

#### Recent Belle results on baryons and charmed baryons

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# **Belle experiment**







CMS energy 10.58 GeV Effective CMS 3-5 GeV



High resolution, general purpose  $4\pi$  spectrometer with particle-id

# The last beam abort of KEKB on June 30, 2010





First physics run on June 2, 1999 Last physics run on June 30, 2010  $L_{peak} = 2.1 \times 10^{34} / cm^2 / s$ L > 1ab<sup>-1</sup>

# **Integrated luminosity of B factories**



> 1 ab<sup>-1</sup> On resonance:  $\Upsilon(5S): 121 \text{ fb}^{-1}$  $\Upsilon(4S): 711 \text{ fb}^{-1}$  $\Upsilon(3S): 3 \text{ fb}^{-1}$  $\Upsilon(2S): 25 \text{ fb}^{-1}$  $\Upsilon(1S): 6 \text{ fb}^{-1}$ Off reson./scan:  $\sim 100 \text{ fb}^{-1}$ 

~ 550 fb<sup>-1</sup> On resonance:  $Y(4S): 433 \text{ fb}^{-1}$  $Y(3S): 30 \text{ fb}^{-1}$  $Y(2S): 14 \text{ fb}^{-1}$ Off resonance: ~ 54 fb<sup>-1</sup>

1998/1 2000/1 2002/1 2004/1 2006/1 2008/1 2010/1 2012/1



Belle II will provide a significantly larger data sample (x50 Belle) that will allow to continue the investigation with a much more powerful instrument

# **The Physics Program**



- → a (Super) B-factory (~ $1.1 \ge 10^9 \text{ BB}$  pairs per ab<sup>-1</sup>);
- a (Super) charm factory (~ $1.3 \ge 10^9 \text{ cc}$  pairs per ab<sup>-1</sup>);
- a (Super)  $\tau$  factory (~0.9 x 10<sup>9</sup>  $\tau^+\tau^-$  pairs per ab<sup>-1</sup>);
- → thanks to the Initial State Radiation, we can effectively scan the range [0.5 – 10] GeV and measure the e<sup>+</sup>e<sup>-</sup> → light hadrons cross-section very precisely;
- → finally we can exploit the clean e<sup>+</sup>e<sup>-</sup> environment to probe the existence of exotic hadrons, dark photons/Higgs, light Dark Matter particles, ...

# **Observation of an excited** $\Omega^-$ baryon

 $\mathcal{R} = \frac{\mathcal{B}(\Omega^{*-} \to \Xi^0 K^-)}{\mathcal{B}(\Omega^{*-} \to \Xi^- \overline{K}^0)} = 1.2 \pm 0.3$ 

Data	Mode	Mass $(MeV/c^2)$	Yield	$\Gamma({\rm MeV})$	$\chi^2$ /d.o.f.	$n_{\sigma}$
$\Upsilon(1S, 2S, 3S)$	$\Xi^0 K^-,  \Xi^- K^0_S$	$2012.4\pm0.7$	$242 \pm 48, \ 279 \pm 71$	$6.4^{+2.5}_{-2.0}$	227/230	8.3
	(simultaneous)					
$\Upsilon(1S, 2S, 3S)$	$\Xi^0 K^-$	$2012.6 \pm 0.8$	$239 \pm 53$	$6.1 \pm 2.6$	115/114	6.9
$\Upsilon(1S,2S,3S)$	$\Xi^- K_S^0$	$2012.0\pm1.1$	$286\pm87$	$6.8\pm3.3$	101/114	4.4
Other	$\Xi^0 K^-$	2012.4 (Fixed)	$209 \pm 63$	6.4 (Fixed)	102/116	3.4
Other	$\Xi^- K^0_S$	2012.4 (Fixed)	$153 \pm 89$	6.4 (Fixed)	133/116	1.7



#### PRL 121, 052003 (2018)

- The gap in the spectrum between the ground state and this excited state ( $\sim$ 340 MeV) is smaller than in other  $\Omega^-$  excited states, which is closer to the negative-parity orbital excitations of many other baryons.
- The narrow width observed implies that the quantum number  $J^P = \frac{3}{2}^-$  is preferable.

# Theoretical interpretation for the $\Omega^{*}(2012)$

It is generally accepted that  $\Omega^*(2012)$  is 1P orbital excitation of the ground state  $\Omega$  baryon with the three strange quarks, whose quantum numbers are  $J^P = \frac{3}{2}^{-1}$ .

Notably, the newly observed  $\Omega^*(2012)$  is revealed as a KE(1530) hadronic molecule. [PRD 98, 054009 (2018), PRD 98, 056013 (2018), arXiv:1807.02145, arXiv:1807.06485, arXiv:1807.06485, .....] The  $K_{\Xi\pi}$  three-body component is largely dominant.

# From PRD 98, 056013 (2018)

FIG. 1: The three-body decays of  $\Omega(2012)$  in the  $K \Xi(1530)$  molecular picture.

	$J^{P} = \frac{3}{2}^{-}$ le $\Omega(2012) \ (K \Xi(1530))$				
W	Vidths (MeV)	Branch Ratio(%)			
KΞ	0.4	14.3			
$K\pi\Xi$	2.4	85.7			
Total	2.8	100.0			

# Search for $\Omega(2012) \rightarrow K\Xi(1530) \rightarrow K\pi\Xi$

We use the same data samples to search for  $\Omega(2012) \rightarrow K\Xi(1530) \rightarrow K\pi\Xi$  in the decay of the narrow resonances  $\Upsilon(1S)$ ,  $\Upsilon(2S)$ , and  $\Upsilon(3S)$ .



No clear  $\Omega(2012)$  signals are observed. We give the upper limits on the ratios of the branching fractions a

ratios of the branching fractions at 90% C.L. as below.

