

# Probing ultralight bosons with compact eccentric binaries

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Supermassive Black Hole and Fundamental Physics

Sep 26, 2021 · ITP CAS

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2107.13527

Weakly coupled & light degrees are well motivated BSM physics  
but are hard to probe

DM: gravity evidences only

Most searching strategies assume beyond gravity couplings

Gravity probe is considered somewhat indirect:  
statistics / collective phenomena

New opportunities from GWs and (SM)BHs  
Precision measurements; more informative

# Ultralight bosons

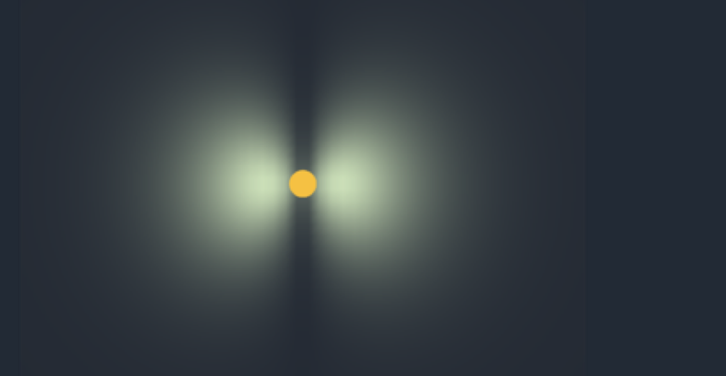
- Gravitational atom

Fast rotating BH / Superradiance  
/ Quasi-stable bound state

$$(\square - \nu^2)\sigma = 0$$

$$\sigma = 2\text{Re}\frac{1}{\sqrt{2\nu}}\psi_{n\ell m}e^{-iE_nt}$$

$|211\rangle$



- Produced through  $\text{Im } E_n$  / Depleted through GWs

$$T_{\text{prod}}(|211\rangle) \sim 10^6 \text{yr} (M/10M_\odot) (0.016/\alpha)^9$$

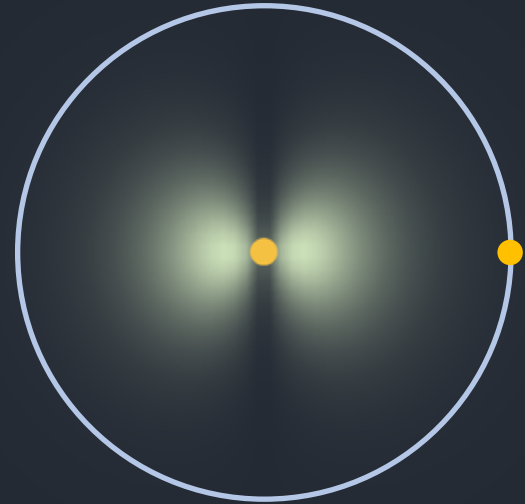
$$T_{\text{depl}}(|211\rangle) \sim 10^8 \text{yr} (M/10M_\odot) (0.062/\alpha)^{15}$$

# “Gravitational collider physics”

- GA-BH binary / GA-PSR binary  
Orbital induced level transition

- Energy transferred between the cloud and the orbit  
Sinking orbit / floating orbit

- Works only when orbital frequency matches energy split  
Several discrete narrow bands in orbital separation
- Many free parameters: challenging for GW observation



# Gravitational collider goes off-shell

- Energy distribution of GA

$$\rho_{211}(\mathbf{x}) = \frac{1}{64\pi} \alpha^5 \nu^5 m_C r^2 e^{-\alpha \nu r} \sin^2 \theta$$

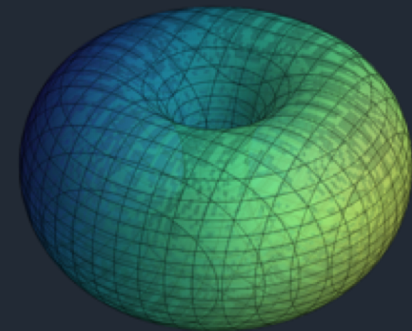
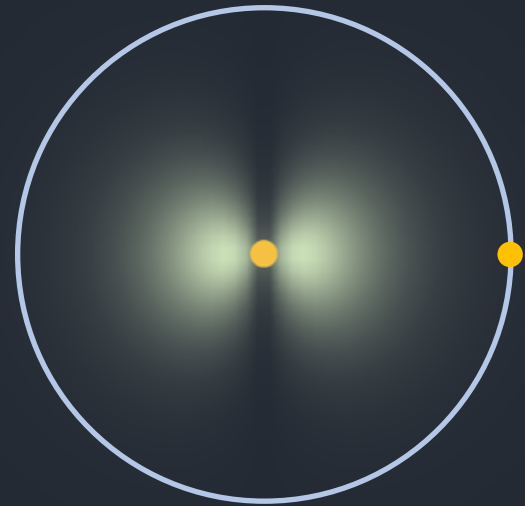
$$m_C \simeq \alpha m_{\text{Kerr}}$$

