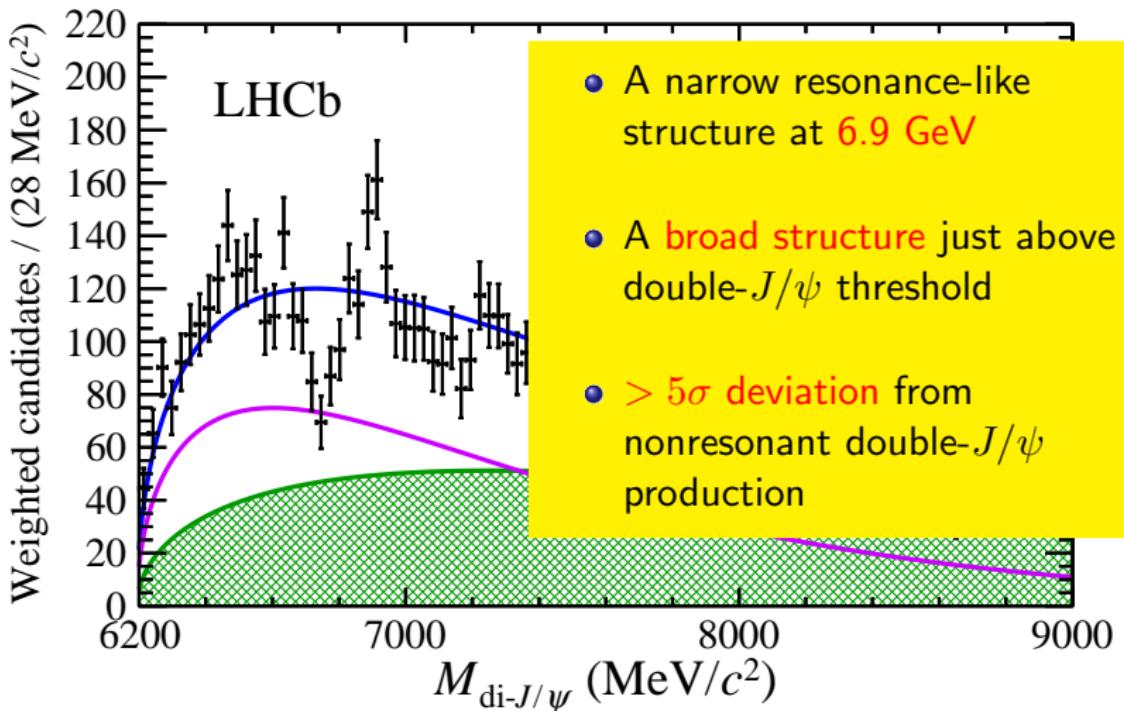
- Di-charmonium production at LHC
- Lattice simulations of heavy di-baryons

Double- J/ψ production @ LHCb (Sc.Bull. 65 (2020) 1983)



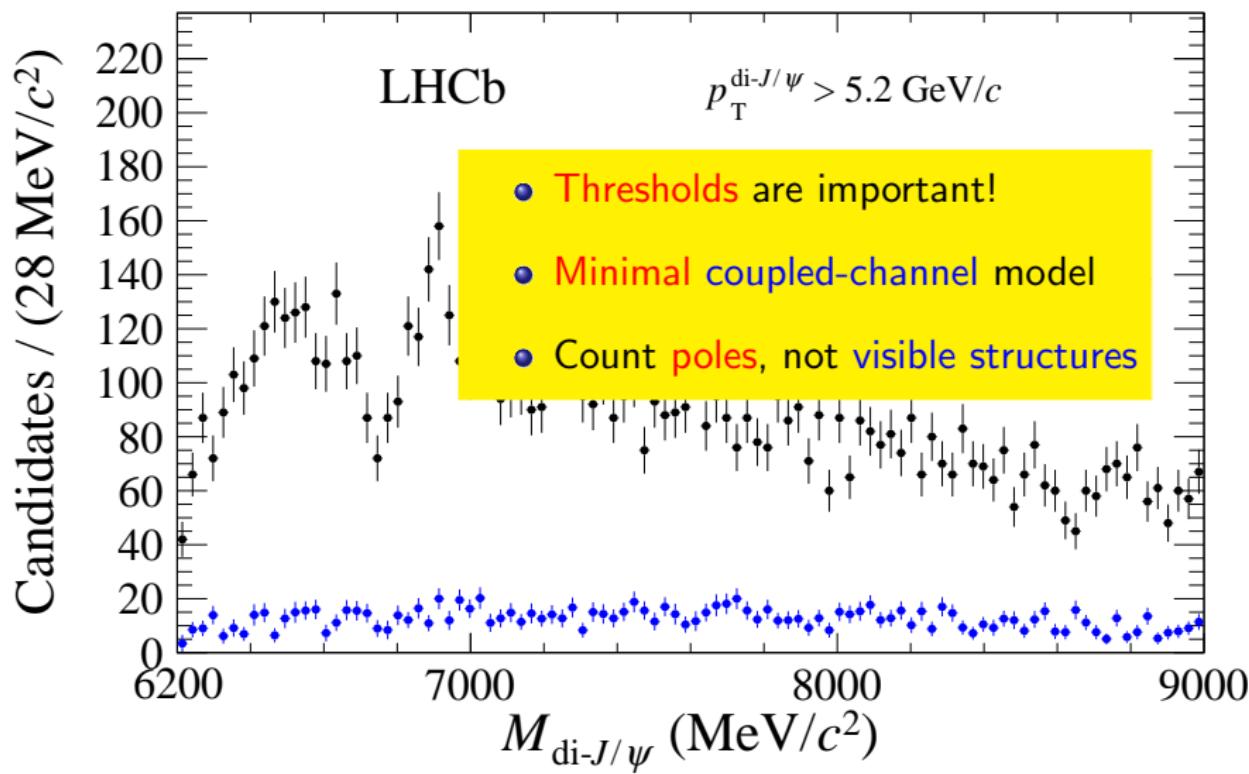
NRSPS=NonResonant Single Parton Scattering
DPS=Double Parton Scattering