$$R_{\Xi^{-}K^{0}}^{\Xi^{-}\pi^{+}K^{-}} = \frac{\mathcal{B}(\Omega \to \Xi(1530)^{0}(\to \Xi^{-}\pi^{+})K^{-})}{\mathcal{B}(\Omega \to \Xi^{-}\overline{K}^{0})} < 9.3\%$$

$$R_{\Xi^{-}K^{0}}^{\Xi^{-}\pi^{0}\overline{K}^{0}} = \frac{\mathcal{B}(\Omega \to \Xi(1530)^{-}(\to \Xi^{-}\pi^{0})\overline{K}^{0})}{\mathcal{B}(\Omega \to \Xi^{-}\overline{K}^{0})} < 81.1\%$$

$$R_{\Xi^{0}K^{-}}^{\Xi^{0}\pi^{-}\overline{K}^{0}} = \frac{\mathcal{B}(\Omega \to \Xi(1530)^{-}(\to \Xi^{0}\pi^{-})\overline{K}^{0})}{\mathcal{B}(\Omega \to \Xi^{0}K^{-})} < 21.3\%$$

$$R_{\Xi^{0}K^{-}}^{\Xi^{0}\pi^{-}} = \frac{\mathcal{B}(\Omega \to \Xi(1530)^{0}(\to \Xi^{0}\pi^{0})K^{-})}{\mathcal{B}(\Omega \to \Xi^{0}K^{-})} < 30.4\%$$

$$R_{\Xi^{0}K^{-}}^{\Xi^{-}\pi^{+}K^{-}} = \frac{\mathcal{B}(\Omega \to \Xi(1530)^{0}(\to \Xi^{-}\pi^{+})K^{-})}{\mathcal{B}(\Omega \to \Xi^{0}K^{-})} < 7.8\%$$

$$R_{\Xi^{-}\overline{K}^{0}}^{\Xi^{0}\pi^{-}\overline{K}^{0}} = \frac{\mathcal{B}(\Omega \to \Xi(1530)^{-}(\to \Xi^{-}\pi^{0})\overline{K}^{0})}{\mathcal{B}(\Omega \to \Xi^{-}\overline{K}^{0})} < 25.6\%$$
S.Jia, \*C.P.Shen et al (Belle)
PRD 100, 032006 (2019)

Search for  $\Omega(2012) \rightarrow K\Xi(1530) \rightarrow K\pi\Xi$ 

A simultaneous fit to all three-body decay modes is performed.



# Evidence for $\Omega_c^0 \to \pi^+ \Omega(2012)^- \to \pi^+ (\overline{K} \Xi)^-$

#### **Motivation:**

- Searching for new production model is very important to understand the nature of Ω(2012)<sup>-</sup>;
- A theoretical study of the  $\Omega(2012)^-$  in the nonleptonic weak decays of  $\Omega_c^0 \to \pi^+ \overline{K} \Xi(1530)(\eta \Omega) \to \pi^+ (\overline{K}\pi \Xi)^-$  and  $(\overline{K}\Xi)^-$  was reported; the authors predicted the clearly  $\Omega(2012)^-$  peak in the  $(\overline{K}\Xi)^-$  invariant mass spectrum of the  $\Omega_c^0 \to \pi^+ (\overline{K}\Xi)^-$ .

[PRD 102, 076009 (2020)]



Evidence for  $\Omega_c^0 \rightarrow \pi^+ \Omega(2012)^- \rightarrow \pi^+ (\overline{K}\Xi)^-$ Phys. Rev. D 104, 052005 (2021)

• To extract the  $\Omega(2012)^-$  signal events from  $\Omega_c^0$  decay, a 2D maximumlikelihood fit is performed to  $M(K^-\Xi^0)/M(K_S^0\Xi^-)$  and  $M(\pi^-\Omega(2012))$ .



• The statistical significances of  $\Omega_c^0 \to \pi^+ \Omega(2012)^- \to \pi^+ K^- \Xi^0$  and  $\Omega_c^0 \to \pi^+ \Omega(2012)^- \to \pi^+ K_S^0 \Xi^-$  decays are 4.0 $\sigma$  and 2.3 $\sigma$ , respectively.

Evidence for  $\Omega_c^0 \rightarrow \pi^+ \Omega(2012)^- \rightarrow \pi^+ (\overline{K}\Xi)^-$ Phys. Rev. D 104, 052005 (2021)

 A 2D un-binned maximum-likelihood simultaneous fit is performed to M((K̄Ξ)<sup>-</sup>) and M(π<sup>+</sup>Ω(2012)<sup>-</sup>) distributions.



 $N_{fit} = 46.6 \pm 12.3$ 

 $\frac{Br(\Omega_c^0 \to \pi^+ \Omega(2012)^-) \times Br(\Omega(2012)^- \to (\overline{K}\Xi)^-)}{Br(\Omega_c^0 \to \pi^+ (\overline{K}\Xi)^-)}$ 

Signal significance:  $4.2\sigma$ (including systematic uncertainties) =  $(6.50 \pm 1.22(\text{stat.}) \pm 0.94(\text{syst.}))\%$ 



However, we realized that (1) the requirement of  $M(\Xi\pi)$  includes large non- $\Xi(1530)$  decay backgrounds; (2) not consider a three-body phase space in  $M(K\pi\Xi)$ , which increases sharply due to the unstable  $\Xi(1530)$  constituent.

The expected M( $\Xi\pi$ ) lineshape and three-body phase space are shown below according to Ref. [PRD 81, 094028 (2010)].



# **The Charmed Baryon Physics**



- The weak decay of charmed baryon has not been understood well.
- Three diagrams contribute in the tree level, but their strengths are not known.
- Ground state charm baryon is a good laboratory for studying strange baryons as decay proceed via  $c \rightarrow s$  transition.
- Belle has collected  $\sim 1 \text{ ab}^{-1} \text{ e}^+\text{e}^-$  data samples (mainly at  $\Upsilon(4S)$ ).
  - $10^9 e^+e^- \rightarrow c\bar{c}$  samples
  - $7.7 \times 10^8 B\overline{B}$  samples
- Huge data sample enable to study various charmed baryons.

# Baryon production at B-factory





Baryons produced via fragmentation

- Charmed baryons rather direct
- Hyperons later stage of fragmentation

#### **Huge statistics**

B is efficiently produced via Y(4s)

Once bottom is produced, it favorably decays into charm.

# Huge statistics, good quality



# Measurements of Branching Fractions of $\Lambda_c^+ \rightarrow p\pi^0$ and $\Lambda_c^+ \rightarrow p\eta$ decays at Belle

#### **Motivation:**

- The weak decay of charmed baryons is very useful for testing many contradictory theoretical models and methods. However, the cognition and exploration of charmed baryon goes pretty slowly.
- The precision of measurement of the decay branching fraction remains poor for many Cabibbo-favored (CF) decays and even worse for some decays dominated by Cabibbosuppressed even though many different experiments like Belle and BESIII have hard work on improving the measurement results of charmed baryons.
- ► In theory, the singly Cabibbo-suppressed (SCS) decays  $\Lambda_c^+ \to p\pi^0$  and  $\Lambda_c^+ \to p\eta$ proceed dominantly through internal W-emission and W-exchange. The measurement of these two decay branching fractions may **be interesting to study the underlying dynamic of charmed baryon decays.**
- ► In experiment, BESIII report the branching fractions of these two SCS decays, which are  $B(\Lambda_c^+ \rightarrow p\pi^0) < 2.7 \times 10^{-4}$  at 90% confidence level and  $B(\Lambda_c^+ \rightarrow p\eta) = (1.24 \pm 0.30) \times 10^{-3}$ .
- ▶ In this analysis, we utilize the much higher statistic sample of  $\Lambda_c^+$  collected by Belle detector to improve the measurement precision.

