

Early Kinetic Decoupling Effect on Forbidden Dark Matter

刘学文 烟台大学

arXiv: [2301.12199](https://arxiv.org/abs/2301.12199)

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Outline

- **E**arly **K**inetic **D**ecoupling (EKD): High Accuracy of Calculations
- Dealing with EKD: **C**oupled **B**oltzmann **E**quations (CBE)
- EKD on **F**orbidden **D**ark **M**atter: One order of magnitude larger
- Summary

1 Early Kinetic Decoupling

Boltzmann Equation

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Boltzmann Equation

- Evolution of Phase space density f_χ

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Boltzmann Equation

$$E\left(\frac{\partial}{\partial t} - H\vec{p} \cdot \frac{\partial}{\partial \vec{p}}\right) f_{\chi} = C[f_{\chi}]$$

- Evolution of Phase space density f_{χ}
- Liouville Operator: Change in spacetime of the DM f_{χ}
- $C[f_{\chi}]$: Collision term, interactions between DM and other particles (including itself) that may alter f_{χ}

Boltzmann Equation: Traditional or “Standard” Treatment

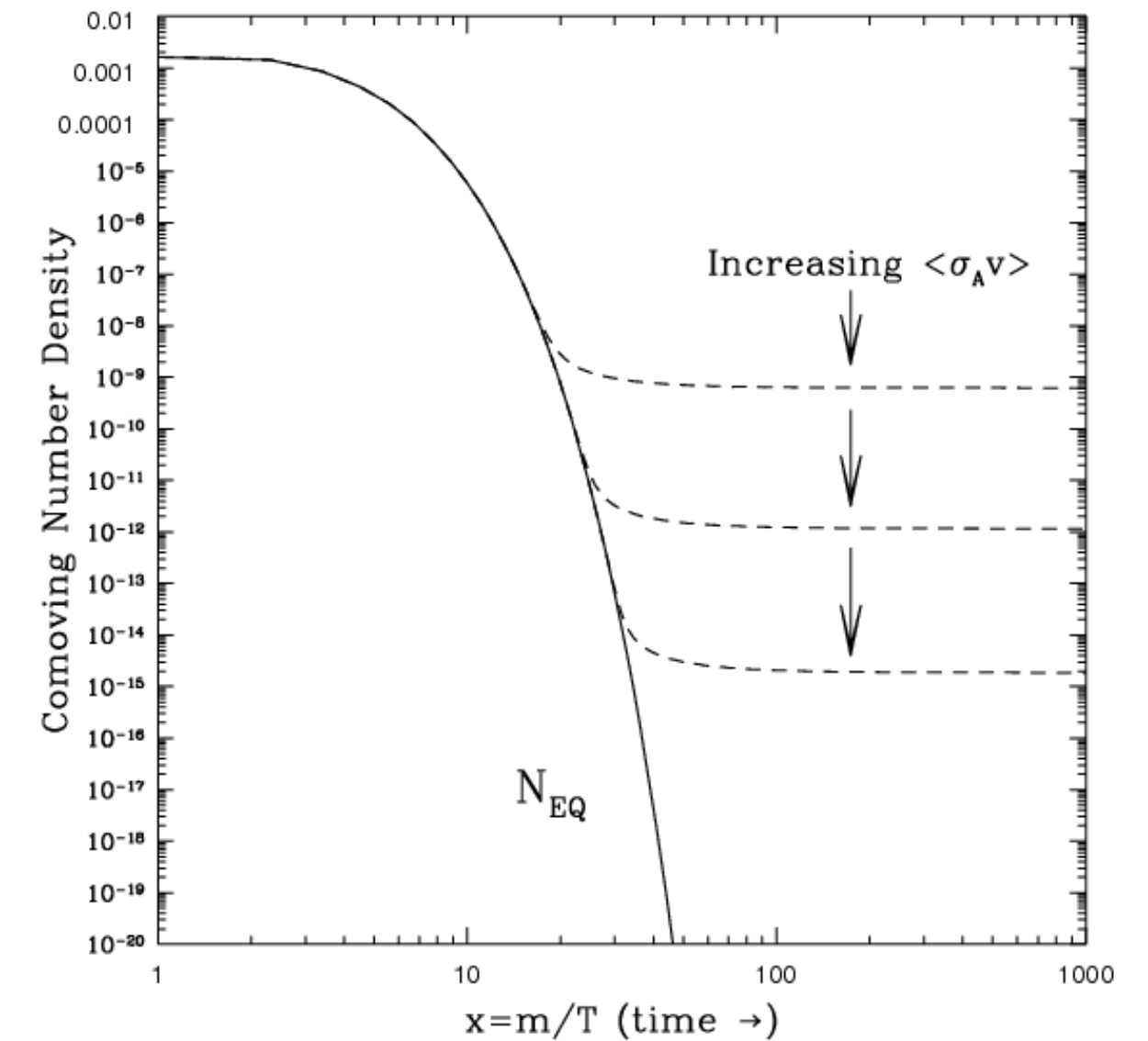
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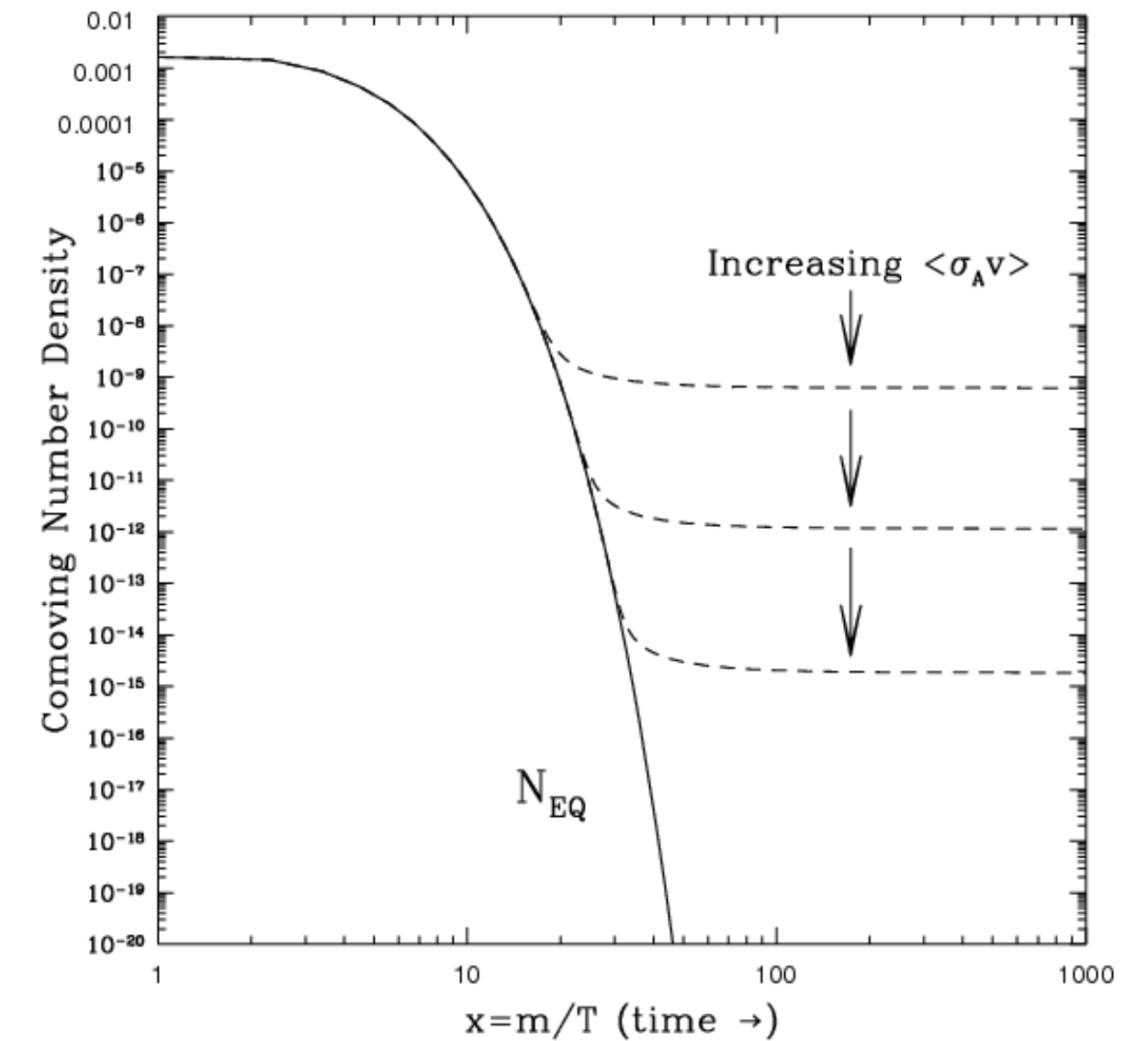


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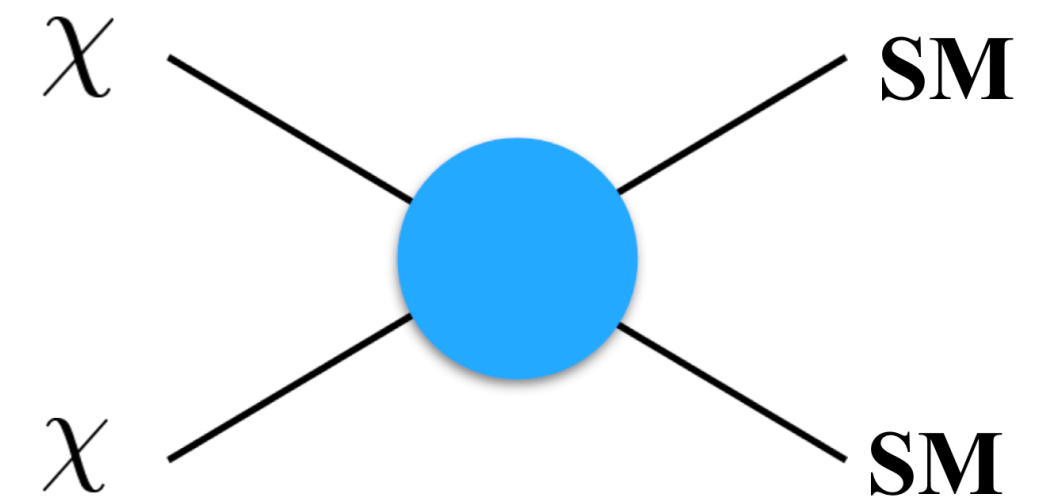
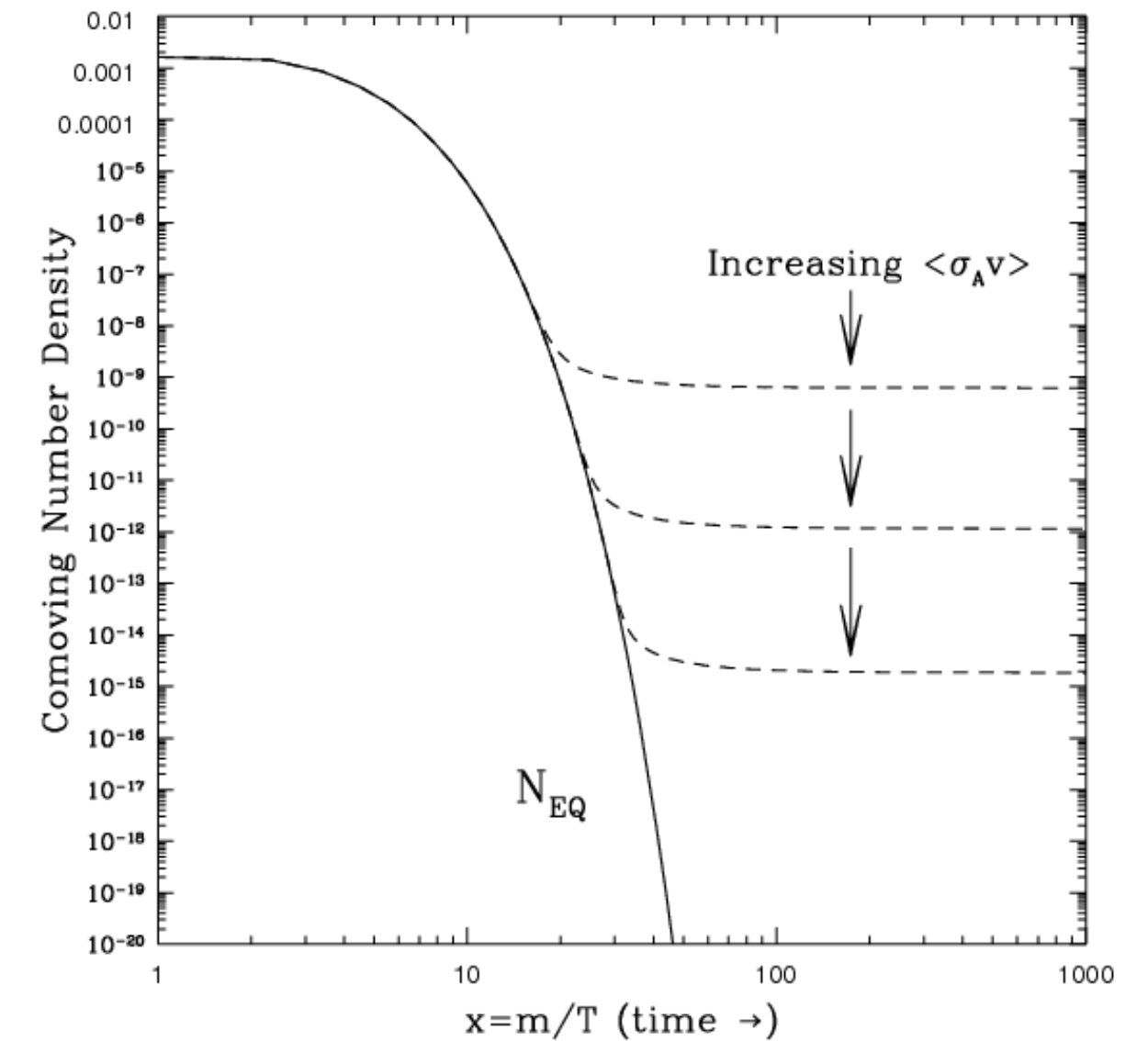
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- Using $C[f_\chi] = C_{\text{ann.}}$, from annihilation process