Double- J/ψ production @ LHCb (Sc.Bull. 65 (2020) 1983)

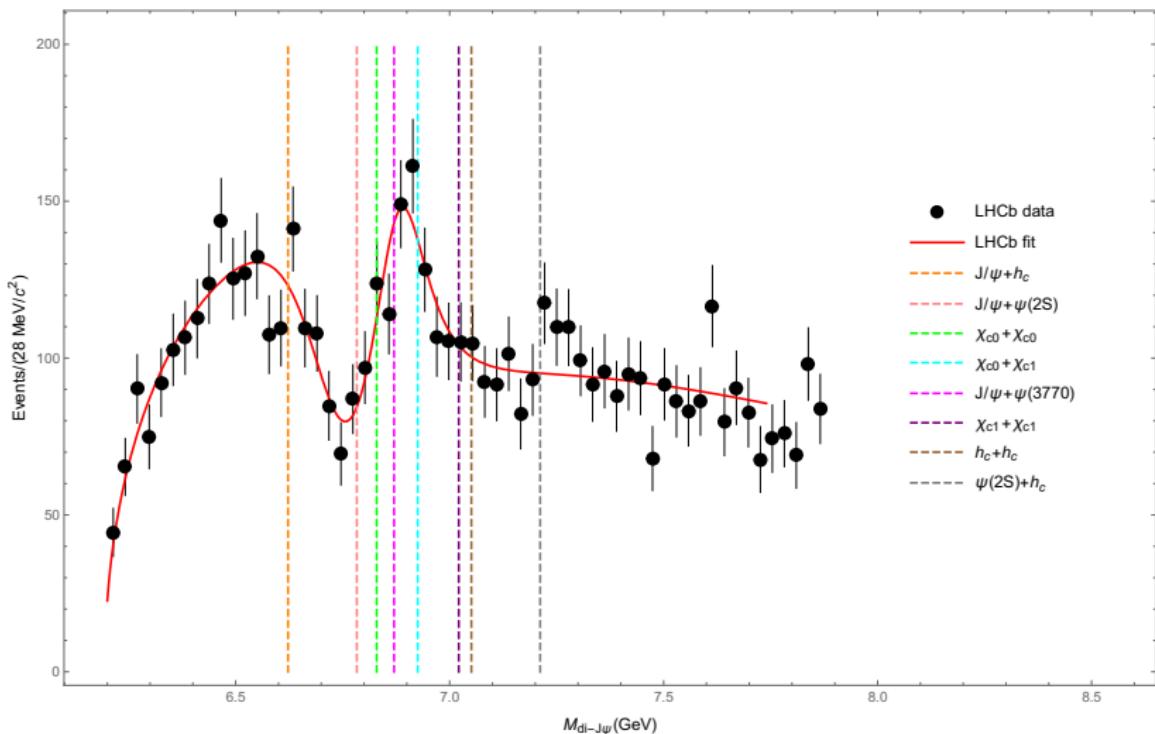


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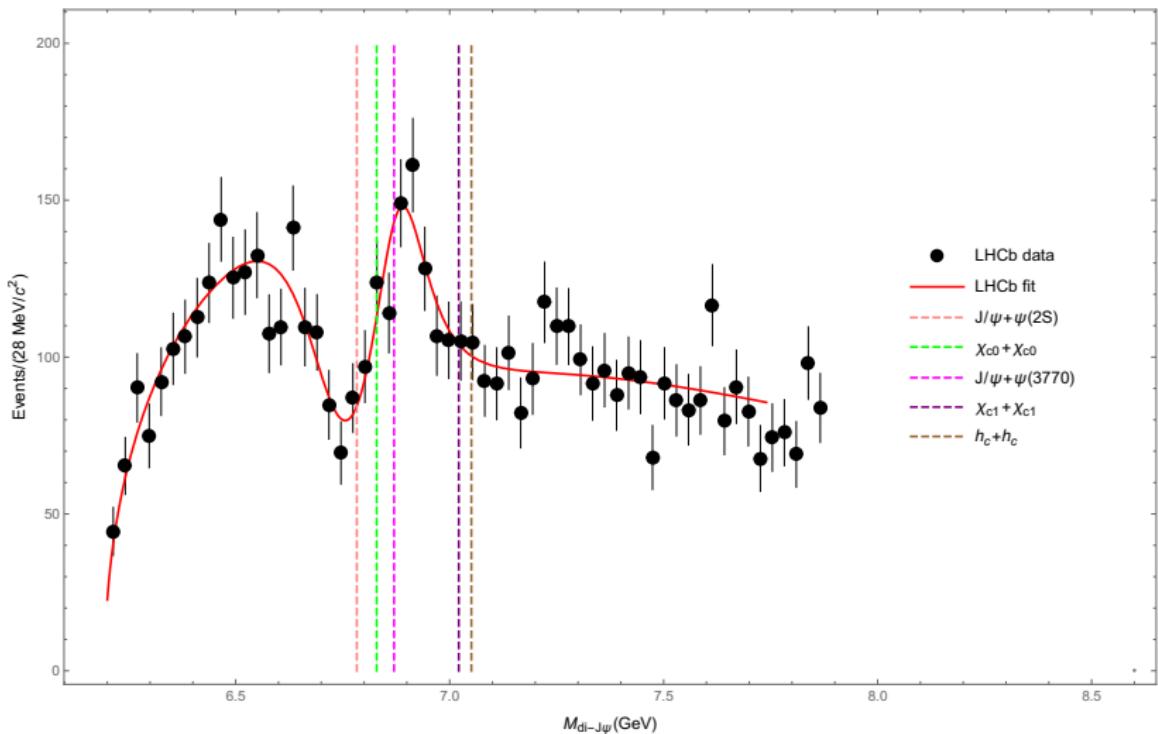
Double- J/ψ production: Theoretical data analysis