$$J_C \simeq J_{\text{Kerr}} \quad Q_C \simeq \alpha^{-3} Q_{\text{Kerr}}$$

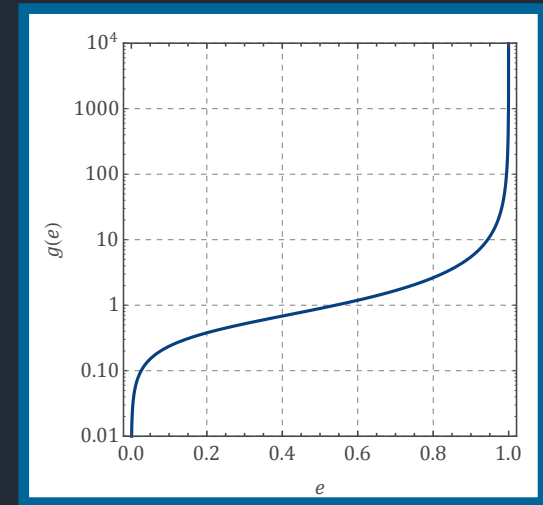
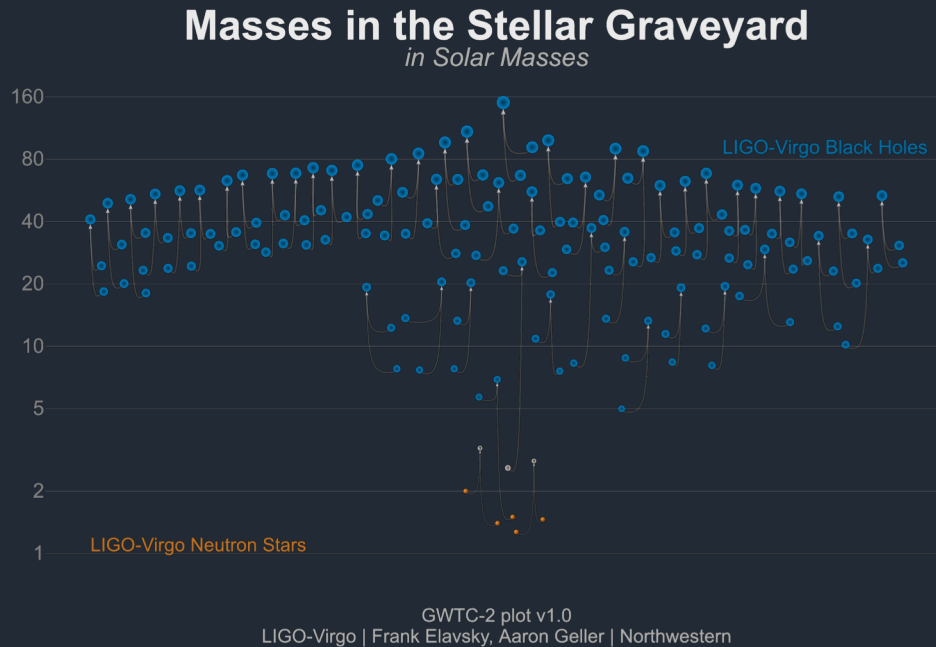
- A small fraction of mass  
but huge quadrupole  $\sim 1/\alpha^3$

1~10 for NS, ~100 for WD

- Significant orbital precession



# Compact eccentric binaries



$$\mathcal{G}(e) \equiv \frac{e^{12/19}}{1 - e^2} \left( 1 + \frac{121}{304} e^2 \right)^{870/2299}$$

