

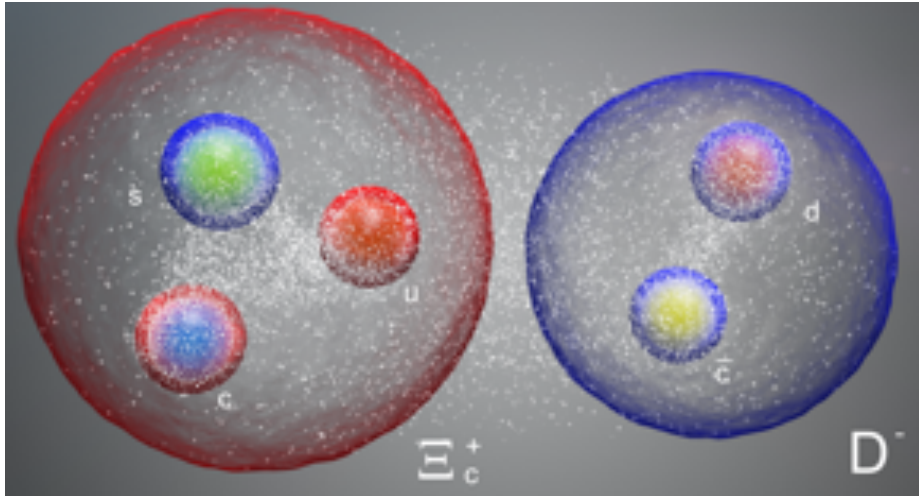
Observation of a neutral pentaquark candidate in B-meson decays

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中国科学院大学

2022年7月8日



Outline



- Observation of first neutral pentaquark in $B^- \rightarrow J/\psi \Lambda \bar{p}$ decays

LHCb-PAPER-2022-031 in preparation

- Evidence for a charged pentaquark in $B_s^0 \rightarrow J/\psi p \bar{p}$ decays

PRL128(2022)062001

Pentaquark in the quark model

A SCHEMATIC MODEL OF BARYONS AND MESONS *

M. GELL-MANN

California Institute of Technology, Pasadena, California

Received 4 January 1964

the triplet as "quarks" q and the members of the anti-triplet as anti-quarks $\bar{q}$. Baryons can now be constructed from quarks by using the combinations (qqq) , $(qqq\bar{q})$, etc., while mesons are made out of $(q\bar{q})$, $(qq\bar{q}\bar{q})$, etc. It is assuming that the lowest

AN SU_3 MODEL FOR STRONG INTERACTION SYMMETRY AND ITS BREAKING
II *)

G. Zweig **)

CERN--Geneva

In general, we would expect that baryons are built not only from the product of three aces, AAA , but also from $\bar{A}AAAA$, $\bar{A}AAAAA$, etc., where $\bar{A}$ denotes an anti-ace. Similarly, mesons could be formed from $\bar{A}A$, $\bar{A}AAA$ etc. For the low mass mesons and baryons we will assume the simplest possibilities, $\bar{A}A$ and AAA , that is, "deuces and treys".

□ Predicted in the quark model

□ Minimal content $qqqq\bar{q}$

Exotic hadron naming convention

LHCb-PUB-2022-013, arXiv:2206.15233

Minimal quark content	Current name	$I^{(G)}, J^{P(C)}$	Proposed name	Reference
$c\bar{c}$	$\chi_{c1}(3872)$	$I^G = 0^+, J^{PC} = 1^{++}$	$\chi_{c1}(3872)$	[24, 25]
$c\bar{c}u\bar{d}$	$Z_c(3900)^+$	$I^G = 1^+, J^P = 1^+$	$T_{\psi 1}^b(3900)^+$	[26–28]
$c\bar{c}u\bar{d}$	$X(4100)^+$	$I^G = 1^-$	$T_{\psi}(4100)^+$	[29]
$c\bar{c}u\bar{d}$	$Z_c(4430)^+$	$I^G = 1^+, J^P = 1^+$	$T_{\psi 1}^b(4430)^+$	[30, 31]
$c\bar{c}(s\bar{s})$	$\chi_{c1}(4140)$	$I^G = 0^+, J^{PC} = 1^{++}$	$\chi_{c1}(4140)$	[32–35]
$c\bar{c}u\bar{s}$	$Z_{cs}(4000)^+$	$I = \frac{1}{2}, J^P = 1^+$	$T_{\psi s 1}^{\theta}(4000)^+$	[7]
$c\bar{c}u\bar{s}$	$Z_{cs}(4220)^+$	$I = \frac{1}{2}, J^P = 1^?$	$T_{\psi s 1}(4220)^+$	[7]
$c\bar{c}c\bar{c}$	$X(6900)$	$I^G = 0^+, J^{PC} = ??^+$	$T_{\psi\psi}(6900)$	[4]
$cs\bar{u}\bar{d}$	$X_0(2900)$	$J^P = 0^+$	$T_{cs 0}(2900)^0$	[5, 6]
$cs\bar{u}\bar{d}$	$X_1(2900)$	$J^P = 1^-$	$T_{cs 1}(2900)^0$	[5, 6]
$cc\bar{u}\bar{d}$	$T_{cc}(3875)^+$		$T_{cc}(3875)^+$	[8, 9]
$b\bar{b}u\bar{d}$	$Z_b(10610)^+$	$I^G = 1^+, J^P = 1^+$	$T_{\Upsilon 1}^b(10610)^+$	[36]
$c\bar{c}uud$	$P_c(4312)^+$	$I = \frac{1}{2}$	$P_{\psi}^N(4312)^+$	[3]
$c\bar{c}uds$	$P_{cs}(4459)^0$	$I = 0$	$P_{\psi s}^{\Lambda}(4459)^0$	[20]

□ P for pentaquarks, T for tetraquarks

□ P states: i.e. $P_{\psi}^N(4312)^+$, $P_{\psi s}^{\Lambda}(4459)^0$

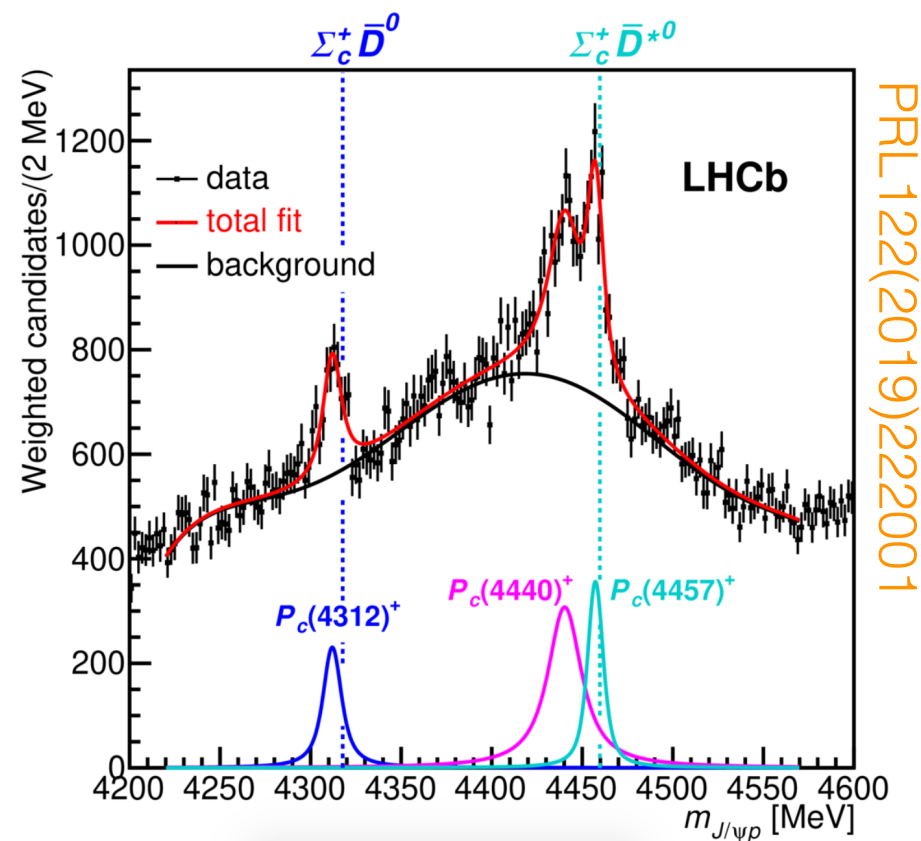
Superscript: denotes isospin: $\Lambda, N, \Sigma, \Delta$ for $I = 0, 1/2, 1, 3/2$

Subscript: Υ, ψ, ϕ for hidden beauty, charm, strangeness; b, c, s for open flavor quantum numbers

Pentaquark candidates in B-baryon decays

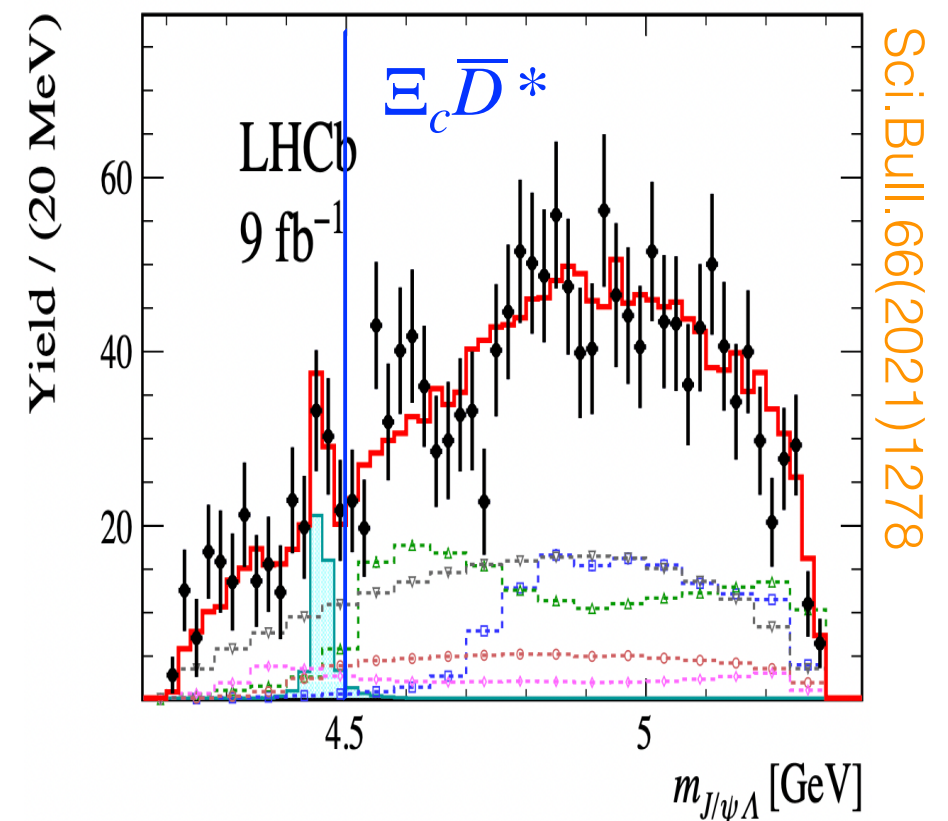
Observation in $\Lambda_b^0 \rightarrow pK^-J/\psi$

$P_\psi^N(4312)^+$, $P_\psi^N(4440)^+$, $P_\psi^N(4457)^+$



Evidence in $\Xi_b^- \rightarrow J/\psi \Lambda K^-$

$P_{\psi s}^\Lambda(4459)^0$

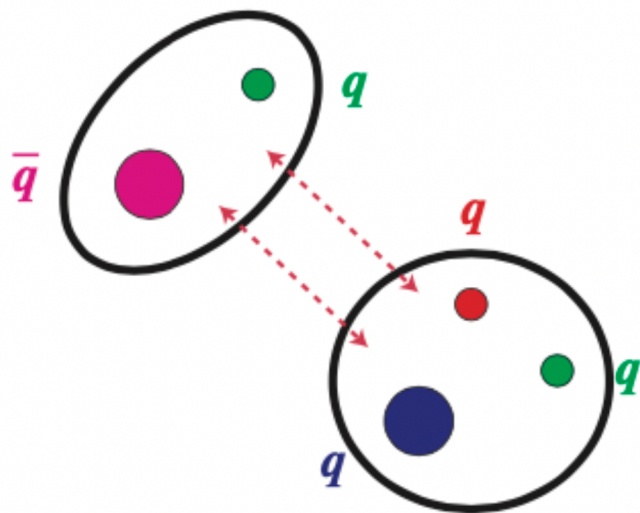


- ❑ Narrow structures require excellent mass resolution
- ❑ Interesting physics related to thresholds
- ❑ Spin and parity not determined