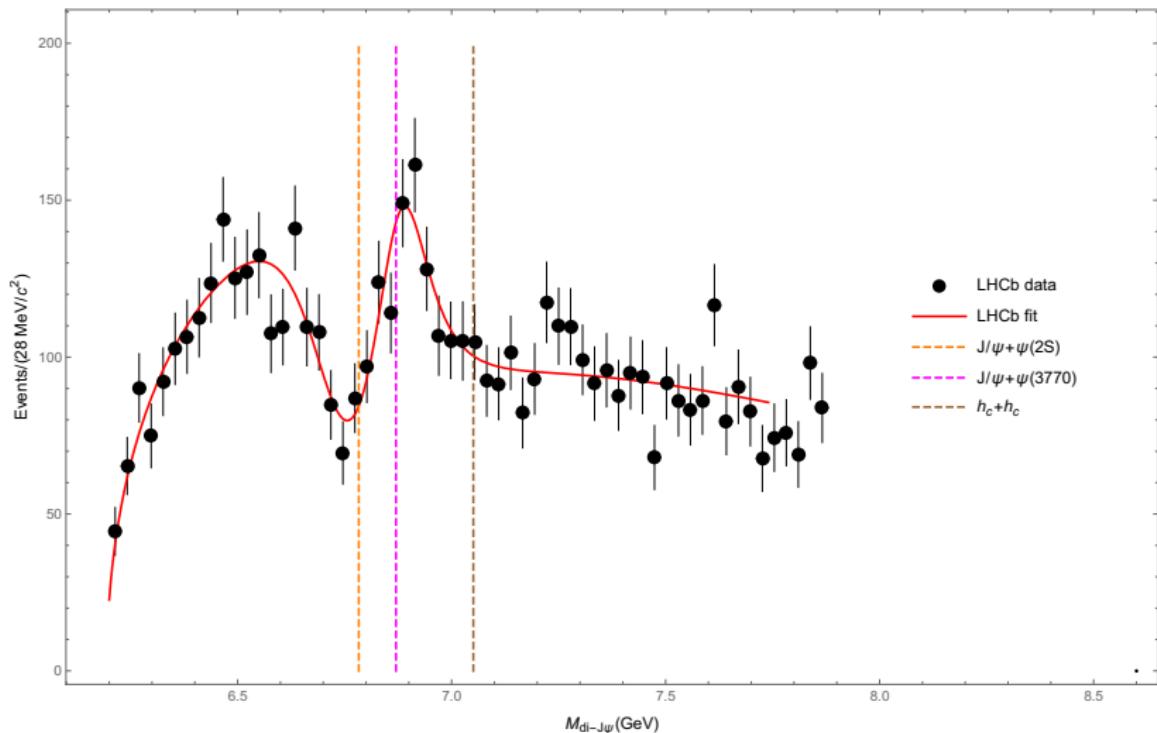
All channels



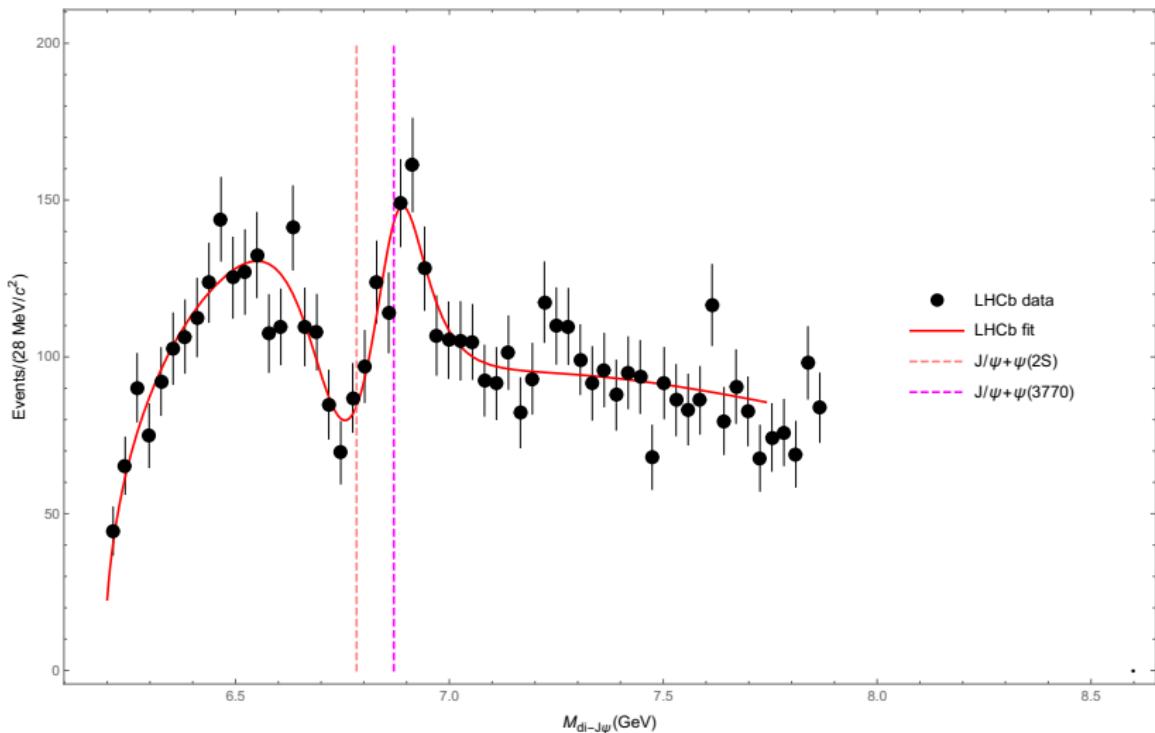
Only S -wave channels (no $J/\psi h_c$, $\psi(2S)h_c$, $\chi_{c0}\chi_{c1}$)