Resolving orbital precession needs observably large orbital eccentricity

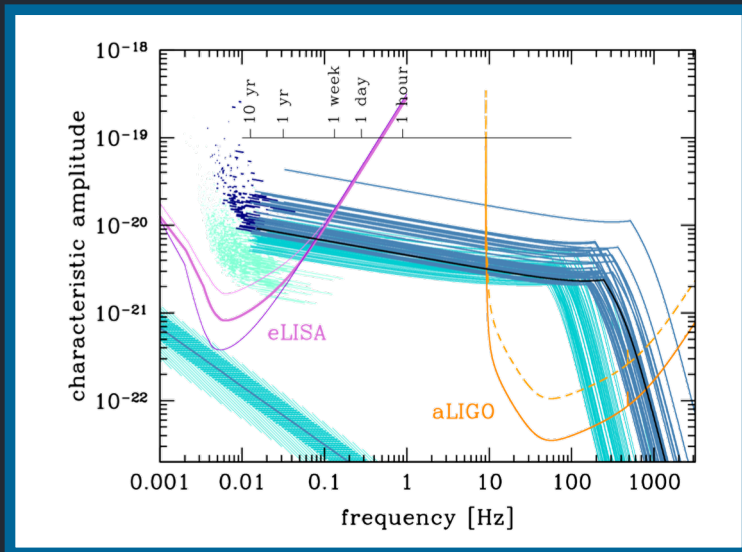
No eccentric binaries confirmed in LIGO, with a good reason

# Compact eccentric binaries

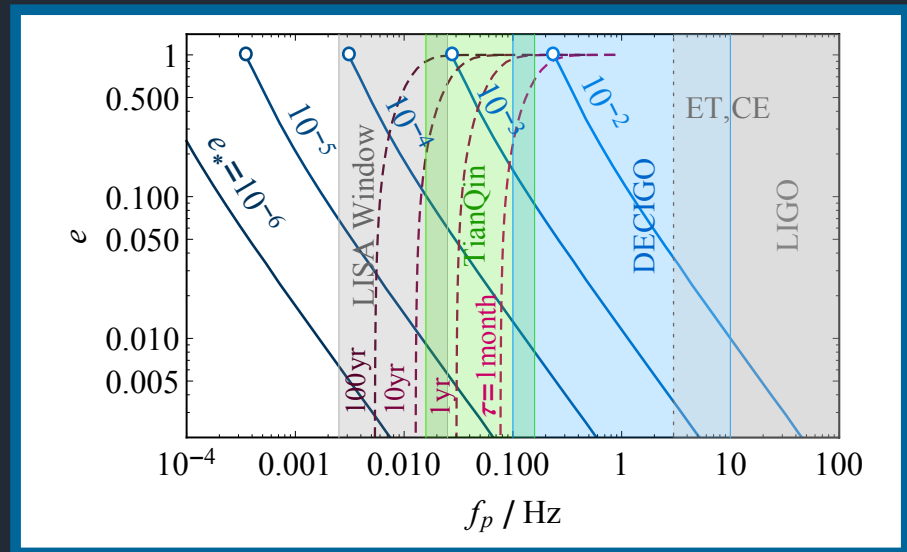
But LIGO-type binaries are expected to be **eccentric**  
at lower frequencies