Interpretations

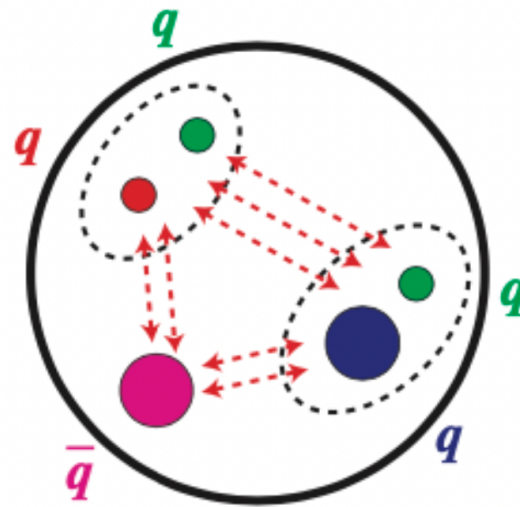
Their nature is still largely unknown, various interpretations including:

Hadronic Molecules



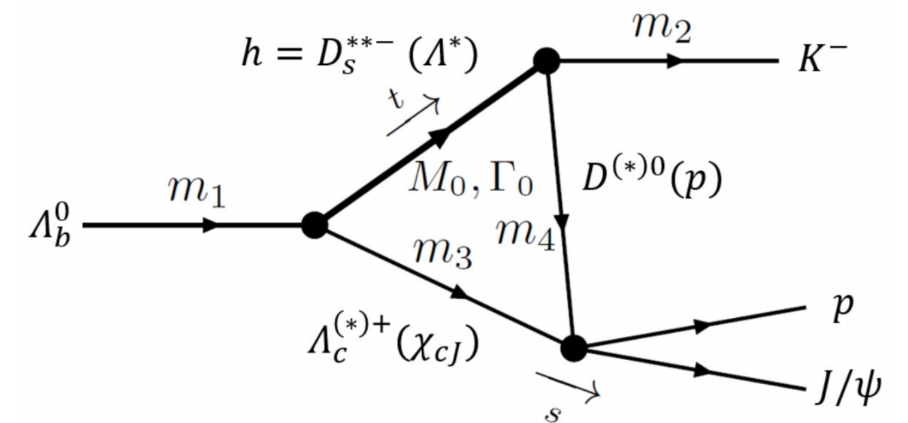
Rev. Mod. Phys. 90(2018)015004
PRD103(2021)112006
Eur.Phys.J.C 82 (2022) 7, 581

Compact pentaquark



Phys. Rept. 668(2017)1
Few Body Syst. 57 (2016)1185

Rescattering effects

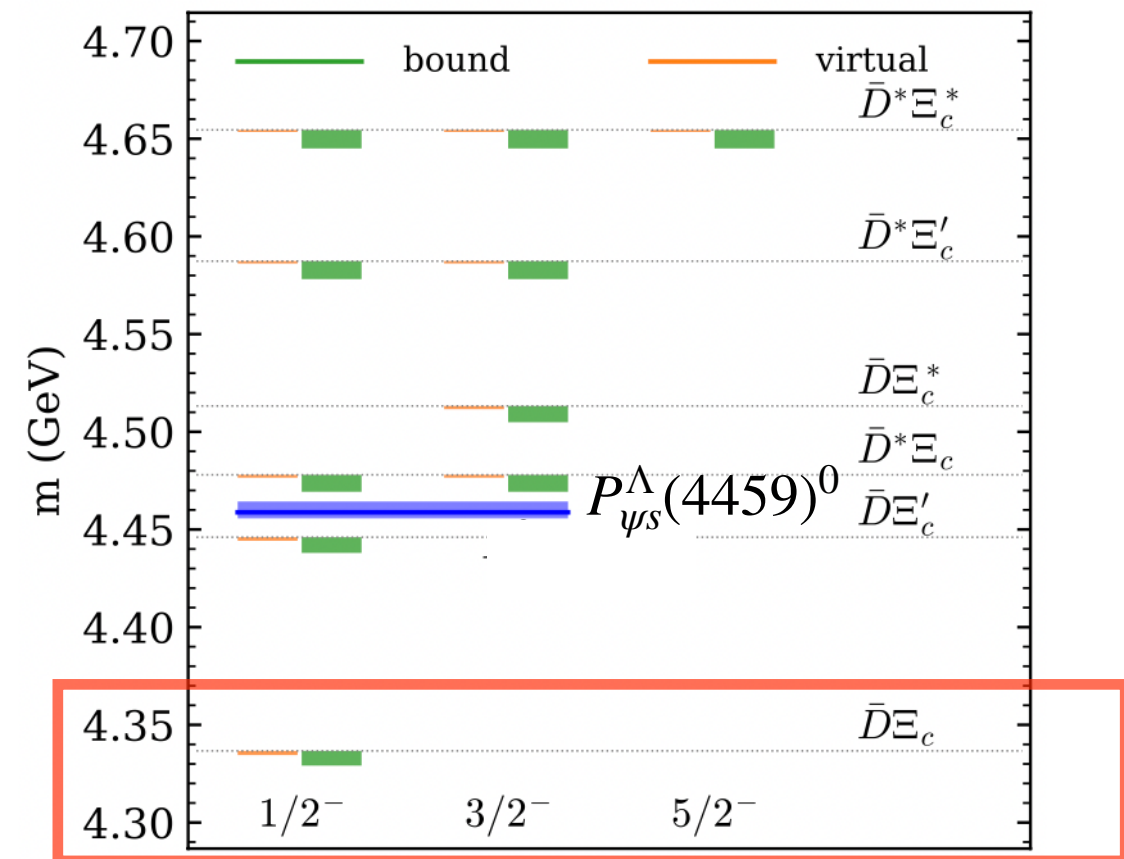
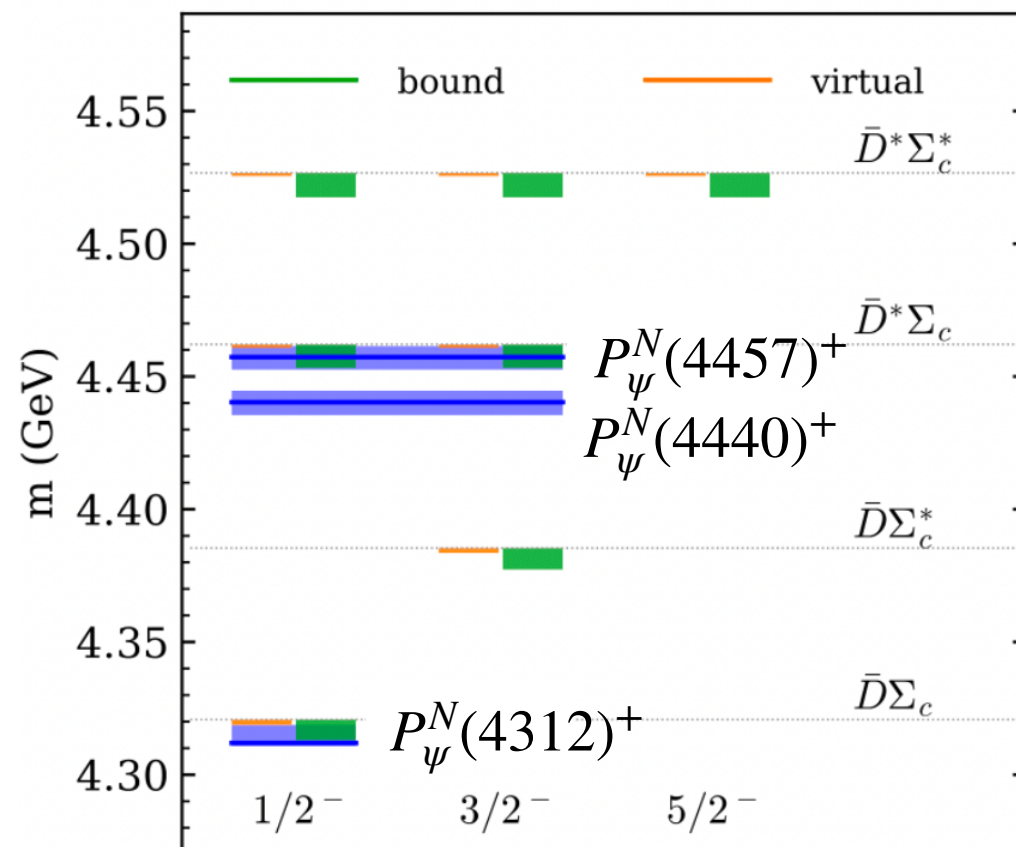


Prog. Part. Nucl. Phys.
112 (2020)103757

More pentaquark states?

contact terms which are resummed to generate poles. It turns out that if a system is attractive near threshold by the light meson exchange, there is a pole close to threshold corresponding to a bound state or a virtual state, depending on the strength of interaction and the cutoff. In total, 229 molecular states are predicted. The observed near-threshold structures with hidden-charm, like the

[Progr.Phys.41\(2021\)65-93](#)