#### Measurement of $\Lambda_c^+ o p K^- \pi^+$ decay

PRD103, 072004 (2021) A method of branching ratio with respect to CF decay  $\Lambda_c^+ \rightarrow pK^-\pi^+$  (reference mode) is applied to measure the branching fractions of two SCS decays.

 $\frac{B(SCS)}{B(CF)} = \frac{N^{obs}(SCS)}{\epsilon^{MC}(SCS)} \times \frac{\epsilon^{MC}(CF)}{N^{obs}(CF)}$ 

#### Signal efficiency estimation: Dalitz method.



Left: Dalitz plot from data; Right: Dalitz plot of efficiency from signal MC.

$$\varepsilon = \sum s_i / \sum_j (s_j / \varepsilon_j) = (\mathbf{14.06} \pm \mathbf{0.01}) \%.$$



#### Measurement of $\Lambda_c^+ \rightarrow p\pi^0 (\rightarrow \gamma\gamma)$ decay PRD103, 072004 (2021)

- The efficiency estimated from signal MC sample is  $(8.891 \pm 0.030)\%$ .
- There is no obvious signal excess in  $M(p\pi^0)$  from data. We set an upper limit on branching fraction of  $B(\Lambda_c^+ \rightarrow p\pi^0) < 8 \times 10^{-5}$  at 90% C.L., reducing the value to more than half of the current best upper limit of 2.7 × 10<sup>-4</sup>.



Left: fit to the invariant mass distribution of  $p\pi^0$  with a fixed signal yield of **1269**. Right: The likelihood distribution changing with the branching fraction with the systematic uncertainty involved.

#### Measurement of $\Lambda_c^+ \rightarrow p\eta (\rightarrow \gamma \gamma)$ decay PRD103, 072004 (2021)

• The efficiency estimated from signal MC sample is  $(8.279 \pm 0.030)\%$ .



Gaussian + CB for signal. Second-order polynomial for background.

Yield: **7734**  $\pm$  **263**  $\chi^2/ndf$ =1.23

• A significant  $\Lambda_c^+$  signal is observed in M( $p\eta$ ) distribution from data. The branching fraction is  $B(\Lambda_c^+ \rightarrow p\eta) = (1.42 \pm 0.05 \pm 0.11) \times 10^{-3}$ , which is consistent with the latest BESIII measured result of  $(1.24 \pm 0.30) \times 10^{-3}$  with much improved precision.

• The measured  $B(\Lambda_c^+ \to p\eta)$  is at least an order of magnitude larger than  $B(\Lambda_c^+ \to p\pi^0)$ , which is consistent with the theoretical prediction of an internal W-emission mechanism involving an s quark in  $\Lambda_c^+ \to p\eta$ .

#### Precise Measurement of $\Lambda_c^+ o p\omega$

PRD 104, 072008 (2021)

• A method to measure the branching fractions of  $\Lambda_c^+ \rightarrow p\omega(\rightarrow \pi^+\pi^-\pi^0)$  decays is:

 $\frac{B(\text{Decay mode})}{B(\Lambda_c^+ \to pK^-\pi^+)} = \frac{y \text{ (Decay mode)}}{B_{\text{PDG}} \times y(\Lambda_c^+ \to pK^-\pi^+)} \text{ (y is the efficiency-corrected yield).}$ 



Simultaneous fit for events in the  $\omega$  signal region and the normalized  $\omega$  sideband regions.

$$\mathcal{B}(\Lambda_c^+ \to p\omega) = (8.27 \pm 0.75 \pm 0.62 \pm 0.42) \times 10^{-4}$$

#### Most precise result to data

- This result is **consistent with** the LHCb result  $(9.4 \pm 3.9) \times 10^{-4}$ .
- This result agrees with the of theoretical predictions within uncertainties based on the SU(3)F flavor symmetry.

#### First Observation of $\Lambda_c^+ o p\eta'$

• A method to measure the branching fractions of  $\Lambda_c^+ \to p\eta' (\to \pi^+ \pi^- \eta)$  decays is:

 $\frac{B(\text{Decay mode})}{B(\Lambda_c^+ \to pK^-\pi^+)} = \frac{y \text{ (Decay mode)}}{B_{\text{PDG}} \times y(\Lambda_c^+ \to pK^-\pi^+)} \text{ (y is the efficiency-corrected yield).}$ 



 $\mathcal{B}(\Lambda_c^+ \to p\eta') = (4.73 \pm 0.82 \pm 0.43 \pm 0.24) \times 10^{-4}$ 

First observation of  $\Lambda_c^+$  in  $\Lambda_c^+ \rightarrow p\eta'$ 

• Our result is **consistent with** the most theoretical calculations based on  $SU(3)_F$  symmetry.

NEV

### Measurements of absolute Brs of $\Xi_c^0$

Summary of the measured branching fractions and the ratios of  $\Xi_c^0$  decays

Y.B.Li, C.P.Shen et al (Belle) PRL122, 082001 (2019)

BF	Result	Theory	PDG
$\mathcal{B}(B^- \to \overline{\Lambda}_c^- \Xi_c^0)$	$(9.51 \pm 2.10 \pm 0.88) \times 10^{-4}$	$\sim 10^{-3}$	
$\mathcal{B}(B^- \to \overline{\Lambda}_c^- \Xi_c^0) \mathcal{B}(\Xi_c^0 \to \Xi^- \pi^+)$	$(1.71\pm 0.28\pm 0.15)\times 10^{-5}$		$(2.4\pm 0.9)\times 10^{-5}$
$\mathcal{B}(B^- \to \overline{\Lambda}_c^- \Xi_c^0) \mathcal{B}(\Xi_c^0 \to \Lambda \mathrm{K}^- \pi^+)$	$(1.11\pm 0.26\pm 0.10)\times 10^{-5}$		$(2.1\pm 0.9)\times 10^{-5}$
$\mathcal{B}(B^- \to \overline{\Lambda}_c^- \Xi_c^0) \mathcal{B}(\Xi_c^0 \to p K^- K^- \pi^+)$	$(5.47 \pm 1.78 \pm 0.57) \times 10^{-6}$		
$\mathcal{B}(\Xi_c^0 \to \Xi^- \pi^+)$	$(1.80 \pm 0.50 \pm 0.14)\%$	1.12% or 0.74%	
$\mathcal{B}(\Xi_c^0  o \Lambda \mathrm{K}^- \pi^+)$	$(1.17\pm 0.37\pm 0.09)\%$		
$\mathcal{B}(\Xi_c^0  o pK^-K^-\pi^+)$	$(0.58\pm 0.23\pm 0.05)\%$		
$\mathcal{B}(\Xi_c^0 \to \Lambda K^- \pi^+) / \mathcal{B}(\Xi_c^0 \to \Xi^- \pi^+)$	$0.65\pm 0.18\pm 0.04$		$\textbf{1.07} \pm \textbf{0.14}$
$\mathcal{B}(\Xi_c^0  o pK^-K^-\pi^+)/\mathcal{B}(\Xi_c^0  o \Xi^-\pi^+)$	$0.32 \pm 0.12 \pm 0.07$		$0.34\pm0.04$