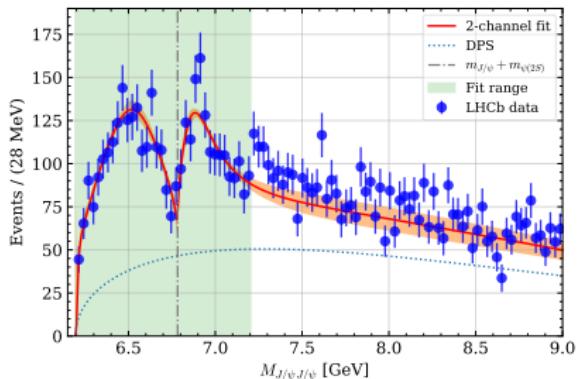
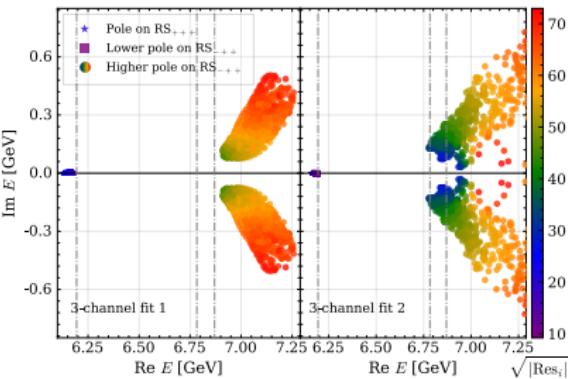
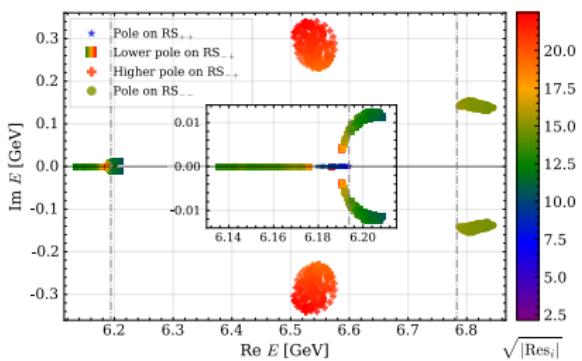
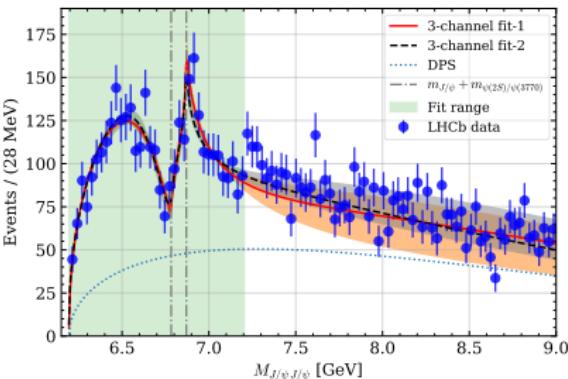
No heavy exchanges (no $\chi_{c0}\chi_{c0}$, $\chi_{c1}\chi_{c1}$)



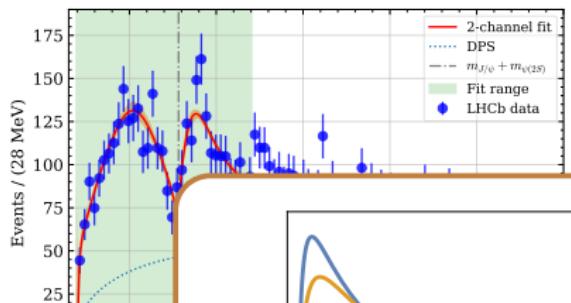
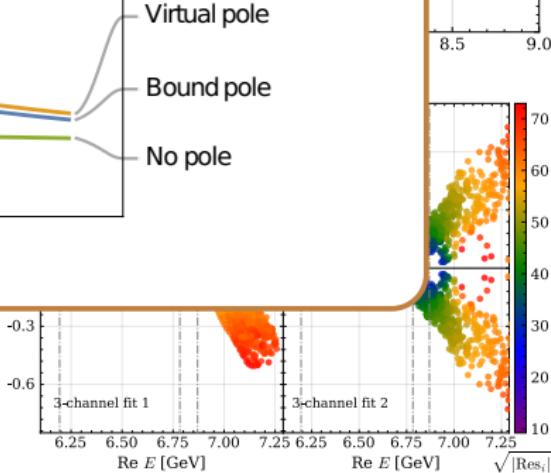
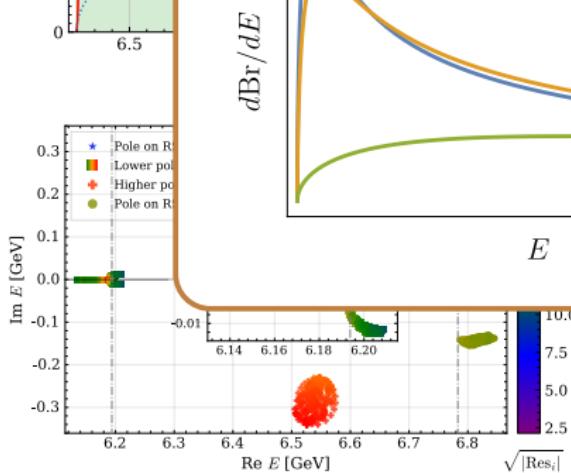
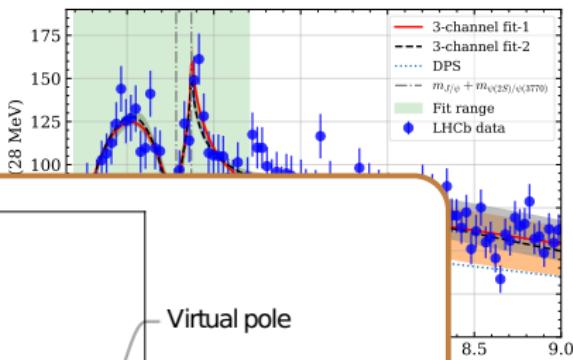
Only HQSS-allowed channels (no $h_c h_c$)