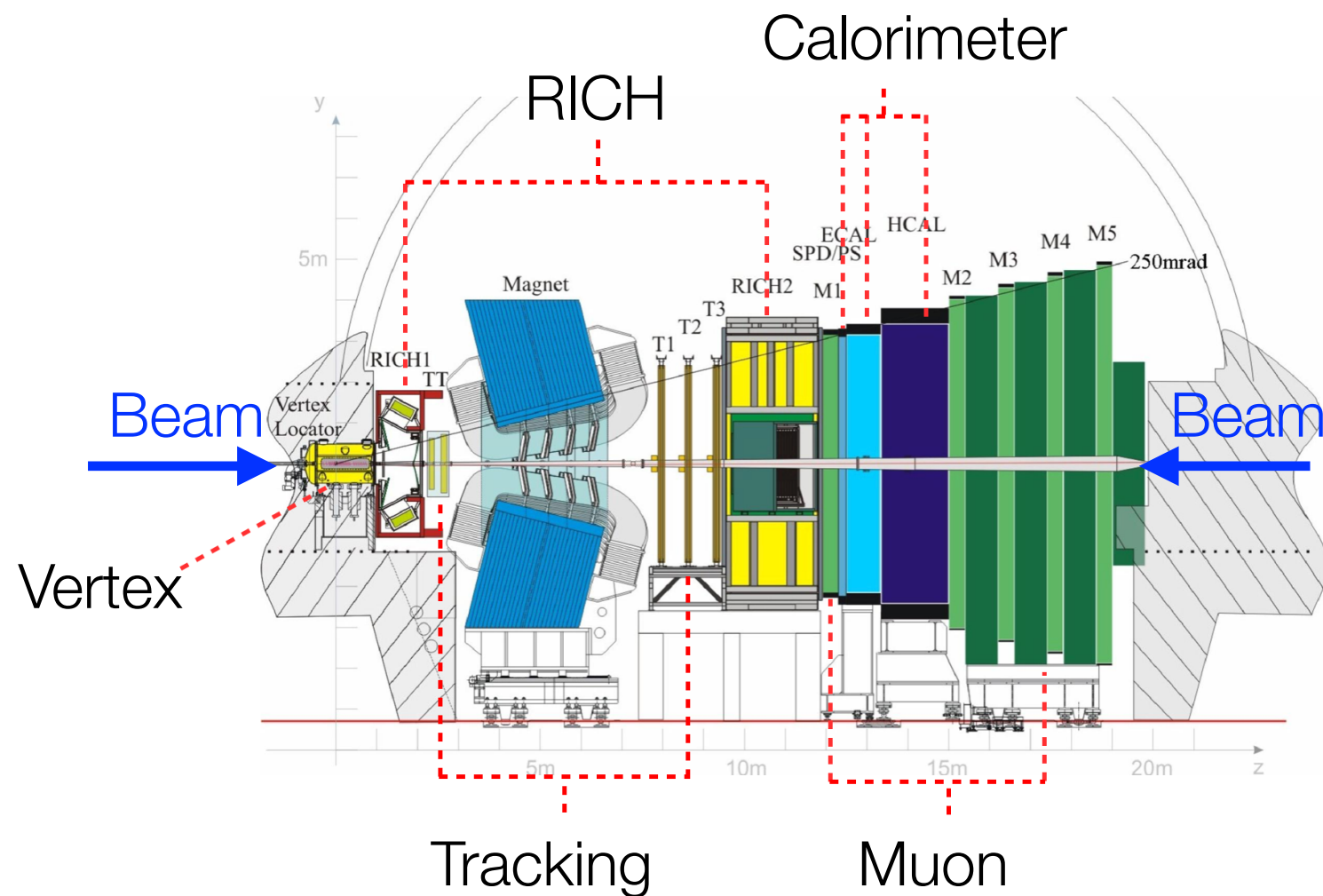


LHCb detector and performance

Int. J. Mod. Phys. A 30 (2015) 1530022

LHCb is a forward spectrometer

JINST3(2008)S08005

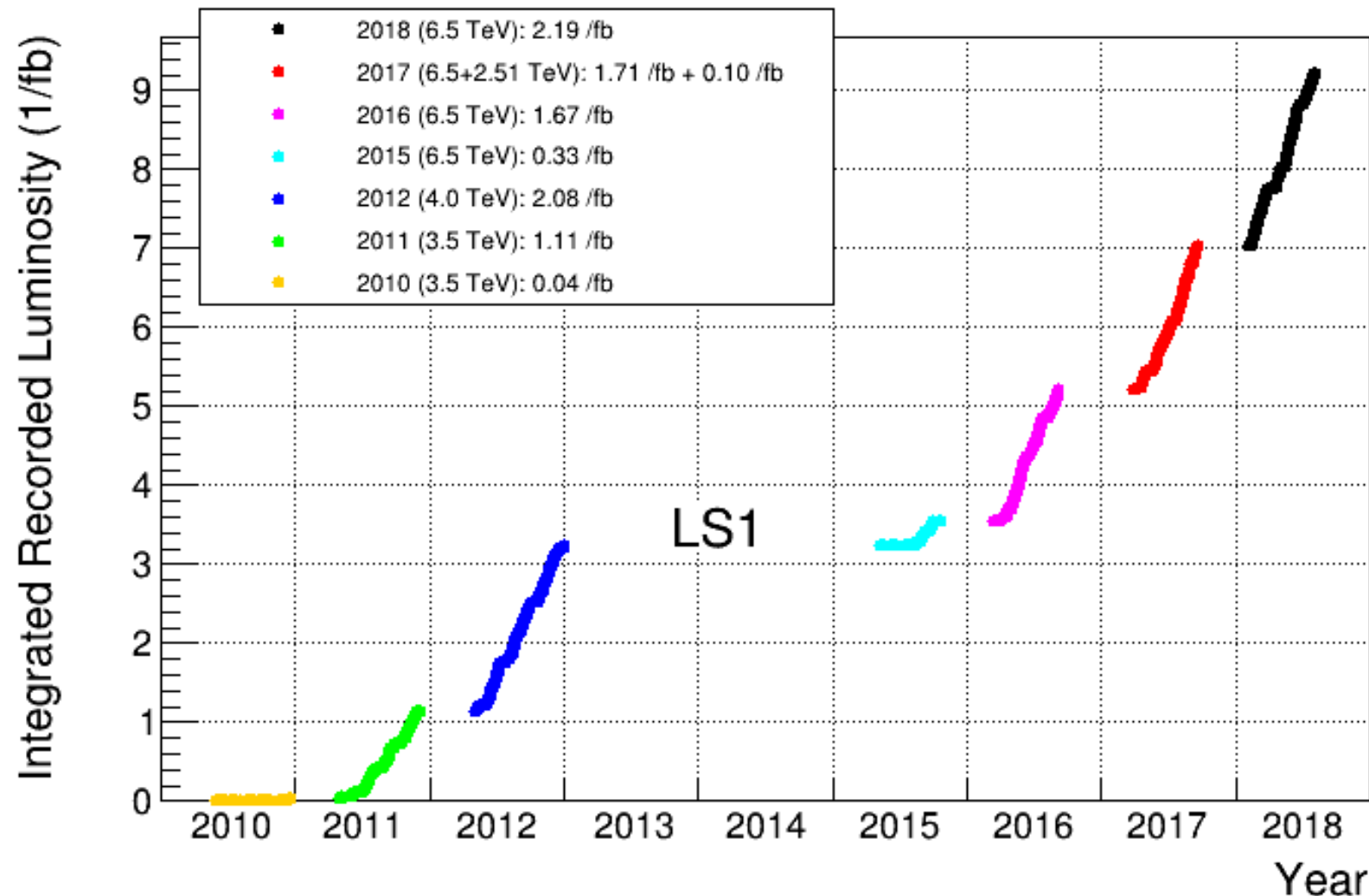


$2 < \eta < 5$, 25% of $b\bar{b}$ pairs
inside LHCb acceptance

Excellent time resolution, IP
resolution, mass resolution,
tracking, PID performance

LHCb collected luminosity

LHCb Cumulative Integrated Recorded Luminosity in pp, 2010-2018



9fb⁻¹ pp collision sample collected in run1+run2

Large b hadrons produced

$$\begin{array}{cccc}
 B^+ & : & B^0 & : & B_s^0 & : & \Lambda_b^0 \\
 (u\bar{b}) & & (d\bar{b}) & & (s\bar{b}) & & (udb) \\
 4 & : & 4 & : & 1 & : & 2
 \end{array}$$

Run1: @7/8 TeV
2010-12, 3fb⁻¹

Run2: @13TeV
2015-18, 6fb⁻¹

Why in B-meson decays

- ❑ Small Q-value, providing excellent mass resolution, allows to search for narrow structures
- ❑ Search for pentaquark and anti-pentaquark states at the same time
- ❑ Sensitive to structures in baryon and anti-baryon system

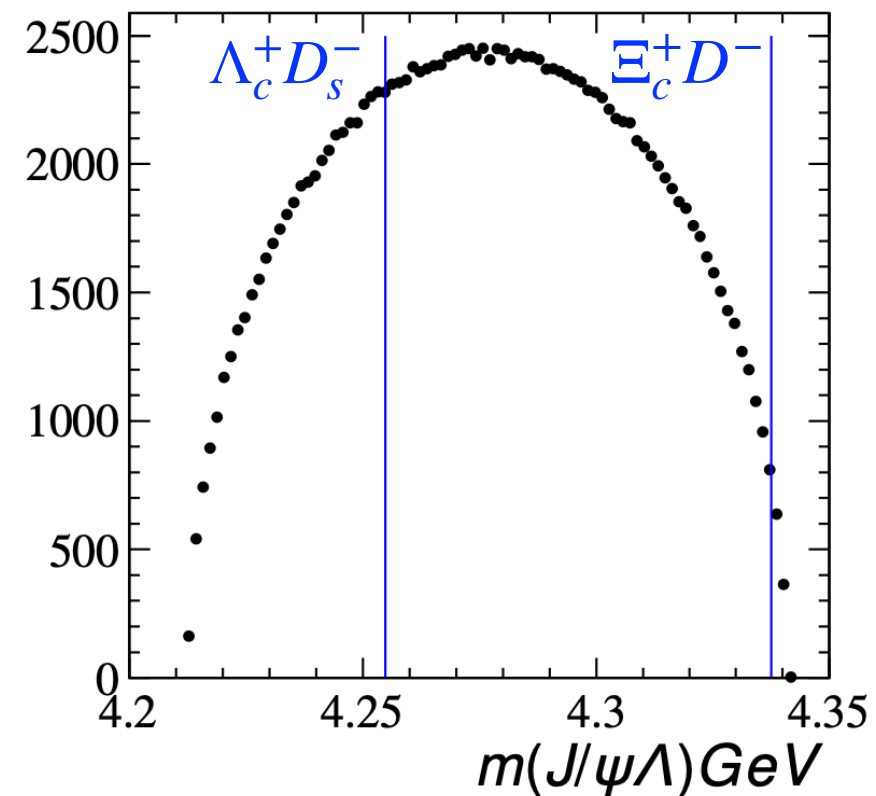
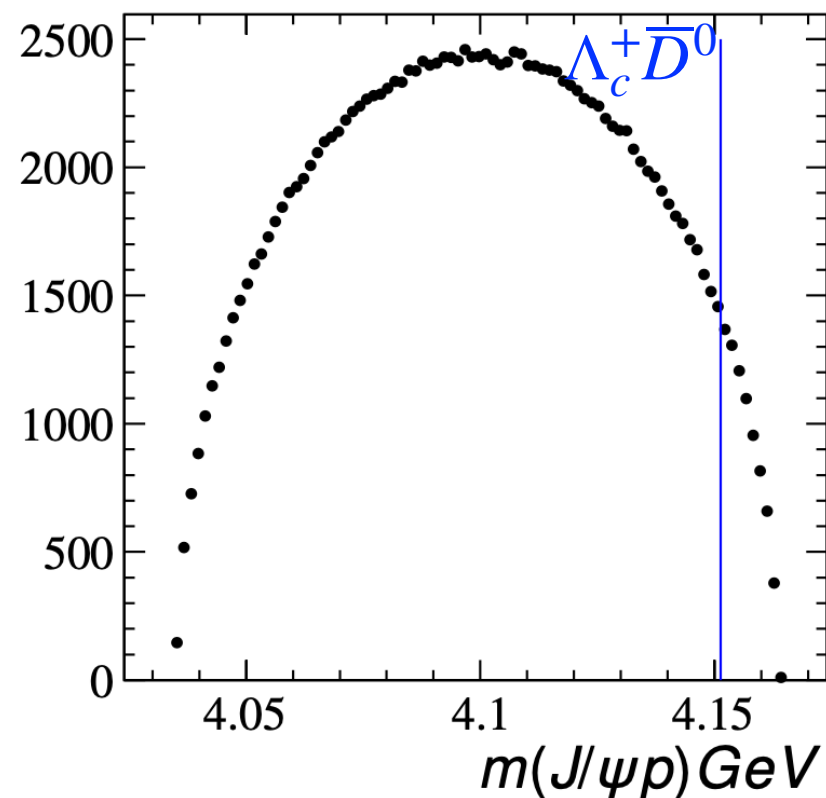
Amplitude analysis in $B^- \rightarrow J/\psi \Lambda \bar{p}$ decays

LHCb-PAPER-2022-031 in preparation

$B^- \rightarrow J/\psi \Lambda \bar{p}$ decays

□ Q-value $\sim 128\text{MeV}$

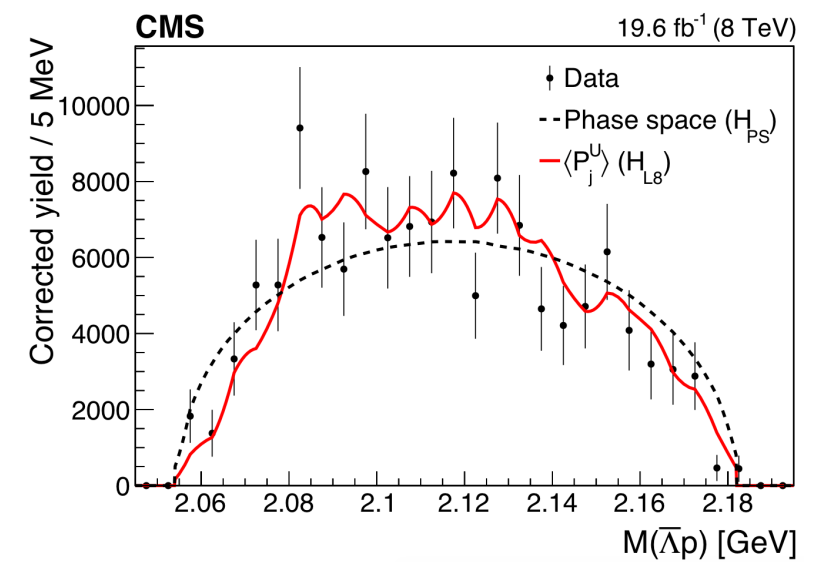
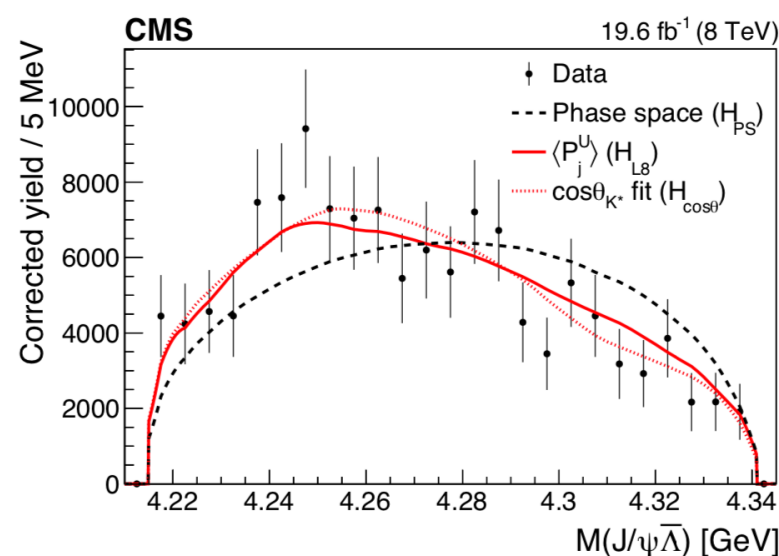
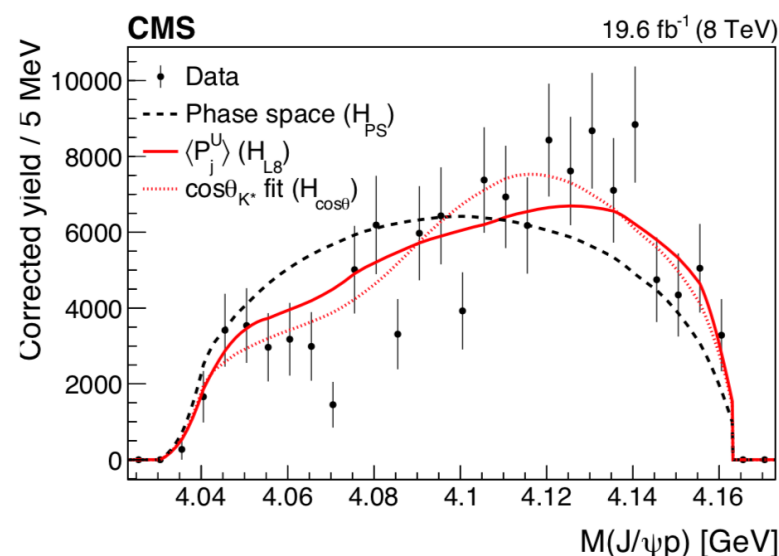
□ Thresholds: $\Lambda_c^+ \bar{D}^0$ in $m(J/\psi \bar{p})$, $\Lambda_c^+ D_s^-$ and $\Xi_c^+ D^-$ in $m(J/\psi \Lambda)$