- We have performed an analysis of  $B^- \to \overline{\Lambda}_c^- \Xi_c^0$  inclusively and exclusively
- First model-independent measurement of absolute Brs of E<sup>0</sup><sub>c</sub> decays
- The branching fraction  $B(B^- \rightarrow \overline{\Lambda}_c^- \Xi_c^0)$  is measured for the first time
- The  $B(\Xi_c^0 \to \Xi^- \pi^+)$  can be used to determine the BR of other  $\Xi_c^0$  decays.

# **Measurement of** $\Xi_c^+$ **absolute BRs**

Y. B. Li. C. P. Shen et al (Belle) PRD 100, 031101 (2019)

BF	Result	Theory	PDG
$\mathcal{B}(\overline{B}{}^0 \to \overline{\Lambda}_c^- \Xi_c^+)$	$(1.16\pm0.42\pm0.15)\times10^{-3}$	$\sim 10^{-3}$	
$\mathcal{B}(\overline{B}{}^0 \to \overline{\Lambda}{}^c \Xi^+_c))\mathcal{B}(\Xi^+_c \to \Xi^- \pi^+ \pi^+)$	$(3.32\pm 0.74\pm 0.33)\times 10^{-5}$		$(1.8 \pm 1.8)  imes 10^{-5}$
$\mathcal{B}(\overline{B}{}^0 \to \overline{\Lambda}{}^c \Xi^+_c) \mathcal{B}(\Xi^+_c \to pK^-\pi^+)$	$(5.27 \pm 1.51 \pm 0.69) \times 10^{-5}$		
$\mathcal{B}(\Xi_c^+ \to \Xi^- \pi^+ \pi^+)$	$({\bf 2.86 \pm 1.21 \pm 0.38})\%$	$(1.47 \pm 0.84)\%$	
$\mathcal{B}(\Xi_c^+ \to pK^-\pi^+)$	$(0.45\pm 0.21\pm 0.07)\%$	$(2.2 \pm 0.8)\%$	
$\mathcal{B}(\Xi_c^+ \to pK^-\pi^+)/\mathcal{B}(\Xi_c^+ \to \Xi^-\pi^+\pi^+)$	$0.16\pm 0.06\pm 0.02$		$0.21\pm0.04$

- First model –independent  $\mathcal{B}(\overline{B}^0 \to \overline{\Lambda}_c^- \Xi_c^+)$  measurement
- $\mathcal{B}(\Xi_c^+ \to \Xi^- \pi^+ \pi^+)$  can be used to determine the BR of other  $\Xi_c^+$  decay

#### **Measurements of Brs and asymmetry parameters of**

 $\Xi_c^0 \rightarrow \Lambda \overline{K}^{*0}$ ,  $\Xi_c^0 \rightarrow \Sigma^0 \overline{K}^{*0}$ , and  $\Xi_c^0 \rightarrow \Sigma^+ K^{*-}_{JHEP \ 06 \ (2021) \ 160}$ 

- There are some difficulties for the theoretical study in the non-leptonic decays of charmed baryons due to the failure of the factorization approach.
- Branching fraction measurements help to distinguish different theoretical models.
- □ The asymmetry parameters of  $\Xi_c^0$  are still not well measured, which is important to test parity violation in charmed-baryon sectors.

Decay branching fractions (%) and asymmetry parameters of the Cabibbo favored  $B_c \rightarrow B_n + V$  decays in QCD and SU(3)<sub>F</sub> approach.

Br	Branching fractions KK [1]		Zen [2] H		HYZ	[3]	GLT [4]		
	$\Xi_c^0\to\Lambda^0\overline{K}{}^{*0}$	1.55		1.15		0.46±0.21		1.37±0.26	
	$\Xi_c^0 \to \Sigma^0 \overline{K}^{*0}$ 0.8		0.85	0.85 0.77		0.27±0.22		0.42±0.23	
	$\Xi_c^0\to \Sigma^+ K^{*-}$		0.54	0.37		0.93±0.29		0.24±0.17	
Asymmetry parameters		KK [1]		Zen [2]		GLT [4]			
	$\Xi_c^0\to\Lambda^0\overline{K}{}^{*0}$		0.58		+0.49		-0.67±0.24		
	$\Xi_c^0\to \Sigma^0\overline{K}{}^{*0}$		-0.87		+0.25		-0.42±0.62		
	$\Xi_c^0\to \Sigma^+ K^{*-}$		-0.60		+0.51		-0	$-0.76^{+0.64}_{-0.24}$	
L] Z. Phys. C 55, 659 (1992) [2] Phys. Rev. D 50, 5787 (1994) [3] Phys. Lett. B 792, 35 (2019)									

[4] Phys. Rev. D 101, 053002 (2020)



#### Asymmetry parameter extractions

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For  $\Xi_c^0 \to \Lambda^0 \overline{K}^{*0}$ ,  $\Xi_c^0 \to \Sigma^0 \overline{K}^{*0}$ , and  $\Xi_c^0 \to \Sigma^+ K^{*-}$ , the differential decay rates [PRD 101, 053002 (2020)] are given by:

$$\frac{dN}{d\cos\theta_{\Lambda}} \propto 1 + \alpha (\Xi_{c}^{0} \to \Lambda \overline{K}^{*0}) \alpha (\Lambda \to p\pi^{-}) \cos\theta_{\Lambda},$$

$$\frac{dN}{d\cos\theta_{\Sigma^{0}}} \propto 1 + \alpha (\Xi_{c}^{0} \to \Sigma^{0} \overline{K}^{*0}) \alpha (\Sigma^{0} \to \Lambda \gamma) \cos\theta_{\Sigma^{0}}, \text{ and}$$

$$\frac{dN}{d\cos\theta_{\Sigma^{+}}} \propto 1 + \alpha (\Xi_{c}^{0} \to \Sigma^{+} K^{*-}) \alpha (\Sigma^{+} \to p\pi^{0}) \cos\theta_{\Sigma^{+}}.$$

Definitions of  $\theta_{\Lambda}$ ,  $\theta_{\Sigma^0}$ , and  $\theta_{\Sigma^+}$ :



This measurement is insensitive to production polarization of \(\mathcal{E}\_c^0\) in B-factory [PRD 63, 111102 (2001)].

