# Probing Ultralight Dark Matter with Space-based Gravitational-Wave Interferometers



### 中国科学院大学

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Contents





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# Motivation

- Standard model is not complete
- Dark Matter and Dark Energy
- Neutrino mass
- Matter-Antimatter asymmetry
- Theoretical Problems
  - Strong CP problem
  - Hierarchy problem
  - ➢ Fermion mass hierarchy
  - Unification of forces

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# Ultralight Bosonic Fields

- > Well-motivated in many physical and cosmological models
- > Popular dark matter candidate, dark energy candidate
- > Topological objects, domain walls, compact objects, …







 $\succ$  Scalar field  $\phi$ 

$$\mathcal{L} = \frac{1}{2} \partial^{\mu} \phi \partial_{\mu} \phi - \frac{1}{2} m_{\phi}^{2} \phi^{2} - C \frac{\phi}{M_{P}} \mathcal{O}_{\mathrm{SM}}, \quad \phi\left(t, \vec{x}\right) = \phi_{\vec{k}} e^{i(\omega t - \vec{k} \cdot \vec{x} + \theta_{0})},$$

> Interaction depending on the underlying theory, e.g.

$$C\frac{\phi}{M_P}m_{\psi}\overline{\psi}\psi \Rightarrow m_{\psi} \to \left(1+C\frac{\phi}{M_P}\right)m_{\psi}, \quad S = -\int m(\phi)\sqrt{-\eta_{\mu\nu}dx^{\mu}dx^{\nu}}.$$
$$\delta x^i(t,\vec{x}) = \mathcal{M}_s\hat{k}^i e^{im_{\phi}(t-v\hat{k}\cdot\vec{x})}, \quad \mathcal{M}_s \propto \phi_{\vec{k}}|\vec{k}|/m_{\phi}^2$$

 $\succ$  Vector field  $A_{\mu}$ 

$$\mathcal{L} = -\frac{1}{4}F^{\mu\nu}F_{\mu\nu} + \frac{1}{2}m_A^2 A^\nu A_\nu - \epsilon_D e J_D^\nu A_\nu, \quad \vec{A}(t, \vec{x}) = |\vec{A}|\hat{e}_A e^{i(\omega t - \vec{k} \cdot \vec{x})},$$
$$\delta x^i(t, \vec{x}) = \mathcal{M}_v \hat{e}_A^i e^{im_A(t - v\hat{k} \cdot \vec{x})}, \mathcal{M}_v \propto \epsilon_D e q_{D,j} |\vec{A}| / m_A M_j$$

DM property

$$\phi_{\vec{k}} = \frac{\sqrt{2\rho_{\rm DM}}}{m_{\phi}}, \qquad |\vec{A}| = \frac{\sqrt{2\rho_{\rm DM}}}{m_{A}}, \qquad v \sim 10^{-3}, \quad \vec{k} \approx m_{\phi} \vec{v} \text{ and } \omega \approx m_{\phi}$$

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# **Physical Effects**

### Atomic physics

- Arvanitaki, Huang & Tilburg (2014), Graham, Kaplan, Mardon, Rajendran & Terrano (2015), Safronova, Budker, DeMille, Kimball, Derevianko & Clark (2018),
- ≻ Stadnik (2022), ……
- Astrophysical physics
  - Pierce, Riles & Zhao (2018), Morisaki & Suyama (2019), Guo, Riles, Yang & Zhao 2019, Grote & Stadnik (2019),
  - An, Huang, Liu & Xue (2021), Chen, Shu, Xue, Yuan & Zhao (2019), Xia, Xu & Zhou (2020), Sun, Yang & Zhang (2021), Wu, Chen, & Huang (2023),
  - ≻ Liu, Lou & Ren (2021), Luu, Liu, Ren, Broadhurst, Yang, Wang & Xie (2023), ……
- Underground searches
  - ➢ Dark Matter Experiments, PandaX, XENONnT, …





# Space-based GW Interferometers

### LISA, Taiji and TianQin, sensitivity band 0.1 mHz ~ 0.1 Hz



Response

Gravitational wave can change the structure of spacetime, and the physical distance between objects

One can measure the phase by laser







 $\succ \text{ Response } \frac{\delta v(t)}{v_0} \equiv y_{BA} = -\frac{1}{2} \frac{n_i n_j}{1 + \vec{k} \cdot \vec{n}} \left[ h_{ij} \left( t - \frac{\vec{k} \cdot \vec{x}_B}{c} \right) - h_{ij} \left( t - \frac{\vec{k} \cdot \vec{x}_A + L}{c} \right) \right]$ 





### Response

- One-way Doppler shift

 $(2(1+\hat{n}_{rs}\cdot\hat{k}))$ 

$$\delta t_{rs} = -\hat{n}_{rs} \cdot \left[\delta \vec{x}(t, \vec{x}_r) - \delta \vec{x}(t - L, \vec{x}_s)\right]$$

$$rac{\delta 
u_{rs}}{
u_0} = rac{
u_{rs} - 
u_0}{
u_0} = -rac{d\,\delta t_{rs}}{dt}.$$