$B^- \rightarrow J/\psi \Lambda \bar{p}$ @CMS

JHEP12(2019)100

- Limited statistics ~450 signals @8TeV
- Pure phase space hypothesis can not describe data
- $K_{2,3,4}^{*+}$ contributions decrease the incompatibility with data



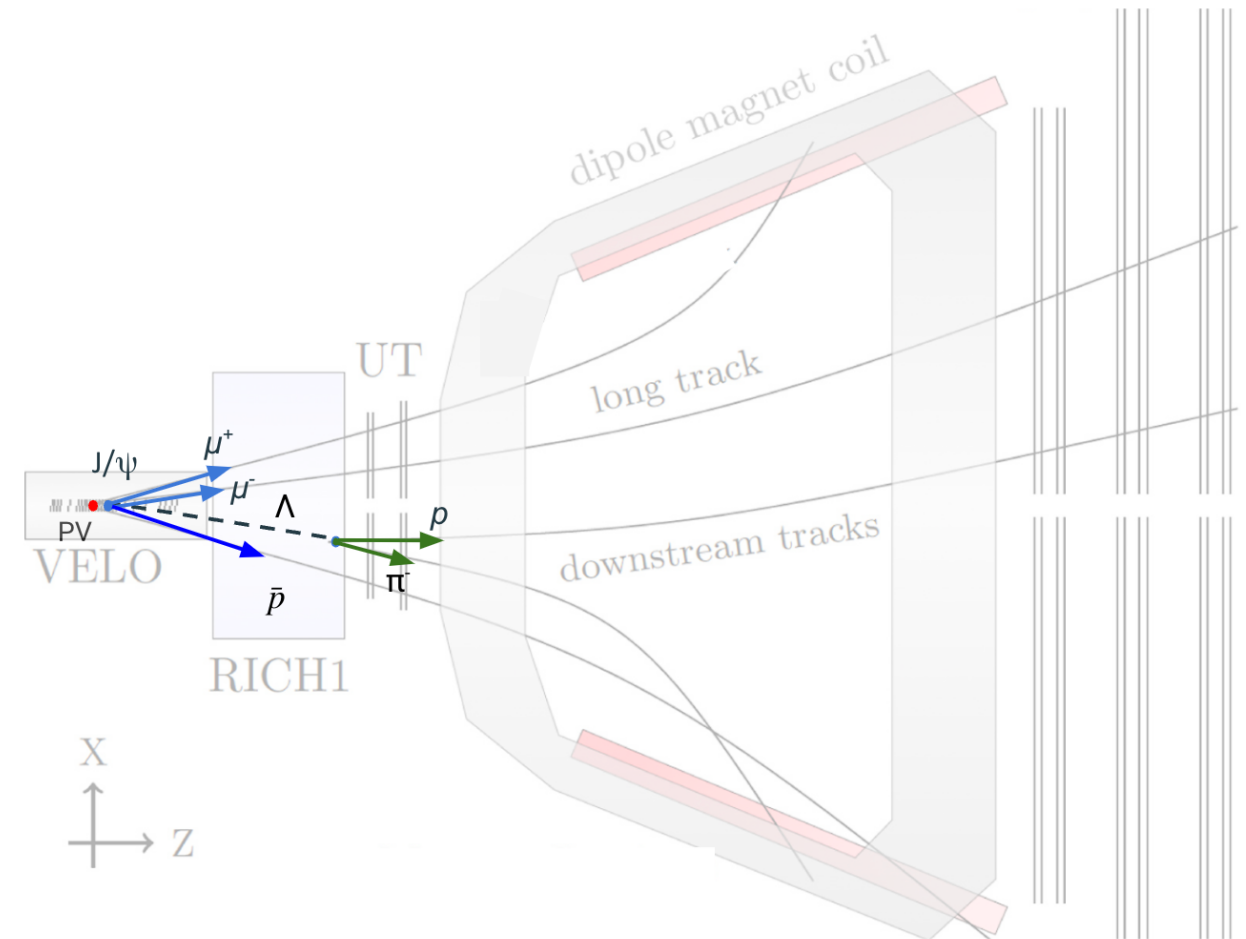
$B^- \rightarrow J/\psi \Lambda \bar{p}$ reconstruction

- Trigger on detached $J/\psi \rightarrow \mu^+ \mu^-$

- $\Lambda \rightarrow p \pi^-$ reconstruction:

 - Λ decay inside VELO

 - Λ decay after VELO



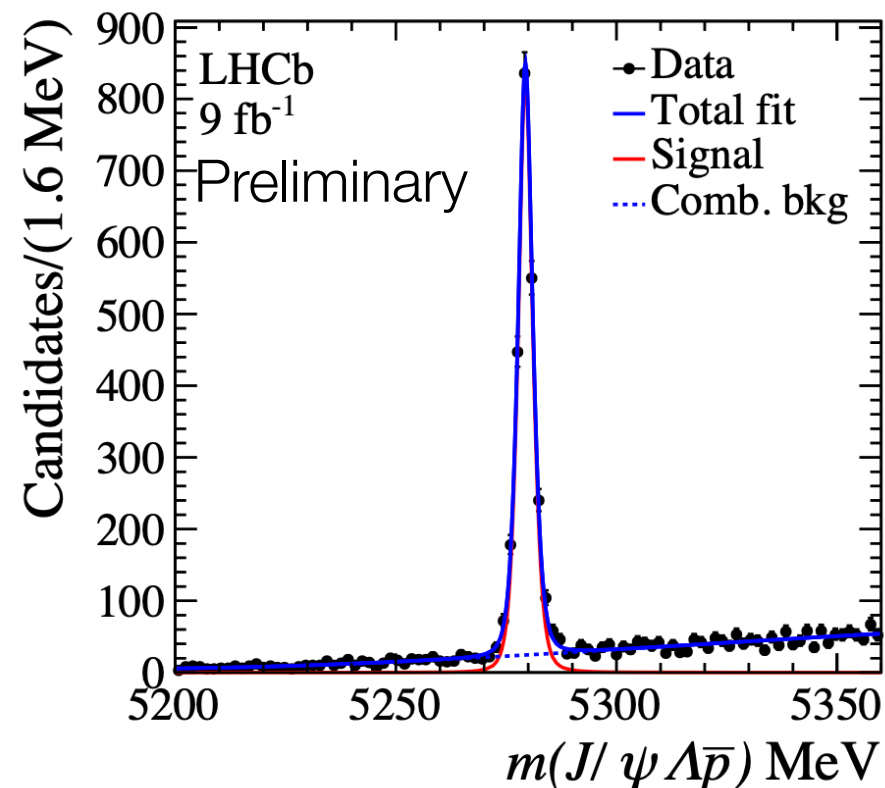
- PID for bachelor anti-proton

- Good-quality vertex for Λ , J/ψ and B^-

- BDT optimization

$B^- \rightarrow J/\psi \Lambda \bar{p}$ candidates

LHCb-PAPER-2022-031 in preparation



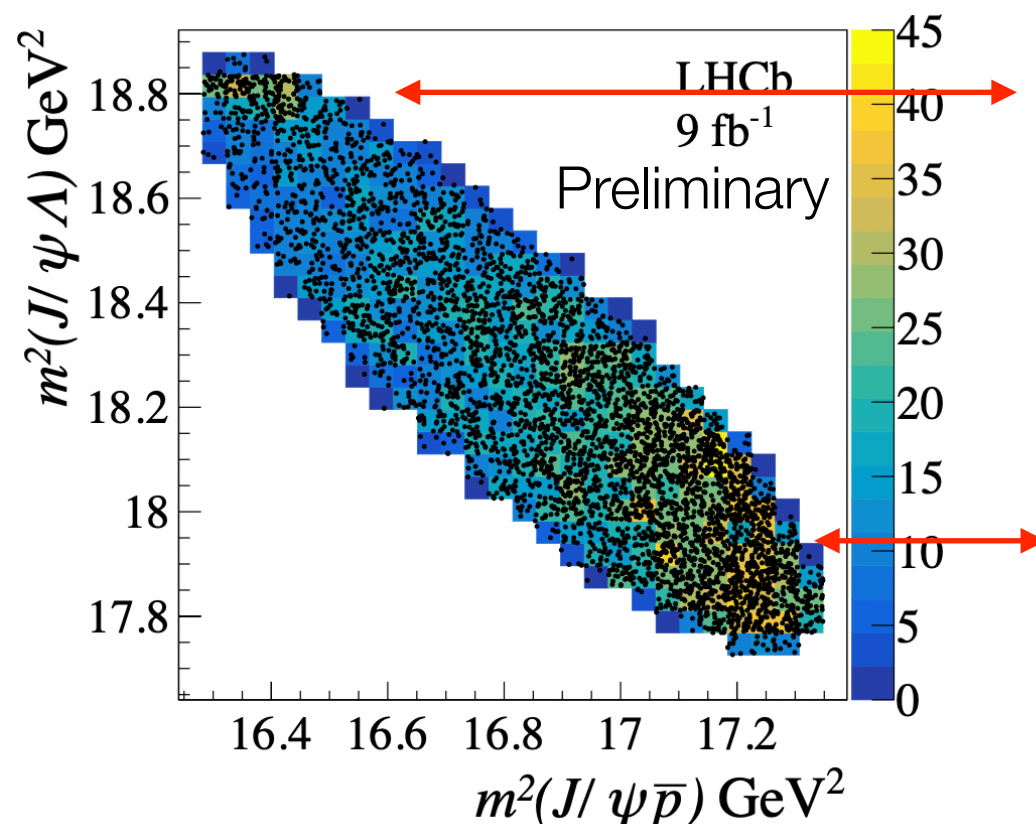
□ LHCb run1 and run2, 9fb⁻¹ dataset

Signal window $\pm 2.5\sigma$

4400 candidates

93% purity

B mass resolution 2MeV



□ Narrow structure in $J/\psi \Lambda$

□ Activity in the $J/\psi p$ system

$B^- \rightarrow J/\psi \Lambda \bar{p}$ amplitude analysis

- Amplitudes developed using helicity formalism, assuming no CPV
- Three decay chains: $J/\psi K^{*-}(\rightarrow \Lambda \bar{p})$, $\Lambda P_c^-(\rightarrow J/\psi \bar{p})$ and $\bar{p} P_{cs}^0(\rightarrow J/\psi \Lambda)$
- Resonant lineshape: Relativistic Breit-Wigner (RBW)
- Non-resonant (NR) lineshape: constant or 2nd order polynomial
- Amplitude fit to data by minimizing

$$-2 \log \mathcal{L}(\vec{\omega}) = -2 \sum_i \log \left[(1 - \beta) \mathcal{P}_{\text{sig}}(m_{\bar{p}\Lambda, i}, \vec{\Omega}_i | \vec{\omega}) + \beta \mathcal{P}_{\text{bkg}}(m_{\bar{p}\Lambda, i}, \vec{\Omega}_i) \right],$$

Model with only K^{*-}

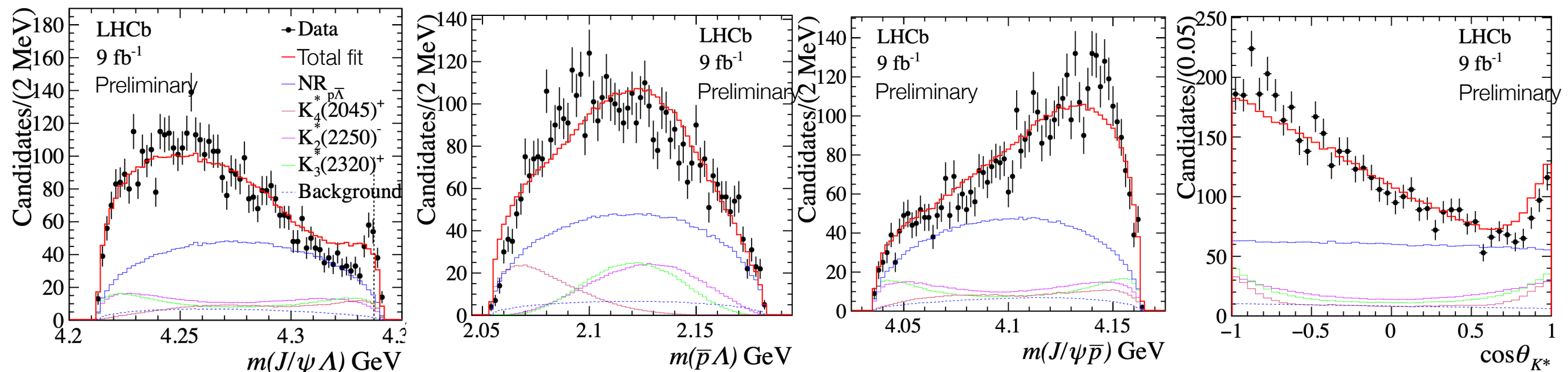
LHCb-PAPER-2022-031 in preparation

□ $K_{2,3,4}^{*-}$ peak out of phsp, and contribution not obvious in $\bar{p}\Lambda$ distribution

Resonance	Mass (MeV)	Natural width (MeV)	J^P
$K_4^*(2045)^+$	2045 ± 9	198 ± 30	4^+
$K_2^*(2250)^+$	2247 ± 17	180 ± 30	2^-
$K_3^*(2320)^+$	2324 ± 24	150 ± 30	3^+

□ Amplitude model with $K_{2,3,4}^{*-} + \text{NR}(\bar{p}\Lambda)$, can not describe data.

$$\chi_{\max}^2 / \text{ndof} = 123/33$$



Baseline and nominal models

Contributions in two amplitude models

Baseline model

$NR(\bar{p}\Lambda)$ constant

$NR(\bar{p}J/\psi)$ 2nd order poly.