The asymmetry parameter  $\alpha(\Sigma^0 \to \Lambda \gamma)$  is expected to be zero due to the case of parity conservation for an electromagnetic decay of  $\Sigma^0 \to \Lambda \gamma$ .

# Asymmetry parameters

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Note that  $\alpha(\Lambda \rightarrow p\pi^-) = 0.747 \pm 0.010$  and  $\alpha(\Sigma^+ \rightarrow p\pi^0) = -0.980 \pm 0.017$  from PDG.

$\alpha(\Xi_c^0 \to \Lambda \bar{K}^{*0}) \alpha(\Lambda \to p\pi^-)$	$0.115 \pm 0.164 ({ m stat.}) \pm 0.038 ({ m syst.})$
$\alpha(\Xi_c^0 \to \Sigma^0 \bar{K}^{*0}) \alpha(\Sigma^0 \to \gamma \Lambda)$	$0.008 \pm 0.072 ({ m stat.}) \pm 0.008 ({ m syst.})$
$\alpha(\Xi_c^0 \to \Sigma^+ K^{*-}) \alpha(\Sigma^+ \to p \pi^0)$	$0.514 \pm 0.295 ({ m stat.}) \pm 0.012 ({ m syst.})$
$\alpha(\Xi_c^0 \to \Lambda \bar{K}^{*0})$	$0.15 \pm 0.22 ({ m stat.}) \pm 0.05 ({ m syst.})$
$\alpha(\Xi_c^0\to\Sigma^+K^{*-})$	$-0.52 \pm 0.30 ({ m stat.}) \pm 0.02 ({ m syst.})$

# Measurements of $Br(\Xi_c^0 \to \Lambda K_S^0 / \Sigma^0 K_S^0 / \Sigma^+ K^-)$

- ➤ The Cabibbo-favored (CF) two-body hadronic weak decays  $\Xi_c^0 \rightarrow B_f$  (light baryons) + *P*(pseudoscalar mesons) have been analyzed based on the dynamical model calculations and  $SU(3)_F$  flavor symmetry methods.
- The predicted branching fractions for the CF decays Ξ<sup>0</sup><sub>c</sub> → ΛK<sup>0</sup>/ Σ<sup>0</sup>K<sup>0</sup>/Σ<sup>+</sup>K<sup>-</sup> based on different models are listed in the below Table.

Mode	dynamical model [PRD 101, 014011 (2020)]	SU(3) I [PLB 794, 19 (2019)]	SU(3) II [3] [JHEP 02 165 (2020)]
$\Xi_c^0 \to \Lambda \overline{K}{}^0$	$13.3 \times 10^{-3}$	$(10.5 \pm 0.6) \times 10^{-3}$	$(8.3 \pm 5.3) \times 10^{-3}$
$\Xi_c^0 \to \Sigma^0 \overline{K}{}^0$	$0.4 \times 10^{-3}$	$(0.8 \pm 0.8) \times 10^{-3}$	$(7.9 \pm 4.8) \times 10^{-3}$
$\Xi_c^0 \to \Sigma^+ K^-$	$7.8 \times 10^{-3}$	$(5.9 \pm 1.1) \times 10^{-3}$	$(22.0 \pm 5.7) \times 10^{-3}$

→ In this analysis, we present the measurements of branching fractions of  $\Xi_c^0 \to \Lambda K_S^0$ ,  $\Xi_c^0 \to \Sigma^0 K_S^0$ , and  $\Xi_c^0 \to \Sigma^+ K^-$  using all Belle data.

#### Measurements of $Br(\Xi_c^0 \rightarrow \Lambda K_S^0 / \Sigma^0 K_S^0 / \Sigma^+ K^-)$ arXiv:2111.08981



Relative branching fractions:  $\frac{Br(\Xi_{c}^{0} \to \Lambda K_{S}^{0})}{Br(\Xi_{c}^{0} \to \Xi^{-}\pi^{+})} = 0.229 \pm 0.008 \pm 0.012$   $\frac{Br(\Xi_{c}^{0} \to \Sigma^{0} K_{S}^{0})}{Br(\Xi_{c}^{0} \to \Xi^{-}\pi^{+})} = 0.038 \pm 0.006 \pm 0.004$   $\frac{Br(\Xi_{c}^{0} \to \Sigma^{+} K^{-})}{Br(\Xi_{c}^{0} \to \Xi^{-}\pi^{+})} = 0.123 \pm 0.007 \pm 0.010$ 

Absolute branching fractions:  $Br(\Xi_{c}^{0} \to \Lambda K_{S}^{0}) = (4.12 \pm 0.14 \pm 0.21 \pm 1.19) \times 10^{-3}$   $Br(\Xi_{c}^{0} \to \Sigma^{0} K_{S}^{0}) = (0.69 \pm 0.10 \pm 0.08 \pm 0.20) \times 10^{-3}$  $Br(\Xi_{c}^{0} \to \Sigma^{+} K^{-}) = (2.21 \pm 0.13 \pm 0.19 \pm 0.64) \times 10^{-3}$ 

The measured ratios of the branching fractions among the three decay modes are consistent with the theoretical predictions based on SU(3) flavor symmetry approaches within the theoretical errors, but contradict those predicted by dynamical model calculations.

# Measurement of $\mathcal{B}(\Omega_c^0 \to \Omega^- l^+ \nu)$ and $\mathcal{B}(\Xi_c^0 \to \Xi^- l \nu)$

Semileptonic decays of charmed baryons:

- Ideal test of QCD in transition region of (non-)perturbative.
- The cleanest processes among charm decays
- > Verify lepton flavor universality (LFU).  $\mathcal{B}(\Lambda_c^+ \rightarrow \Lambda e^+ \nu_e) = (3.6 \pm 0.4)\%$  PRL 115, 221805(2015)

**Experimentally:** 

- •BESIII measured the  $\mathcal{B}(\Lambda_c^+ \to \Lambda l^+ \nu)$
- •ARGUS and CLEOII measured  $\mathcal{B}(\Xi_c \to \Xi \ l^+ \nu)$  | arge uncertainty •CLEO measured  $\mathcal{B}(\Omega_c^0 \to \Omega^- e^+ \nu)$  |



 $\frac{\beta_{0}}{\beta_{0}} = \frac{\beta_{0}}{(1): \text{ CLEO}} + \frac{\beta_{120302-003}}{\beta_{0}} + \frac{\beta_{120302-003}}{\beta_{0}} + \frac{\beta_{120302-003}}{\beta_{120302-003}} + \frac{$ 