## Important targets of space GW telescopes

PSR-BH are expected to be eccentric, too  
Could be seen in the future

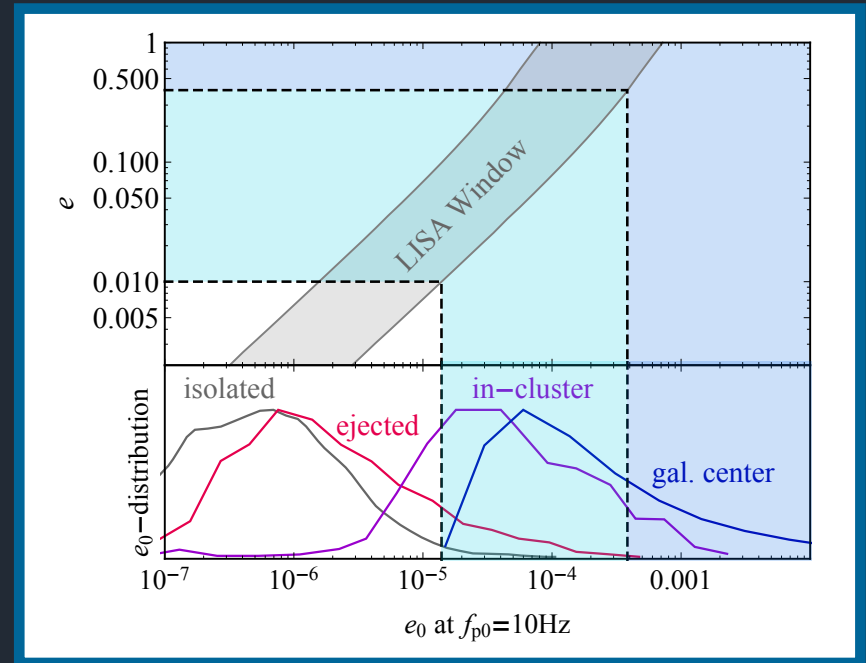
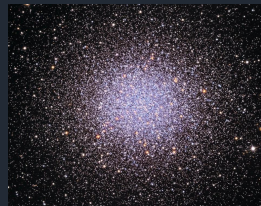
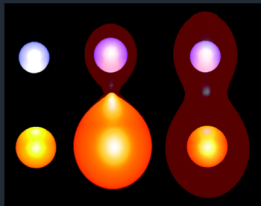
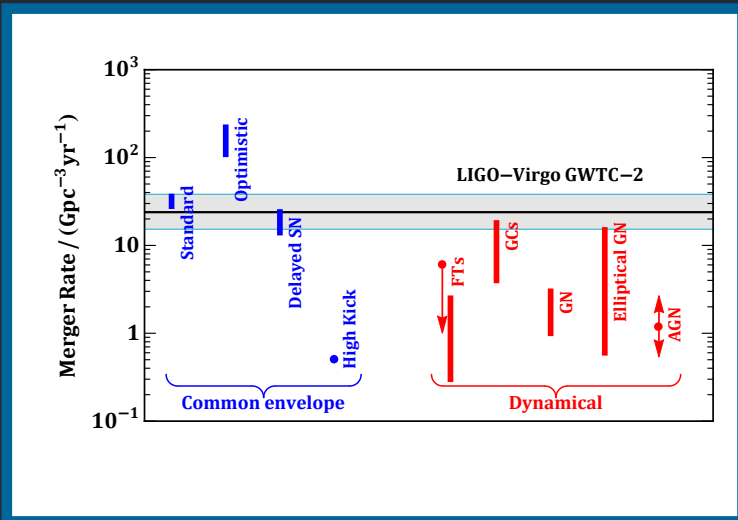


## Sesana 16



Randall &amp; ZZX 19

# Compact eccentric binaries



Randall & ZZX 17, 18

LIGO/Virgo: 2010.14533; CE: 1405.7016; Field Triples: 1608.07642,  
1703.06614; GCs: 1809.01152GN: 1706.09896; Elliptical GN:  
1705.05848; AGN: 1602.03831



# Cloud-induced precessions

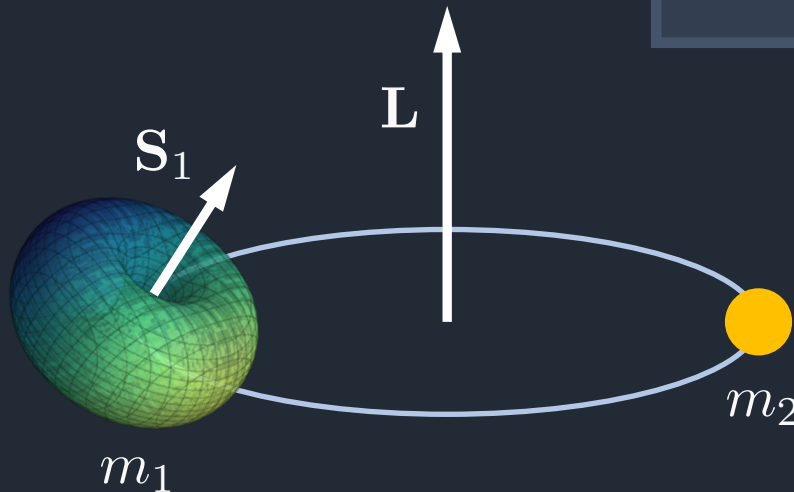
Interaction potential:  
orbital, **quadrupole-orbit**, and spin-orbit