Boltzmann Equation: An assumption

- **Kinetic equilibrium:** DM and SM share the same temperature, as the elastic scattering is frequency .
- Kinetic equilibrium determines the form of the phase space density

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- Kinetic equilibrium determines the form of the phase space density

$$f_{\chi} = \frac{n_{\chi}}{n_{\chi}^{\text{eq}}} f_{\chi}^{\text{eq}} (T_{\text{SM}})$$

Boltzmann Equation: Problem

- Kinetic equilibrium is **NOT** satisfied at some scenarios
- DM kinetic decouples from the thermal bath in an **EARLY** stage
- i.e. **Early Kinetic Decoupling**
- DM has its own temperature than SM, different phase space density

Boltzmann Equation: Problem

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$$f_{\chi}(T_{\text{SM}}) \rightarrow f_{\chi}(T_{\text{DM}})$$

2 Coupled Boltzmann Equations (CBE)

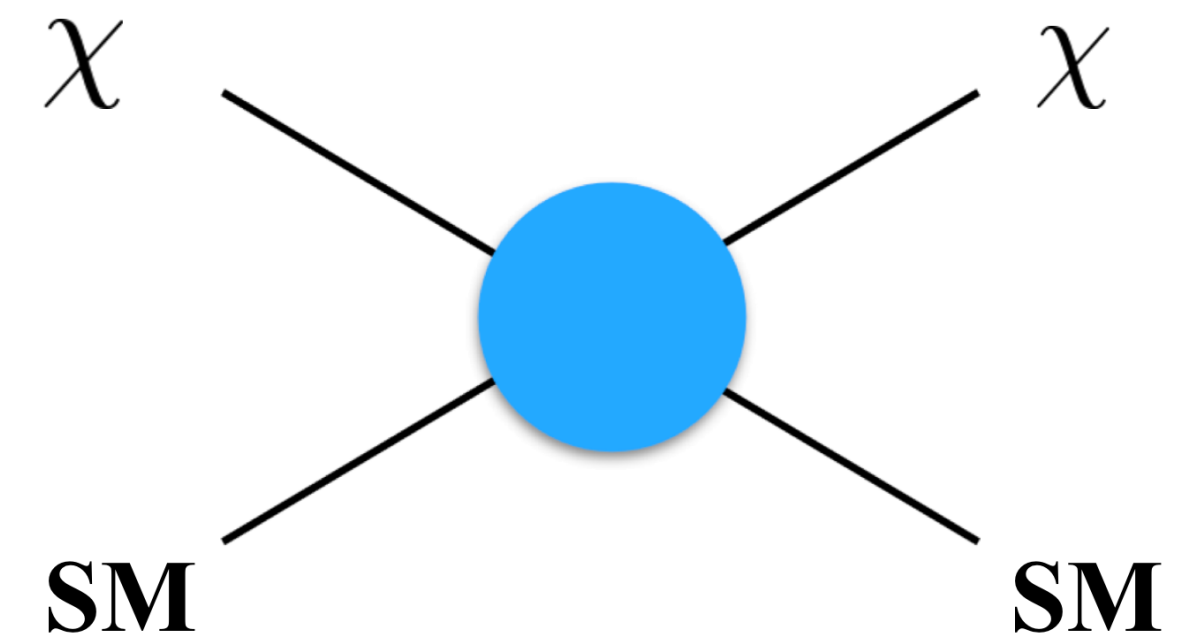
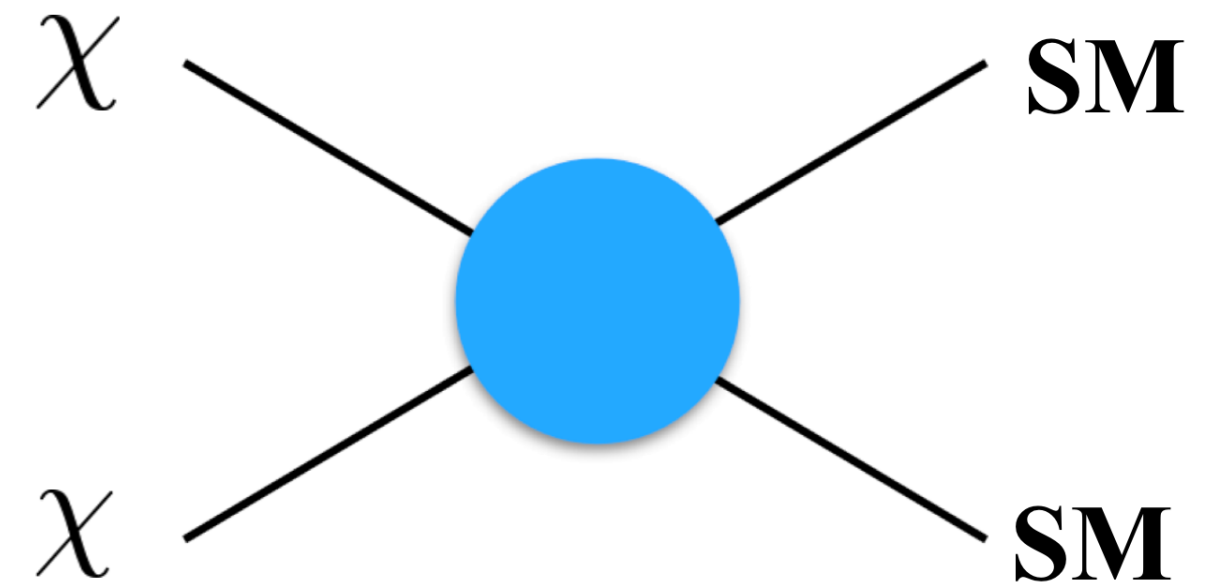
Dealing with Early Kinetic Decoupling

- Including a NEW collision term from elastic scattering

$$C[f_\chi] = C_{\text{ann.}}[f_\chi] + C_{\text{el.}}[f_\chi]$$

- Define the temperature for DM:

$$T_\chi = \frac{g_\chi}{3n_\chi} \int \frac{d^3p}{(2\pi)^3} \frac{\vec{p}^2}{E} f_\chi(\vec{p})$$