 $\mathcal{B}(\Lambda_c^+ \to \Lambda \,\mu^+ \nu_e) = (3.5 \pm 0.4)\%$  PLB 767, 42 (2017)

Measurement of 
$$\mathcal{B}(\Omega_c^0 \to \Omega^- l^+ \nu)$$
 and  $\mathcal{B}(\Xi_c^0 \to \Xi^- l \nu)$ 



# Measurement of $\mathcal{B}(\Xi_c^0 \to \Xi^- l \nu)$





#### PRL 127, 121803 (2021)

#### Fit component:

- Signal : True signal histogram
- BKG1:  $\Xi^-$  sideband
- BKG2:  $\Xi^- \ell^- \Xi^-$  sideband
- BKG3:  $\Xi_c^0 \to \Xi^- \pi^+ \ell \nu$  histogram
- BKG5:  $\Xi_c^0 \rightarrow \Xi^- \pi + h$  histogram
- BKG4: Bkg histogram from B decay

$$B(\Xi_c^0 \to \Xi^- e^+ \nu_e) = (1.31 \pm 0.39)\%$$

Previous:  $(2.34 \pm 1.59)\%$ 

 $\mathcal{B}(\mathcal{Z}_c^0 \rightarrow \mathcal{Z}^- \mu^+ \nu_\mu) = (1.27 \pm 0.39)\%$ 

Consistent with LFU

# Measurement of $\mathcal{B}(\Omega_c^0 \to \Omega^- l^+ \nu)$

#### $\Omega_c^0 \to \Omega^- \pi^+$ : Fragmentation Function extraction



Peterson's fragmentation function

$$\frac{dN}{dx_p} \approx \frac{1}{x_p} \cdot \frac{1}{\left(1 - \frac{1}{x_p} - \frac{\epsilon_p}{1 - x_p}\right)^2}$$

• data with  $p^*_{\Omega\ell(\pi)}/p^*_{\max}$ >0.5 used in fit •  $\epsilon_p = 0.1160 \pm 0.014$ 

preliminary

 $\Omega_c^0 \to \Omega^- \mu^+ \nu$ : signal extraction Similar Data-driven method used in  $\Xi_c^0 \to \Xi^- l \nu$ 



 $\frac{\mathcal{B}(\Omega_c^0 \to \Omega^- e^+ \nu)}{\mathcal{B}(\Omega_c^0 \to \Omega^- \pi^+)} = 1.98 \pm 0.15$ Previous: 2.4 ± 1.2  $\frac{\mathcal{B}(\Omega_c^0 \to \Omega^- \mu^+ \nu)}{\mathcal{B}(\Omega_c^0 \to \Omega^- \pi^+)} = 1.94 \pm 0.21$ 

Consistent with LFU

# $\Xi_c$ worklist:

#### 1. Measurement of absolute decay branching fractions

 $\mu? \left\{ \begin{array}{l} \mathcal{B}(\Xi_c^0 \to \Xi^- \pi^+) &= (1.80 \pm 0.52)\% \text{PRL 122 082001} \\ \mathcal{B}(\Xi_c^+ \to \Xi^- \pi^+ \pi^+) &= (2.86 \pm 1.27)\% \text{ PRD 100 031101} \\ \mathcal{B}(\Xi_c^0 \to \Xi^- e^+ \nu_e) &= (1.8 \pm 1.2)\% \text{ PDG} \\ \mathcal{B}(\Xi_c^+ \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratios to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratio to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratio to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratio to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratio to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratio to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratio to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0 e^+ \nu_e) &= (1.8^{+0.7}_{-0.8})\% \text{ PDG, ratio to} \\ \mathcal{B}(\Xi_c^- \to \Xi^0$ 

#### 2. Find more decay modes:

#### 3.Decay parameter measurement:



# Charmed baryon spectra at BelleII

- LHCb recently reconstructed  $K^+\Lambda_c^-$  and found  $E_c(2923)^0$ ,  $E_c(2939)^0$  and  $E_c(2965)^0$
- $E_c(2930)^0$  is the sum of them ?



Phys. Rev. Lett. 124, 222001

We use LHCb's results as input, refit our plot and find the fit result is good
 More data needed to extract signal yield with low uncertainty
 Notice Ξ<sub>c</sub>(2923)<sup>+</sup> → Λ<sup>+</sup><sub>c</sub>K<sup>-</sup>π<sup>+</sup>, can we find more decay channel?
 How about Σ<sub>c</sub>(2800)? Is Σ<sub>c</sub>(2800) the overlap of several Σ<sub>c</sub>?

▶ ....



- Although Belle has stopped data taking for >10 years ago, we are still producing exciting results [China group has made great contributions].
- ▶ Belle II started data taking on 25 March 2019 with its full detector.
- The Belle II experiment at SuperKEKB aims to find New Physics beyond the SM with ultimate precision measurement (a few %, typically) of heavy flavor decays.
- SuperKEKB has achieved Lpeak =2.4 x 10<sup>34</sup>cm<sup>-2</sup>s<sup>-1</sup> (world highest luminosity)
- ▶ Belle II is performing as expected, and obtained early physics results.
- SuperKEKB/Belle II aims at accumulate 50ab<sup>-1</sup> by ~2030, by further improving the luminosity performance.





# 感谢您的批评指正 <sub>沈成平</sub> shencp@fudan.edu.cn

# **BackUp : Semileptonic decays**

=

 Semileptonic decays of charmed baryons:
 be calculated with QCD factorization approach
 The cleanest processes among charm decays
 Decay rate depend on CKM matrix element |V<sub>cq</sub>| and strong interaction (form factor)
 Verify lepton universality.

$$\mathcal{A}(\mathbf{B_{c}} \to \mathbf{B_{n}}\ell^{+}\nu_{\ell}) = \frac{G_{F}}{\sqrt{2}}V_{cq}\langle \mathbf{B_{n}} | J_{\mu}^{(V-A)} | \mathbf{B_{c}} \rangle \bar{u}_{\nu}\gamma^{\mu} \frac{(1-\gamma_{5})}{2}v_{\ell}$$

$$\overset{\mathbf{B_{c}} = (\Xi_{c}^{0}, -\Xi_{c}^{+}, \Lambda_{c}),}{\underset{\mathbf{B}_{n} = \begin{pmatrix} \frac{1}{\sqrt{6}}\Lambda + \frac{1}{\sqrt{2}}\Sigma^{0} & \Sigma^{+} & p \\ \Sigma^{-} & \frac{1}{\sqrt{6}}\Lambda - \frac{1}{\sqrt{2}}\Sigma^{0} & n \\ \Xi^{-} & \Xi^{0} & -\sqrt{\frac{2}{3}}\Lambda \end{pmatrix}} J_{\mu}^{(V-A)} = \bar{q}\gamma_{\mu}(1-\gamma_{5})/2c$$

$$\overline{u}_{\nu} \text{ and } v_{\ell} \text{: Dirac bispinors } \mathbf{q} = \mathbf{s}, \mathbf{d}$$

$$G_{F} \text{: Fermi constant}$$

# BackUp: $\Xi_c \to \Xi(\Lambda) l \nu$

## $\Xi_c \rightarrow \Xi (\Lambda) l \nu$ have been well studied theoretically.

- Total branching fraction(B),
- forward-backward asymmetries  $(\mathcal{A}_{FB})$ ,
- differential decay rates  $(d\Gamma/dq^2)$

are observable values that uncover the underline dynamics of QCD.