Spin-spin always negligible

$$V_O = -\frac{Gm\mu}{r} + \text{1PN terms}$$

$$V_Q = \frac{Gm\mu}{2r^3} q_1 (1 - 3 \cos^2 n_1)$$

$$V_{SO} = \frac{2GL}{r^3} (m_1 + \frac{3}{4}m_2) s_1 \cos n_1$$



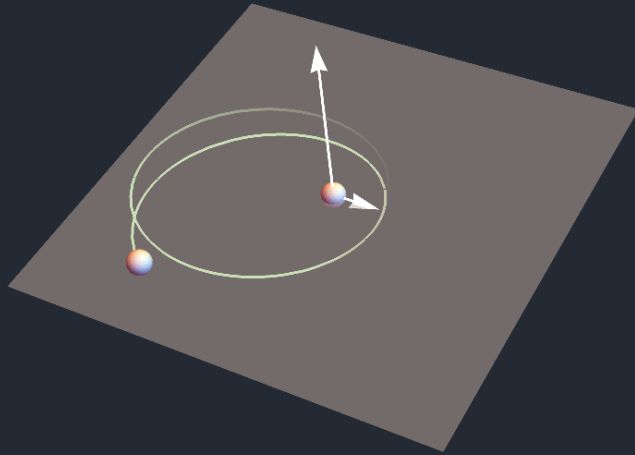
$$\cos n_1 = \mathbf{S}_1 \cdot \mathbf{L}$$

$$q_1 = Q_1/m_1$$

$$s_1 = S_1/m_1$$

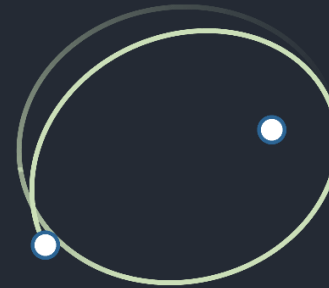
# Cloud-induced precessions

## Nodal precession



Confined in a small cone by  
angular momentum conservation,  
Unimportant for GW

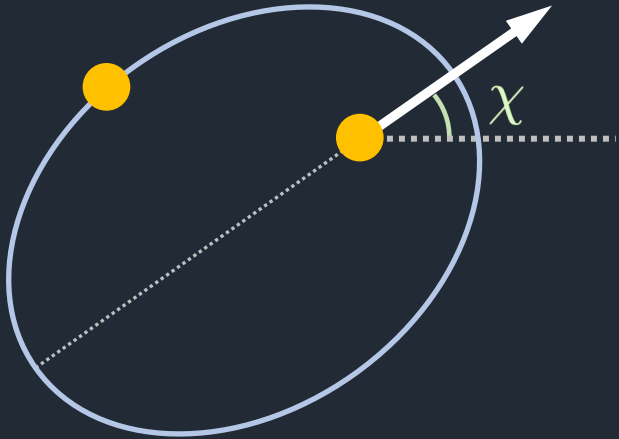
## Apsidal precession



Typically very significant,  
**main signal**

\* Also spin precessions, but happen to be unobservable

# Apsidal precession rate



$$\dot{\chi} = \dot{\chi}_{\text{GR}} + \dot{\chi}_{\text{Q}} + \dot{\chi}_{\text{S}}$$

$$V_{\text{O}} = -\frac{Gm\mu}{r} + \text{1PN terms}$$

$$V_{\text{Q}} = \frac{Gm\mu}{2r^3} q_1 (1 - 3 \cos^2 n_1)$$

$$V_{\text{SO}} = \frac{2GL}{r^3} (m_1 + \frac{3}{4}m_2) s_1 \cos n_1$$

$$\dot{\chi}_{\text{GR}} = \frac{3Gm\omega}{a(1 - e^2)}$$

$$\dot{\chi}_{\text{Q}} = \frac{3\omega}{4a^2(1 - e^2)^2} q_1 (1 - 3 \cos^2 n_1)$$

$$\dot{\chi}_{\text{S}} = \frac{-3\omega^2}{(1 - e^2)^{3/2}} s_1 \left(1 + \frac{m_1}{3m}\right) \cos n_1$$

# Parameter space

$$\frac{\dot{\chi}_Q}{\dot{\chi}_{GR}} \simeq \frac{3u_1}{2\alpha^3} \frac{Gm_1}{a(1-e^2)}$$