Fits & poles

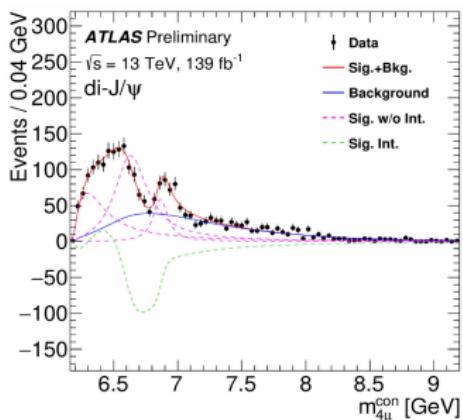
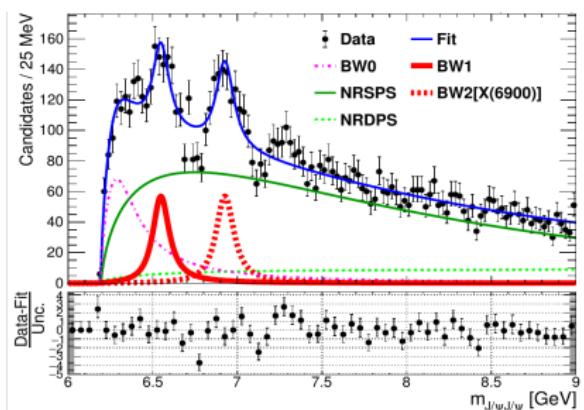
 $J/\psi J/\psi$ & $J/\psi\psi(2S)$ (2+5 params) $J/\psi J/\psi$, $J/\psi\psi(2S)$ & $J/\psi\psi(3770)$ (2+6 params)

Fits & poles

 $J/\psi J/\psi$ & $J/\psi\psi(2S)$ (2+5 params) $J/\psi J/\psi$, $J/\psi\psi(2S)$ & $J/\psi\psi(3770)$ (2+6 params)

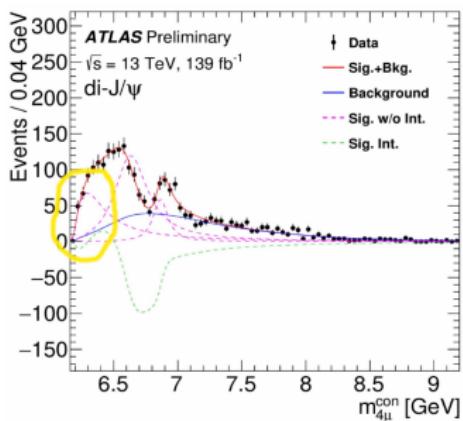
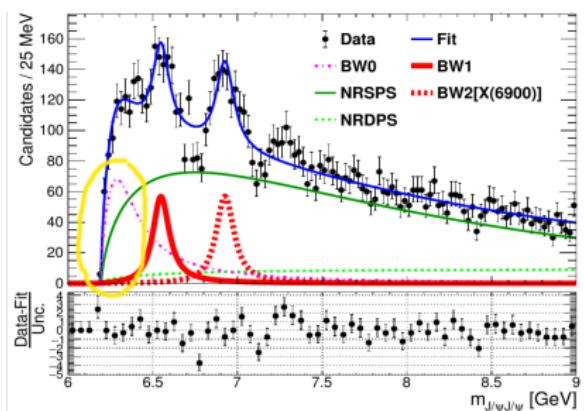
Data from CMS and ATLAS