#### Theories such as:

SU(3) symmetry arXiv:1901.05610 PRD, 97, 073006 (2018) PLB, 792, 214 (2019); Quasipotential approach relativistic quark model EPJC, 79, 695 (2019); light-front quark modelCPC, 42,093101(2018); light-cone QCD sum rules EPJC, 79, 695 (2019) calculated:

$$\mathcal{B}(\Xi_c^0 \to \Xi^- l^+ \nu_e) = 1.35 \sim 7.26\%$$
  $\mathcal{B}(\Xi_c^+ \to \Xi^0 l^+ \nu_e) = 3.3 \sim 28.6\%$ 

 $\mathcal{B}(\Xi_c^+\to\Lambda l^+\nu_e)=0.082{\sim}0.17\%$ 

Also EPJC, 79, 695 (2019) predict:

$$\frac{6\Gamma(\Xi_c \to \Lambda e\nu_e)}{|V_{cd}|^2} = \frac{\Gamma(\Xi_c \to \Xi e\nu_e)}{|V_{cs}|^2}$$

# **BackUp:** Excited states:

#### By John Yelton

https://hflav-eos.web.cern.ch/hflav-

eos/charm/baryons/Excited\_Apr19/baryons\_April19.html

Charmed Barvon	Mode	Mass	Natural Width	$J^{P}$	Status and Comments
Excited State	mode	$(MeV/c^2)$	$(MeV/c^2)$	0	
$\Lambda_{c}(2595)^{+}$	$\Lambda_c^+\pi^+\pi^-, \Sigma_c\pi$	$2592.25 \pm 0.28$	$2.59 \pm 0.30 \pm 0.47$	$1/2^{-}$	well established, most precise mmeasurement by CDF 1
$A_c(2625)^+$	$\Lambda_c^+\pi^+\pi^-$	$2628.11 \pm 0.19$	< 1.9	3/2-	well established, most precise measurements by CDF 1
$\Lambda_{c}(2765)^{+}$	$\Lambda_c^+\pi^+\pi^-, \Sigma_c\pi$	$2766.6 \pm 2.4$	50	??	discovered by CLEO, seen by Belle, but parameters not measured 2
$A_{c}(2880)^{+}$	$\Lambda^+_{\alpha}\pi^+\pi^-, \Sigma_c\pi,$	$2881.53 \pm 0.35$	$5.8 \pm 1.1$	5/2+	well established and seen in more than one mode 24
	$\Sigma_c(2520)\pi, D^0p$			(experimental evidence)	
$\Lambda_{c}(2940)^{+}$	$D^0p, \Sigma_c\pi$	$2939.3^{+1.4}_{-1.5}$	$17^{+8}_{-6}$	??	Seen by both BaBar 4 and BelleMizuk
$\Sigma_c(2455)^{++}$	$\Lambda_c^+\pi^+$	$167.510 \pm 0.17$	$1.89\pm^{+0.09}_{-0.18}$	1/2+	well established, most precise measurements by Belle 5
$\Sigma_{c}(2455)^{+}$	$\Lambda_c^+\pi^+$	$166.4\pm0.4$	< 4.6 @ 90% CL	$1/2^+$	well established, but parameters not measured precisely
$\Sigma_{c}(2455)^{0}$	$\Lambda_c^+\pi^+$	$167.29\pm0.17$	$1.83^{+0.11}_{-0.19}$	$1/2^+$	well established, most precise measurements by Belle 5
$\Sigma_c(2520)^{++}$	$\Lambda_c^+ \pi^+$	$231.95^{+0.17}_{-0.12}$	$14.78\pm +0.30_{-0.40}$	3/2+	well etablished, most precise measurements by Belle 5
$\Sigma_{c}(2520)^{+}$	$\Lambda_c^+\pi^+$	$231.0\pm2.3$	< 17 @ 90% CL	$3/2^+$	fairly well established, awaits precise measurement
$\Sigma_{c}(2520)^{0}$	$\Lambda_c^+\pi^+$	$232.02\substack{+0.15\\-0.14}$	$15.3_{-0.5}^{+0.4}$	$3/2^+$	well established, most precise measurements by Belle 5
$\Sigma_c(2800)^{++}$	$\Lambda_c^+\pi^+$	$514^{+4}_{-6}$	$75^{+18+12}_{-13-11}$	tentatively identified	observed by Belle 6 - should be confirmed
$\Sigma_{c}(2800)^{+}$	$\Lambda_c^+ \pi^0$	$505^{+15}_{-5}$	$62^{+37+52}_{-23-38}$	as members of the predicted	
$\Sigma_{c}(2800)^{0}$	$\Lambda_c^+\pi^-$	$519^{+5}_{-7}$	$72^{+22}_{-15}$	$\Sigma_{c2} 3/2^{-}$ isospin triplet?	same states as that below?
	$\Lambda_c^+\pi^-$	$560\pm8\pm10$	$86^{+33}_{-22}$		seen by Babar 7 in resonant substructure of B decays - needs confirmation
$\Xi_c^{\prime+}$	$\Xi_c^+\gamma$	$110.5\pm0.4$		1/2+	well established
$\Xi_c^{\prime 0}$	$\Xi_c^0 \gamma$	$108.3\pm0.4$		$1/2^+$	well established
$\Xi_c(2645)^+$	$\Xi_c^0 \pi^+$	$178.5\pm0.1$	$2.1\pm0.2$	$3/2^+$	well established, widths recently measured by Belle 8
$\Xi_c(2645)^0$	$\Xi_c^+\pi^-$	$174.7\pm0.1$	$2.4\pm0.2$	$3/2^+$	
$\Xi_c(2790)^+$	$\Xi_{c}^{\prime 0}\pi^{+}$	$320.7\pm0.5$	$9\pm1$	1/2-	well established, widths recently measured by Belle 8
$\Xi_c(2790)^0$	$\Xi_c^{\prime +}\pi^-$	$323.8\pm0.5$	$10 \pm 1$	$1/2^{-}$	
$\Xi_c(2815)^+$	$\Xi_c(2645)^0\pi^+$	$348.8\pm0.1$	$2.43\pm0.23$	3/2-	well established, widths recently measured by Belle 8
$\Xi_c(2815)^0$	$\Xi_{c}(2645)^{+}\pi^{-}$	$349.4\pm0.1$	$2.54\pm0.23$	$3/2^{-}$	
$\Xi_c(2930)^+$	$\Lambda_c^+ K_S^0$	$2942.3 \pm 4.4 \pm 1.5$	$14.8 \pm 8.8 \pm 2.5$	??	"evidence" recently reported by Belle 9
$\Xi_c(2930)^0$	$\Lambda_c^+ K^-$	$2928.9 \pm 3.0 \substack{+0.9 \\ -12.0}$	$19.5 \pm 8.4^{+5.9}_{-7.9}$	??	originally reported by BaBar 11, confirmed by Belle 10
$\Xi_c(2970)^+$	$\Lambda_c^+ K^- \pi^+, \Sigma_c^{++} K^-, \Xi_c(2645)^0 \pi^+$	$2967.2\pm0.8$	$21 \pm 3$	??	well established, but parameters in different modes and experiments differ
$\Xi_c(2970)^0$	$arepsilon_c(2645)^+\pi^-$	$2970.4\pm0.8$	$28 \pm 3$	??	well established, but parameters in different modes and experiments differ
$\Xi_c(3055)^+$	$\Sigma_c^{++}K^-, \Lambda D$	$3055.7\pm0.4$	$8.0 \pm 1.9$	??	seen by Belle and BaBar 12 14
$\Xi_c(3055)^0$	$\Lambda D$	$3059.0\pm0.8$	$6.2 \pm 2.4$	??	newly observed by Belle 14
$\Xi_c(3080)^+$	$\Lambda_{c}^{+}K^{-}\pi^{+}, \Sigma_{c}^{++}K^{-}, \Sigma_{c}(2520)^{++}K^{-}, \Lambda D$	$3077.8\pm0.3$	$3.6\pm0.7$	??	seen by Belle and BaBar 12 15
$\Xi_c(3080)^0$	$\Lambda_{c}^{+}K_{S}^{0}\pi^{-},  \Sigma_{c}^{0}K_{S}^{0},  \Sigma_{c}(2520)^{0}K_{S}^{0}$	$3079.9 \pm 1.0$	$5.6 \pm 2.2$	??	seen by Belle and BaBar 12 14 15
$\Omega_{c}(2770)^{0}$	$\Omega_c^0 \gamma$	$2765.9\pm2.0$	0	3/2+	seen by BaBar 16 and Belle 17
$\Omega_c(3000)^0$	$\Xi_c^+ K^-$	$3000.4 \pm 0.2 \pm 0.1^{+0.3}_{-0.5}$	$4.5\pm0.6\pm0.3$	??	LHĊb 18
$\Omega_{c}(3050)^{0}$	$\Xi_c^+ K^-$	$3050.2 \pm 0.1 \pm 0.1 \pm 0.1 \substack{+0.3 \\ -0.5}$	< 1.2,95%CL	??	LHCb 18
$\Omega_{c}(3066)^{0}$	$\Xi_c^+ K^-$	$3065.6 \pm 0.1 \pm 0.3^{+0.3}_{-0.5}$	$3.5\pm0.4\pm0.2$	??	LHCb 18
$\Omega_{c}(3090)^{0}$	$\Xi_c^+ K^-$	$3090.2 \pm 0.3 \pm 0.5^{+0.3}_{-0.5}$	$8.7\pm1.0\pm0.8$	??	LHCb 18
$\Omega_{c}(3119)^{0}$	$\Xi_c^+ K^-$	$3119.1 \pm 0.3 \pm 0.9 \substack{+0.3 \\ -0.5}$	$1.1\pm0.8\pm0.4$	??	LHCb 18
$\Omega_{c}(3118)^{0}$	$\Xi_c^+ K^-$	$3188 \pm 5 \pm 13$	$60\pm15\pm11$	??	Reported by LHCb [18], not clear if it is several resonances