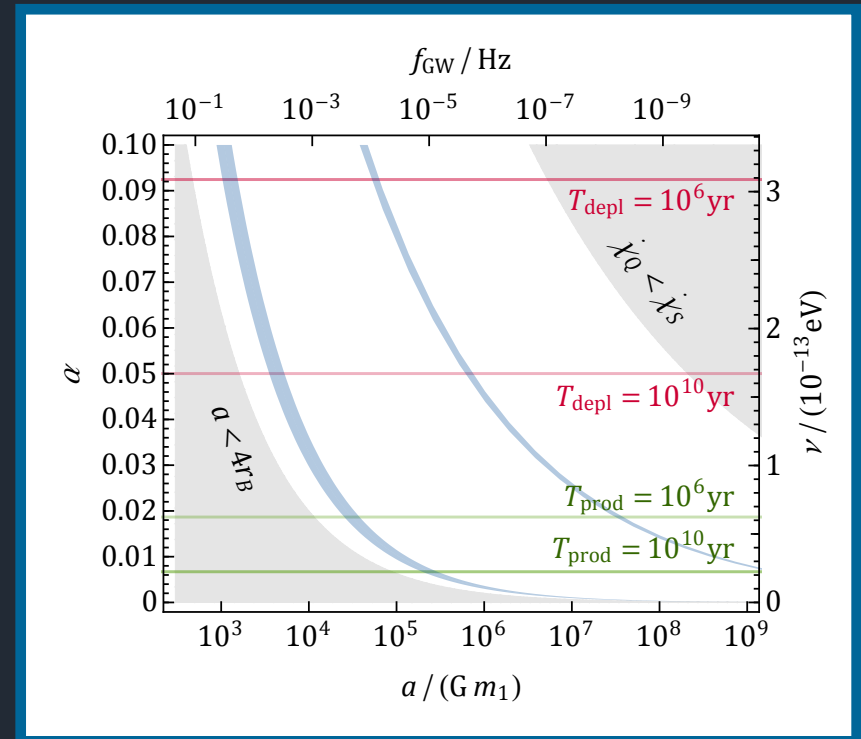
$$\frac{\dot{\chi}_Q}{\dot{\chi}_S} \simeq \frac{3u_1}{2\alpha^3(1+u_1/3)} \left[ \frac{Gm_1}{a(1-e^2)} \right]^{1/2}$$

$$u_1 \equiv m_1/m$$

GR precession can be subtracted  
up to measurement errors

SO precession as contamination

Much widened parameter space



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# GW signals

Perturbativity pushes the system to mHz ~ dHz window

$$f_2 = \frac{1}{\pi} \sqrt{\frac{Gm}{a^3}} < \frac{1}{\pi} \sqrt{\frac{Gm}{(4r_B)^3}} \sim 0.14 \text{Hz} \left( \frac{\alpha}{0.05} \right)^3 \left( \frac{10 M_\odot}{m_1} \right)$$

GWs trace orbital motion

$$h^{ij} = \frac{2G}{d} \Lambda^{ij,kl} \ddot{M}^{kl} \quad M^{ij} = \mu r^i r^j \quad r_i = r(\cos \bar{\psi}, \sin \bar{\psi})$$
$$r = \frac{a(1 - e^2)}{1 + e \cos \psi} \quad \bar{\psi} = \psi + \chi$$

Eccentricity: nonuniform motion (Kepler's 2nd law)