- Analyse CMS and ATLAS data in the **double- J/ψ channel**



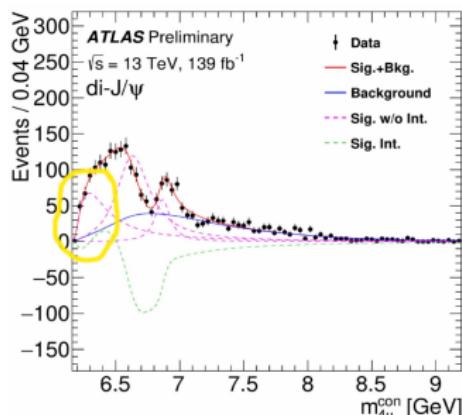
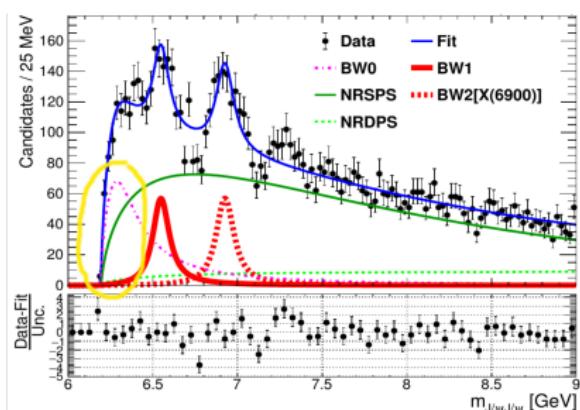
Data from CMS and ATLAS

- Analyse CMS and ATLAS data in the **double- J/ψ channel**



Data from CMS and ATLAS

- Analyse CMS and ATLAS data in the double- J/ψ channel



- Analyse data in the complimentary $\psi(2S)J/\psi$ channel
- Combined analysis of all data sets (?)

Multipole expansion in QED

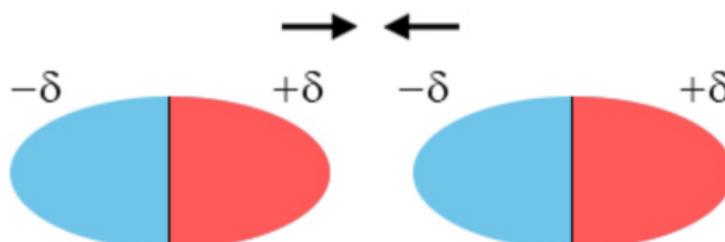
Field produced by charge distribution at $R \gg r$

$$\Phi(R) = \frac{Q}{R} + \text{Dipole} + \text{Quadrupole} + \dots$$



If $Q = 0$, **dipole** term provides leading contribution

Attraction

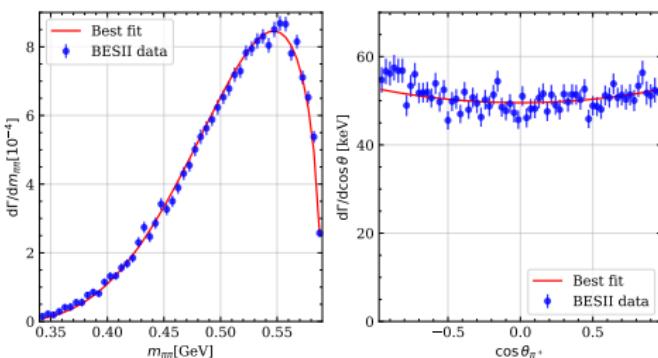


Polarisability measures how easily neutral object is **distorted** by electric field

Multipole expansion for charmonia

Multipole expansion for soft-gluon exchanges ($r_{\bar{Q}Q} \ll \Lambda_{\text{QCD}}^{-1}$): $H_{\text{int}} \approx -\frac{1}{2}\zeta_a \mathbf{r} \cdot \mathbf{E}^a$
 (Gottfried'1977, Voloshin'1978, Peskin'1979, ... Voloshin&Sibirtsev'2005)

$$\beta_{\psi\psi'} = \frac{1}{9} \langle \psi' | \mathbf{r} \frac{1}{\hat{H}_O - M} \mathbf{r} | \psi \rangle \implies \left\{ \begin{array}{l} \mathcal{M}(\psi(2S) \rightarrow J/\psi \pi\pi) = \beta_{\psi(2S)J/\psi} \langle \pi\pi | \mathbf{E}^a \cdot \mathbf{E}^a | 0 \rangle \\ V_{J/\psi J/\psi}(r) \propto \beta_{J/\psi J/\psi}^2 \propto \xi^2 \end{array} \right. \quad \left(\xi = \frac{\beta_{J/\psi J/\psi}}{\beta_{\psi(2S)J/\psi}} \right)$$



$$|\beta_{\psi(2S)J/\psi}| \approx 1.81 \text{ GeV}^{-3}$$

Chromopolarisability of J/ψ

- Simple estimate

$$\xi \sim I_{J/\psi J/\psi} / I_{\psi(2S)J/\psi} \sim 10 \quad I_{\psi' \psi} = \langle \psi' | e^{-i \Delta \mathbf{q}_c \cdot \mathbf{r} / 2} | \psi \rangle$$

(Dong et al.'2021)

- More advanced estimate

$$a_0^{\text{th}}(J/\psi \pi \rightarrow J/\psi \pi) \approx 0.0036 \xi \text{ fm}$$

(Dong et al.'2021)

$$|a_0^{\text{lat}}(J/\psi \pi \rightarrow J/\psi \pi)| \sim 0.01 \text{ fm} \implies \xi \simeq 3$$

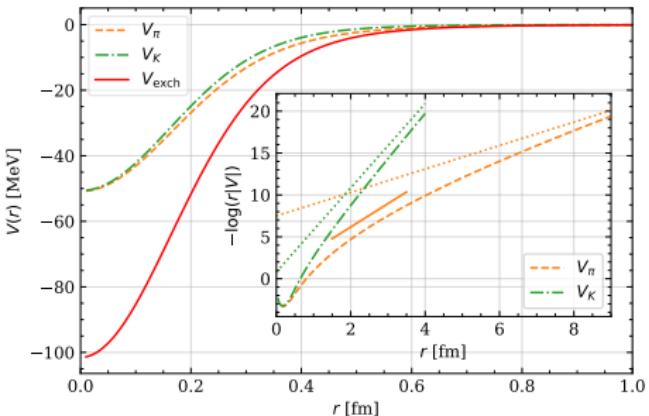
(Yokokawa et al.'2006, Liu et al.'2008)