$P_{\psi s}^{\Lambda}(J/\psi\Lambda)$ RBW

Nominal model

Fit results of nominal model

LHCb-PAPER-2022-031 in preparation

First neutral pentaquark candidate $P_{\psi S}^{\Lambda}(4338)^0$ near threshold of $\Xi_c^+ D^-$

Significance: $>10\sigma$ wrt baseline model

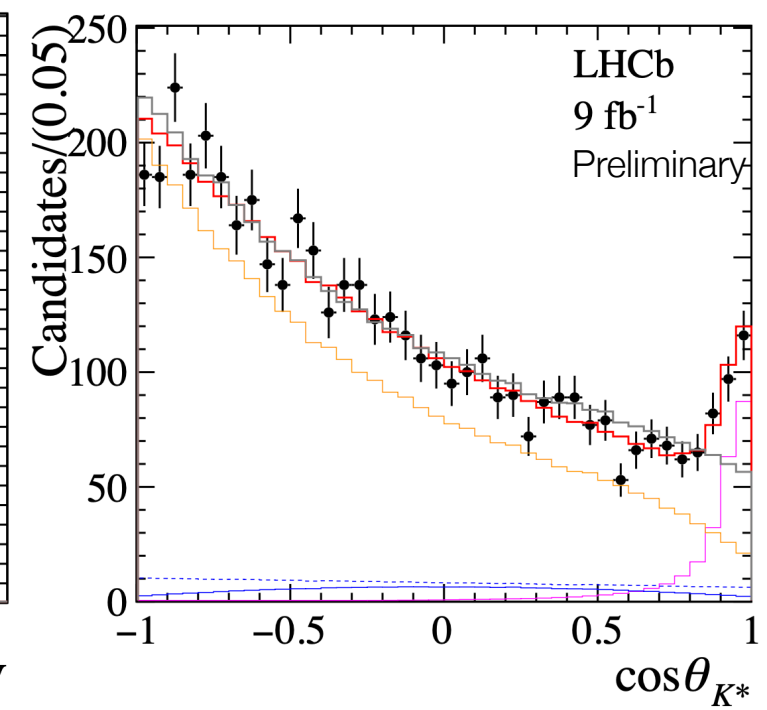
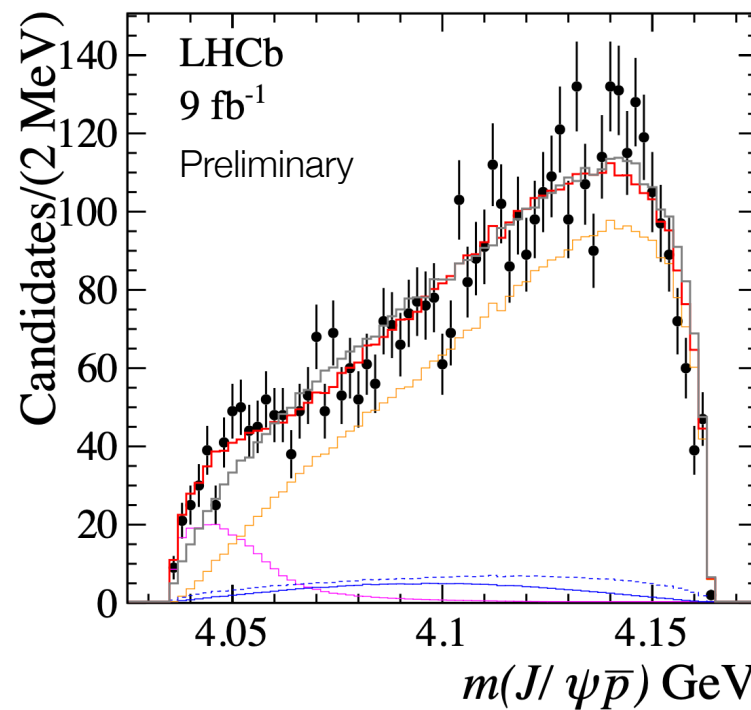
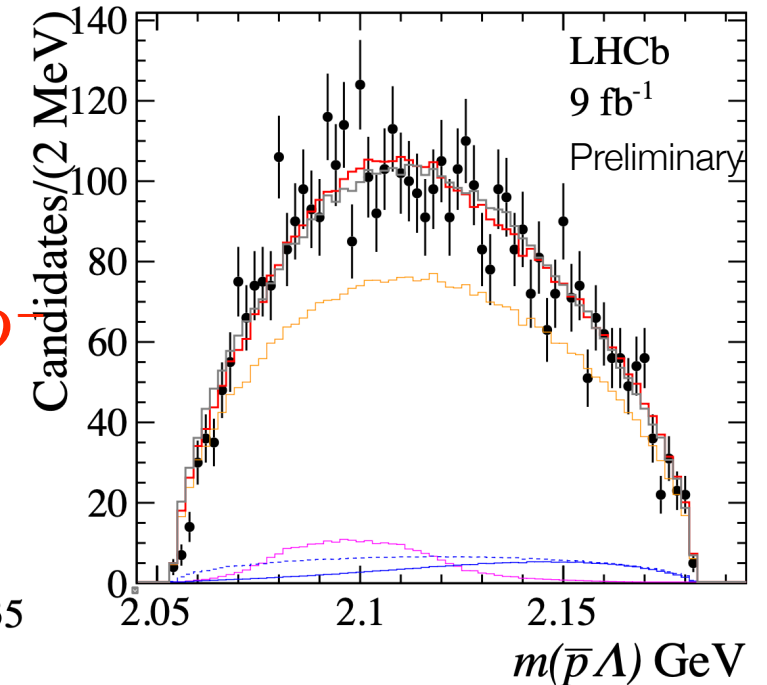
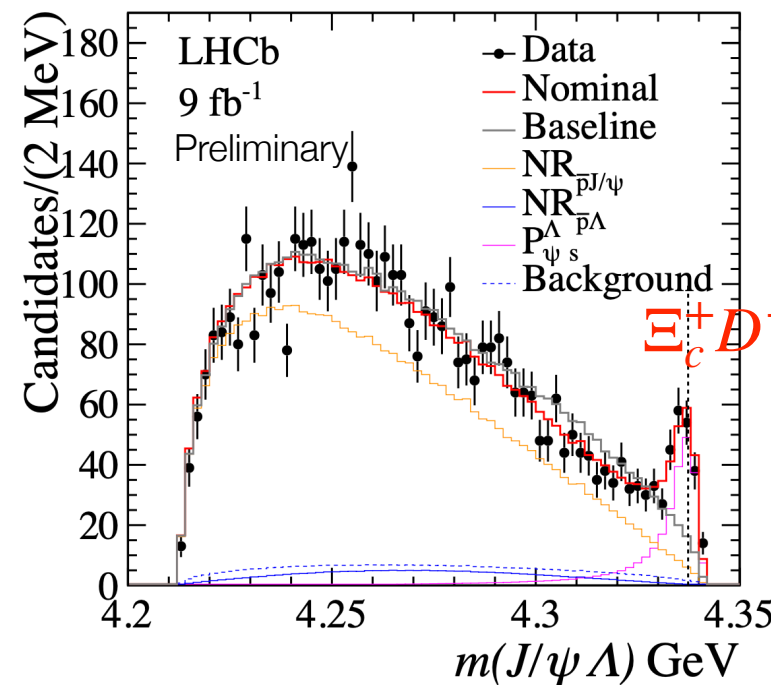
$$m(P_{\psi S}^{\Lambda}) = 4338.3 \pm 0.7 \pm 0.4 \text{ MeV}$$

$$\Gamma(P_{\psi S}^{\Lambda}) = 7.0 \pm 1.2 \pm 1.3 \text{ MeV}$$

Spin 1/2 assigned

$1/2^-$ preferred

$1/2^+$ excluded at 90% C.L.



Structure around $\Lambda_c^+ \bar{D}^0$

LHCb-PAPER-2022-031 in preparation

Amplitude fit with $\text{RBW}(P_{\psi}^N)$
for possible resonance
around $\Lambda_c^+ \bar{D}^0$