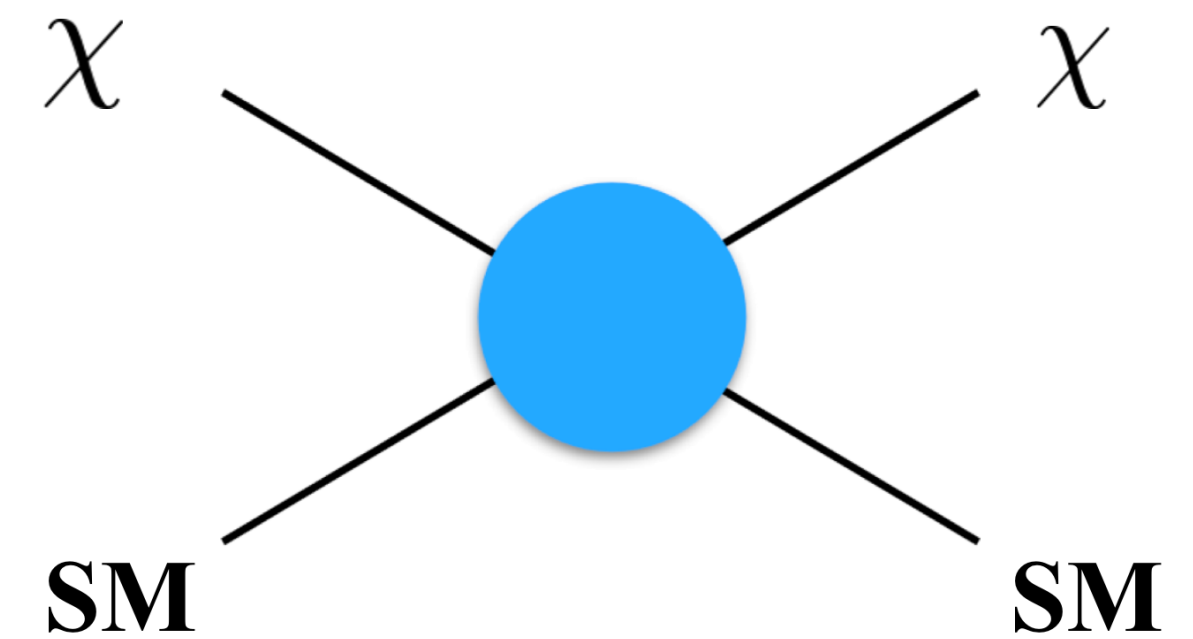
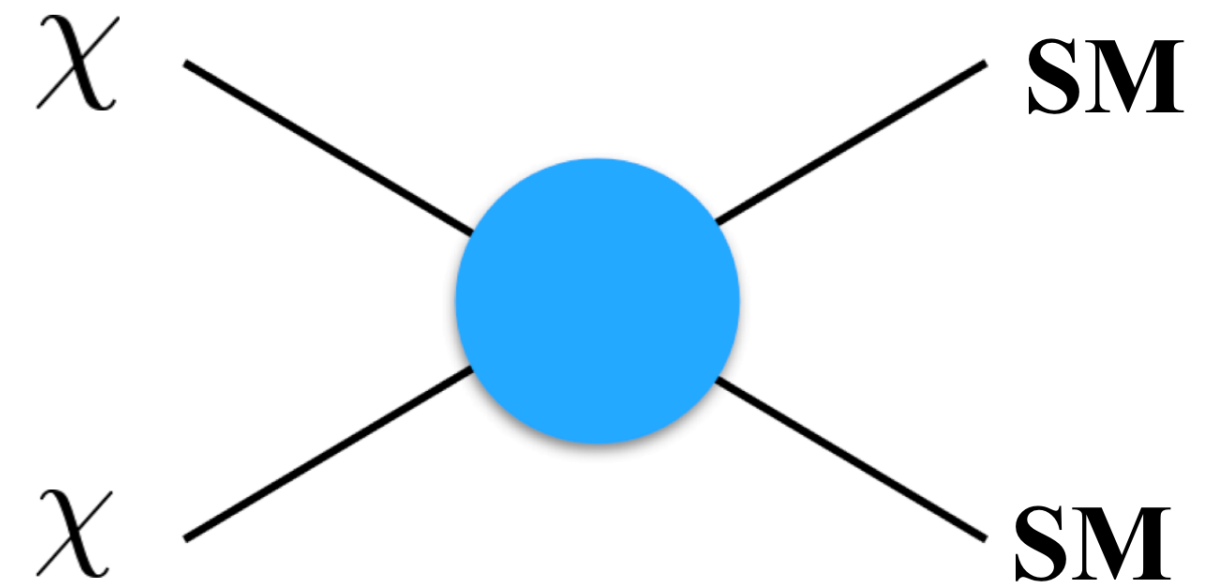
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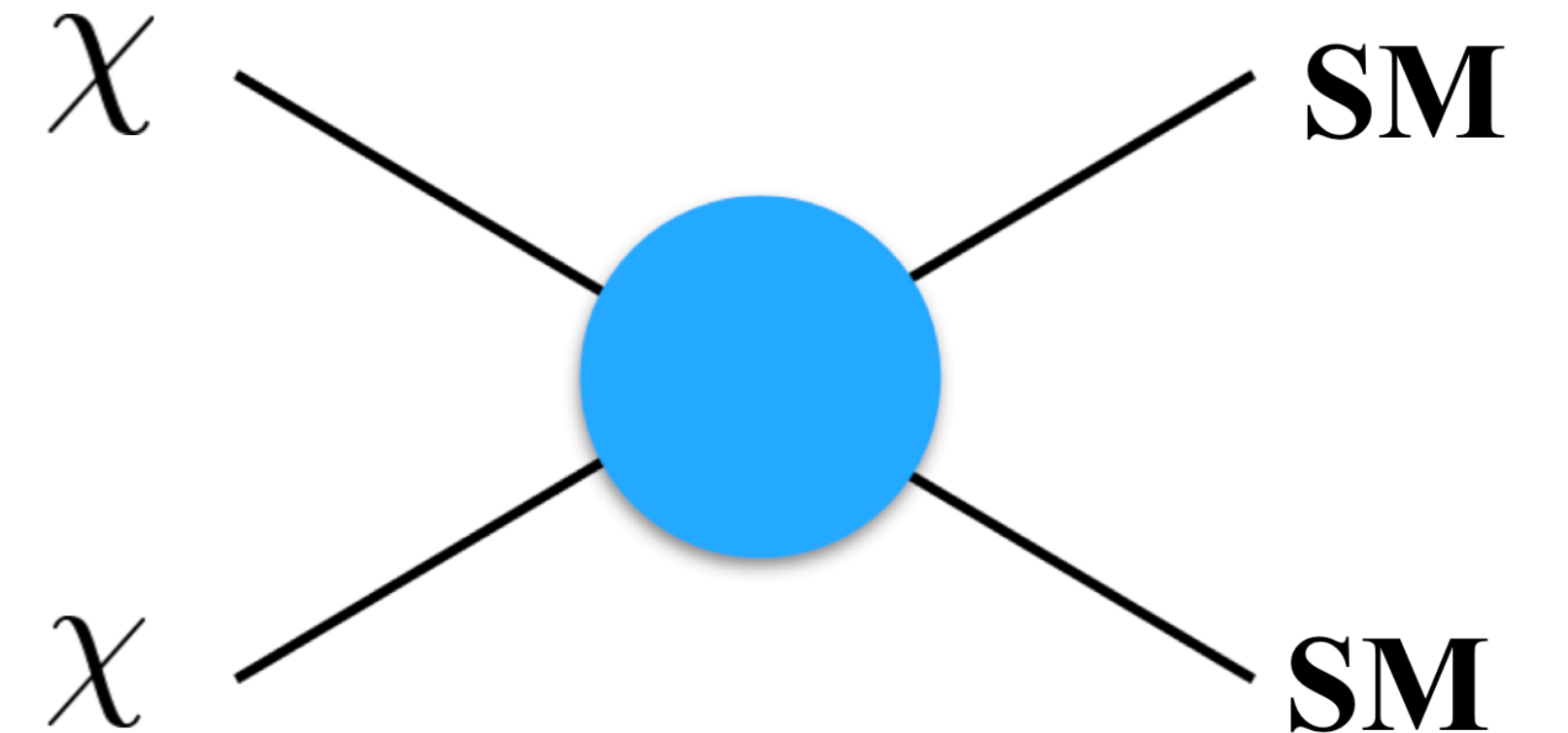
- Define the temperature for DM:

$$T_\chi = \frac{g_\chi}{3n_\chi} \int \frac{d^3p}{(2\pi)^3} \frac{\vec{p}^2}{E} f_\chi(\vec{p}) \equiv \frac{s^{2/3}}{m_\chi} y$$



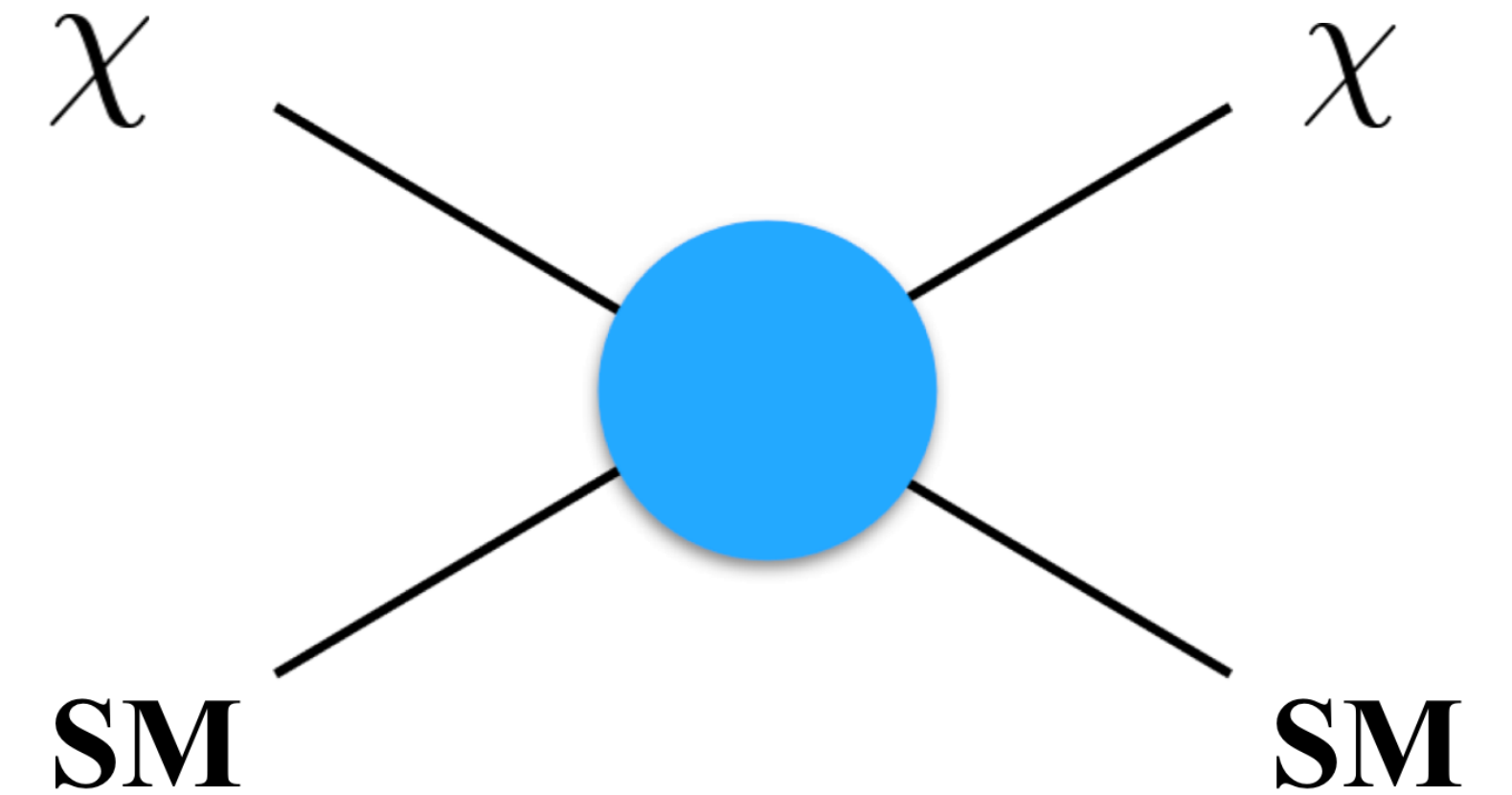
Collision term from annihilation

$$\begin{aligned}
 C_{\text{ann.}} = & \frac{1}{2g_\chi} \sum \int \frac{d^3 p'}{(2\pi)^3 2E_{p'}} \int \frac{d^3 k}{(2\pi)^3 2E_k} \int \frac{d^3 k'}{(2\pi)^3 2E_{k'}} \\
 & \times (2\pi)^4 \delta^4(p + p' - k - k') \\
 & \times [- |\mathcal{M}_{\chi\chi \rightarrow \mathcal{B}\mathcal{B}'}|^2 f_\chi(\vec{p}) f_\chi(\vec{p}') (1 \pm f_{\mathcal{B}}^{\text{eq}}(\vec{k})) (1 \pm f_{\mathcal{B}'}^{\text{eq}}(\vec{k}')) \\
 & + |\mathcal{M}_{\mathcal{B}\mathcal{B}' \rightarrow \chi\chi}|^2 f_{\mathcal{B}}^{\text{eq}}(\vec{k}) f_{\mathcal{B}'}^{\text{eq}}(\vec{k}') (1 \pm f_\chi(\vec{p})) (1 \pm f_\chi(\vec{p}'))]
 \end{aligned}$$



NEW collision term from elastic scattering

$$\begin{aligned}
 C_{\text{el.}} = & \frac{1}{2g_\chi} \sum \int \frac{d^3 p'}{(2\pi)^3 2E_{p'}} \int \frac{d^3 k}{(2\pi)^3 2E_k} \int \frac{d^3 k'}{(2\pi)^3 2E_{k'}} \\
 & \times (2\pi)^4 \delta^4(p + p' - k - k') \\
 & \times \left[- |\mathcal{M}_{\chi\mathcal{B} \rightarrow \chi\mathcal{B}}|^2 f_\chi(\vec{p}) f_{\mathcal{B}}^{\text{eq}}(\vec{k}) (1 \pm f_\chi(\vec{p}')) (1 \pm f_{\mathcal{B}}^{\text{eq}}(\vec{k}')) \right. \\
 & \left. + |\mathcal{M}_{\chi\mathcal{B} \rightarrow \chi\mathcal{B}}|^2 f_\chi(\vec{p}') f_{\mathcal{B}}^{\text{eq}}(\vec{k}') (1 \pm f_\chi(\vec{p})) (1 \pm f_{\mathcal{B}}^{\text{eq}}(\vec{k})) \right]
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Early Kinetic Decoupling and Coupled Boltzmann Equation

- Integrating the Boltzmann Equation [\[2103.01944\]](#)

Early Kinetic Decoupling and Coupled Boltzmann Equation

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$$g_\chi \int \frac{d^3p}{(2\pi)^3} \frac{1}{E} \quad \text{AND} \quad g_\chi \int \frac{d^3p}{(2\pi)^3} \frac{1}{E} \frac{\vec{p}^2}{E} \left(E \left(\frac{\partial}{\partial t} - H \vec{p} \cdot \frac{\partial}{\partial \vec{p}} \right) f_\chi = C[f_\chi] \right)$$

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$$(1) \quad \frac{Y'}{Y} = \frac{sY}{x\tilde{H}} \left[\frac{Y_{\text{eq}}^2}{Y^2} \langle \sigma v \rangle_T - \langle \sigma v \rangle_{T_\chi} \right]$$