Apsidal precession: two sets of frequencies in radial and angular directions

# GW signals

- Eccentricity: harmonics

$$\omega_n = n\omega$$

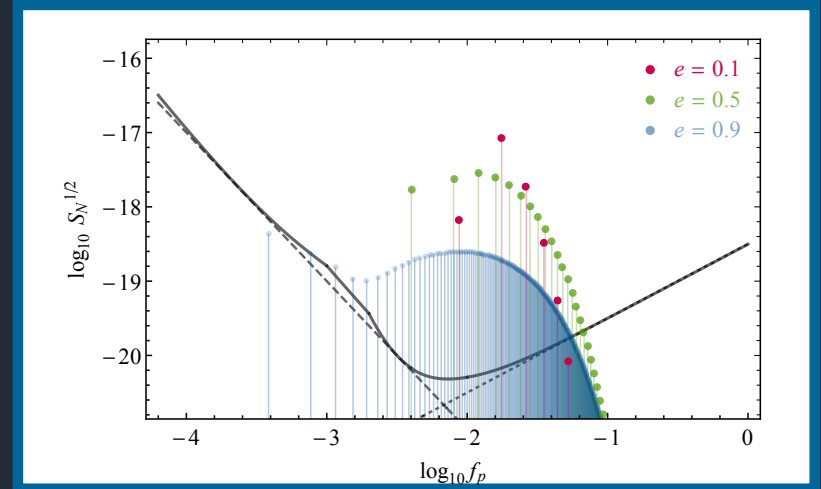
- Precession: triple split

$$(\omega_n, \omega_{n\pm}) \equiv (n\omega, n\omega \pm 2\dot{\chi})$$

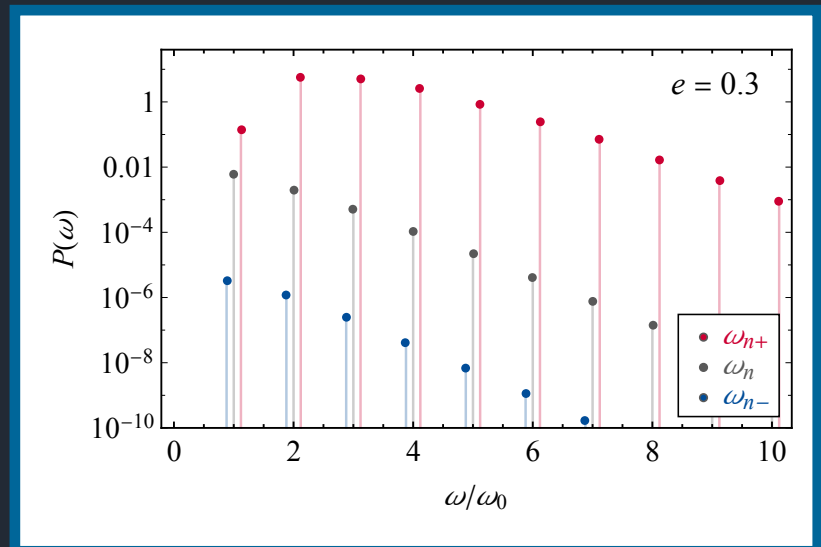
The highest component almost always dominates; the other two practically unobservable

- Observable: anharmonic overtone

absent for circular orbit  
at least at Newtonian level



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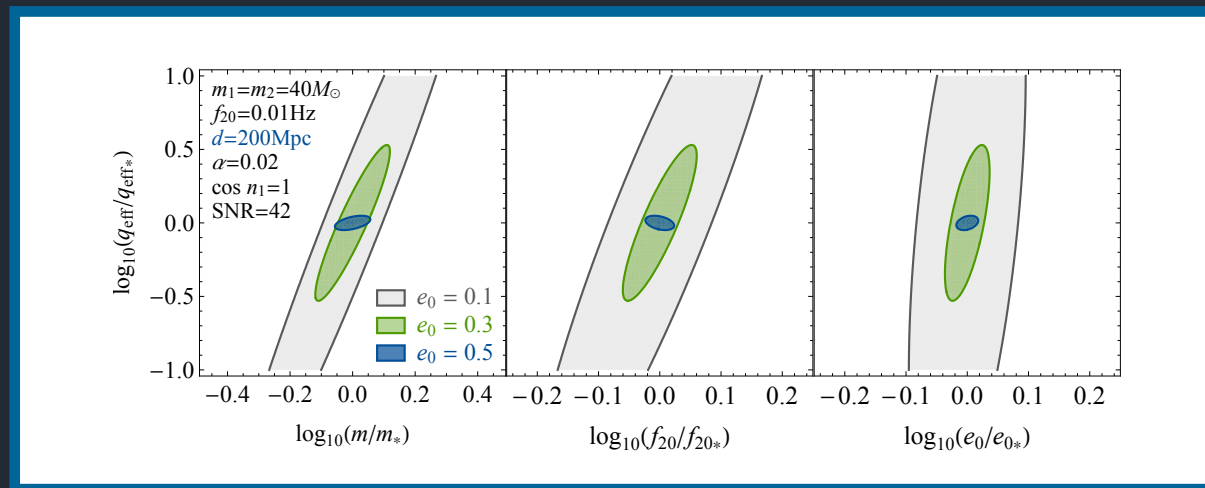


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# GW signals

- Observability of anharmonic overtones relies on the frequency resolution
- 1yr data @ 0.01Hz  $10^{-5} < q_1/a^2 < \alpha$