$$m(P_{\psi S}^{\Lambda}) = 4338.8 \pm 1.1 \text{ MeV}$$

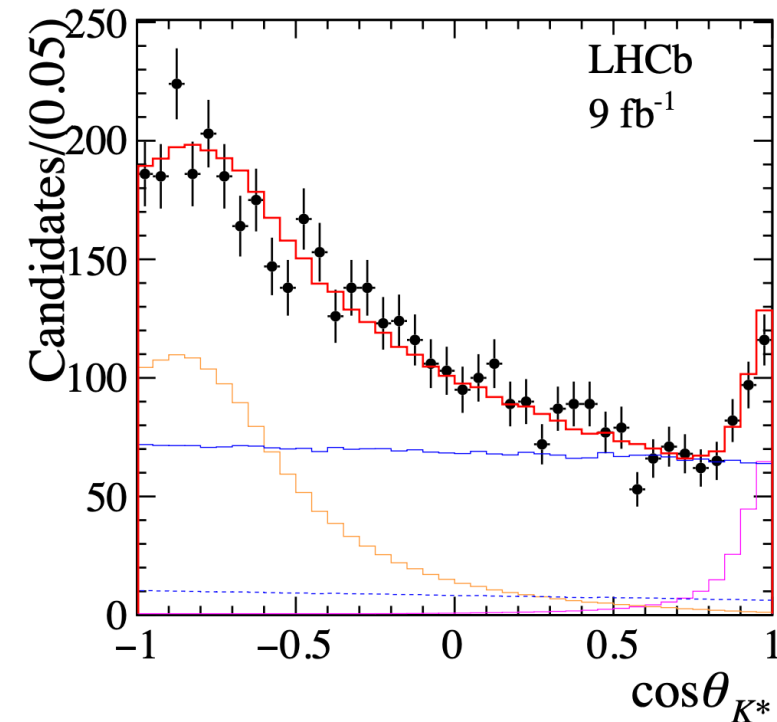
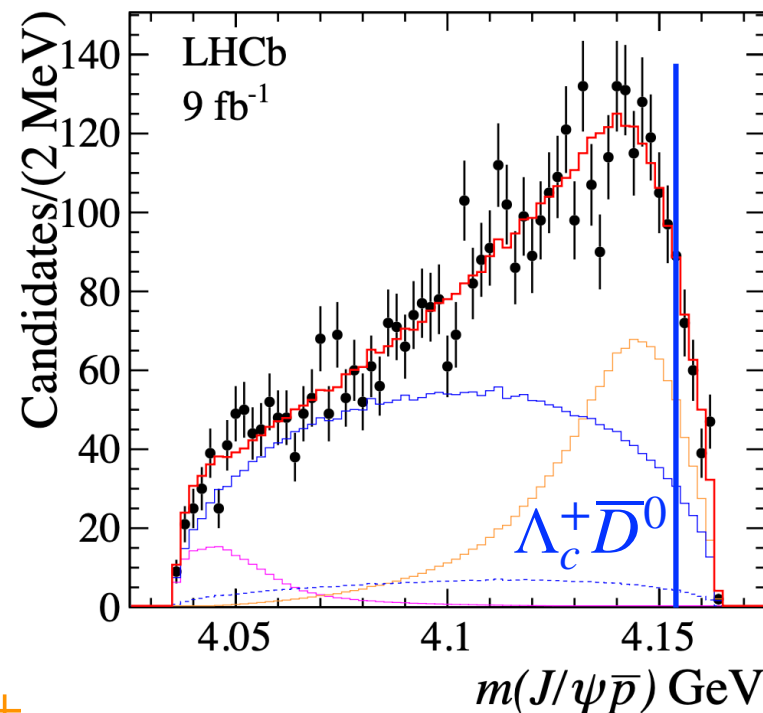
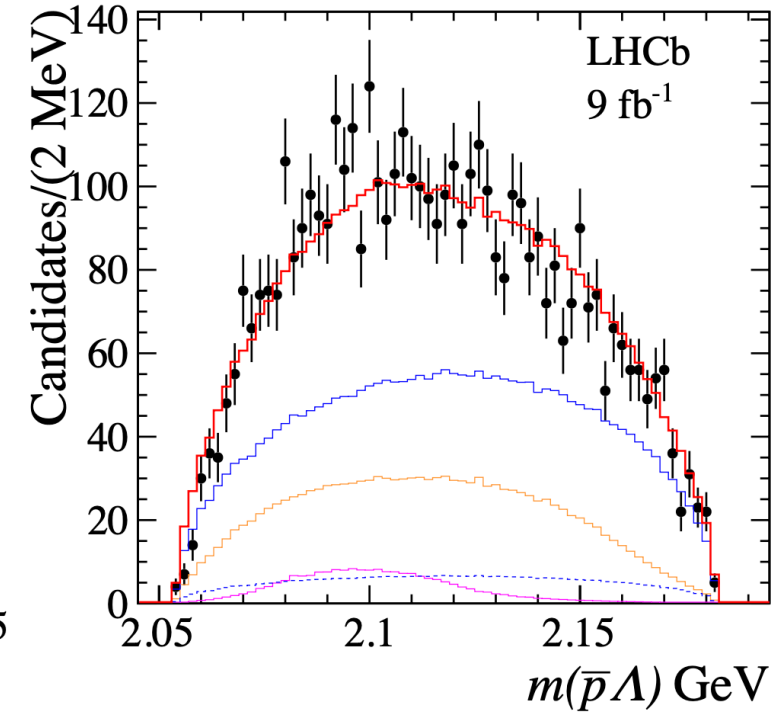
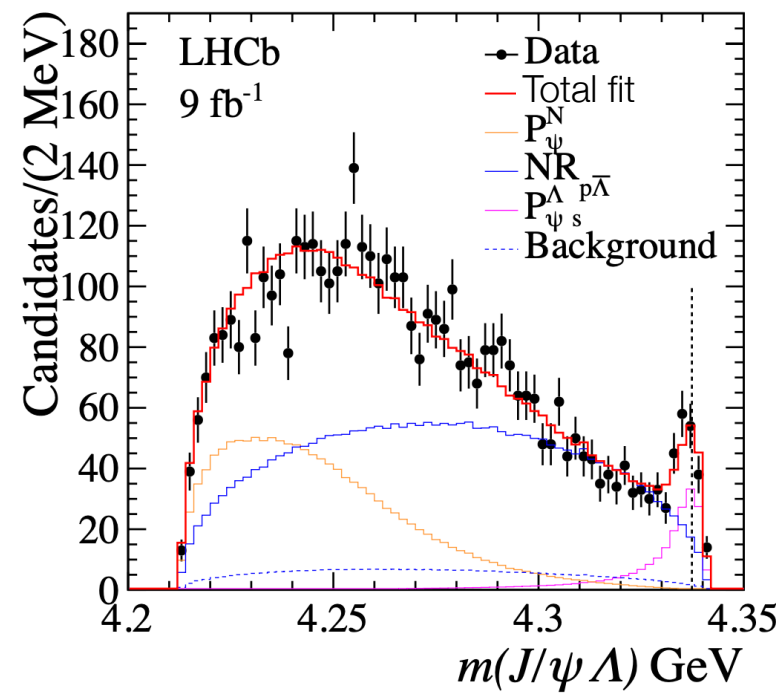
$$\Gamma(P_{\psi S}^{\Lambda}) = 8.4 \pm 1.6 \text{ MeV}$$

$$m(P_{\psi}^N) = 4152.3 \pm 2.0 \text{ MeV}$$

$$\Gamma(P_{\psi}^N) = 41.8 \pm 6.0 \text{ MeV}$$

BUT, $-2\Delta \log L \sim 80$, worse
than nominal fit
2nd polynomial is preferred

=> no evidence for $P_{\psi}^N(4152)^+$



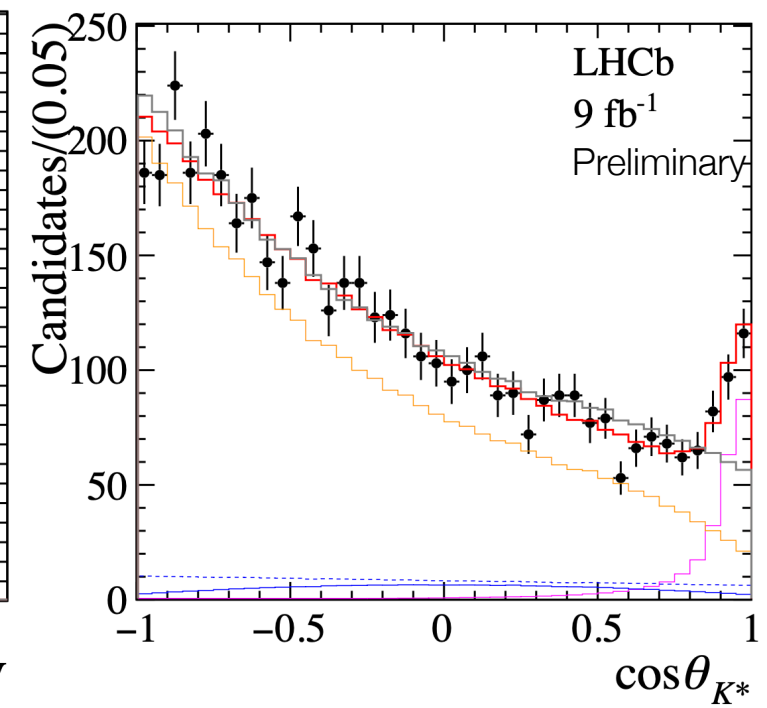
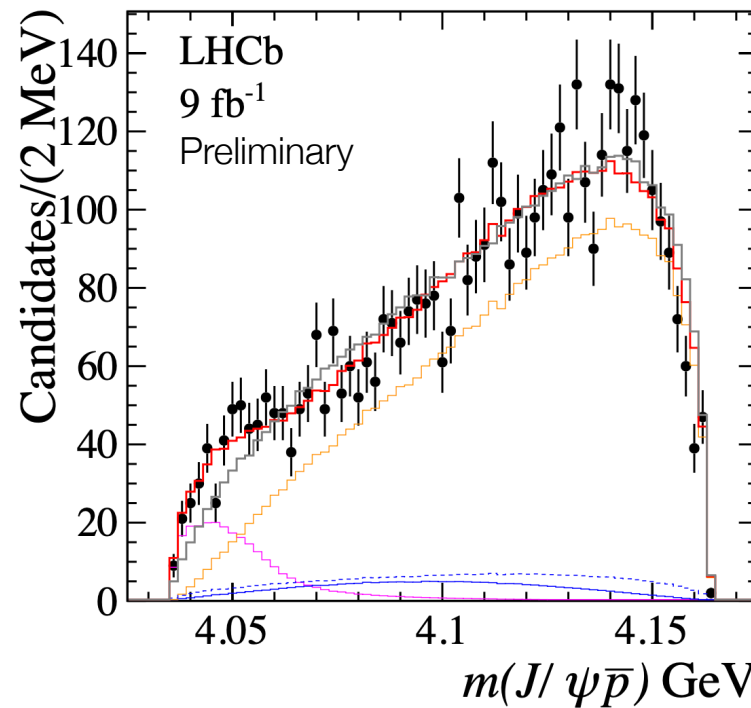
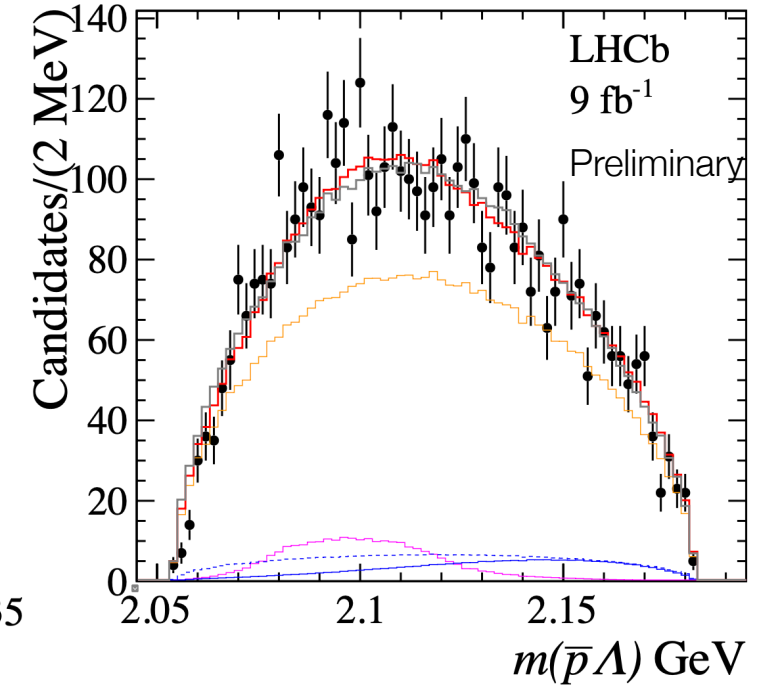
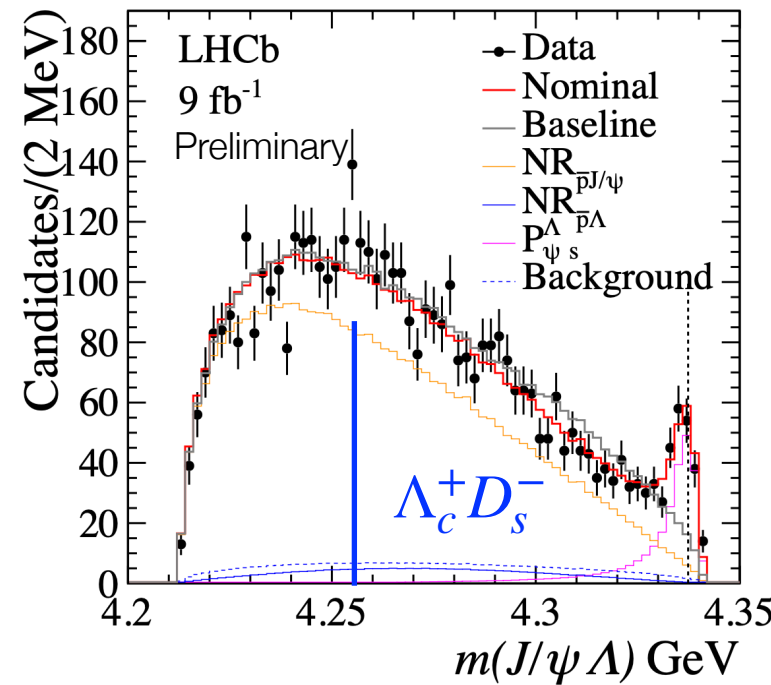
Structure around $\Lambda_c^+ D_s^-$

LHCb-PAPER-2022-031 in preparation

□ Few events excess around threshold of $\Lambda_c^+ D_s^-$

□ A second $\text{RBW}(P_{\psi S}^{\Lambda 0})$ added to the nominal model

□ A p -value of 20% determined from toys
 $\Rightarrow$ no evidence for $P_{\psi S}^{\Lambda}(4255)^0$