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- Evolution of DM number density ($Y = n_\chi/s$) and temperature T_χ (y)

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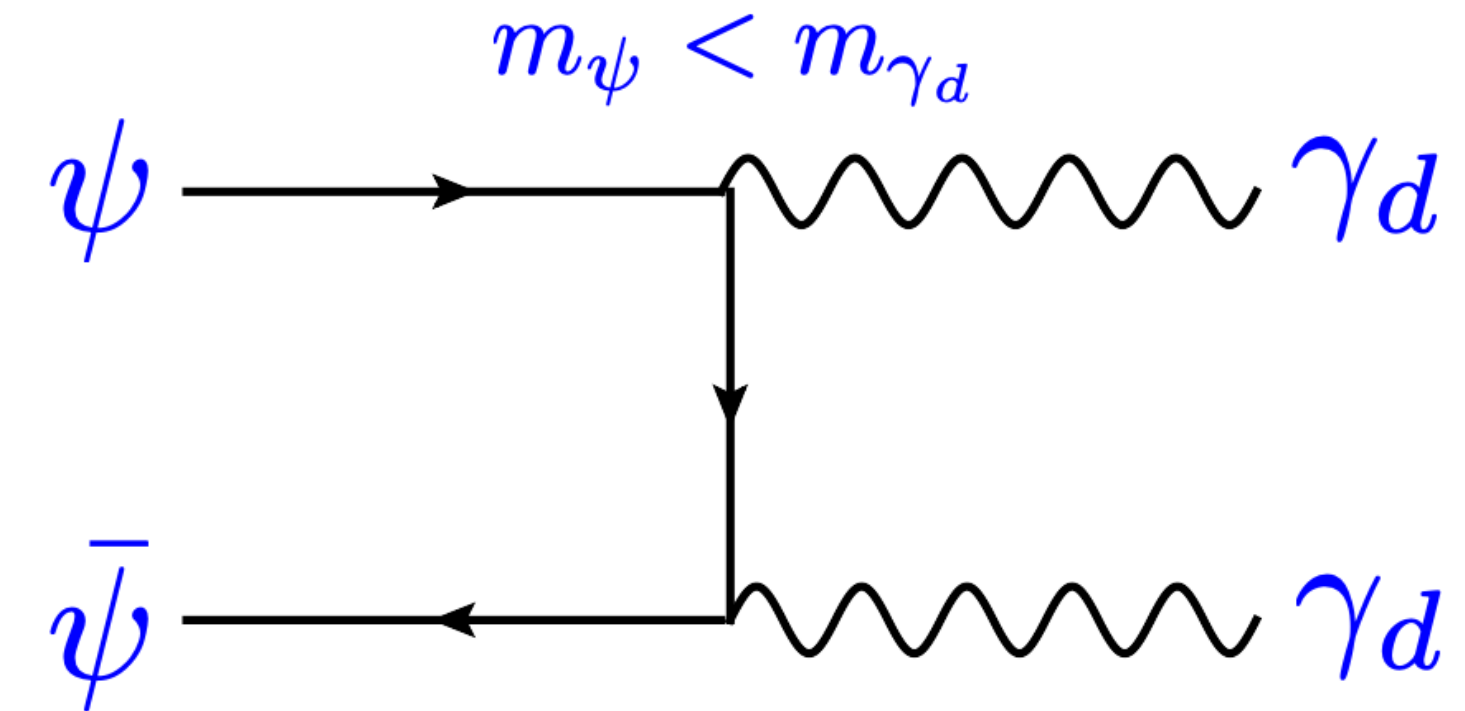
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3 EKD on Forbidden Dark Matter

Forbidden Dark Matter

K. Griest, D. Seckel,
Phys.Rev.D 43 (1991) 3191-3203

- Forbidden: $m_\chi < m_{\text{SM}}$
- Annihilations are forbidden at zero temperature
- Nonzero at finite temperature.



D'Agnolo, Ruderman,
[1505.07107]

Our Forbidden Model

$$-\mathcal{L} \supset g_{ij}\phi\bar{l}_i l_j + g_{ij}^A\phi\bar{l}_i\gamma_5 l_j + g_\chi\phi\bar{\chi}\chi + g_\chi^A\phi\bar{\chi}\gamma_5\chi$$

- Particles: Fermionic DM χ , scalar ϕ , $l_i = \mu, \tau$
- Case: $l_i = e$ excluded by BBN and CMB
- Forbidden: $\Delta \equiv (m_\ell - m_\chi)/m_\chi > 0$
- Annihilation cross section exponentially suppressed

$$\langle\sigma v\rangle_{\chi\chi\rightarrow\ell\ell} = \langle\sigma v\rangle_{\ell\ell\rightarrow\chi\chi}\frac{(n_\ell^{\text{eq}})^2}{(n_\chi^{\text{eq}})^2} \simeq \langle\sigma v\rangle_{\ell\ell\rightarrow\chi\chi}e^{-2\Delta x}$$

Scattering term

- Non-relativistic DM, Fokker-Planck operator [1205.1914, 2004.10041]

$$C_{\text{el.}} \simeq \frac{E}{2} \gamma(T) \left[TE \partial_p^2 + \left(2T \frac{E}{p} + p + T \frac{p}{E} \right) \partial_p + 3 \right] f_\chi$$

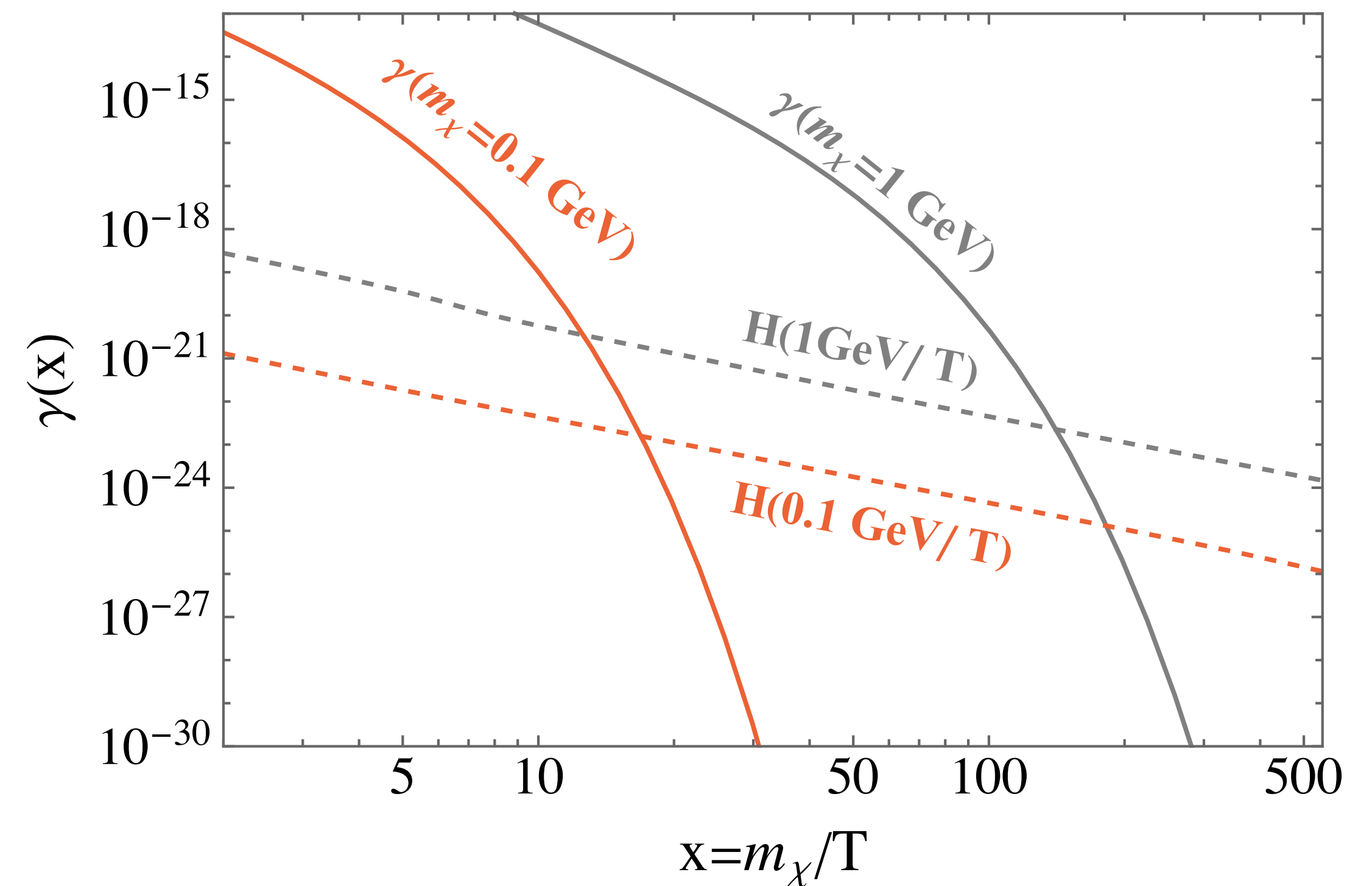
- $\gamma(T)$ is the momentum transfer rate:

$$\gamma(T) = \frac{1}{3g_\chi m_\chi T} \int \frac{d^3k}{(2\pi)^3} f_{\mathcal{B}}^{\text{eq}}(E_k) \left[1 \mp f_{\mathcal{B}}^{\text{eq}}(E_k) \right] \int_{-4k_{\text{cm}}^2}^0 dt (-t) \frac{d\sigma}{dt} v \propto e^{-(\Delta+1)x}$$

$\Delta + 1 > 1$ implies the scattering frequency experiences a strong suppression