- Quark model calculation (approximation: J/ψ = Coulombic system)

$$\begin{aligned} \beta_{J/\psi J/\psi} &= \frac{1}{9} \langle J/\psi | \mathbf{r} \frac{1}{\hat{H}_O - M_{J/\psi}} \mathbf{r} | J/\psi \rangle = \frac{1}{9} \int \frac{d^3 p}{(2\pi)^3} \int_0^\infty r dr \frac{u_{J/\psi}(r) v_{p,l=1}(r)}{2m_c + \frac{p^2}{m_c} - M_{J/\psi}} \\ &= \frac{0.93}{\alpha_s^4 m_c^3} = 19^{+15}_{-14} \text{ GeV}^{-3} \implies 3 \lesssim \xi \lesssim 19 \end{aligned}$$

(Brambilla et al.'2016, Dong et al.'2023)

Interaction potential in di- J/ψ system



$$V_{\text{tot}}(r, \Lambda) = V_\pi(r, \Lambda) + V_K(r, \Lambda) = V_{\text{CT}}(r, \Lambda) + V_{\text{exch}}(r, \Lambda)$$

$$V_{\text{exch}}(r, \Lambda) = -\frac{1}{4\pi M_{J/\psi}^2} \int \frac{d^3 q}{(2\pi)^3} e^{i\mathbf{q}\cdot\mathbf{r}} \int_{4m_\pi^2}^\infty d\mu^2 \frac{\text{Im}\mathcal{M}_{J/\psi J/\psi}(\mu^2)}{\mu^2 + q^2} F\left(\frac{q^2 + \mu^2}{\Lambda^2}\right)$$

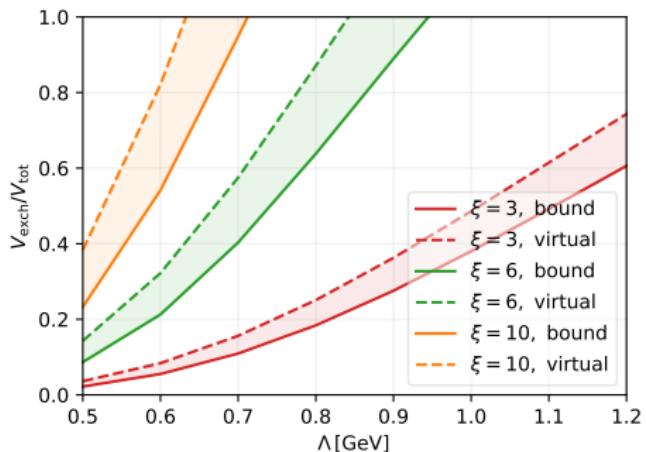
Exchange of charmonia is suppressed as $\Lambda_{\text{QCD}}^2/m_c^2$

$\Rightarrow V_{CT}$ mainly comes from pion/kaon exchanges

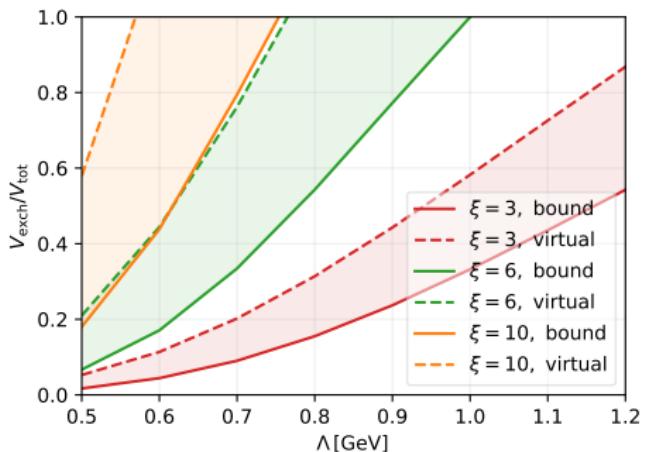
\implies For natural $\Lambda \sim 1$ GeV $V_{\text{CT}} \sim V_{\text{exch}}$ $\implies V_{\text{exch}}/V_{\text{tot}} \gtrsim \frac{1}{2}$