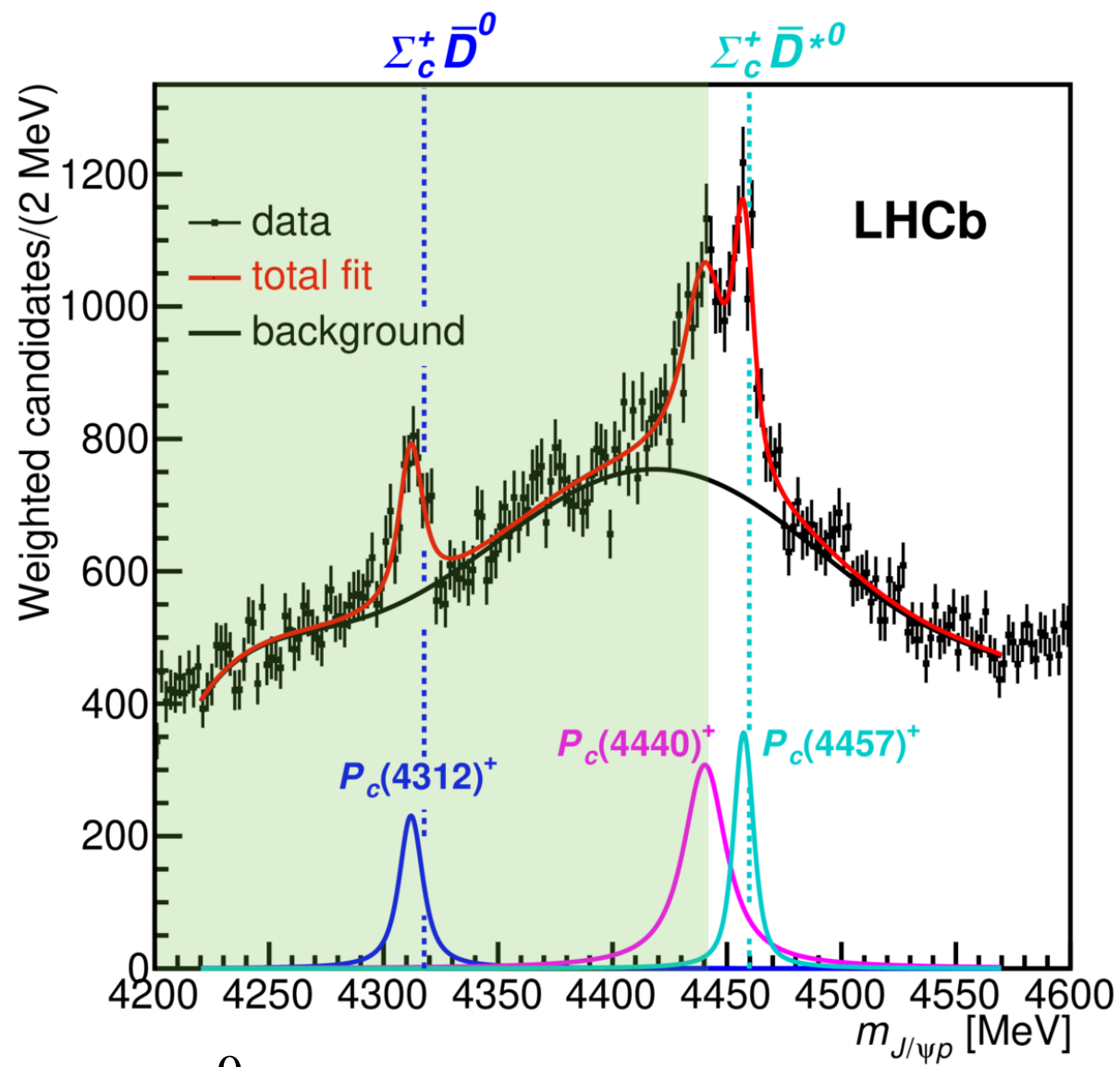


Amplitude analysis in $B_s^0 \rightarrow J/\psi p \bar{p}$ decays

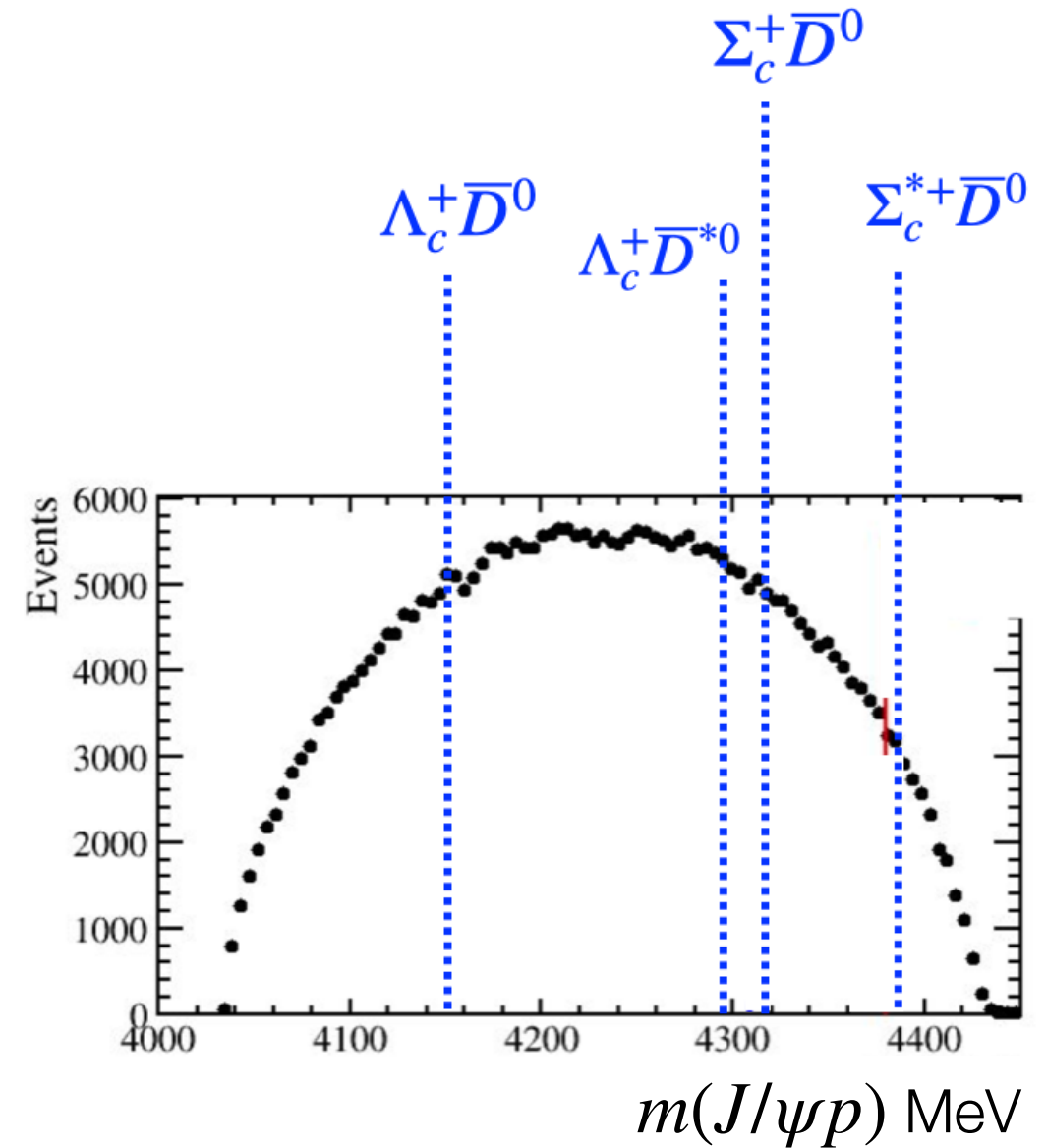
[PRL128\(2022\)062001](#)

$B_s^0 \rightarrow J/\psi p \bar{p}$ decays

- Q-value $\sim 393\text{MeV}$ and four thresholds
- Check $P_\psi^N(4312)^+$, $P_\psi^N(4440)^+$



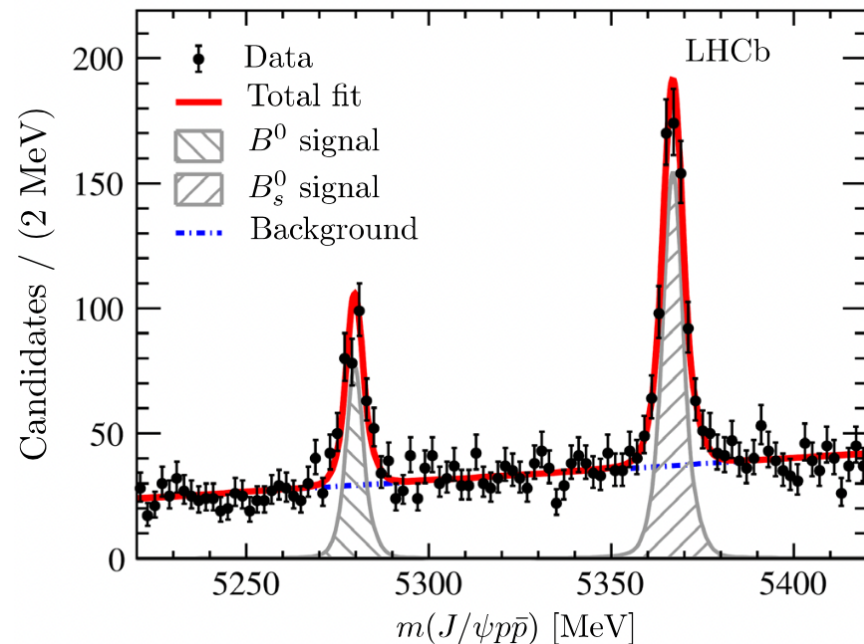
$$\Lambda_b^0 \rightarrow p K^- J/\psi$$



$$B_s^0 \rightarrow J/\psi p \bar{p}$$

$B_s^0 \rightarrow J/\psi p \bar{p}$ candidates

PRL 122(2019)191804



First observed at LHCb in 2019

$$\mathcal{B}(B_s^0 \rightarrow J/\psi p \bar{p})$$

$$= [3.58 \pm 0.19(\text{stat}) \pm 0.39(\text{syst})] \times 10^{-6}$$

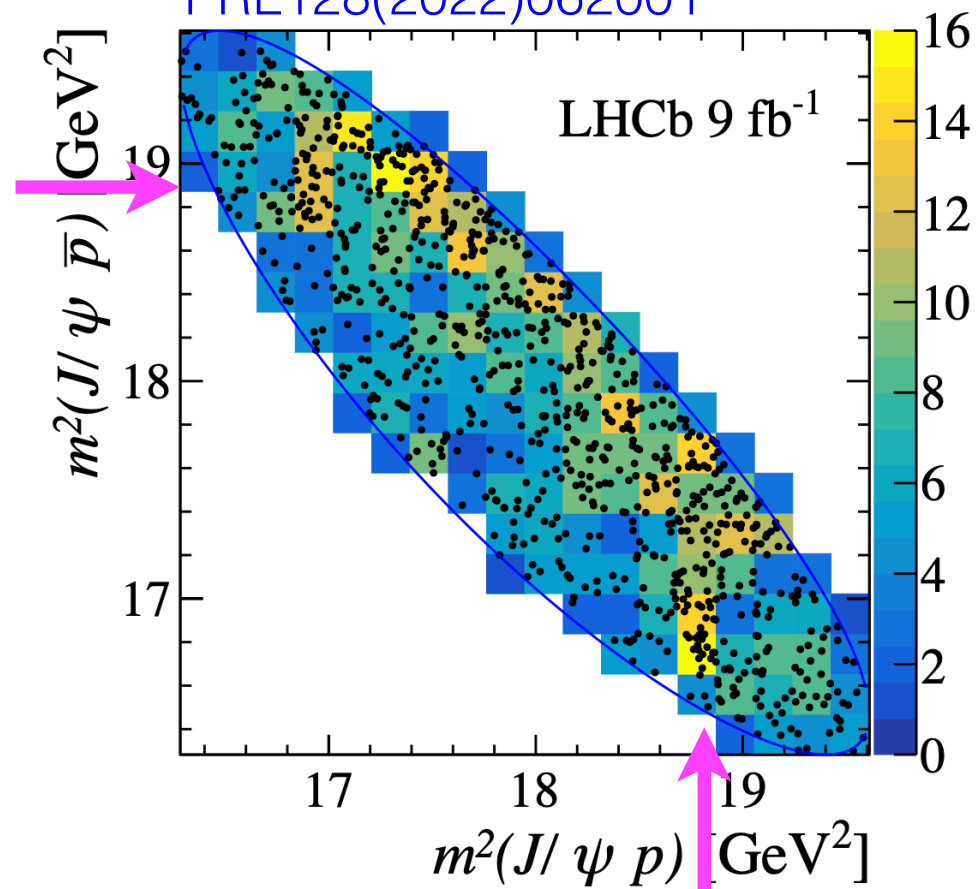
enhanced by 2 orders w.r.t. estimation w/o resonant contributions

800 candidates

85% purity

B mass resolution 3.5 MeV

PRL 128(2022)062001



Hints of structures in Dalitz plot

Full amplitude analysis using 9fb⁻¹ data

Amplitude fit

- Amplitudes developed using helicity formalism

- Assuming no CPV

- For untagged B decay, total amplitude

$$|\bar{\mathcal{M}}|^2 = \frac{1}{2} (|\mathcal{M}(B^0)|^2 + |\mathcal{M}(\bar{B}^0)|^2)$$