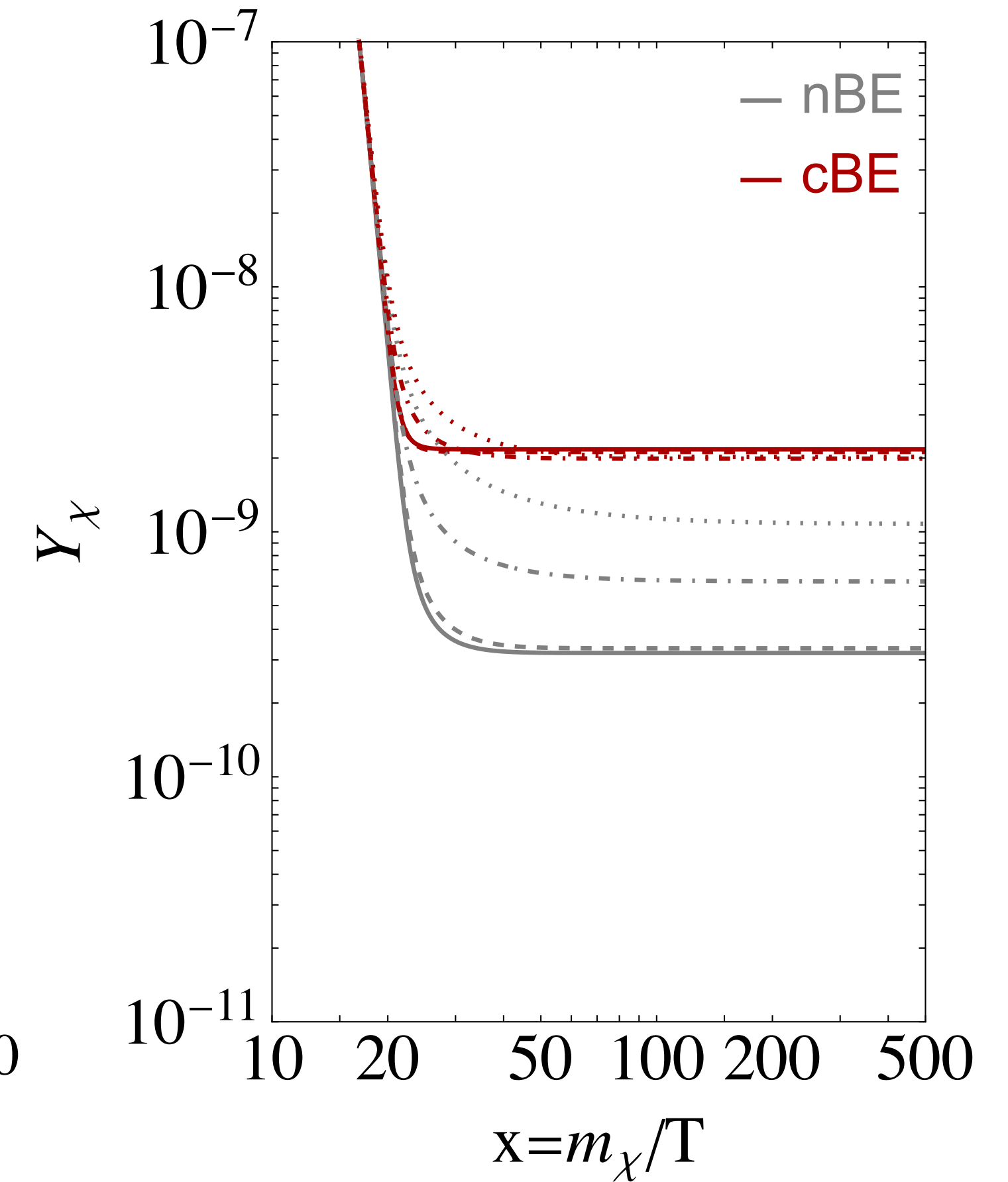
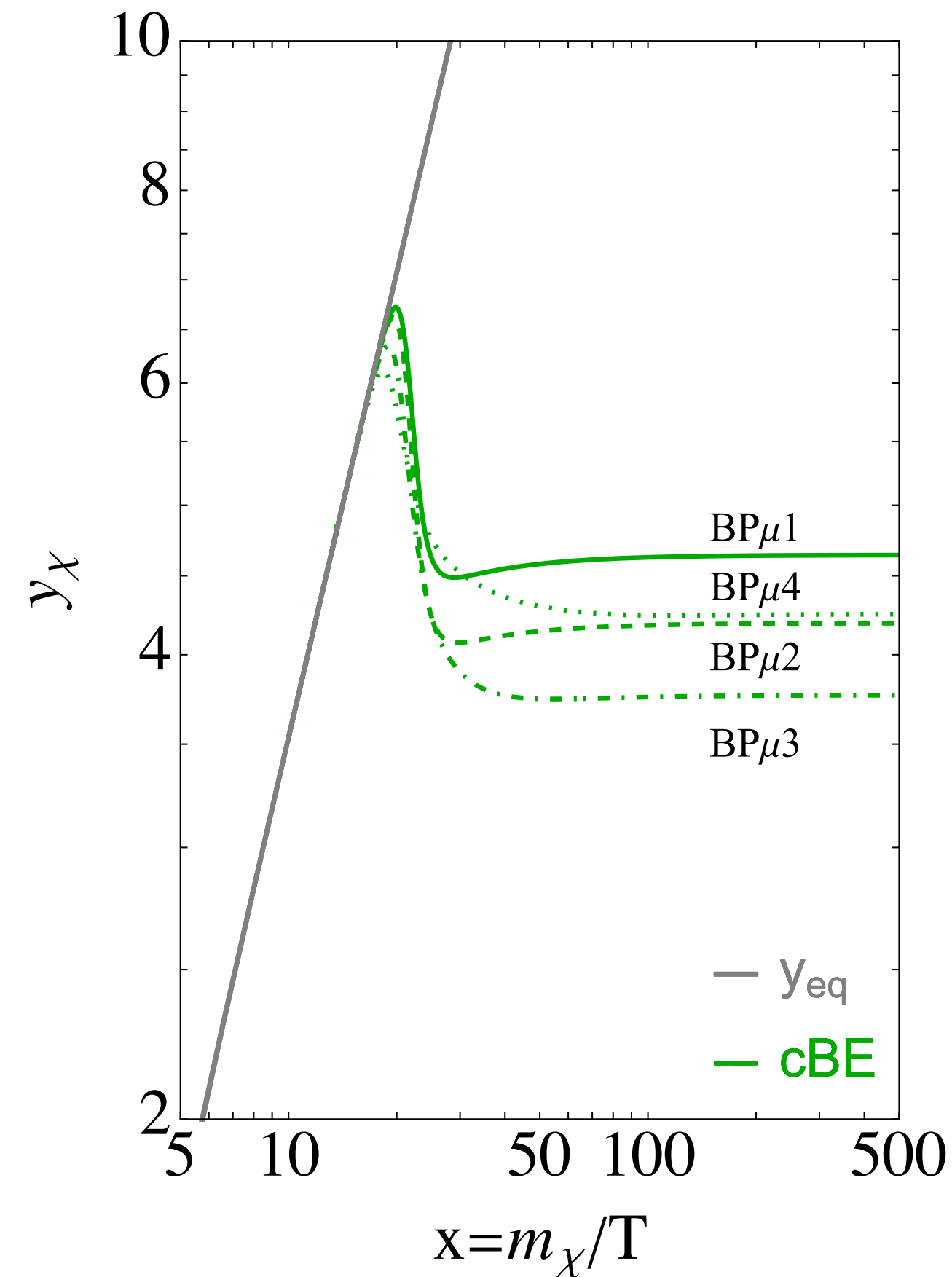
EKD happens in Forbidden Dark Matter

- Hubble rate: $H(T) = \frac{\pi\sqrt{g_{\text{eff}}}}{\sqrt{90}} \frac{T^2}{M_{\text{Pl}}}$
- $\gamma(T)$ becoming comparable with $H(T)$ happens much earlier than traditional WIMP
- The kinetic decoupling starts at $x \sim 20$ (freeze-out time)



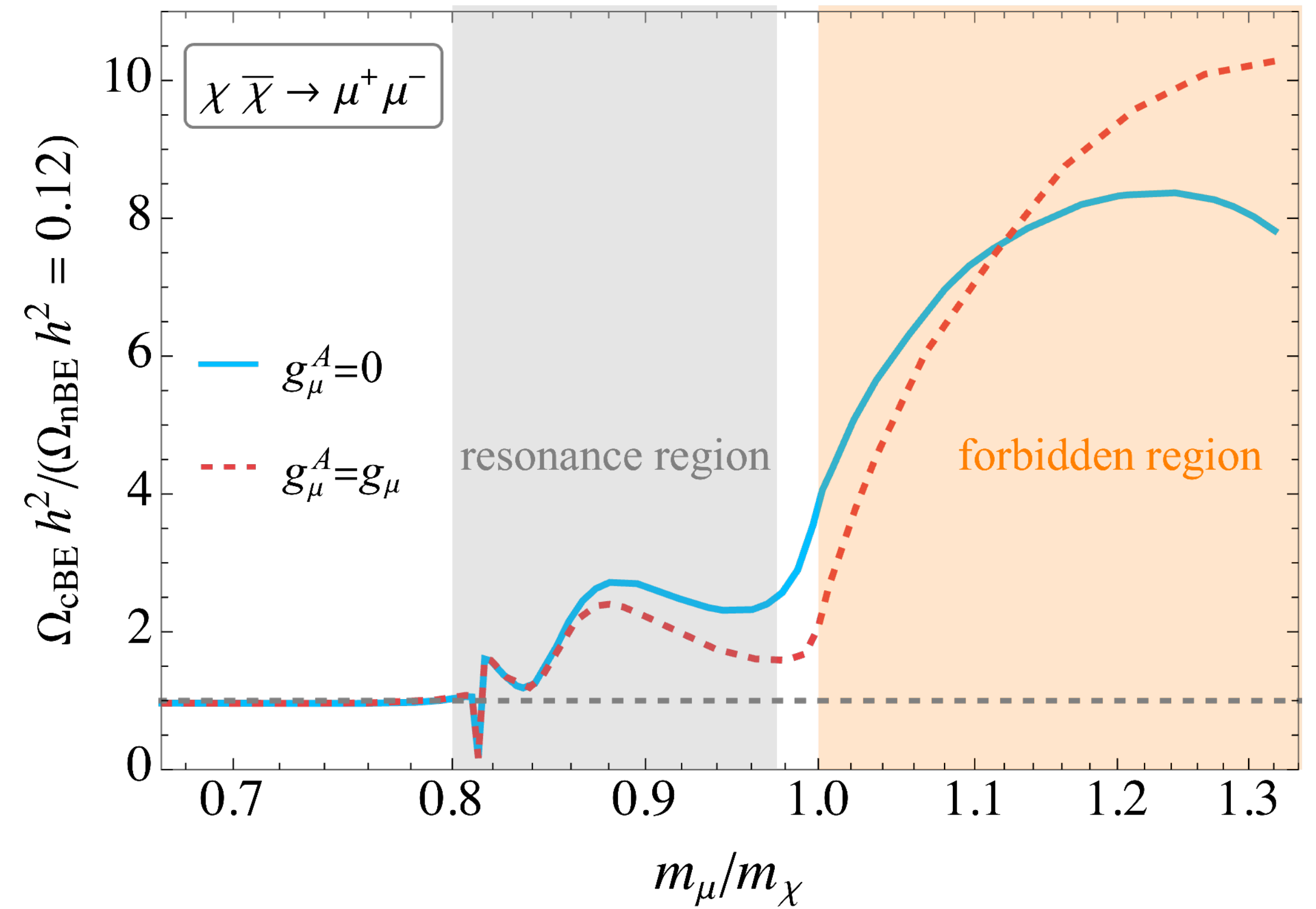
EKD on Forbidden Dark Matter

- A drop of the temperature and a large DM yield for CBE treatment
- there is a cooling phase during the evolution
- the decreased kinetic energy of DM particles suppresses the forbidden annihilation rate, which gives rise to larger abundances.



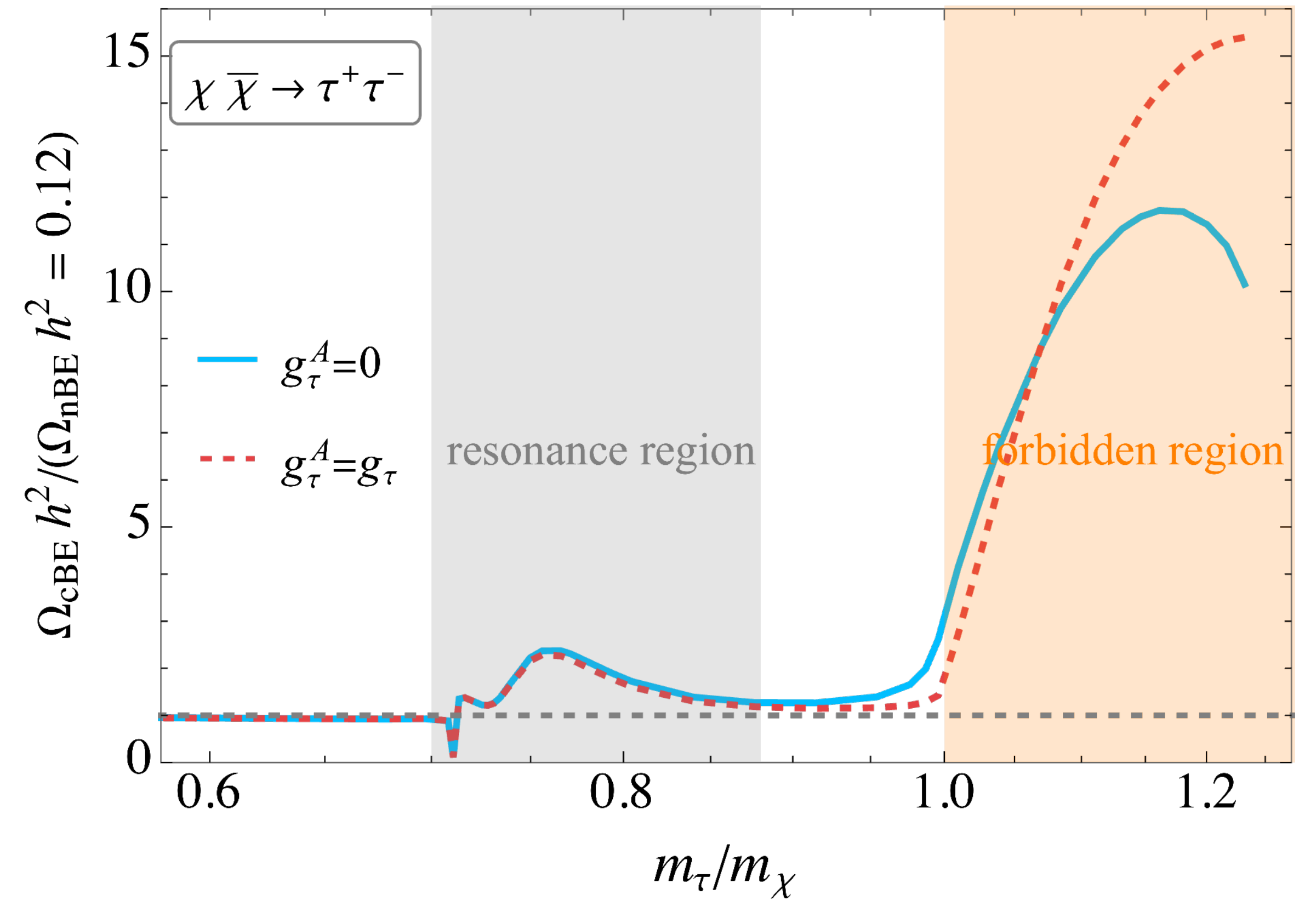
EKD effect on Forbidden Dark Matter

- $\mu\mu$ channel
- Forbidden region: $m_\mu/m_\chi > 1$
- $\Omega_{\text{cBE}} h^2 / \Omega_{\text{nBE}} h^2 \sim 2 - 10$
- Resonance region: see [1706.07433](#), [2004.10041](#)