Fractional frequency change

$$y_{rs}(t) \equiv \frac{\delta \nu_{rs}}{\nu_0} = \mu_{rs} \left[ h(t, \vec{x_r}) - h(t - L, \vec{x_s}) \right], \ h(t, \vec{x}) \propto e^{im(t - v\hat{k} \cdot \vec{x_s})}$$
$$\mu_{rs} = \begin{cases} \hat{k} \cdot \hat{n}_{rs} & \text{for scalar field,} \\ \hat{e}_A \cdot \hat{n}_{rs} & \text{for vector field,} \\ \frac{\hat{n}_{rs}^i \hat{n}_{rs}^j e_{ij}(\hat{k}, \psi)}{\hat{n}} & \text{for gravitational wave,} \end{cases}$$







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# Time-Delay Interferometry

- > The arm lengths are not equal
- Laser frequency noises dominate

$$\begin{split} X(t) &\equiv \left[ \Delta y_{PD1}(t) - \Delta y_{PD2}(t) \right] - \left[ \Delta y_{PD1} \left( t - T_2 \right) - \Delta y_{PD2} \left( t - T_1 \right) \right] \\ &= \left[ H_1(t) - H_2(t) + p \left( t - T_1 \right) - p \left( t - T_2 \right) \right] \\ &- \left[ H_1 \left( t - T_2 \right) - H_2 \left( t - T_1 \right) + p \left( t - T_1 \right) - p \left( t - T_2 \right) \right] \\ &= H_1(t) - H_2(t) - H_1 \left( t - T_2 \right) + H_2 \left( t - T_1 \right), \end{split}$$

➤ Michelson interferometry  $X(t) \equiv [\Delta y_{PD2} (t - T_1) + \Delta y_{PD1}(t)] - [\Delta y_{PD1} (t - T_2) + \Delta y_{PD2}(t)]$ 



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#### Input for TDI

$$\begin{split} \eta_{i'} &\equiv s_{i'} + \frac{\varepsilon_{i'} - \tau_{i'}}{2} + D_{i+1'} \frac{\varepsilon_{i-1} - \tau_{i-1}}{2} + \frac{\tau_i - \tau_{i'}}{2} \\ \eta_i &\equiv s_i + \frac{\varepsilon_i - \tau_i}{2} + D_{i-1} \frac{\varepsilon_{i+1'} - \tau_{i+1'}}{2} - D_{i-1} \frac{\tau_{i+1} - \tau_{i+1'}}{2} \end{split}$$

L3

$$\begin{split} \eta_{1'} &\sim D_{2'} p_3 - p_1, \, \eta_1 \sim D_3 p_2 - p_1, \\ \eta_{2'} &\sim D_{3'} p_1 - p_2, \, \eta_2 \sim D_1 p_3 - p_2, \\ \eta_{3'} &\sim D_{1'} p_2 - p_3, \, \eta_3 \sim D_2 p_1 - p_3. \end{split}$$

phasemeter

clock





downlink

# Time-Delay Interferometry

- There are multiple combinations
- Michelson channels
- $$\begin{split} X(t) &= (\eta_{2':\mathbf{322'}} + \eta_{1:\mathbf{22'}} + \eta_{3:\mathbf{2'}} + \eta_{1'}) (\eta_{3:\mathbf{2'3'3}} + \eta_{1':\mathbf{3'3}} + \eta_{2':\mathbf{3}} + \eta_{1}), \\ Y(t) &= (\eta_{3':\mathbf{133'}} + \eta_{2:\mathbf{33'}} + \eta_{1:\mathbf{3'}} + \eta_{2'}) (\eta_{1:\mathbf{3'1'1}} + \eta_{2':\mathbf{1'1}} + \eta_{3':\mathbf{1}} + \eta_{2}), \\ Z(t) &= (\eta_{1':\mathbf{211'}} + \eta_{3:\mathbf{11'}} + \eta_{2:\mathbf{1'}} + \eta_{3'}) (\eta_{2:\mathbf{1'2'2}} + \eta_{3':\mathbf{2'2}} + \eta_{1':\mathbf{2}} + \eta_{3}). \end{split}$$

### Sagnac channels

$$\begin{aligned} \alpha(t) &= (\eta_{2':\mathbf{1'2'}} + \eta_{3':\mathbf{2'}} + \eta_{1'}) - (\eta_{3:\mathbf{13}} + \eta_{2:\mathbf{3}} + \eta_{1}), \\ \beta(t) &= (\eta_{3':\mathbf{2'3'}} + \eta_{1':\mathbf{3'}} + \eta_{2'}) - (\eta_{1:\mathbf{21}} + \eta_{3:\mathbf{1}} + \eta_{2}), \\ \gamma(t) &= (\eta_{1':\mathbf{3'1'}} + \eta_{2':\mathbf{1'}} + \eta_{3'}) - (\eta_{2:\mathbf{32}} + \eta_{1:\mathbf{2}} + \eta_{3}). \end{aligned}$$