A partial time-domain Fisher forecast



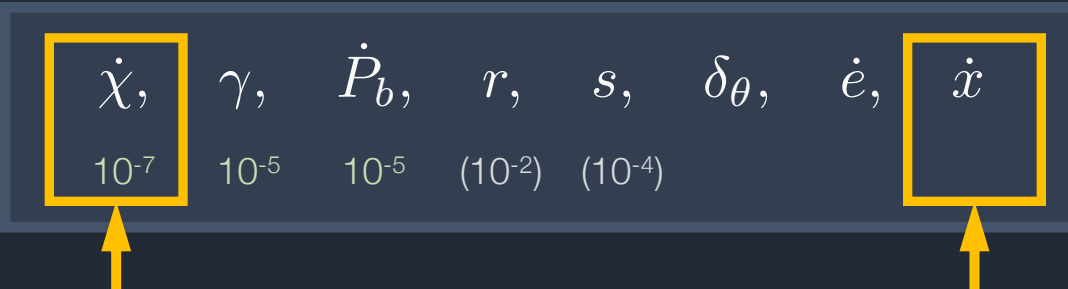
# Pulsar-GA binaries

- 5 Keplerian parameters:

$P_b$	$T_0$	$x = a \sin I$	$e$	$\chi_0$
$10^{-12}$	$10^{-13}$	$10^{-8}$	$10^{-7}$	$10^{-8}$

- Post-Keplerian parameters:

$\dot{\chi},$	$\gamma,$	$\dot{P}_b,$	$r,$	$s,$	$\delta_\theta,$	$\dot{e},$	$\dot{x}$
$10^{-7}$	$10^{-5}$	$10^{-5}$	$(10^{-2})$	$(10^{-4})$			



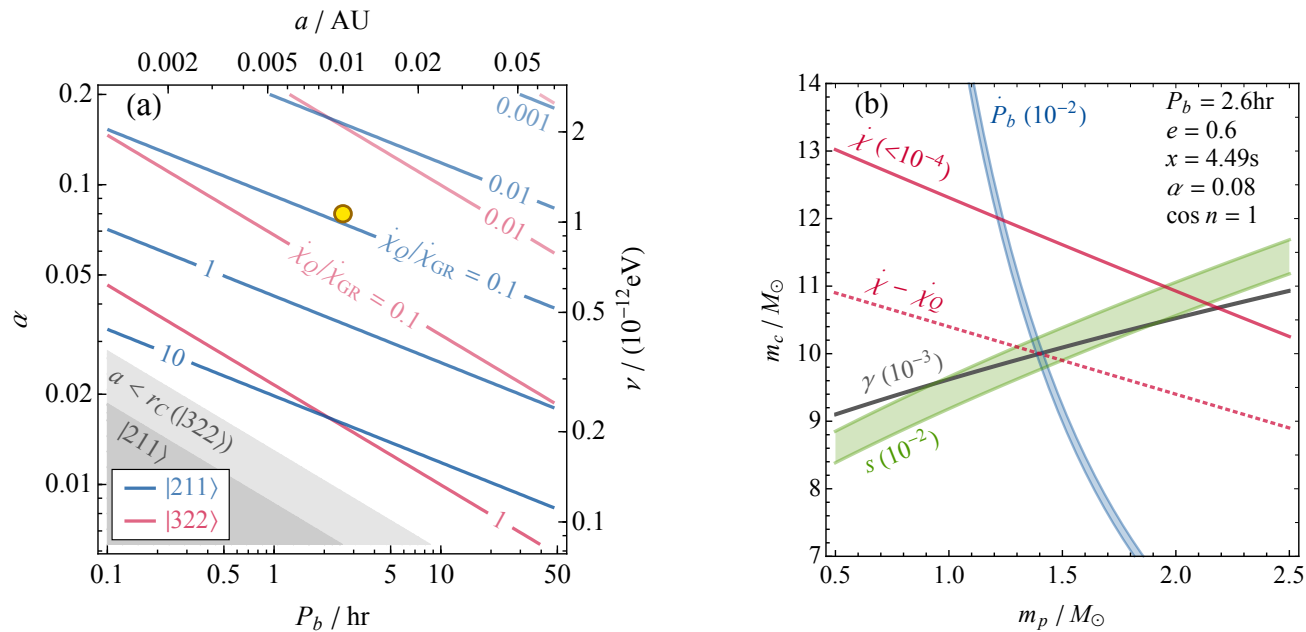
Strongly affected by mass quadrupole

- 2 PKs to fix the system (two masses)
- more PKs as consistency check



# Pulsar-GA binaries

GA induces dramatic misfit in apsidal precession



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# Summary

Astro BHs typically born with large spin  
natural factories and habitat for ultralight bosons

GA-borne binaries typically exhibit large orbital precessions  
much larger than other known contributions

Future GW and pulsar timing data could discover them

Also applicable to binaries with SMBHs