EKD effect on Forbidden Dark Matter

- Similar results for $\tau\tau$ channel
- $\Omega_{\text{cBE}} h^2 / \Omega_{\text{nBE}} h^2 \sim 2 - 15$



Parameter space of $\chi\chi \rightarrow \mu^+\mu^-$

- Conclusion:

Viable parameter space shrinks by EKD

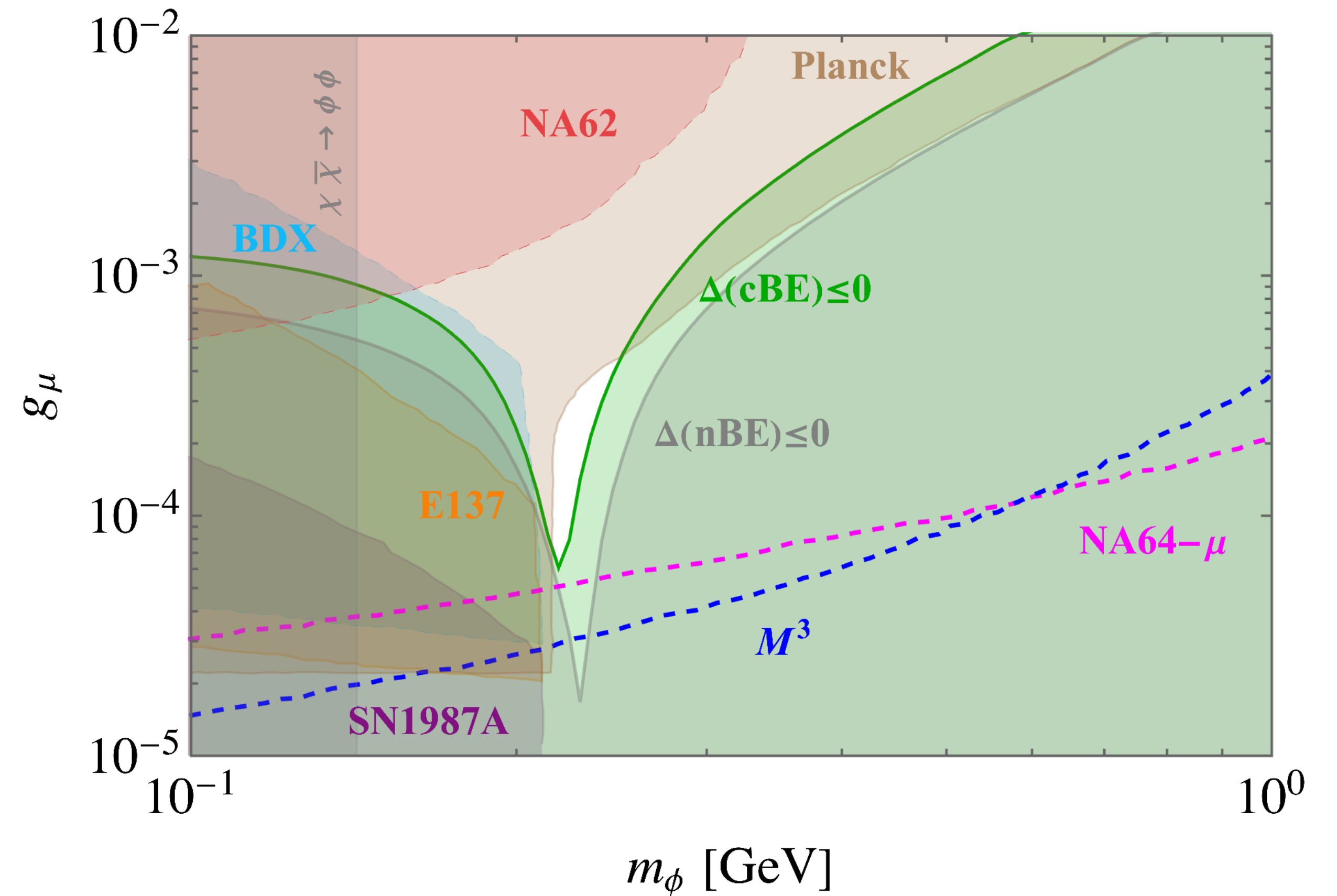
- Constraints:

(1) Planck (CMB) : $\chi\chi \rightarrow \gamma\gamma$

(2) SN 1987A: $\gamma p \rightarrow p\phi$ energy loss process

(3) SLAC E137 and JLab BDX: NA62 : $K \rightarrow \mu\nu\phi$

(4) NA64- μ and M^3 : $\mu N \rightarrow \mu N + \phi$



Parameter space of $\chi\chi \rightarrow \tau^+\tau^-$

- Conclusion:

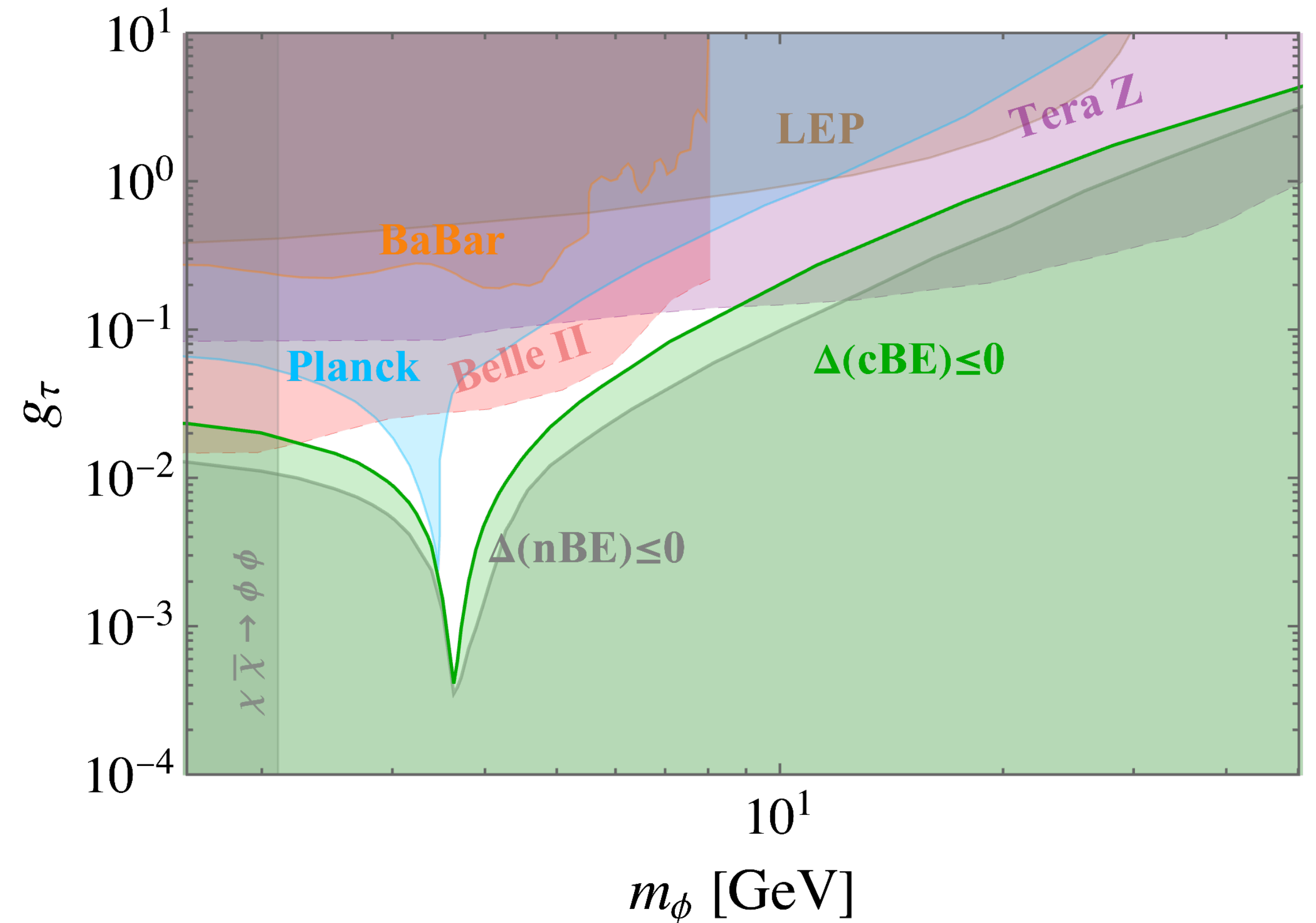
Same conclusion on relic density

- Constraints:

(1) Planck (CMB) : $\chi\chi \rightarrow \gamma\gamma$

(2) Babar and Belle II $(50 \text{ ab}^{-1})e^+e^- \rightarrow \gamma + \text{inv.}$

(3) LEP and Tera : $Z \rightarrow \tau^+\tau^- + \phi$



Summary

- Early Kinetic Decoupling effect is an indispensable ingredient.

- Effects on Forbidden DM:

The decreased kinetic energy of DM particles suppresses the forbidden annihilation rate, which gives rise to larger abundances, up to **an order of magnitude larger**

- EKD on other scenarios: Resonant annihilation of dark matter, Sommerfeld-enhanced annihilation (new work to appear), etc.