### $\succ \zeta$ channel

 $\zeta(t) = (\eta_{1':1'} + \eta_{2':2'} + \eta_{3':3'}) - (\eta_{1:1} + \eta_{2:2} + \eta_{3:3}).$ 





# Transfer Function

- Fourier transform
- One-way single link

$$h(t) = \frac{\sqrt{T}}{2\pi} \int_0^\infty \tilde{h}(\omega) e^{i\omega t} d\omega$$
  
$$y_{rs}(t) = \mu_{rs} \frac{\sqrt{T}}{2\pi} \int_0^\infty d\omega \ \tilde{h}(\omega) e^{i\omega t} \left[ e^{-i\vec{k}\cdot\vec{x}_r} - e^{-i(\tau + \vec{k}\cdot\vec{x}_s)} \right],$$
  
$$\tilde{y}_{rs}(\omega) = \mu_{rs} \ \tilde{h}(\omega) \left[ e^{-i(\vec{k}\cdot\vec{x}_r)} - e^{-i(\tau + \vec{k}\cdot\vec{x}_s)} \right].$$

Transfer function, sky and polarization averaged



#### $\succ$ DM is also different from gravitational wave, velocity effect, $\cdots$ α $10^{0}$ 10<sup>0</sup> $10^{-2}$ $10^{-2}$ $10^{-4}$ $10^{-4}$ **∠** 10<sup>-6</sup> $10^{-8}$ $10^{-6}$ $10^{-10}$ GW GW $10^{-8}$ Scalar $10^{-12}$ Scalar Vector Vector $10^{-14}$ 10<sup>-5</sup> 10<sup>-2</sup> $10^{-4}$ $10^{-3}$ $10^{-1}$ $10^{-4}$ 10<sup>-3</sup> 10-2 $10^{0}$ $10^{-1}$ 10<sup>0</sup> f(Hz) f(Hz) 14

- **Transfer Functions** 
  - Different channels have different transfer functions

### Sensitivity

 $> \text{ Defined by } S_O(f) = \frac{N_O(f)}{R_O(f)}, \quad N_X = 16\sin^2(\tau) \left\{ [3 + \cos(2\tau)] S_{acc} + S_{oms} \right\}, \quad \tau = 2\pi f L \\ S_{oms}(f) = \left( s_{oms} \frac{2\pi f}{c} \right)^2 \left[ 1 + \left( \frac{2 \times 10^{-3} \text{ Hz}}{f} \right)^4 \right] \frac{1}{\text{Hz}}, \qquad \qquad \text{LISA: } s_{oms} = 15 \times 10^{-12} \text{ m}, s_{acc} = 3 \times 10^{-15} \text{ m/s}^2, \\ S_{acc}(f) = \left( \frac{s_{acc}}{2\pi fc} \right)^2 \left[ 1 + \left( \frac{0.4 \times 10^{-3} \text{ Hz}}{f} \right)^2 \right] \left[ 1 + \left( \frac{f}{8 \times 10^{-3} \text{ Hz}} \right)^4 \right] \frac{1}{\text{Hz}}, \qquad \qquad \text{TianQin: } s_{oms} = 1 \times 10^{-12} \text{ m}, s_{acc} = 1 \times 10^{-15} \text{ m/s}^2.$ 

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Sensitivity

> Optimal channels  $A = \frac{1}{\sqrt{2}} [Z - X], E = \frac{1}{\sqrt{6}} [X - 2Y + Z], T = \frac{1}{\sqrt{3}} [X + Y + Z].$   $\frac{1}{S_{\eta}} = \frac{1}{S_A} + \frac{1}{S_E} + \frac{1}{S_T}$ 

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## Sensitivity on scalar DM





assuming  $d_m = 0$  and  $d_g^* \approx 0.9 d_g$ .

Equivalence principle is violated.

➢ MICROSCOPE



Sensitivity on vector DM For example, vector fields couple to baryon number *B*, or *B-L*, effectively neutron number. Sensitivity on ratio  $\epsilon_D = e_D/e$ 



# Summary

- Ultralight bosonic fields (ULBFs) are motivated and predicted in many physical and cosmological theories
- ULBFs can also be dark matter candidates, ULDM
- The tiny coupling between ULBFs and standard model particles can induce observable physical effects
- We use the space-based gravitational-wave interferometers to probe ULBFs and ULDM, taking the various time-delay interferometry channels into account

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- Good sensitivity can be obtained in some parameter region
- Caveat: Stochastic effect to be implemented, stay tuned

### Thanks for your attention!