# BackUp: Current charmed baryon status

$\Lambda_c^+$	$1/2^+$	****
$arLambda_c(2595)^+$	$1/2^-$	***
$arLambda_c(2625)^+$	$3/2^-$	***
$arLambda_c(2765)^+$ or $arLambda_c(2765)$		*
$arLambda_c(2860)^+$	$3/2^+$	***
$arLambda_c(2880)^+$	$5/2^+$	***
$arLambda_c(2940)^+$	$3/2^-$	***
$\Sigma_c(2455)$	$1/2^+$	****
$\Sigma_c(2520)$	$3/2^+$	***
$\Sigma_c(2800)$		***
$\Xi_c^+$	$1/2^+$	***
$\Xi_c^0$	$1/2^+$	****
$\Xi_c^{\prime+}$	$1/2^+$	***
$ec{\Xi}_{c}^{\prime 0}$	$1/2^+$	***
$arepsilon_c(2645)$	$3/2^+$	***
$\Xi_c(2790)$	$1/2^{-}$	***
$\Xi_c(2815)$	$3/2^-$	***
$arepsilon_c(2930)$		**
$arepsilon_c(2970)$		***
was $arepsilon_c(2980)$		
$arepsilon_c(3055)$		***
$arepsilon_c(3080)$		***
$arepsilon_c(3123)$		*
$arOmega_c^0$	$1/2^{+}$	***
${\it \Omega_c(2770)}^0$	$3/2^+$	***
$arOmega_c(3000)^0$		***
$arOmega_c(3050)^0$		***
$arOmega_c(3065)^0$		***
$\Omega_c(3090)^0$		***
$\Omega_c(3120)^0$		***
- ( )		

\*\*\*\*: Existence is certain, properties fairly explored.

- \*\*\*:Existence is very likely or certain, further confirmation is desirable and/or quantum numbers, branching fractions, etc. are not well determined.
  - \*\*:Evidence of existence is only fair.
    - \*:Evidence of existence is poor.

Only  $\Lambda_c^+$ ,  $\Xi_c^0$  and  $\Sigma_c(2455)$  in \*\*\*\* status



#### 发现了新粲重子激发态E<sub>c</sub>(2930)

 相继发现带电的和中性的E<sub>c</sub>(2930),具有奇特性质:唯一在B介 子衰变过程中发现,但未在正负电子直接对撞过程中观测到的E<sub>c</sub> 激发态,引起广泛关注。



Y.B.Li, \*C.P.Shen *et al* (Belle) EPJC 78, 928 (2018); Y.B.Li, \*C.P.Shen *et al* (Belle) EPJC 78, 252 (2018)