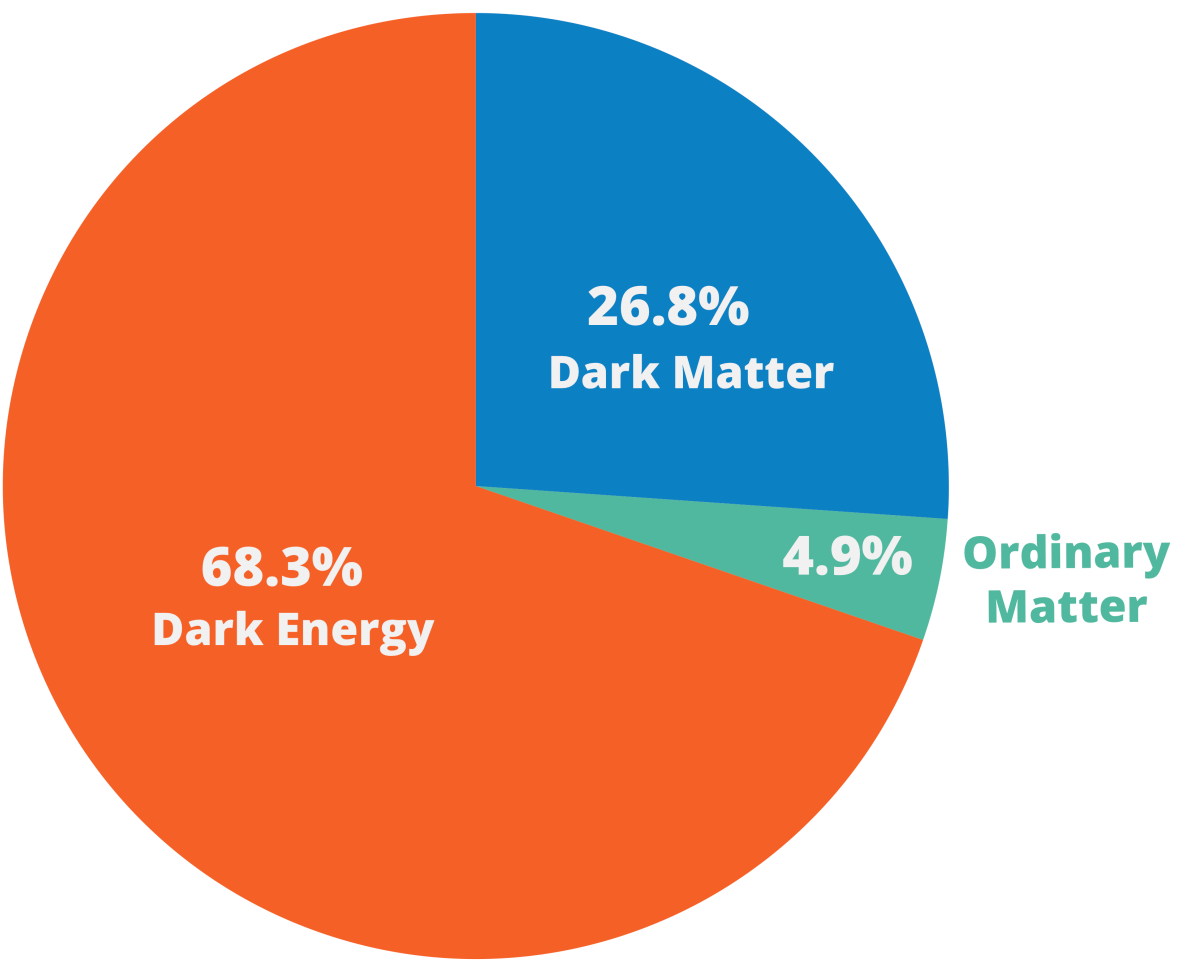
Thanks

BACKUP SLIDES

Motivation

Motivation

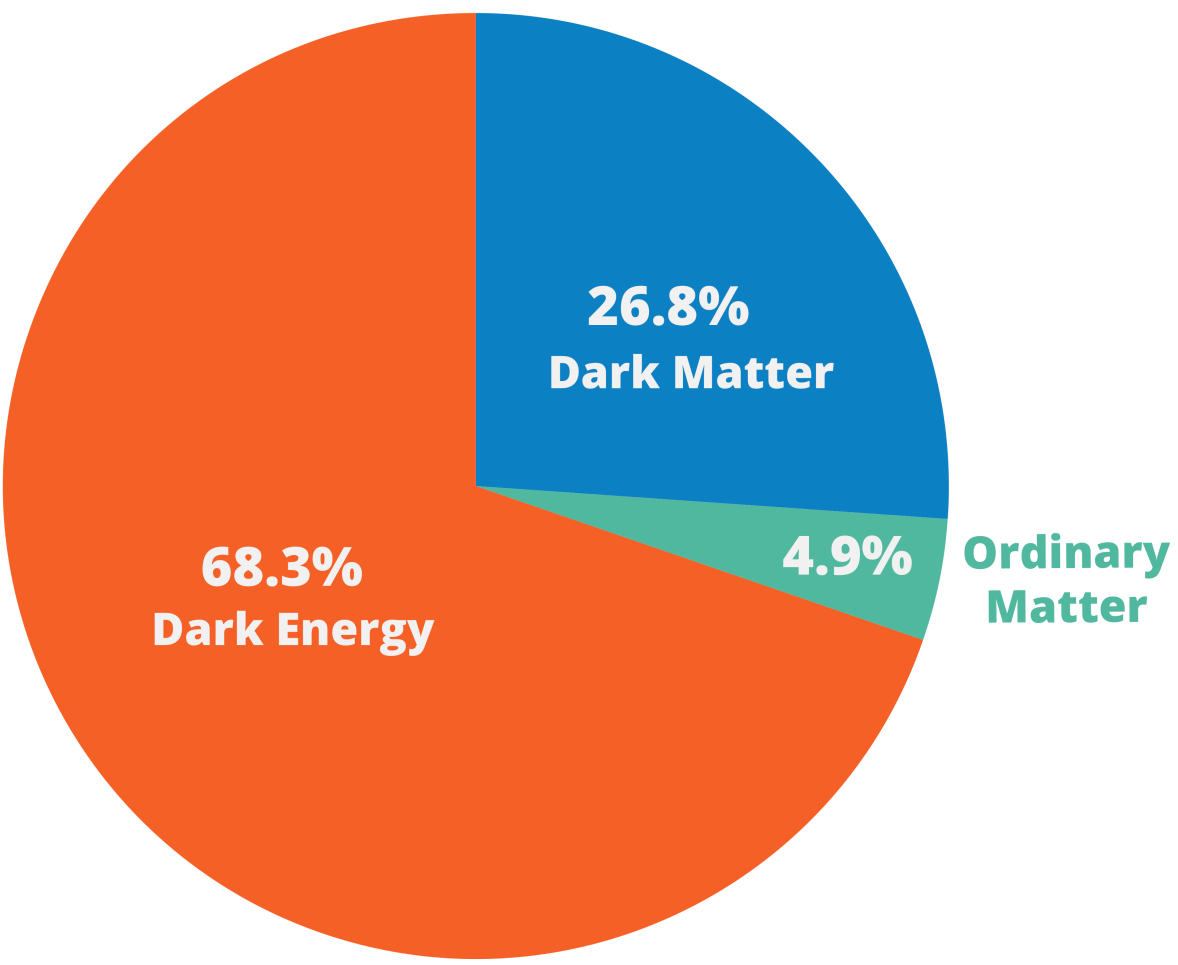
Relic Density measurement from Planck2018



Parameter	TT+lowE 68% limits	TE+lowE 68% limits	EE+lowE 68% limits	TT,TE,EE+lowE 68% limits	TT,TE,EE+lowE+lensing 68% limits
$\Omega_b h^2$	0.02212 ± 0.00022	0.02249 ± 0.00025	0.0240 ± 0.0012	0.02236 ± 0.00015	0.02237 ± 0.00015
$\Omega_c h^2$	0.1206 ± 0.0021	0.1177 ± 0.0020	0.1158 ± 0.0046	0.1202 ± 0.0014	0.1200 ± 0.0012
$100\theta_{MC}$	1.04077 ± 0.00047	1.04139 ± 0.00049	1.03999 ± 0.00089	1.04090 ± 0.00031	1.04092 ± 0.00031
τ	0.0522 ± 0.0080	0.0496 ± 0.0085	0.0527 ± 0.0090	$0.0544^{+0.0070}_{-0.0081}$	0.0544 ± 0.0073

Motivation

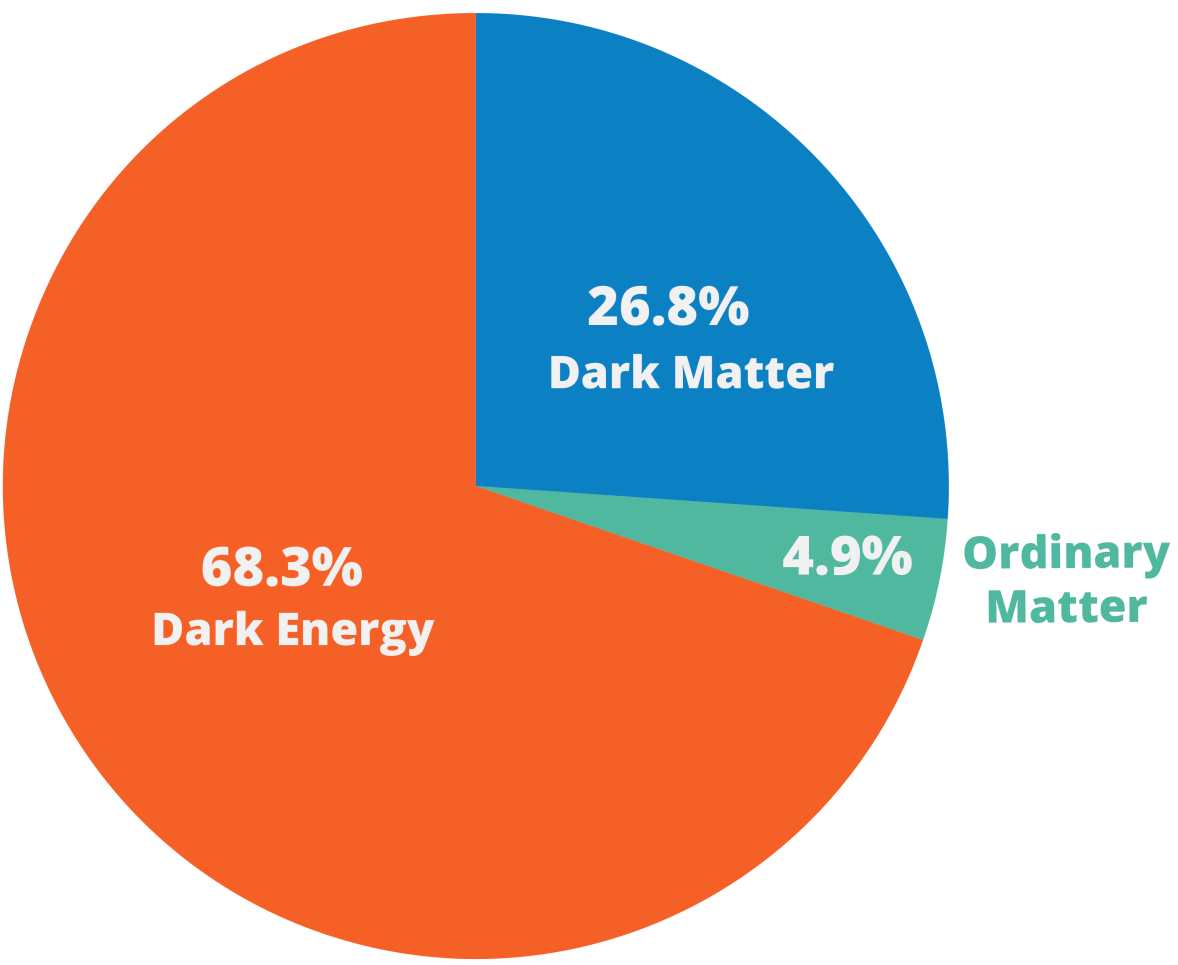
Relic Density measurement from Planck2018



Parameter	TT+lowE 68% limits	TE+lowE 68% limits	EE+lowE 68% limits	TT,TE,EE+lowE 68% limits	TT,TE,EE+lowE+lensing 68% limits
$\Omega_b h^2$	0.02212 ± 0.00022	0.02249 ± 0.00025	0.0240 ± 0.0012	0.02236 ± 0.00015	0.02237 ± 0.00015
$\Omega_c h^2$	0.1206 ± 0.0021	0.1177 ± 0.0020	0.1158 ± 0.0046	0.1202 ± 0.0014	0.1200 ± 0.0012
$100\theta_{MC}$	1.04077 ± 0.00047	1.04139 ± 0.00049	1.03999 ± 0.00089	1.04090 ± 0.00031	1.04092 ± 0.00031
τ	0.0522 ± 0.0080	0.0496 ± 0.0085	0.0527 ± 0.0090	$0.0544^{+0.0070}_{-0.0081}$	0.0544 ± 0.0073