$X(6200)$ as double- J/ψ molecule

$E_B = 1$ MeV

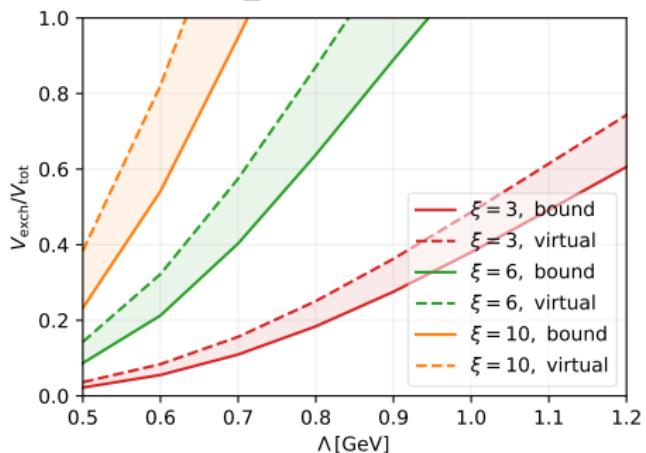


$E_B = 5$ MeV

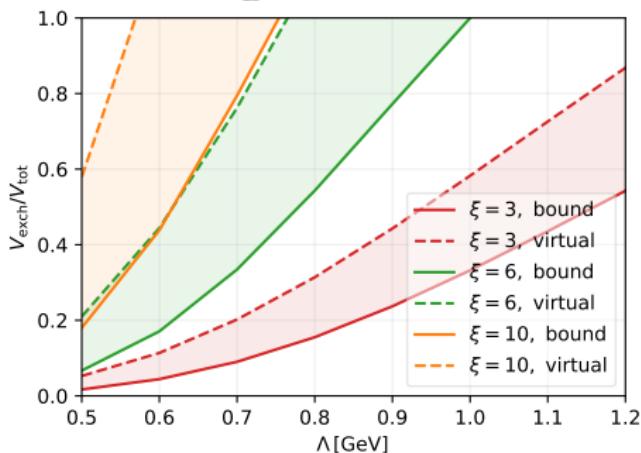


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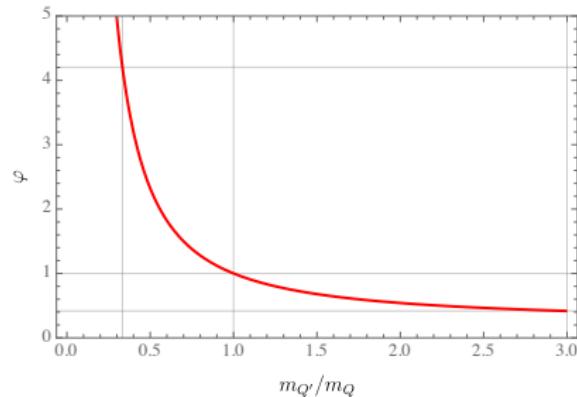
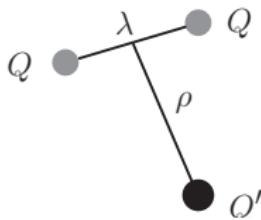
$E_B = 5$ MeV



Conclusion: Existence of a molecular pole near the double- J/ψ threshold is consistent with our knowledge on hadron-hadron interactions

Chromopolarisability of QQQ' baryon

$$3 \otimes 3 \otimes 3 = 1 \oplus 8 \oplus 8 \oplus 10$$



- $Q' = Q$

$$\beta_\Omega = \beta_\Omega^{(\lambda)} + \beta_\Omega^{(\rho)} = \frac{C_\Omega}{m_Q^3 \alpha_s^4} \quad C_\Omega = C_\Omega^{(\rho)} + C_\Omega^{(\lambda)}$$

$$C_\Omega^{(\rho)} \approx 1.1 \quad C_\Omega^{(\lambda)} \approx 1.3$$

- $Q' \neq Q$

$$\beta_\Omega = \beta_\Omega^{(\lambda)} + \beta_\Omega^{(\rho)} = \frac{C_\Omega}{m_Q^3 \alpha_s^4} \varphi(m_{Q'}/m_Q)$$

Di-heavy-baryon molecules?

- Quarkonium $\psi = \bar{Q}Q$

$$\beta_\psi = \frac{C_\psi}{\alpha_s^4 m_Q^3} \quad C_\psi \approx 0.93$$

(Brambilla et al.'2016, Dong et al'2023)

- Baryon $\Omega = QQQ$

$$\beta_\Omega = \frac{C_\Omega}{\alpha_s^4 m_Q^3} \quad C_\Omega \approx 2.4 \approx 2.6 C_\psi$$

(Dong et al'2023)

- Uncertainty from hadron being **not Coulombic**

$$\frac{\delta\beta}{\beta} \sim \frac{\langle \sigma r \rangle}{\langle \alpha_s/r \rangle} \sim \frac{\Lambda_{\text{QCD}}^2}{\alpha_s^3 m_Q^2} \sim \begin{cases} 100\% & Q = c \\ 10\% & Q = b \end{cases}$$

Di-heavy-baryon molecules?

- Quarkonium $\psi = \bar{Q}Q$

$$\beta_{\psi} = \frac{C_{\psi}}{C_{J/\psi}} \quad C_{\psi} \approx 0.93$$

Conclusion: If $X(6200)$ is a di- J/ψ molecule, then
di- Ω_{ccc} and di- Ω_{bbb} molecules are also likely to exist

Supported by recent lattice results:

- $E_B(\text{di}-\Omega_{ccc}) \simeq 5..6 \text{ MeV}$ (Lyu et al.'2021)
- $E_B(\text{di}-\Omega_{bbb}) \simeq 80..100 \text{ MeV}$ (Mathur et al.'2022)
- Uncertainty from hadron being NOT Coulombic

$$\frac{\delta\beta}{\beta} \sim \frac{\langle \sigma r \rangle}{\langle \alpha_s/r \rangle} \sim \frac{\Lambda_{\text{QCD}}^2}{\alpha_s^3 m_Q^2} \sim \begin{cases} 100\% & Q = c \\ 10\% & Q = b \end{cases}$$

Conclusions

- Discovery of $X(3872)$ started new era in hadronic physics of heavy quarks
- LHC studies of double- J/ψ production opened new chapter in this book
- Data collected are analysed
 - using (minimal but realistic) coupled-channel scheme
 - preserving unitarity of multichannel scattering amplitude
 - respecting (approximate but accurate) heavy quark spin symmetry
- Existence of a pole at $J/\psi J/\psi$ threshold is predicted from data analysis
- Conjecture of molecular nature of $X(6200)$ is plausible and consistent with our knowledge of hadron-hadron interactions
- Near-threshold state in double- J/ψ channel may imply existence of double-heavy-baryon molecules