- Same fit strategy as $B^- \rightarrow J/\psi \Lambda \bar{p}$ analysis

Amplitude fit

PRL128(2022)062001

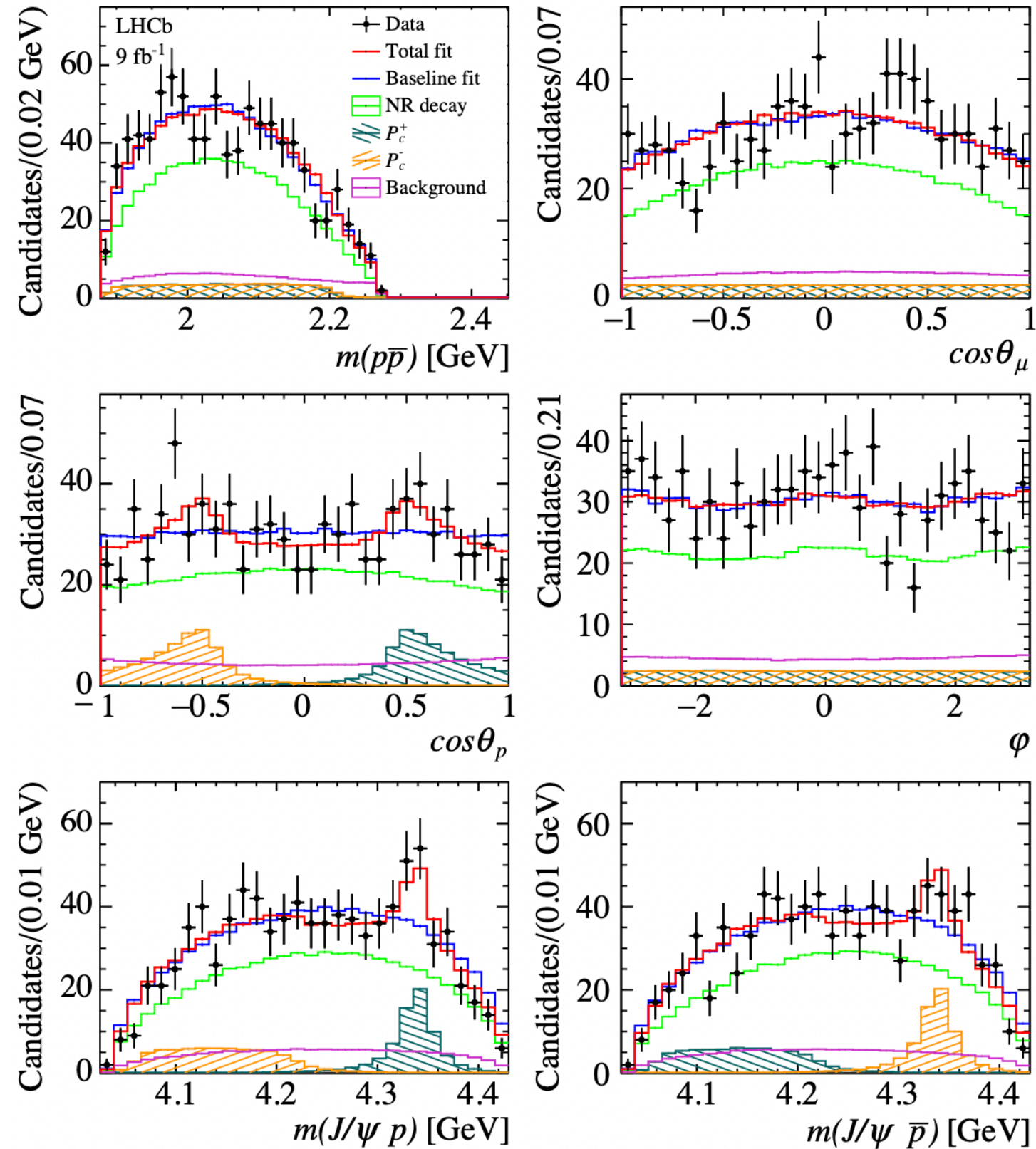
- Two amplitude models

Baseline: $\text{NR}(p\bar{p})$

With $P_{\psi}^{N\pm}$: $\text{NR}(p\bar{p}) + \text{RBW}(P_{\psi}^{N\pm})$

- The same mass, width and couplings for $P_{\psi}^{N\pm}$

- Improvement in mass and helicity distributions



Amplitude fit

PRL128(2022)062001

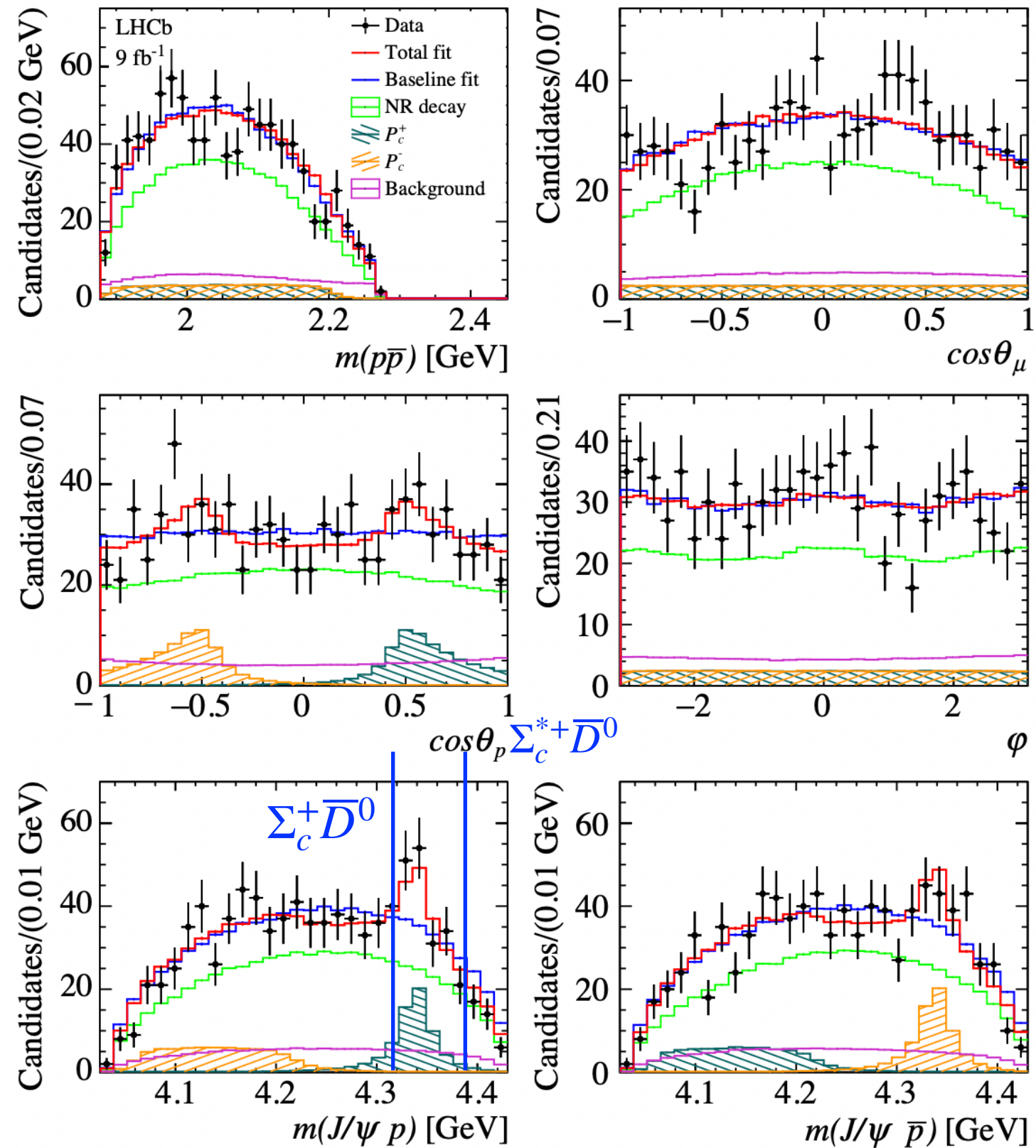
Evidence for a charged pentaquark candidate $P_{\psi}^N(4337)^{\pm}$

Significance: $3.1 \sim 3.7\sigma$ for $J^P(1/2^{\pm}, 3/2^{\pm})$

$$M_{P_c} = 4337_{-4}^{+7}(\text{stat})_{-2}^{+2}(\text{syst}) \text{ MeV}$$

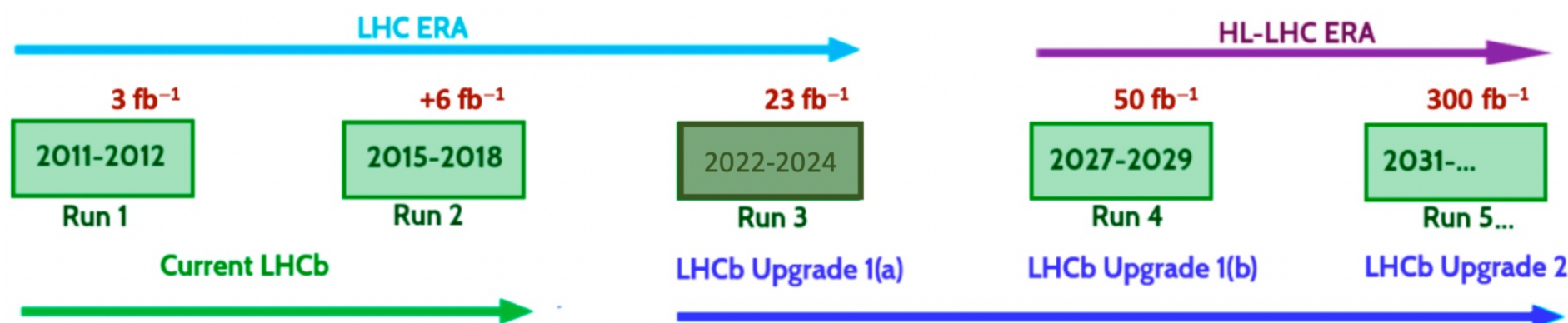
$$\Gamma_{P_c} = 29_{-12}^{+26}(\text{stat})_{-14}^{+14}(\text{syst}) \text{ MeV}$$

No evidence for $P_{\psi}^N(4312)^+$ and $P_{\psi}^N(4440)^+$



Summary

- ❑ Observation of first neutral pentaquark candidate $P_{\psi s}^{\Lambda}(4338)^0$ with strangeness in $B^- \rightarrow J/\psi \Lambda \bar{p}$ decays
narrow state close to $\Xi_c^+ D^-$ threshold
spin 1/2 assigned, negative parity preferred
- ❑ Evidence for a charged pentaquark candidate $P_{\psi}^N(4337)^+$ in $B_s^0 \rightarrow J/\psi p \bar{p}$ decays
- ❑ No evidence for structures around thresholds of $\Lambda_c^+ \bar{D}^0$ and $\Lambda_c^+ D_s^-$
- ❑ Run3 is coming, expect to have more structures in following years



Thank you!

Couplings and lineshapes

- Helicity couplings expressed in LS couplings

$$H_{\lambda_B, \lambda_C}^{A \rightarrow BC} = \sum_L \sum_S \sqrt{\frac{2L+1}{2J_A+1}} B_{L,S} \left(\begin{array}{cc|c} J_B & J_C & S \\ \lambda_B & -\lambda_C & \lambda_B - \lambda_C \end{array} \right) \times \left(\begin{array}{cc|c} L & S & J_A \\ 0 & \lambda_B - \lambda_C & \lambda_B - \lambda_C \end{array} \right)$$

- Helicity couplings of Λ decays are fixed according to decay parameters (PDG)

$$H_{1/2}^{\bar{\Lambda} \rightarrow \bar{p} \pi^+} = (1,0) \quad H_{-1/2}^{\bar{\Lambda} \rightarrow \bar{p} \pi^+} = \sqrt{\frac{1-\alpha_+}{1+\alpha_+}} e^{-i\phi_+} \quad \alpha_+ = -0.758 \pm 0.012$$

$$\phi_+ = +6.5 \pm 3.5^\circ$$

- Line shape of resonant decay

$$\left(\frac{p}{p_0}\right)^L B'_L(p, p_0, d) \times \left(\frac{q}{q_0}\right)^l B'_l(q, q_0, d) \text{BW}(m|m_0, \Gamma_0)$$

$$\text{BW}(m|m_0, \Gamma_0) = \frac{1}{m_0^2 - m^2 - im_0\Gamma(m)}$$

$$\Gamma(m) = \Gamma_0 \left(\frac{q}{q_0}\right)^{2l+1} \left(\frac{m_0}{m}\right) B'_l(q, q_0, d)^2$$

For Non-Resonant (NR), BW replaced by

$$\text{NR} = \begin{cases} 1 & \text{for NR}(\bar{p}\Lambda) \\ c_0 + c_1(m - m_0) + c_2(m - m_0)^2 & \text{for NR}(\bar{p}J/\psi) \end{cases}$$