Motivation

Relic Density measurement from Planck2018

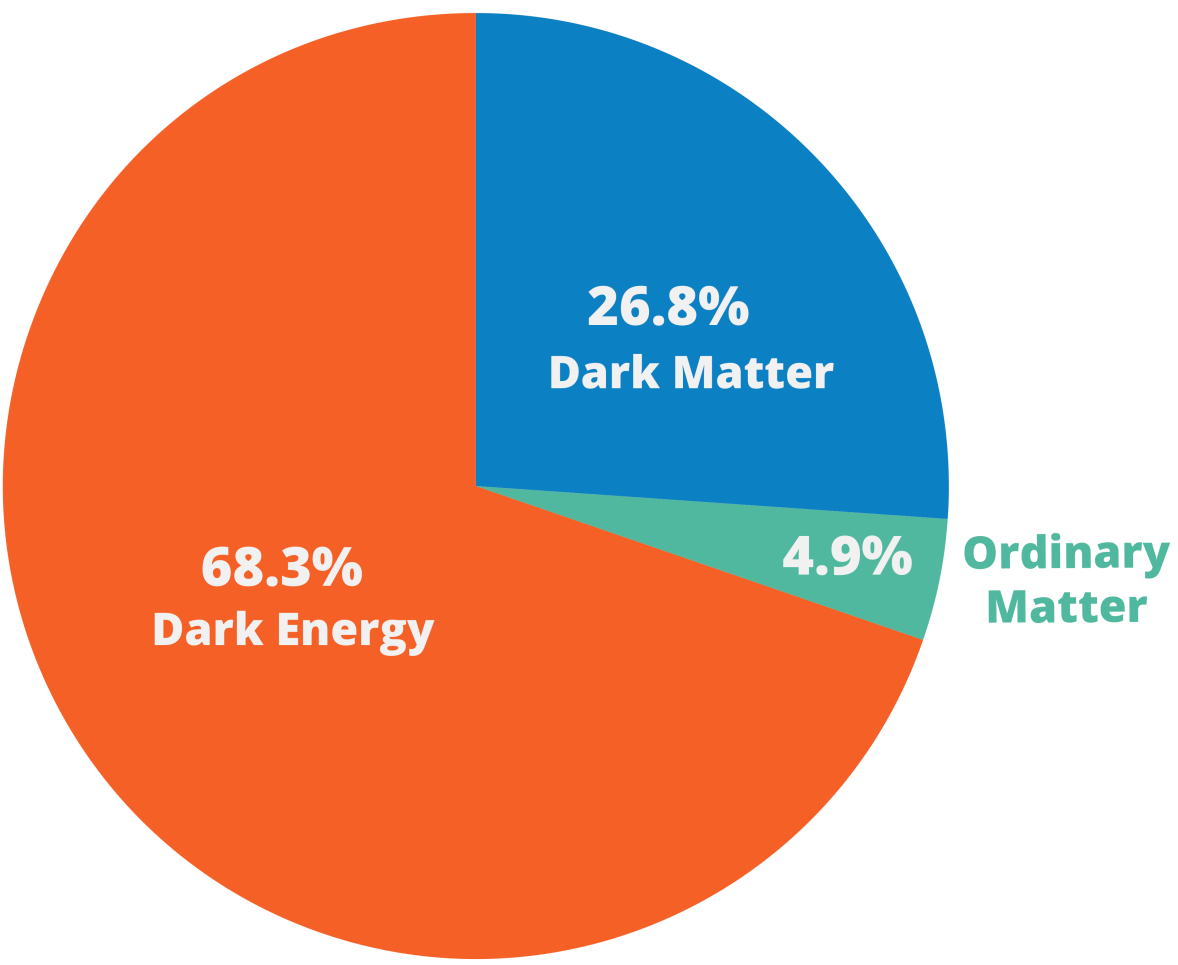


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High precision of Measurements

Motivation

Relic Density measurement from Planck2018

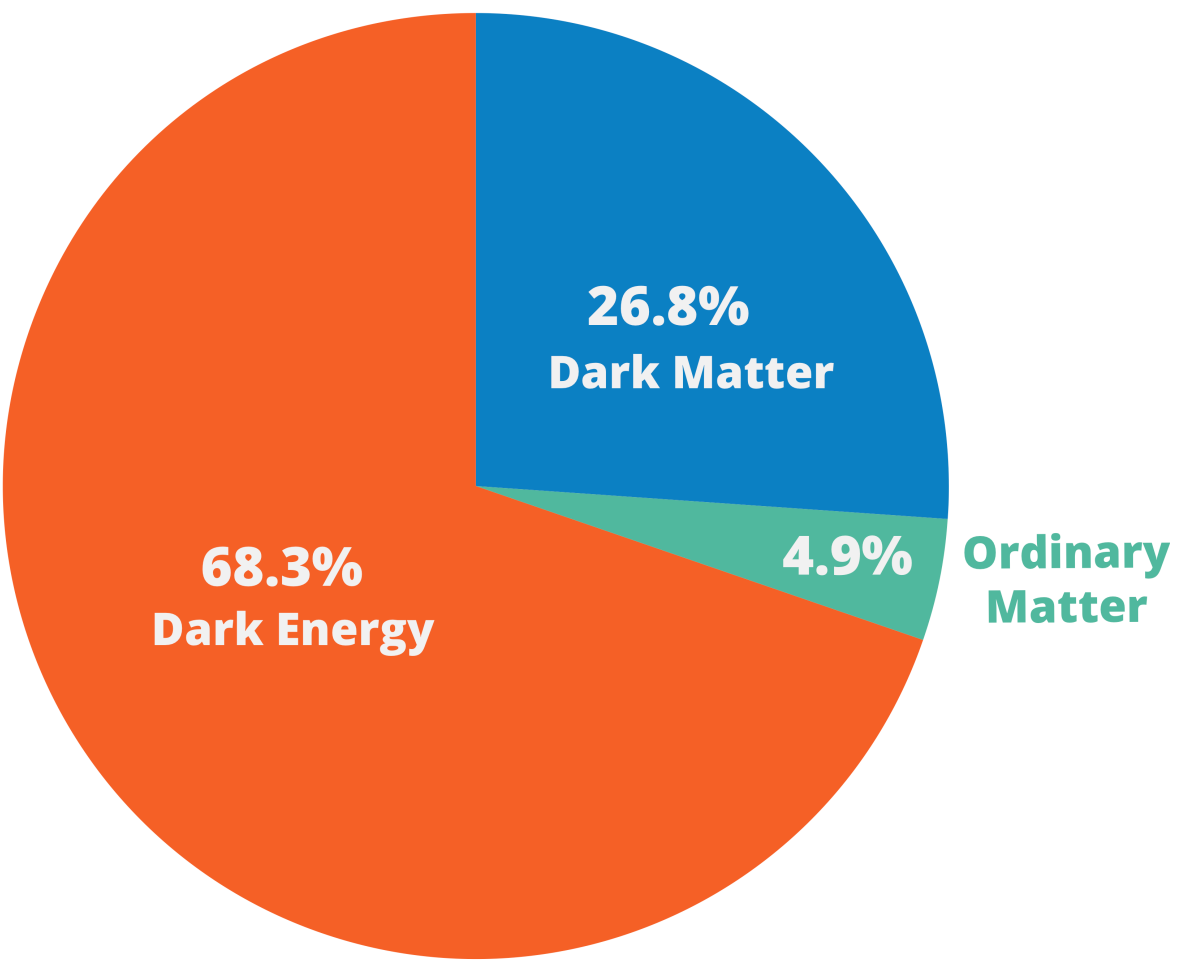


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High precision of Measurements ➡

Motivation

Relic Density measurement from Planck2018

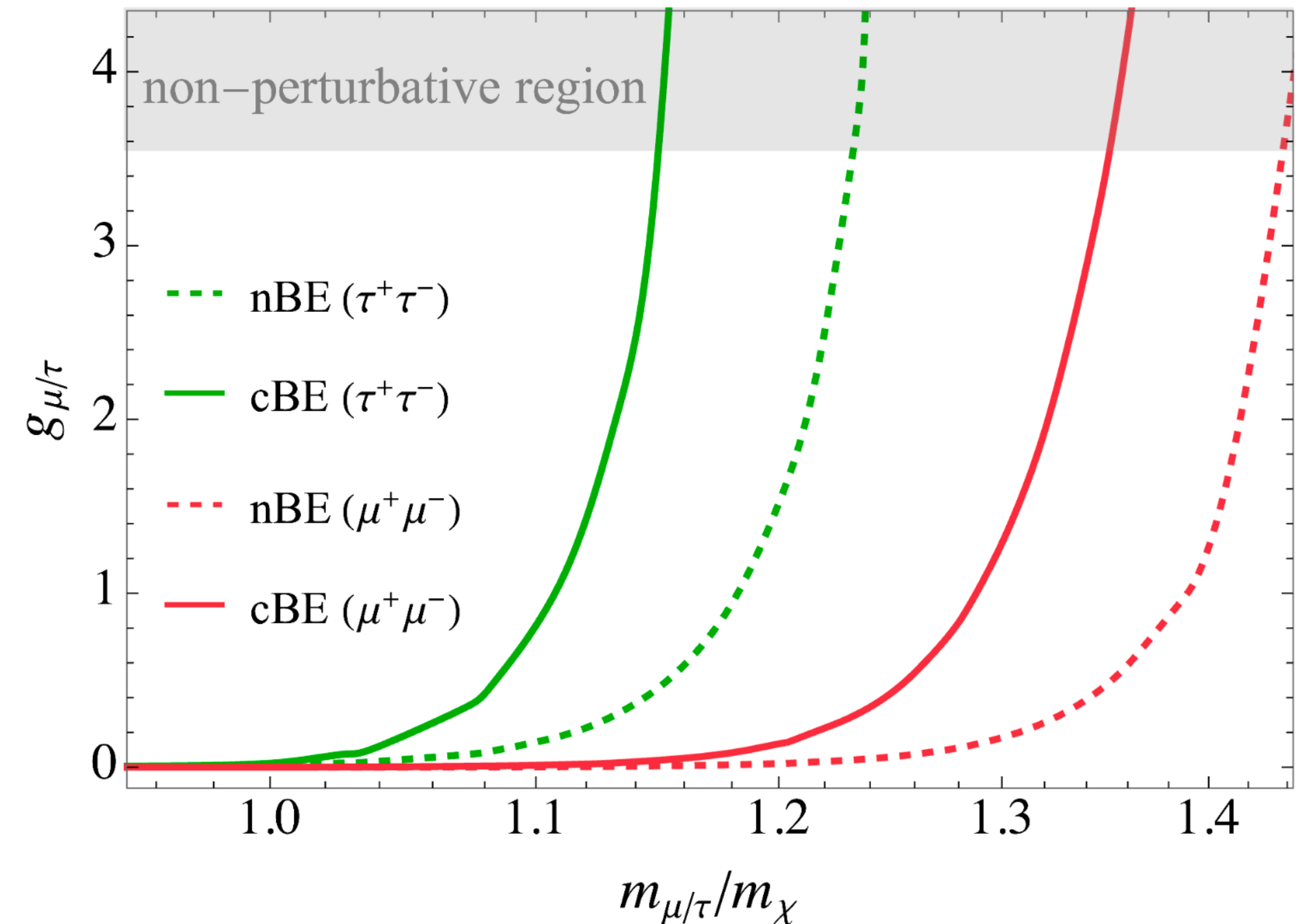


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High precision of Measurements ➡ High Accuracy of Calculations

Upper Boundary on mass ratio

- Fulfill $\Omega h^2 = 0.12$
- Gray region: perturbativity $g_\ell < \sqrt{4\pi}$
- Smaller m_μ/m_χ ratios are allowed from perturbativity requirement compared to nBE



Literature on the EKD

- Resonant annihilation of dark matter: [1706.07433], [2004.10041], [2103.01944]
- Sommerfeld-enhanced annihilation: [2103.01944]
- Sub-threshold annihilation (also known as the forbidden annihilation): [2103.01944]
- Others: [1910.01549, 1912.02870, 2104.03238, 2104.05684, 2106.01956, 2111.01267, 2201.06456, 2204.07078]

